

9 49. Plots of cross sections and related quantities

Table 49.2D: Total hadronic cross section. For data descriptions the HPR_1R_2 model of highest COMPETE-rank¹ is used. As in RPP-2012 it is modified to save the universality of the rising part in new form that explicitly includes dependence of the s_M and H on the initial state mass parameters and the new scale parameter M .

$$\sigma^{a\bar{b}} = H \log^2(s/s_M^{ab}) + P^{ab} + R_1^{ab} (s/s_M^{ab})^{-\eta_1} \pm R_2^{ab} (s/s_M^{ab})^{-\eta_2}$$

$$\rho^{a\bar{b}} = \frac{1}{\sigma^{a\bar{b}}} \left[\pi H \log(s/s_M^{ab}) - R_1^{ab} (s/s_M^{ab})^{-\eta_1} \tan\left(\frac{\eta_1\pi}{2}\right) \pm R_2^{ab} (s/s_M^{ab})^{-\eta_2} \cot\left(\frac{\eta_2\pi}{2}\right) \right]$$

where upper signs in formulas are for particles and lower signs for antiparticles. Adjustable parameters are as follows:

$H = \pi \frac{(\hbar c)^2}{M^2}$ in mb, where notation H^2 is after Heisenberg(1952,1975);

P^{ab} in mb, are Pomeranchuk's(1958) constant terms;

R_i^{ab} in mb are the intensities of the effective secondary Regge pole contributions named after Regge-Gribov(1961);

$s, s_M^{ab} = (m_a + m_b + M)^2$ are in GeV² ;

$m_a, m_b, (m_{\gamma^*} = m_{\rho(\tau\tau_0)})$ are the masses of initial state particles, and M – the mass parameter defining the rate of universal rise of the cross sections are all in GeV.

Parameters M, η_1 and η_2 are universal for all collisions considered. In addition to total cross sections σ , the measured ratios of the real-to-imaginary parts of the forward scattering amplitudes $\rho = \text{Re}(T)/\text{Im}(T)$ are included in the fits by using s to u crossing symmetry and differential dispersion relations.

Exact factorization hypothesis was used for both $H \log^2(s/s_M^{ab})$ and P^{ab} to extend the universal rise of the total hadronic cross sections to the $\gamma(p, d) \rightarrow \text{hadrons}$ and $\gamma\gamma \rightarrow \text{hadrons}$ collisions. This results in one additional adjustable parameter δ with substitutions:

$$\pi H \log^2(s/s_M^{\gamma(p,d)}) + P^{\gamma(p,d)} \Rightarrow \delta \left[\pi(1, \lambda) H \log^2(s/s_M^{\gamma(p,d)}) + P^{p(p,d)} \right],$$

$$\pi H \log^2(s/s_M^{\gamma\gamma}) + P^{\gamma\gamma} \Rightarrow \delta^2 \left[\pi H \log^2(s/s_M^{\gamma\gamma}) + P^{pp} \right].$$

Three fits were made to the 2013-updated data above $\sqrt{s_{min}} = 5, 6, 7$ GeV for all collisions, listed in the ‘‘Beam/Target’’ column of the table . Fit results are shown only for fit above 7 GeV with the best fit quality $FQ = \chi^2/NDF$ and uniformity of data descriptions by groups (last column). The total number of adjusted parameters is **35**.

HPR₁R₂ at $\sqrt{s} \geq 7$ GeV	$M = 2076 \pm 0.016$ [GeV] $\eta_1 = 0.412 \pm 0.017$ $\delta = (3.112 \pm 0.027) \times 10^{-3}$		$H = 0.2838 \pm 0.0045$ [mb] $\eta_2 = 0.5626 \pm 0.0092$ $\lambda = 1.456 \pm 0.058$		FQ_{Int} = 0.86 FQ_{Ext} = 0.87
P [mb]	R_1 [mb]	R_2 [mb]	Beam/Target	Npt=836 npt	χ^2/npt by groups
33.73±0.33	13.67±0.33	7.77±0.18	$(\bar{p})p / p$	219	1.09
33.77±0.38	14.05±0.63	6.93±0.29	$(\bar{p})p / n$	48	0.39
33.20±3.90	-14.±47.	-15.±52.	Σ^- / n	8	0.41
18.08±0.29	10.44±0.32	1.977±0.078	π^\mp / p	137	0.91
15.84±0.20	5.12±0.28	3.538±0.095	K^\mp / p	85	0.76
15.73±0.22	4.81±0.40	1.86±0.13	K^\mp / n	48	0.56
	0.0132±0.0023		γ / p	34	0.56
	(-6.0±3.3)×10⁻⁵		γ / γ	31	0.68
	0.0256±0.0044		γ / d	3	0.31
64.79±0.75	27.06±0.85	15.46±0.37	$(\bar{p})p / d$	75	0.97
36.66±0.62	17.89±0.82	0.38±0.14	π^\mp / d	81	0.71
32.28±0.46	7.02±0.71	5.74±0.16	K^\mp / d	67	0.67

All fits are stable: in all cases the $FQ < 1.0$, Hessian matrices are positive definite. Quoted parameter uncertainties are obtained by direct MC-propagation of the uncertainties in each input experimental data point assuming independent normal distributions. Computer-readable data files are available at <http://pdg.lbl.gov/current/xsect/>. (Courtesy of the COMPAS group, IHEP, Protvino, September 2013)

¹Physical motivations for the model and related references see in mini-review. All plots give details for Figures 49.8 and 49.9 of the mini-review.

²For collisions with deuteron target $H_d = \lambda H$ where dimensionless parameter λ is introduced to test the universality of the Heisenberg rise for particle–nuclear and nuclear–nuclear collisions