

# WIMPs and Other Particles Searches for

OMITTED FROM SUMMARY TABLE

## WIMPS AND OTHER PARTICLE SEARCHES

Revised August 2013 by K. Hikasa (Tohoku University).

We collect here those searches which do not appear in any of the above search categories. These are listed in the following order:

1. Galactic WIMP (weakly-interacting massive particle) searches
2. Concentration of stable particles in matter
3. General new physics searches
4. Limits on jet-jet resonance in hadron collisions
5. Limits on neutral particle production at accelerators
6. Limits on charged particles in  $e^+e^-$  collisions
7. Limits on charged particles in hadron reactions
8. Limits on charged particles in cosmic rays
9. Searches for quantum black hole production

Note that searches appear in separate sections elsewhere for Higgs bosons (and technipions), other heavy bosons (including  $W_R$ ,  $W'$ ,  $Z'$ , leptoquarks, axigluons), axions (including pseudo-Goldstone bosons, Majorons, familons), heavy leptons, heavy neutrinos, free quarks, monopoles, supersymmetric particles, and compositeness. We include specific WIMP searches in the appropriate sections when they yield limits on hypothetical particles such as supersymmetric particles, axions, massive neutrinos, monopoles, *etc.*

We omit papers on CHAMP's, millicharged particles, and other exotic particles. We no longer list for limits on tachyons and centauros. See our 1994 edition for these limits.

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## GALACTIC WIMP SEARCHES

These limits are for weakly-interacting stable particles that may constitute the invisible mass in the galaxy. Unless otherwise noted, a local mass density of  $0.3 \text{ GeV}/\text{cm}^3$  is assumed; see each paper for velocity distribution assumptions. In the papers the limit is given as a function of the  $X^0$  mass. Here we list limits only for typical mass values of 20 GeV, 100 GeV, and 1 TeV. Specific limits on supersymmetric dark matter particles may be found in the Supersymmetry section.

### ———— Limits for Spin-Independent Cross Section ———— ———— of Dark Matter Particle ( $X^0$ ) on Nucleon ————

Isoscalar coupling is assumed to extract the limits from those on  $X^0$ -nuclei cross section.

#### For $m_{X^0} = 20 \text{ GeV}$

For limits from  $X^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$<7.3 \times 10^{-7}$	90	AGNES	16	DS50 Ar
$<1 \times 10^{-5}$	90	AGNES	15	DSID Ar
$<1.5 \times 10^{-6}$	90	<sup>1</sup> AGNESE	15A	CDM2 Ge
$<1.5 \times 10^{-7}$	90	<sup>2</sup> AGNESE	15B	CDM2 Ge
$<2 \times 10^{-6}$	90	<sup>3</sup> AMOLE	15	PICO $\text{C}_3\text{F}_8$
$<1.2 \times 10^{-5}$	90	CHOI	15	SKAM H, solar $\nu$ ( $b\bar{b}$ )
$<1.19 \times 10^{-6}$	90	CHOI	15	SKAM H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<2 \times 10^{-8}$	90	<sup>4</sup> XIAO	15	PANX Xe
$<2.0 \times 10^{-7}$	90	<sup>5</sup> AGNESE	14	SCDM Ge
$<3.7 \times 10^{-5}$	90	<sup>6</sup> AGNESE	14A	SCDM Ge
$<1 \times 10^{-9}$	90	<sup>7</sup> AKERIB	14	LUX Xe
$<2 \times 10^{-6}$	90	<sup>8</sup> ANGLOHER	14	CRES $\text{CaWO}_4$
$<5 \times 10^{-6}$	90	FELIZARDO	14	SMPL $\text{C}_2\text{ClF}_5$
$<8 \times 10^{-6}$	90	<sup>9</sup> LEE	14A	KIMS CsI
$<2 \times 10^{-4}$	90	<sup>10</sup> LIU	14A	CDEX Ge
$<1 \times 10^{-5}$	90	<sup>11</sup> YUE	14	CDEX Ge
$<1.08 \times 10^{-4}$	90	<sup>12</sup> AARTSEN	13	ICCB H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<1.5 \times 10^{-5}$	90	<sup>13</sup> ABE	13B	XMAS Xe
$<3.1 \times 10^{-6}$	90	<sup>14</sup> AGNESE	13	CDM2 Si
$<3.4 \times 10^{-6}$	90	<sup>15</sup> AGNESE	13A	CDM2 Si
$<2.2 \times 10^{-6}$	90	<sup>16</sup> AGNESE	13A	CDM2 Si
$<5 \times 10^{-5}$	90	<sup>17</sup> LI	13B	TEXO Ge
		<sup>18</sup> ZHAO	13	CDEX Ge
$<1.2 \times 10^{-7}$	90	AKIMOV	12	ZEP3 Xe
		<sup>19</sup> ANGLOHER	12	CRES $\text{CaWO}_4$
$<8 \times 10^{-6}$	90	<sup>20</sup> ANGLOHER	12	CRES $\text{CaWO}_4$
$<7 \times 10^{-9}$	90	<sup>21</sup> APRILE	12	X100 Xe
		<sup>22</sup> ARCHAMBAU..	12	PICA F ( $\text{C}_4\text{F}_{10}$ )
$<7 \times 10^{-7}$	90	<sup>23</sup> ARMENGAUD	12	EDE2 Ge
		<sup>24</sup> BARRETO	12	DMIC CCD

$<2 \times 10^{-6}$	90	BEHNKE	12	COUP	CF <sub>3</sub> I
$<7 \times 10^{-6}$		25 FELIZARDO	12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
$<1.5 \times 10^{-6}$	90	KIM	12	KIMS	CsI
$<5 \times 10^{-5}$	90	26 AALSETH	11	CGNT	Ge
		27 AALSETH	11A	CGNT	Ge
$<5 \times 10^{-7}$	90	28 AHMED	11	CDM2	Ge, inelastic
$<2.7 \times 10^{-7}$	90	29 AHMED	11A	RVUE	Ge
		30 AHMED	11B	CDM2	Ge, low threshold
$<3 \times 10^{-6}$	90	31 ANGLE	11	XE10	Xe
$<7 \times 10^{-8}$	90	32 APRILE	11	X100	Xe
		33 APRILE	11A	X100	Xe, inelastic
$<2 \times 10^{-8}$	90	21 APRILE	11B	X100	Xe
		34 HORN	11	ZEP3	Xe
$<2 \times 10^{-7}$	90	AHMED	10	CDM2	Ge
$<1 \times 10^{-5}$	90	35 AKERIB	10	CDM2	Si, Ge, low threshold
$<1 \times 10^{-7}$	90	APRILE	10	X100	Xe
$<2 \times 10^{-6}$	90	ARMENGAUD	10	EDE2	Ge
$<4 \times 10^{-5}$	90	FELIZARDO	10	SMPL	C <sub>2</sub> ClF <sub>3</sub>
$<1.5 \times 10^{-7}$	90	36 AHMED	09	CDM2	Ge
$<2 \times 10^{-4}$	90	37 LIN	09	TEXO	Ge
		38 AALSETH	08	CGNT	Ge

<sup>1</sup> AGNESE 15A reanalyse AHMED 11B low threshold data. See their Fig. 12 (left) for improved limits extending down to 5 GeV.

<sup>2</sup> AGNESE 15B reanalyse AHMED 10 data.

<sup>3</sup> See their Fig. 7 for limits extending down to 4 GeV.

<sup>4</sup> See their Fig. 13 for limits extending down to 5 GeV.

<sup>5</sup> This limit value is provided by the authors. See their Fig. 4 for limits extending down to  $m_{\chi^0} = 3.5$  GeV.

<sup>6</sup> This limit value is provided by the authors. AGNESE 14A result is from CDMSlite mode operation with enhanced sensitivity to low mass  $m_{\chi^0}$ . See their Fig. 3 for limits extending down to  $m_{\chi^0} = 3.5$  GeV (see also Fig. 4 in AGNESE 14).

<sup>7</sup> See their Fig. 5 for limits extending down to  $m_{\chi^0} = 5.5$  GeV.

<sup>8</sup> See their Fig. 5 for limits extending down to  $m_{\chi^0} = 1$  GeV.

<sup>9</sup> See their Fig. 5 for limits extending down to  $m_{\chi^0} = 5$  GeV.

<sup>10</sup> LIU 14A result is based on prototype CDEX-0 detector. See their Fig. 13 for limits extending down to  $m_{\chi^0} = 2$  GeV.

<sup>11</sup> See their Fig. 4 for limits extending down to  $m_{\chi^0} = 4.5$  GeV.

<sup>12</sup> AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $\chi^0$  trapped by the sun in data taken between June 2010 and May 2011.

<sup>13</sup> See their Fig. 8 for limits extending down to  $m_{\chi^0} = 7$  GeV.

<sup>14</sup> This limit value is provided by the authors. AGNESE 13 use data taken between Oct. 2006 and July 2007. See their Fig. 4 for limits extending down to  $m_{\chi^0} = 7$  GeV.

<sup>15</sup> This limit value is provided by the authors. AGNESE 13A use data taken between July 2007 and Sep. 2008. Three candidate events are seen. Assuming these events are real, the best fit parameters are  $m_{\chi^0} = 8.6$  GeV and  $\sigma = 1.9 \times 10^{-5}$  pb.

<sup>16</sup> This limit value is provided by the authors. Limit from combined data of AGNESE 13 and AGNESE 13A. See their Fig. 4 for limits extending down to  $m_{\chi^0} = 5.5$  GeV.

<sup>17</sup> See their Fig. 4 for limits extending down to  $m_{\chi^0} = 4$  GeV.

- <sup>18</sup> See their Fig. 5 for limits for  $m_{\chi^0} = 4\text{--}12$  GeV.
- <sup>19</sup> ANGLOHER 12 observe excess events above the expected background which are consistent with  $\chi^0$  with mass  $\sim 25$  GeV (or 12 GeV) and spin-independent  $\chi^0$ -nucleon cross section of  $2 \times 10^{-6}$  pb (or  $4 \times 10^{-5}$  pb).
- <sup>20</sup> Reanalysis of ANGLOHER 09 data with all three nuclides. See also BROWN 12.
- <sup>21</sup> See also APRILE 14A.
- <sup>22</sup> See their Fig. 7 for cross section limits for  $m_{\chi^0}$  between 4 and 12 GeV.
- <sup>23</sup> See their Fig. 4 for limits extending down to  $m_{\chi^0} = 7$  GeV.
- <sup>24</sup> See their Fig. 13 for cross section limits for  $m_{\chi^0}$  between 1.2 and 10 GeV.
- <sup>25</sup> See also DAHL 12 for a criticism.
- <sup>26</sup> See their Fig. 4 for limits extending to  $m_{\chi^0} = 3.5$  GeV.
- <sup>27</sup> AALSETH 11A find indications of annual modulation of the data, the energy spectrum being compatible with  $\chi^0$  mass around 8 GeV. See also AALSETH 13.
- <sup>28</sup> AHMED 11 search for  $\chi^0$  inelastic scattering. See their Fig. 8–10 for limits. The inelastic cross section reduces to the elastic cross section at the limit of zero mass splitting (Fig. 8, left).
- <sup>29</sup> AHMED 11A combine CDMS II and EDELWEISS data.
- <sup>30</sup> AHMED 11B give limits on spin-independent  $\chi^0$ -nucleon cross section for  $m_{\chi^0} = 4\text{--}12$  GeV in the range  $10^{-3}\text{--}10^{-5}$  pb. See their Fig. 3.
- <sup>31</sup> See their Fig. 3 for limits down to  $m_{\chi^0} = 4$  GeV.
- <sup>32</sup> APRILE 11 reanalyze APRILE 10 data.
- <sup>33</sup> APRILE 11A search for  $\chi^0$  inelastic scattering. See their Fig. 2 and 3 for limits. See also APRILE 14A.
- <sup>34</sup> HORN 11 perform detector calibration by neutrons. Earlier results are only marginally affected.
- <sup>35</sup> See their Fig. 10 and 12 for limits extending to  $\chi^0$  mass of 1 GeV.
- <sup>36</sup> Superseded by AHMED 10.
- <sup>37</sup> See their Fig. 6(a) for cross section limits for  $m_{\chi^0}$  extending down to 2 GeV.
- <sup>38</sup> See their Fig. 2 for cross section limits for  $m_{\chi^0}$  between 4 and 10 GeV.

### For $m_{\chi^0} = 100$ GeV

For limits from  $\chi^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$<2.0 \times 10^{-8}$	90	AGNES	16	DS50 Ar
$<6 \times 10^{-8}$	90	AGNES	15	DSID Ar
$<4 \times 10^{-8}$	90	<sup>1</sup> AGNESE	15B	CDM2 Ge
$<7.13 \times 10^{-6}$	90	CHOI	15	SKAM H, solar $\nu$ ( $b\bar{b}$ )
$<6.26 \times 10^{-7}$	90	CHOI	15	SKAM H, solar $\nu$ ( $W^+ W^-$ )
$<2.76 \times 10^{-7}$	90	CHOI	15	SKAM H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<1.5 \times 10^{-8}$	90	XIAO	15	PANX Xe
$<1 \times 10^{-9}$	90	AKERIB	14	LUX Xe
$<4.0 \times 10^{-6}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $W^+ W^-$ )
$<1.0 \times 10^{-4}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $b\bar{b}$ )
$<1.6 \times 10^{-6}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $\tau^+ \tau^-$ )

$<5 \times 10^{-6}$	90	FELIZARDO	14	SMPL	$C_2ClF_5$
$<6.01 \times 10^{-7}$	90	<sup>3</sup> AARTSEN	13	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
$<3.30 \times 10^{-5}$	90	<sup>3</sup> AARTSEN	13	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
$<1.9 \times 10^{-6}$	90	<sup>4</sup> ADRIAN-MAR..13	13	ANTR	H, solar $\nu$ ( $W^+ W^-$ )
$<1.2 \times 10^{-4}$	90	<sup>4</sup> ADRIAN-MAR..13	13	ANTR	H, solar $\nu$ ( $b\bar{b}$ )
$<7.6 \times 10^{-7}$	90	<sup>4</sup> ADRIAN-MAR..13	13	ANTR	H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<2 \times 10^{-6}$	90	<sup>5</sup> AGNESE	13	CDM2	Si
$<1.6 \times 10^{-6}$	90	<sup>6</sup> BOLIEV	13	BAKS	H, solar $\nu$ ( $W^+ W^-$ )
$<1.9 \times 10^{-5}$	90	<sup>6</sup> BOLIEV	13	BAKS	H, solar $\nu$ ( $b\bar{b}$ )
$<7.1 \times 10^{-7}$	90	<sup>6</sup> BOLIEV	13	BAKS	H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<1.67 \times 10^{-6}$	90	<sup>7</sup> ABBASI	12	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
$<1.07 \times 10^{-4}$	90	<sup>7</sup> ABBASI	12	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
$<4 \times 10^{-8}$	90	AKIMOV	12	ZEP3	Xe
$<1.4 \times 10^{-6}$	90	<sup>8</sup> ANGLOHER	12	CRES	$CaWO_4$
$<3 \times 10^{-9}$	90	<sup>9</sup> APRILE	12	X100	Xe
$<3 \times 10^{-7}$	90	BEHNKE	12	COUP	$CF_3I$
$<7 \times 10^{-6}$		FELIZARDO	12	SMPL	$C_2ClF_5$
$<2.5 \times 10^{-7}$	90	<sup>10</sup> KIM	12	KIMS	CsI
$<2 \times 10^{-4}$	90	AALSETH	11	CGNT	Ge
		<sup>11</sup> AHMED	11	CDM2	Ge, inelastic
$<3.3 \times 10^{-8}$	90	<sup>12</sup> AHMED	11A	RVUE	Ge
		<sup>13</sup> AJELLO	11	FLAT	
$<3 \times 10^{-8}$	90	<sup>14</sup> APRILE	11	X100	Xe
		<sup>15</sup> APRILE	11A	X100	Xe, inelastic
$<1 \times 10^{-8}$	90	<sup>9</sup> APRILE	11B	X100	Xe
$<5 \times 10^{-8}$	90	<sup>16</sup> ARMENGAUD	11	EDE2	Ge
		<sup>17</sup> HORN	11	ZEP3	Xe
$<4 \times 10^{-8}$	90	AHMED	10	CDM2	Ge
$<9 \times 10^{-6}$	90	AKERIB	10	CDM2	Si, Ge, low threshold
		<sup>18</sup> AKIMOV	10	ZEP3	Xe, inelastic
$<5 \times 10^{-8}$	90	APRILE	10	X100	Xe
$<1 \times 10^{-7}$	90	ARMENGAUD	10	EDE2	Ge
$<3 \times 10^{-5}$	90	FELIZARDO	10	SMPL	$C_2ClF_3$
$<5 \times 10^{-8}$	90	<sup>19</sup> AHMED	09	CDM2	Ge
		<sup>20</sup> ANGLE	09	XE10	Xe, inelastic
$<3 \times 10^{-4}$	90	LIN	09	TEXO	Ge
		<sup>21</sup> GIULIANI	05	RVUE	

<sup>1</sup> AGNESE 15B reanalyse AHMED 10 data.

<sup>2</sup> AVRORIN 14 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun in data taken between 1998 and 2003. See their Table 1 for limits assuming annihilation into neutrino pairs.

<sup>3</sup> AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between June 2010 and May 2011.

<sup>4</sup> ADRIAN-MARTINEZ 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between Jan. 2007 and Dec. 2008.

<sup>5</sup> AGNESE 13 use data taken between Oct. 2006 and July 2007.

<sup>6</sup> BOLIEV 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken from 1978 to 2009. See also SUVOROVA 13 for an older analysis of the same data.

- <sup>7</sup> ABBASI 12 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.
- <sup>8</sup> Reanalysis of ANGLOHER 09 data with all three nuclides. See also BROWN 12.
- <sup>9</sup> See also APRILE 14A.
- <sup>10</sup> See their Fig. 6 for a limit on inelastically scattering  $X^0$  for  $m_{X^0} = 70$  GeV.
- <sup>11</sup> AHMED 11 search for  $X^0$  inelastic scattering. See their Fig. 8–10 for limits.
- <sup>12</sup> AHMED 11A combine CDMS and EDELWEISS data.
- <sup>13</sup> AJELLO 11 search for  $e^\pm$  flux from  $X^0$  annihilations in the Sun. Models in which  $X^0$  annihilates into an intermediate long-lived weakly interacting particles or  $X^0$  scatters inelastically are constrained. See their Fig. 6–8 for limits.
- <sup>14</sup> APRILE 11 reanalyze APRILE 10 data.
- <sup>15</sup> APRILE 11A search for  $X^0$  inelastic scattering. See their Fig. 2 and 3 for limits. See also APRILE 14A.
- <sup>16</sup> Supersedes ARMENGAUD 10. A limit on inelastic cross section is also given.
- <sup>17</sup> HORN 11 perform detector calibration by neutrons. Earlier results are only marginally affected.
- <sup>18</sup> AKIMOV 10 give cross section limits for inelastically scattering dark matter. See their Fig. 4.
- <sup>19</sup> Superseded by AHMED 10.
- <sup>20</sup> ANGLE 09 search for  $X^0$  inelastic scattering. See their Fig. 4 for limits.
- <sup>21</sup> GIULIANI 05 analyzes the spin-independent  $X^0$ -nucleon cross section limits with both isoscalar and isovector couplings. See their Fig. 3 and 4 for limits on the couplings.

### For $m_{X^0} = 1$ TeV

For limits from  $X^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$<8.6 \times 10^{-8}$	90	AGNES	16	DS50 Ar
$<2 \times 10^{-7}$	90	AGNES	15	DSID Ar
$<2 \times 10^{-7}$	90	<sup>1</sup> AGNESE	15B	CDM2 Ge
$<1 \times 10^{-8}$	90	AKERIB	14	LUX Xe
$<2.2 \times 10^{-6}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $W^+ W^-$ )
$<5.5 \times 10^{-5}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $b\bar{b}$ )
$<6.8 \times 10^{-7}$	90	<sup>2</sup> AVRORIN	14	BAIK H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<3.46 \times 10^{-7}$	90	<sup>3</sup> AARTSEN	13	ICCB H, solar $\nu$ ( $W^+ W^-$ )
$<7.75 \times 10^{-6}$	90	<sup>3</sup> AARTSEN	13	ICCB H, solar $\nu$ ( $b\bar{b}$ )
$<6.9 \times 10^{-7}$	90	<sup>4</sup> ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $W^+ W^-$ )
$<1.5 \times 10^{-5}$	90	<sup>4</sup> ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $b\bar{b}$ )
$<1.8 \times 10^{-7}$	90	<sup>4</sup> ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<4.3 \times 10^{-6}$	90	<sup>5</sup> BOLIEV	13	BAKS H, solar $\nu$ ( $W^+ W^-$ )
$<3.4 \times 10^{-5}$	90	<sup>5</sup> BOLIEV	13	BAKS H, solar $\nu$ ( $b\bar{b}$ )
$<1.2 \times 10^{-6}$	90	<sup>5</sup> BOLIEV	13	BAKS H, solar $\nu$ ( $\tau^+ \tau^-$ )
$<2.12 \times 10^{-7}$	90	<sup>6</sup> ABBASI	12	ICCB H, solar $\nu$ ( $W^+ W^-$ )
$<6.56 \times 10^{-6}$	90	<sup>6</sup> ABBASI	12	ICCB H, solar $\nu$ ( $b\bar{b}$ )
$<4 \times 10^{-7}$	90	AKIMOV	12	ZEP3 Xe
$<1.1 \times 10^{-5}$	90	<sup>7</sup> ANGLOHER	12	CRES $\text{CaWO}_4$
$<2 \times 10^{-8}$	90	<sup>8</sup> APRILE	12	X100 Xe

$<2 \times 10^{-6}$	90	BEHNKE	12	COUP	CF <sub>3</sub> I
$<4 \times 10^{-6}$		FELIZARDO	12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
$<1.5 \times 10^{-6}$	90	KIM	12	KIMS	CsI
		<sup>9</sup> AHMED	11	CDM2	Ge, inelastic
$<1.5 \times 10^{-7}$	90	<sup>10</sup> AHMED	11A	RVUE	Ge
$<2 \times 10^{-7}$	90	<sup>11</sup> APRILE	11	X100	Xe
$<8 \times 10^{-8}$	90	<sup>8</sup> APRILE	11B	X100	Xe
$<2 \times 10^{-7}$	90	<sup>12</sup> ARMENGAUD	11	EDE2	Ge
		<sup>13</sup> HORN	11	ZEP3	Xe
$<2 \times 10^{-7}$	90	AHMED	10	CDM2	Ge
$<4 \times 10^{-7}$	90	APRILE	10	X100	Xe
$<6 \times 10^{-7}$	90	ARMENGAUD	10	EDE2	Ge
$<3.5 \times 10^{-7}$	90	<sup>14</sup> AHMED	09	CDM2	Ge

<sup>1</sup> AGNESE 15B reanalyse AHMED 10 data.

<sup>2</sup> AVRORIN 14 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun in data taken between 1998 and 2003. See their Table 1 for limits assuming annihilation into neutrino pairs.

<sup>3</sup> AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between June 2010 and May 2011.

<sup>4</sup> ADRIAN-MARTINEZ 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between Jan. 2007 and Dec. 2008.

<sup>5</sup> BOLIEV 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken from 1978 to 2009. See also SUVOROVA 13 for an older analysis of the same data.

<sup>6</sup> ABBASI 12 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.

<sup>7</sup> Reanalysis of ANGLOHER 09 data with all three nuclides. See also BROWN 12.

<sup>8</sup> See also APRILE 14A.

<sup>9</sup> AHMED 11 search for  $X^0$  inelastic scattering. See their Fig. 8–10 for limits.

<sup>10</sup> AHMED 11A combine CDMS and EDELWEISS data.

<sup>11</sup> APRILE 11 reanalyze APRILE 10 data.

<sup>12</sup> Supersedes ARMENGAUD 10. A limit on inelastic cross section is also given.

<sup>13</sup> HORN 11 perform detector calibration by neutrons. Earlier results are only marginally affected.

<sup>14</sup> Superseded by AHMED 10.

### Limits for Spin-Dependent Cross Section of Dark Matter Particle ( $X^0$ ) on Proton

#### For $m_{X^0} = 20$ GeV

For limits from  $X^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

VALUE (pb)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$< 1.2 \times 10^{-3}$	90	AMOLE	15	PICO C <sub>3</sub> F <sub>8</sub>
$< 1.43 \times 10^{-3}$	90	CHOI	15	SKAM H, solar $\nu$ ( $b\bar{b}$ )
$< 1.42 \times 10^{-4}$	90	CHOI	15	SKAM H, solar $\nu$ ( $\tau^+ \tau^-$ )
$< 5 \times 10^{-3}$	90	FELIZARDO	14	SMPL C <sub>2</sub> ClF <sub>5</sub>
$< 1.29 \times 10^{-2}$	90	<sup>1</sup> AARTSEN	13	ICCB H, solar $\nu$ ( $\tau^+ \tau^-$ )
$< 3.17 \times 10^{-2}$	90	<sup>2</sup> APRILE	13	X100 Xe
$< 3 \times 10^{-2}$	90	ARCHAMBAU..12	PICA	F (C <sub>4</sub> F <sub>10</sub> )

< 6 × 10 <sup>-2</sup>	90	BEHNKE	12	COUP	CF <sub>3</sub> I
< 20	90	DAW	12	DRFT	F (CF <sub>4</sub> )
< 7 × 10 <sup>-3</sup>		FELIZARDO	12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
< 0.15	90	KIM	12	KIMS	CsI
< 1 × 10 <sup>5</sup>	90	<sup>3</sup> AHLEN	11	DMTP	F (CF <sub>4</sub> )
< 0.1	90	<sup>3</sup> BEHNKE	11	COUP	CF <sub>3</sub> I
< 1.5 × 10 <sup>-2</sup>	90	<sup>4</sup> TANAKA	11	SKAM	H, solar ν ( <i>b</i> $\bar{b}$ )
< 0.2	90	ARCHAMBAU..09		PICA	F
< 4	90	LEBEDENKO	09A	ZEP3	Xe
< 0.6	90	ANGLE	08A	XE10	Xe
<100	90	ALNER	07	ZEP2	Xe
< 1	90	LEE	07A	KIMS	CsI
< 20	90	<sup>5</sup> AKERIB	06	CDMS	<sup>73</sup> Ge, <sup>29</sup> Si
< 2	90	SHIMIZU	06A	CNTR	F (CaF <sub>2</sub> )
< 0.5	90	ALNER	05	NAIA	NaI
< 1.5	90	BARNABE-HE..05		PICA	F (C <sub>4</sub> F <sub>10</sub> )
< 1.5	90	GIRARD	05	SMPL	F (C <sub>2</sub> ClF <sub>5</sub> )
< 35	90	MIUCHI	03	BOLO	LiF
< 30	90	TAKEDA	03	BOLO	NaF

<sup>1</sup>AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between June 2010 and May 2011.

<sup>2</sup>The value has been provided by the authors. APRILE 13 note that the proton limits on Xe are highly sensitive to the theoretical model used. See also APRILE 14A.

<sup>3</sup>Use a direction-sensitive detector.

<sup>4</sup>TANAKA 11 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.

<sup>5</sup>See also AKERIB 05.

### For $m_{X^0} = 100$ GeV

For limits from  $X^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

VALUE (pb)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 1.5 × 10 <sup>-3</sup>	90	AMOLE	15	PICO C <sub>3</sub> F <sub>8</sub>
< 3.19 × 10 <sup>-3</sup>	90	CHOI	15	SKAM H, solar ν ( <i>b</i> $\bar{b}$ )
< 2.80 × 10 <sup>-4</sup>	90	CHOI	15	SKAM H, solar ν ( $W^+ W^-$ )
< 1.24 × 10 <sup>-4</sup>	90	CHOI	15	SKAM H, solar ν ( $\tau^+ \tau^-$ )
< 8 × 10 <sup>2</sup>	90	<sup>1</sup> NAKAMURA	15	NAGE CF <sub>4</sub>
< 1.7 × 10 <sup>-3</sup>	90	<sup>2</sup> AVRORIN	14	BAIK H, solar ν ( $W^+ W^-$ )
< 4.5 × 10 <sup>-2</sup>	90	<sup>2</sup> AVRORIN	14	BAIK H, solar ν ( <i>b</i> $\bar{b}$ )
< 7.1 × 10 <sup>-4</sup>	90	<sup>2</sup> AVRORIN	14	BAIK H, solar ν ( $\tau^+ \tau^-$ )
< 6 × 10 <sup>-3</sup>	90	FELIZARDO	14	SMPL C <sub>2</sub> ClF <sub>5</sub>
< 2.68 × 10 <sup>-4</sup>	90	<sup>3</sup> AARTSEN	13	ICCB H, solar ν ( $W^+ W^-$ )
< 1.47 × 10 <sup>-2</sup>	90	<sup>3</sup> AARTSEN	13	ICCB H, solar ν ( <i>b</i> $\bar{b}$ )
< 8.5 × 10 <sup>-4</sup>	90	<sup>4</sup> ADRIAN-MAR..13		ANTR H, solar ν ( $W^+ W^-$ )
< 5.5 × 10 <sup>-2</sup>	90	<sup>4</sup> ADRIAN-MAR..13		ANTR H, solar ν ( <i>b</i> $\bar{b}$ )
< 3.4 × 10 <sup>-4</sup>	90	<sup>4</sup> ADRIAN-MAR..13		ANTR H, solar ν ( $\tau^+ \tau^-$ )
< 1.00 × 10 <sup>-2</sup>	90	<sup>5</sup> APRILE	13	X100 Xe
< 7.1 × 10 <sup>-4</sup>	90	<sup>6</sup> BOLIEV	13	BAKS H, solar ν ( $W^+ W^-$ )
< 8.4 × 10 <sup>-3</sup>	90	<sup>6</sup> BOLIEV	13	BAKS H, solar ν ( <i>b</i> $\bar{b}$ )



< $3.1 \times 10^{-4}$	90	<sup>6</sup> BOLIEV	13	BAKS	H, solar $\nu$ ( $\tau^+ \tau^-$ )
< $7.07 \times 10^{-4}$	90	<sup>7</sup> ABBASI	12	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
< $4.53 \times 10^{-2}$	90	<sup>7</sup> ABBASI	12	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
< $7 \times 10^{-2}$	90	ARCHAMBAU..	12	PICA	F ( $C_4F_{10}$ )
< $1 \times 10^{-2}$	90	BEHNKE	12	COUP	CF <sub>3</sub> I
< 1.8	90	DAW	12	DRFT	F (CF <sub>4</sub> )
< $9 \times 10^{-3}$		FELIZARDO	12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
< $2 \times 10^{-2}$	90	KIM	12	KIMS	CsI
< $2 \times 10^3$	90	<sup>1</sup> AHLEN	11	DMTP	F (CF <sub>4</sub> )
< $7 \times 10^{-2}$	90	BEHNKE	11	COUP	CF <sub>3</sub> I
< $2.7 \times 10^{-4}$	90	<sup>8</sup> TANAKA	11	SKAM	H, solar $\nu$ ( $W^+ W^-$ )
< $4.5 \times 10^{-3}$	90	<sup>8</sup> TANAKA	11	SKAM	H, solar $\nu$ ( $b\bar{b}$ )
		<sup>9</sup> FELIZARDO	10	SMPL	C <sub>2</sub> ClF <sub>3</sub>
< $6 \times 10^3$	90	<sup>1</sup> MIUCHI	10	NAGE	CF <sub>4</sub>
< 0.4	90	ARCHAMBAU..	09	PICA	F
< 0.8	90	LEBEDENKO	09A	ZEP3	Xe
< 1.0	90	ANGLE	08A	XE10	Xe
< 15	90	ALNER	07	ZEP2	Xe
< 0.2	90	LEE	07A	KIMS	CsI
< $1 \times 10^4$	90	<sup>1</sup> MIUCHI	07	NAGE	F (CF <sub>4</sub> )
< 5	90	<sup>10</sup> AKERIB	06	CDMS	<sup>73</sup> Ge, <sup>29</sup> Si
< 2	90	SHIMIZU	06A	CNTR	F (CaF <sub>2</sub> )
< 0.3	90	ALNER	05	NAIA	NaI
< 2	90	BARNABE-HE..	05	PICA	F ( $C_4F_{10}$ )
<100	90	BENOIT	05	EDEL	<sup>73</sup> Ge
< 1.5	90	GIRARD	05	SMPL	F (C <sub>2</sub> ClF <sub>5</sub> )
< 0.7		<sup>11</sup> GIULIANI	05A	RVUE	
		<sup>12</sup> GIULIANI	04	RVUE	
		<sup>13</sup> GIULIANI	04A	RVUE	
< 35	90	MIUCHI	03	BOLO	LiF
< 40	90	TAKEDA	03	BOLO	NaF

<sup>1</sup> Use a direction-sensitive detector.

<sup>2</sup> AVRORIN 14 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun in data taken between 1998 and 2003. See their Table 1 for limits assuming annihilation into neutrino pairs.

<sup>3</sup> AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between June 2010 and May 2011.

<sup>4</sup> ADRIAN-MARTINEZ 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between Jan. 2007 and Dec. 2008.

<sup>5</sup> The value has been provided by the authors. APRILE 13 note that the proton limits on Xe are highly sensitive to the theoretical model used. See also APRILE 14A.

<sup>6</sup> BOLIEV 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken from 1978 to 2009. See also SUVOROVA 13 for an older analysis of the same data.

<sup>7</sup> ABBASI 12 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.

<sup>8</sup> TANAKA 11 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.

<sup>9</sup> See their Fig. 3 for limits on spin-dependent proton couplings for  $X^0$  mass of 50 GeV.

<sup>10</sup> See also AKERIB 05.

<sup>11</sup> GIULIANI 05A analyze available data and give combined limits.

<sup>12</sup>GIULIANI 04 reanalyze COLLAR 00 data and give limits for spin-dependent  $X^0$ -proton coupling.

<sup>13</sup>GIULIANI 04A give limits for spin-dependent  $X^0$ -proton couplings from existing data.

### For $m_{X^0} = 1$ TeV

For limits from  $X^0$  annihilation in the Sun, the assumed annihilation final state is shown in parenthesis in the comment.

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 1 × 10 <sup>-2</sup>	90	AMOLE 15	PICO	C <sub>3</sub> F <sub>8</sub>
< 1.5 × 10 <sup>3</sup>	90	NAKAMURA 15	NAGE	CF <sub>4</sub>
< 2.7 × 10 <sup>-3</sup>	90	1 AVRORIN 14	BAIK	H, solar $\nu$ ( $W^+ W^-$ )
< 6.9 × 10 <sup>-2</sup>	90	1 AVRORIN 14	BAIK	H, solar $\nu$ ( $b\bar{b}$ )
< 8.4 × 10 <sup>-4</sup>	90	1 AVRORIN 14	BAIK	H, solar $\nu$ ( $\tau^+ \tau^-$ )
< 4.48 × 10 <sup>-4</sup>	90	2 AARTSEN 13	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
< 1.00 × 10 <sup>-2</sup>	90	2 AARTSEN 13	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
< 8.9 × 10 <sup>-4</sup>	90	3 ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $W^+ W^-$ )
< 2.0 × 10 <sup>-2</sup>	90	3 ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $b\bar{b}$ )
< 2.3 × 10 <sup>-4</sup>	90	3 ADRIAN-MAR..13	ANTR	H, solar $\nu$ ( $\tau^+ \tau^-$ )
< 7.57 × 10 <sup>-2</sup>	90	4 APRILE 13	X100	Xe
< 5.4 × 10 <sup>-3</sup>	90	5 BOLIEV 13	BAKS	H, solar $\nu$ ( $W^+ W^-$ )
< 4.2 × 10 <sup>-2</sup>	90	5 BOLIEV 13	BAKS	H, solar $\nu$ ( $b\bar{b}$ )
< 1.5 × 10 <sup>-3</sup>	90	5 BOLIEV 13	BAKS	H, solar $\nu$ ( $\tau^+ \tau^-$ )
< 2.50 × 10 <sup>-4</sup>	90	6 ABBASI 12	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
< 7.86 × 10 <sup>-3</sup>	90	6 ABBASI 12	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
< 8 × 10 <sup>-2</sup>	90	BEHNKE 12	COUP	CF <sub>3</sub> I
< 8	90	DAW 12	DRFT	F (CF <sub>4</sub> )
< 6 × 10 <sup>-2</sup>		FELIZARDO 12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
< 8 × 10 <sup>-2</sup>	90	KIM 12	KIMS	CsI
< 8 × 10 <sup>3</sup>	90	7 AHLEN 11	DMTP	F (CF <sub>4</sub> )
< 0.4	90	BEHNKE 11	COUP	CF <sub>3</sub> I
< 2 × 10 <sup>-3</sup>	90	8 TANAKA 11	SKAM	H, solar $\nu$ ( $b\bar{b}$ )
< 2 × 10 <sup>-2</sup>	90	8 TANAKA 11	SKAM	H, solar $\nu$ ( $W^+ W^-$ )
< 1 × 10 <sup>-3</sup>	90	9 ABBASI 10	ICCB	KK dark matter
< 2 × 10 <sup>4</sup>	90	7 MIUCHI 10	NAGE	CF <sub>4</sub>
< 8.7 × 10 <sup>-4</sup>	90	ABBASI 09B	ICCB	H, solar $\nu$ ( $W^+ W^-$ )
< 2.2 × 10 <sup>-2</sup>	90	ABBASI 09B	ICCB	H, solar $\nu$ ( $b\bar{b}$ )
< 3	90	ARCHAMBAU..09	PICA	F
< 6	90	LEBEDENKO 09A	ZEP3	Xe
< 9	90	ANGLE 08A	XE10	Xe
<100	90	ALNER 07	ZEP2	Xe
< 0.8	90	LEE 07A	KIMS	CsI
< 4 × 10 <sup>4</sup>	90	7 MIUCHI 07	NAGE	F (CF <sub>4</sub> )
< 30	90	10 AKERIB 06	CDMS	<sup>73</sup> Ge, <sup>29</sup> Si
< 1.5	90	ALNER 05	NAIA	NaI
< 15	90	BARNABE-HE..05	PICA	F (C <sub>4</sub> F <sub>10</sub> )
<600	90	BENOIT 05	EDEL	<sup>73</sup> Ge
< 10	90	GIRARD 05	SMPL	F (C <sub>2</sub> ClF <sub>5</sub> )
<260	90	MIUCHI 03	BOLO	LiF
<150	90	TAKEDA 03	BOLO	NaF

- <sup>1</sup> AVRORIN 14 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun in data taken between 1998 and 2003. See their Table 1 for limits assuming annihilation into neutrino pairs.
- <sup>2</sup> AARTSEN 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between June 2010 and May 2011.
- <sup>3</sup> ADRIAN-MARTINEZ 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken between Jan. 2007 and Dec. 2008.
- <sup>4</sup> The value has been provided by the authors. APRILE 13 note that the proton limits on Xe are highly sensitive to the theoretical model used. See also APRILE 14A.
- <sup>5</sup> BOLIEV 13 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the sun in data taken from 1978 to 2009. See also SUVOROVA 13 for an older analysis of the same data.
- <sup>6</sup> ABBASI 12 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.
- <sup>7</sup> Use a direction-sensitive detector.
- <sup>8</sup> TANAKA 11 search for neutrinos from the Sun arising from the pair annihilation of  $X^0$  trapped by the Sun. The amount of  $X^0$  depends on the  $X^0$ -proton cross section.
- <sup>9</sup> ABBASI 10 search for  $\nu_\mu$  from annihilations of Kaluza-Klein photon dark matter in the Sun.
- <sup>10</sup> See also AKERIB 05.

**Limits for Spin-Dependent Cross Section**
  

**of Dark Matter Particle ( $X^0$ ) on Neutron**

**For  $m_{X^0} = 20$  GeV**

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.09	90	FELIZARDO	14	SMPL C <sub>2</sub> ClF <sub>5</sub>
< 8	90	<sup>1</sup> UCHIDA	14	XMAS <sup>129</sup> Xe, inelastic
< 1.13 × 10 <sup>-3</sup>	90	<sup>2</sup> APRILE	13	X100 Xe
< 0.02	90	AKIMOV	12	ZEP3 Xe
		<sup>3</sup> AHMED	11B	CDM2 Ge, low threshold
< 0.06	90	AHMED	09	CDM2 Ge
< 0.04	90	LEBEDENKO	09A	ZEP3 Xe
< 50		<sup>4</sup> LIN	09	TEXO Ge
< 6 × 10 <sup>-3</sup>	90	ANGLE	08A	XE10 Xe
< 0.5	90	ALNER	07	ZEP2 Xe
< 25	90	LEE	07A	KIMS CsI
< 0.3	90	<sup>5</sup> AKERIB	06	CDMS <sup>73</sup> Ge, <sup>29</sup> Si
< 30	90	SHIMIZU	06A	CNTR F (CaF <sub>2</sub> )
< 60	90	ALNER	05	NAIA NaI
< 20	90	BARNABE-HE.	05	PICA F (C <sub>4</sub> F <sub>10</sub> )
< 10	90	BENOIT	05	EDEL <sup>73</sup> Ge
< 4	90	KLAPDOR-K...	05	HDMS <sup>73</sup> Ge (enriched)
< 600	90	TAKEDA	03	BOLO NaF

<sup>1</sup> Derived limit from search for inelastic scattering  $X^0 + ^{129}\text{Xe} \rightarrow X^0 + ^{129}\text{Xe}^*$  (39.58 keV).

<sup>2</sup> The value has been provided by the authors. See also APRILE 14A.

<sup>3</sup> AHMED 11B give limits on spin-dependent  $X^0$ -neutron cross section for  $m_{X^0} = 4$ –12 GeV in the range 10<sup>-3</sup>–10 pb. See their Fig. 3.

<sup>4</sup> See their Fig. 6(b) for cross section limits for  $m_{X^0}$  extending down to 2 GeV.

<sup>5</sup> See also AKERIB 05.

**For  $m_{\chi^0} = 100$  GeV**

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.1	90	FELIZARDO	14	SMPL C <sub>2</sub> F <sub>5</sub>
< 0.05	90	<sup>1</sup> UCHIDA	14	XMAS <sup>129</sup> Xe, inelastic
< $4.68 \times 10^{-4}$	90	<sup>2</sup> APRILE	13	X100 Xe
< 0.01	90	AKIMOV	12	ZEP3 Xe
		<sup>3</sup> FELIZARDO	10	SMPL C <sub>2</sub> F <sub>3</sub>
< 0.02	90	AHMED	09	CDM2 Ge
< 0.01	90	LEBEDENKO	09A	ZEP3 Xe
<100	90	LIN	09	TEXO Ge
< 0.01	90	ANGLE	08A	XE10 Xe
< 0.05	90	<sup>4</sup> BEDNYAKOV	08	RVUE Ge
< 0.08	90	ALNER	07	ZEP2 Xe
< 6	90	LEE	07A	KIMS Csl
< 0.07	90	<sup>5</sup> AKERIB	06	CDMS <sup>73</sup> Ge, <sup>29</sup> Si
< 30	90	SHIMIZU	06A	CNTR F (CaF <sub>2</sub> )
< 10	90	ALNER	05	NAIA NaI
< 30	90	BARNABE-HE.	05	PICA F (C <sub>4</sub> F <sub>10</sub> )
< 0.7	90	BENOIT	05	EDEL <sup>73</sup> Ge
< 0.2		<sup>6</sup> GIULIANI	05A	RVUE
< 1.5	90	KLAPDOR-K...	05	HDMS <sup>73</sup> Ge (enriched)
		<sup>7</sup> GIULIANI	04	RVUE
		<sup>8</sup> GIULIANI	04A	RVUE
		<sup>9</sup> MIUCHI	03	BOLO LiF
<800	90	TAKEDA	03	BOLO NaF

<sup>1</sup> Derived limit from search for inelastic scattering  $X^0 + ^{129}\text{Xe}^* \rightarrow X^0 + ^{129}\text{Xe}^*$  (39.58 keV).

<sup>2</sup> The value has been provided by the authors. See also APRILE 14A.

<sup>3</sup> See their Fig. 3 for limits on spin-dependent neutron couplings for  $X^0$  mass of 50 GeV.

<sup>4</sup> BEDNYAKOV 08 reanalyze KLAPDOR-KLEINGROTHAUS 05 and BAUDIS 01 data.

<sup>5</sup> See also AKERIB 05.

<sup>6</sup> GIULIANI 05A analyze available data and give combined limits.

<sup>7</sup> GIULIANI 04 reanalyze COLLAR 00 data and give limits for spin-dependent  $X^0$ -neutron coupling.

<sup>8</sup> GIULIANI 04A give limits for spin-dependent  $X^0$ -neutron couplings from existing data.

<sup>9</sup> MIUCHI 03 give model-independent limit for spin-dependent  $X^0$ -proton and neutron cross sections. See their Fig. 5.

**For  $m_{\chi^0} = 1$  TeV**

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.07	90	FELIZARDO	14	SMPL C <sub>2</sub> F <sub>5</sub>
< 0.2	90	<sup>1</sup> UCHIDA	14	XMAS <sup>129</sup> Xe, inelastic
< $3.64 \times 10^{-3}$	90	<sup>2</sup> APRILE	13	X100 Xe
< 0.08	90	AKIMOV	12	ZEP3 Xe
< 0.2	90	AHMED	09	CDM2 Ge
< 0.1	90	LEBEDENKO	09A	ZEP3 Xe
< 0.1	90	ANGLE	08A	XE10 Xe
< 0.25	90	<sup>3</sup> BEDNYAKOV	08	RVUE Ge

< 0.6	90	ALNER	07	ZEP2	Xe
< 30	90	LEE	07A	KIMS	CsI
< 0.5	90	<sup>4</sup> AKERIB	06	CDMS	<sup>73</sup> Ge, <sup>29</sup> Si
< 40	90	ALNER	05	NAIA	NaI
<200	90	BARNABE-HE.	05	PICA	F (C <sub>4</sub> F <sub>10</sub> )
< 4	90	BENOIT	05	EDEL	<sup>73</sup> Ge
< 10	90	KLAPDOR-K...	05	HDMS	<sup>73</sup> Ge (enriched)
< 4 × 10 <sup>3</sup>	90	TAKEDA	03	BOLO	NaF

<sup>1</sup> Derived limit from search for inelastic scattering  $X^0 + {}^{129}\text{Xe}^* \rightarrow X^0 + {}^{129}\text{Xe}^*$  (39.58 keV).

<sup>2</sup> The value has been provided by the authors. See also APRILE 14A.

<sup>3</sup> BEDNYAKOV 08 reanalyze KLAPDOR-KLEINGROTHAUS 05 and BAUDIS 01 data.

<sup>4</sup> See also AKERIB 05.

## Cross-Section Limits for Dark Matter Particles ( $X^0$ ) on Nuclei

### For $m_{X^0} = 20$ GeV

VALUE (nb)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.03	90	<sup>1</sup> UCHIDA	14	XMAS <sup>129</sup> Xe, inelastic
< 0.08	90	<sup>2</sup> ANGLOHER	02	CRES Al
		<sup>3</sup> BENOIT	00	EDEL Ge
< 0.04	95	<sup>4</sup> KLIMENKO	98	CNTR <sup>73</sup> Ge, inel.
< 0.8		ALESSAND...	96	CNTR O
< 6		ALESSAND...	96	CNTR Te
< 0.02	90	<sup>5</sup> BELLI	96	CNTR <sup>129</sup> Xe, inel.
		<sup>6</sup> BELLI	96C	CNTR <sup>129</sup> Xe
< 4 × 10 <sup>-3</sup>	90	<sup>7</sup> BERNABEI	96	CNTR Na
< 0.3	90	<sup>7</sup> BERNABEI	96	CNTR I
< 0.2	95	<sup>8</sup> SARSA	96	CNTR Na
< 0.015	90	<sup>9</sup> SMITH	96	CNTR Na
< 0.05	95	<sup>10</sup> GARCIA	95	CNTR Natural Ge
< 0.1	95	QUENBY	95	CNTR Na
<90	90	<sup>11</sup> SNOWDEN-...	95	MICA <sup>16</sup> O
< 4 × 10 <sup>3</sup>	90	<sup>11</sup> SNOWDEN-...	95	MICA <sup>39</sup> K
< 0.7	90	BACCI	92	CNTR Na
< 0.12	90	<sup>12</sup> REUSSER	91	CNTR Natural Ge
< 0.06	95	CALDWELL	88	CNTR Natural Ge

<sup>1</sup> UCHIDA 14 limit is for inelastic scattering  $X^0 + {}^{129}\text{Xe}^* \rightarrow X^0 + {}^{129}\text{Xe}^*$  (39.58 keV).

<sup>2</sup> ANGLOHER 02 limit is for spin-dependent WIMP-Aluminum cross section.

<sup>3</sup> BENOIT 00 find four event categories in Ge detectors and suggest that low-energy surface nuclear recoils can explain anomalous events reported by UKDMC and Saclay NaI experiments.

<sup>4</sup> KLIMENKO 98 limit is for inelastic scattering  $X^0 {}^{73}\text{Ge} \rightarrow X^0 {}^{73}\text{Ge}^*$  (13.26 keV).

<sup>5</sup> BELLI 96 limit for inelastic scattering  $X^0 {}^{129}\text{Xe} \rightarrow X^0 {}^{129}\text{Xe}^*$  (39.58 keV).

<sup>6</sup> BELLI 96C use background subtraction and obtain  $\sigma < 150$  pb ( $< 1.5$  fb) (90% CL) for spin-dependent (independent)  $X^0$ -proton cross section. The confidence level is from R. Bernabei, private communication, May 20, 1999.

<sup>7</sup> BERNABEI 96 use pulse shape discrimination to enhance the possible signal. The limit here is from R. Bernabei, private communication, September 19, 1997.

- <sup>8</sup>SARSA 96 search for annual modulation of WIMP signal. See SARSA 97 for details of the analysis. The limit here is from M.L. Sarsa, private communication, May 26, 1997.
- <sup>9</sup>SMITH 96 use pulse shape discrimination to enhance the possible signal. A dark matter density of  $0.4 \text{ GeV cm}^{-3}$  is assumed.
- <sup>10</sup>GARCIA 95 limit is from the event rate. A weaker limit is obtained from searches for diurnal and annual modulation.
- <sup>11</sup>SNOWDEN-IFFT 95 look for recoil tracks in an ancient mica crystal. Similar limits are also given for  $^{27}\text{Al}$  and  $^{28}\text{Si}$ . See COLLAR 96 and SNOWDEN-IFFT 96 for discussion on potential backgrounds.
- <sup>12</sup>REUSSER 91 limit here is changed from published (0.04) after reanalysis by authors. J.L. Vuilleumier, private communication, March 29, 1996.

### For $m_{\chi_0} = 100 \text{ GeV}$

<u>VALUE (nb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$< 3 \times 10^{-3}$	90	1 UCHIDA	14 XMAS	$^{129}\text{Xe}$ , inelastic
$< 0.3$	90	2 ANGLOHER	02 CRES	Al
		3 BELLI	02 RVUE	
		4 BERNABEI	02C DAMA	
		5 GREEN	02 RVUE	
		6 ULLIO	01 RVUE	
		7 BENOIT	00 EDEL	Ge
$< 4 \times 10^{-3}$	90	8 BERNABEI	00D	$^{129}\text{Xe}$ , inel.
		9 AMBROSIO	99 MCRO	
		10 BRHLIK	99 RVUE	
$< 8 \times 10^{-3}$	95	11 KLIMENKO	98 CNTR	$^{73}\text{Ge}$ , inel.
$< 0.08$	95	12 KLIMENKO	98 CNTR	$^{73}\text{Ge}$ , inel.
$< 4$		ALESSAND...	96 CNTR	O
$< 25$		ALESSAND...	96 CNTR	Te
$< 6 \times 10^{-3}$	90	13 BELLI	96 CNTR	$^{129}\text{Xe}$ , inel.
		14 BELLI	96C CNTR	$^{129}\text{Xe}$
$< 1 \times 10^{-3}$	90	15 BERNABEI	96 CNTR	Na
$< 0.3$	90	15 BERNABEI	96 CNTR	I
$< 0.7$	95	16 SARSA	96 CNTR	Na
$< 0.03$	90	17 SMITH	96 CNTR	Na
$< 0.8$	90	17 SMITH	96 CNTR	I
$< 0.35$	95	18 GARCIA	95 CNTR	Natural Ge
$< 0.6$	95	QUENBY	95 CNTR	Na
$< 3$	95	QUENBY	95 CNTR	I
$< 1.5 \times 10^2$	90	19 SNOWDEN-...	95 MICA	$^{16}\text{O}$
$< 4 \times 10^2$	90	19 SNOWDEN-...	95 MICA	$^{39}\text{K}$
$< 0.08$	90	20 BECK	94 CNTR	$^{76}\text{Ge}$
$< 2.5$	90	BACCI	92 CNTR	Na
$< 3$	90	BACCI	92 CNTR	I
$< 0.9$	90	21 REUSSER	91 CNTR	Natural Ge
$< 0.7$	95	CALDWELL	88 CNTR	Natural Ge

- 1 UCHIDA 14 limit is for inelastic scattering  $\chi^0 + {}^{129}\text{Xe}^* \rightarrow \chi^0 + {}^{129}\text{Xe}^*$  (39.58 keV).
- 2 ANGLOHER 02 limit is for spin-dependent WIMP-Aluminum cross section.
- 3 BELLI 02 discuss dependence of the extracted WIMP cross section on the assumptions of the galactic halo structure.
- 4 BERNABEI 02c analyze the DAMA data in the scenario in which  $\chi^0$  scatters into a slightly heavier state as discussed by SMITH 01.
- 5 GREEN 02 discusses dependence of extracted WIMP cross section limits on the assumptions of the galactic halo structure.
- 6 ULLIO 01 disfavor the possibility that the BERNABEI 99 signal is due to spin-dependent WIMP coupling.
- 7 BENOIT 00 find four event categories in Ge detectors and suggest that low-energy surface nuclear recoils can explain anomalous events reported by UKDMC and Saclay NaI experiments.
- 8 BERNABEI 00D limit is for inelastic scattering  $\chi^0 {}^{129}\text{Xe} \rightarrow \chi^0 {}^{129}\text{Xe}$  (39.58 keV).
- 9 AMBROSIO 99 search for upgoing muon events induced by neutrinos originating from WIMP annihilations in the Sun and Earth.
- 10 BRHLIK 99 discuss the effect of astrophysical uncertainties on the WIMP interpretation of the BERNABEI 99 signal.
- 11 KLIMENKO 98 limit is for inelastic scattering  $\chi^0 {}^{73}\text{Ge} \rightarrow \chi^0 {}^{73}\text{Ge}^*$  (13.26 keV).
- 12 KLIMENKO 98 limit is for inelastic scattering  $\chi^0 {}^{73}\text{Ge} \rightarrow \chi^0 {}^{73}\text{Ge}^*$  (66.73 keV).
- 13 BELLI 96 limit for inelastic scattering  $\chi^0 {}^{129}\text{Xe} \rightarrow \chi^0 {}^{129}\text{Xe}^*$  (39.58 keV).
- 14 BELLI 96c use background subtraction and obtain  $\sigma < 0.35$  pb ( $< 0.15$  fb) (90% CL) for spin-dependent (independent)  $\chi^0$ -proton cross section. The confidence level is from R. Bernabei, private communication, May 20, 1999.
- 15 BERNABEI 96 use pulse shape discrimination to enhance the possible signal. The limit here is from R. Bernabei, private communication, September 19, 1997.
- 16 SARSA 96 search for annual modulation of WIMP signal. See SARSA 97 for details of the analysis. The limit here is from M.L. Sarsa, private communication, May 26, 1997.
- 17 SMITH 96 use pulse shape discrimination to enhance the possible signal. A dark matter density of  $0.4 \text{ GeV cm}^{-3}$  is assumed.
- 18 GARCIA 95 limit is from the event rate. A weaker limit is obtained from searches for diurnal and annual modulation.
- 19 SNOWDEN-IFFT 95 look for recoil tracks in an ancient mica crystal. Similar limits are also given for  ${}^{27}\text{Al}$  and  ${}^{28}\text{Si}$ . See COLLAR 96 and SNOWDEN-IFFT 96 for discussion on potential backgrounds.
- 20 BECK 94 uses enriched  ${}^{76}\text{Ge}$  (86% purity).
- 21 REUSSER 91 limit here is changed from published (0.3) after reanalysis by authors. J.L. Vuilleumier, private communication, March 29, 1996.

### For $m_{\chi^0} = 1 \text{ TeV}$

<u>VALUE (nb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.03	90	1 UCHIDA 14	XMAS	${}^{129}\text{Xe}$ , inelastic
< 3	90	2 ANGLOHER 02	CRES	Al
		3 BENOIT 00	EDEL	Ge
		4 BERNABEI 99D	CNTR	SIMP
		5 DERBIN 99	CNTR	SIMP
< 0.06	95	6 KLIMENKO 98	CNTR	${}^{73}\text{Ge}$ , inel.
< 0.4	95	7 KLIMENKO 98	CNTR	${}^{73}\text{Ge}$ , inel.
< 40		ALESSAND...	96	CNTR O
< 700		ALESSAND...	96	CNTR Te

< 0.05	90	8 BELLI	96	CNTR	$^{129}\text{Xe}$ , inel.
< 1.5	90	9 BELLI	96	CNTR	$^{129}\text{Xe}$ , inel.
		10 BELLI	96C	CNTR	$^{129}\text{Xe}$
< 0.01	90	11 BERNABEI	96	CNTR	Na
< 9	90	11 BERNABEI	96	CNTR	I
< 7	95	12 SARSA	96	CNTR	Na
< 0.3	90	13 SMITH	96	CNTR	Na
< 6	90	13 SMITH	96	CNTR	I
< 6	95	14 GARCIA	95	CNTR	Natural Ge
< 8	95	QUENBY	95	CNTR	Na
< 50	95	QUENBY	95	CNTR	I
< 700	90	15 SNOWDEN-...	95	MICA	$^{16}\text{O}$
< 1 $\times 10^3$	90	15 SNOWDEN-...	95	MICA	$^{39}\text{K}$
< 0.8	90	16 BECK	94	CNTR	$^{76}\text{Ge}$
< 30	90	BACCI	92	CNTR	Na
< 30	90	BACCI	92	CNTR	I
< 15	90	17 REUSSER	91	CNTR	Natural Ge
< 6	95	CALDWELL	88	CNTR	Natural Ge

<sup>1</sup> UCHIDA 14 limit is for inelastic scattering  $\chi^0 + ^{129}\text{Xe}^* \rightarrow \chi^0 + ^{129}\text{Xe}^*$  (39.58 keV).

<sup>2</sup> ANGLOHER 02 limit is for spin-dependent WIMP-Aluminum cross section.

<sup>3</sup> BENOIT 00 find four event categories in Ge detectors and suggest that low-energy surface nuclear recoils can explain anomalous events reported by UKDMC and Saclay NaI experiments.

<sup>4</sup> BERNABEI 99D search for SIMPs (Strongly Interacting Massive Particles) in the mass range  $10^3$ – $10^{16}$  GeV. See their Fig. 3 for cross-section limits.

<sup>5</sup> DERBIN 99 search for SIMPs (Strongly Interacting Massive Particles) in the mass range  $10^2$ – $10^{14}$  GeV. See their Fig. 3 for cross-section limits.

<sup>6</sup> KLIMENKO 98 limit is for inelastic scattering  $\chi^0 \ ^{73}\text{Ge} \rightarrow \chi^0 \ ^{73}\text{Ge}^*$  (13.26 keV).

<sup>7</sup> KLIMENKO 98 limit is for inelastic scattering  $\chi^0 \ ^{73}\text{Ge} \rightarrow \chi^0 \ ^{73}\text{Ge}^*$  (66.73 keV).

<sup>8</sup> BELLI 96 limit for inelastic scattering  $\chi^0 \ ^{129}\text{Xe} \rightarrow \chi^0 \ ^{129}\text{Xe}^*$  (39.58 keV).

<sup>9</sup> BELLI 96 limit for inelastic scattering  $\chi^0 \ ^{129}\text{Xe} \rightarrow \chi^0 \ ^{129}\text{Xe}^*$  (236.14 keV).

<sup>10</sup> BELLI 96C use background subtraction and obtain  $\sigma < 0.7$  pb ( $< 0.7$  fb) (90% CL) for spin-dependent (independent)  $\chi^0$ -proton cross section. The confidence level is from R. Bernabei, private communication, May 20, 1999.

<sup>11</sup> BERNABEI 96 use pulse shape discrimination to enhance the possible signal. The limit here is from R. Bernabei, private communication, September 19, 1997.

<sup>12</sup> SARSA 96 search for annual modulation of WIMP signal. See SARSA 97 for details of the analysis. The limit here is from M.L. Sarsa, private communication, May 26, 1997.

<sup>13</sup> SMITH 96 use pulse shape discrimination to enhance the possible signal. A dark matter density of  $0.4 \text{ GeV cm}^{-3}$  is assumed.

<sup>14</sup> GARCIA 95 limit is from the event rate. A weaker limit is obtained from searches for diurnal and annual modulation.

<sup>15</sup> SNOWDEN-IFFT 95 look for recoil tracks in an ancient mica crystal. Similar limits are also given for  $^{27}\text{Al}$  and  $^{28}\text{Si}$ . See COLLAR 96 and SNOWDEN-IFFT 96 for discussion on potential backgrounds.

<sup>16</sup> BECK 94 uses enriched  $^{76}\text{Ge}$  (86% purity).

<sup>17</sup> REUSSER 91 limit here is changed from published (5) after reanalysis by authors. J.L. Vuilleumier, private communication, March 29, 1996.



## Miscellaneous Results from Underground Dark Matter Searches

VALUE	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●			
	<sup>1</sup> APRILE	15 X100	Event rate modulation
	<sup>2</sup> APRILE	15A X100	Electron scattering
<sup>1</sup> APRILE 15 search for periodic variation of electronic recoil event rate in the data between Feb. 2011 and Mar. 2012. No significant modulation is found for periods up to 500 days.			
<sup>2</sup> APRILE 15A search for $X^0$ scattering off electrons. See their Fig. 4 for limits on cross section through axial-vector coupling for $m_{X^0}$ between 0.6 GeV and 1 TeV. For $m_{X^0} = 2$ GeV, $\sigma < 60$ pb (90%CL) is obtained.			

## $X^0$ Annihilation Cross Section

Limits are on  $\sigma v$  for  $X^0$  pair annihilation at threshold.

VALUE ( $\text{cm}^3\text{s}^{-1}$ )	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		<sup>1</sup> AARTSEN	15C ICCB	$\nu$ , Galactic halo
		<sup>2</sup> AARTSEN	15E ICCB	$\nu$ , Galactic center
		<sup>3</sup> ABRAMOWSKI15	HESS	Galactic center
		<sup>4</sup> ACKERMANN 15	FLAT	monochromatic $\gamma$
		<sup>5</sup> ACKERMANN 15A	FLAT	isotropic $\gamma$ background
		<sup>6</sup> ACKERMANN 15B	FLAT	Satellite galaxy
		<sup>7</sup> ADRIAN-MAR.15	ANTR	$\nu$ , Galactic center
$< 2.90 \times 10^{-26}$	95	<sup>8,9</sup> ACKERMANN 14	FLAT	Satellite galaxy, $m = 10$ GeV
$< 1.84 \times 10^{-25}$	95	<sup>8,10</sup> ACKERMANN 14	FLAT	Satellite galaxy, $m = 100$ GeV
$< 1.75 \times 10^{-24}$	95	<sup>8,10</sup> ACKERMANN 14	FLAT	Satellite galaxy, $m = 1$ TeV
$< 4.52 \times 10^{-24}$	95	<sup>11</sup> ALEKSIC 14	MGIC	Segue 1, $m = 1.35$ TeV
		<sup>12</sup> AARTSEN	13C ICCB	Galaxies
		<sup>13</sup> ABRAMOWSKI13	HESS	Central Galactic Halo
		<sup>14</sup> ACKERMANN 13A	FLAT	Galaxy
		<sup>15</sup> ABRAMOWSKI12	HESS	Fornax Cluster
		<sup>16</sup> ACKERMANN 12	FLAT	Galaxy
		<sup>17</sup> ACKERMANN 12	FLAT	Galaxy
		<sup>18</sup> ALIU 12	VRTS	Segue 1
$< 1 \times 10^{-22}$	90	<sup>19</sup> ABBASI 11C	ICCB	Galactic halo, $m=1$ TeV
$< 3 \times 10^{-25}$	95	<sup>20</sup> ABRAMOWSKI11	HESS	Near Galactic center, $m=1$ TeV
$< 1 \times 10^{-26}$	95	<sup>21</sup> ACKERMANN 11	FLAT	Satellite galaxy, $m=10$ GeV
$< 1 \times 10^{-25}$	95	<sup>21</sup> ACKERMANN 11	FLAT	Satellite galaxy, $m=100$ GeV
$< 1 \times 10^{-24}$	95	<sup>21</sup> ACKERMANN 11	FLAT	Satellite galaxy, $m=1$ TeV

<sup>1</sup> AARTSEN 15C search for neutrinos from  $X^0$  annihilation in the Galactic halo. See their Figs. 16 and 17, and Table 5 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 100 GeV and 100 TeV.

<sup>2</sup> AARTSEN 15E search for neutrinos from  $X^0$  annihilation in the Galactic center. See their Figs. 7 and 9, and Table 3 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 30 GeV and 10 TeV.

<sup>3</sup> ABRAMOWSKI 15 search for  $\gamma$  from  $X^0$  annihilation in the Galactic center. See their Fig. 4 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 250 GeV and 10 TeV.

<sup>4</sup> ACKERMANN 15 search for monochromatic  $\gamma$  from  $X^0$  annihilation in the Galactic halo. See their Fig. 8 and Tables 2–4 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 0.2 GeV and 500 GeV.

- <sup>5</sup> ACKERMANN 15A search for  $\gamma$  from  $X^0$  annihilation (both Galactic and extragalactic) in the isotropic  $\gamma$  background. See their Fig. 7 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 10 GeV and 30 TeV.
- <sup>6</sup> ACKERMANN 15B search for  $\gamma$  from  $X^0$  annihilation in 15 dwarf spheroidal satellite galaxies of the Milky Way. See their Figs. 1 and 2 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 2 GeV and 10 TeV.
- <sup>7</sup> ADRIAN-MARTINEZ 15 search for neutrinos from  $X^0$  annihilation in the Galactic center. See their Figs. 10 and 11 and Tables 1 and 2 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 25 GeV and 10 TeV.
- <sup>8</sup> ACKERMANN 14 search for  $\gamma$  from  $X^0$  annihilation in 25 dwarf spheroidal satellite galaxies of the Milky Way. See their Tables II–VII for limits assuming annihilation into  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ ,  $u\bar{u}$ ,  $b\bar{b}$ , and  $W^+W^-$ , for  $X^0$  mass ranging from 2 GeV to 10 TeV.
- <sup>9</sup> Limit assuming  $X^0$  pair annihilation into  $b\bar{b}$ .
- <sup>10</sup> Limit assuming  $X^0$  pair annihilation into  $W^+W^-$ .
- <sup>11</sup> ALEKSIC 14 search for  $\gamma$  from  $X^0$  annihilation in the dwarf spheroidal galaxy Segue 1. The listed limit assumes annihilation into  $W^+W^-$ . See their Figs. 6, 7, and 16 for limits on  $\sigma \cdot v$  for annihilation channels  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ ,  $b\bar{b}$ ,  $t\bar{t}$ ,  $\gamma\gamma$ ,  $\gamma Z$ ,  $W^+W^-$ ,  $ZZ$  for  $X^0$  mass between  $10^2$  and  $10^4$  GeV.
- <sup>12</sup> AARTSEN 13C search for neutrinos from  $X^0$  annihilation in nearby galaxies and galaxy clusters. See their Figs. 5–7 for limits on  $\sigma \cdot v$  for  $X^0 X^0 \rightarrow \nu\bar{\nu}$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ , and  $W^+W^-$  for  $X^0$  mass between 300 GeV and 100 TeV.
- <sup>13</sup> ABRAMOWSKI 13 search for monochromatic  $\gamma$  from  $X^0$  annihilation in the Milky Way halo in the central region. Limit on  $\sigma \cdot v$  between  $10^{-28}$  and  $10^{-25}$   $\text{cm}^3 \text{s}^{-1}$  (95% CL) is obtained for  $X^0$  mass between 500 GeV and 20 TeV for  $X^0 X^0 \rightarrow \gamma\gamma$ .  $X^0$  density distribution in the Galaxy by Einasto is assumed. See their Fig. 4.
- <sup>14</sup> ACKERMANN 13A search for monochromatic  $\gamma$  from  $X^0$  annihilation in the Milky Way. Limit on  $\sigma \cdot v$  for the process  $X^0 X^0 \rightarrow \gamma\gamma$  in the range  $10^{-29}$ – $10^{-27}$   $\text{cm}^3 \text{s}^{-1}$  (95% CL) is obtained for  $X^0$  mass between 5 and 300 GeV. The limit depends slightly on the assumed density profile of  $X^0$  in the Galaxy. See their Tables VII–X and Fig. 10. Supersedes ACKERMANN 12.
- <sup>15</sup> ABRAMOWSKI 12 search for  $\gamma$ 's from  $X^0$  annihilation in the Fornax galaxy cluster. See their Fig. 7 for limits on  $\sigma \cdot v$  for  $X^0$  mass between 0.1 and 100 TeV for the annihilation channels  $\tau^+\tau^-$ ,  $b\bar{b}$ , and  $W^+W^-$ .
- <sup>16</sup> ACKERMANN 12 search for monochromatic  $\gamma$  from  $X^0$  annihilation in the Milky Way. Limit on  $\sigma \cdot v$  in the range  $10^{-28}$ – $10^{-26}$   $\text{cm}^3 \text{s}^{-1}$  (95% CL) is obtained for  $X^0$  mass between 7 and 200 GeV if  $X^0$  annihilates into  $\gamma\gamma$ . The limit depends slightly on the assumed density profile of  $X^0$  in the Galaxy. See their Table III and Fig. 15.
- <sup>17</sup> ACKERMANN 12 search for  $\gamma$  from  $X^0$  annihilation in the Milky Way in the diffuse  $\gamma$  background. Limit on  $\sigma \cdot v$  of  $10^{-24}$   $\text{cm}^3 \text{s}^{-1}$  or larger is obtained for  $X^0$  mass between 5 GeV and 10 TeV for various annihilation channels including  $W^+W^-$ ,  $b\bar{b}$ ,  $gg$ ,  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ . The limit depends slightly on the assumed density profile of  $X^0$  in the Galaxy. See their Figs. 17–20.
- <sup>18</sup> ALIU 12 search for  $\gamma$ 's from  $X^0$  annihilation in the dwarf spheroidal galaxy Segue 1. Limit on  $\sigma \cdot v$  in the range  $10^{-24}$ – $10^{-20}$   $\text{cm}^3 \text{s}^{-1}$  (95% CL) is obtained for  $X^0$  mass between 10 GeV and 2 TeV for annihilation channels  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ ,  $b\bar{b}$ , and  $W^+W^-$ . See their Fig. 3.
- <sup>19</sup> ABBASI 11C search for  $\nu_\mu$  from  $X^0$  annihilation in the outer halo of the Milky Way. The limit assumes annihilation into  $\nu\nu$ . See their Fig. 9 for limits with other annihilation channels.
- <sup>20</sup> ABRAMOWSKI 11 search for  $\gamma$  from  $X^0$  annihilation near the Galactic center. The limit assumes Einasto DM density profile.

- <sup>21</sup> ACKERMANN 11 search for  $\gamma$  from  $X^0$  annihilation in ten dwarf spheroidal satellite galaxies of the Milky Way. The limit for  $m = 10$  GeV assumes annihilation into  $b\bar{b}$ , the others  $W^+W^-$ . See their Fig. 2 for limits with other final states. See also GERINGER-SAMETH 11 for a different analysis of the same data.

## Dark Matter Particle ( $X^0$ ) Production in Hadron Collisions

Searches for  $X^0$  production in association with observable particles ( $\gamma$ , jets, ...) in high energy hadron collisions. If a specific form of effective interaction Lagrangian is assumed, the limits may be translated into limits on  $X^0$ -nucleon scattering cross section.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •			
1	AAD 15AS	ATLS	$b(\bar{b}) + \cancel{E}_T, t\bar{t} + \cancel{E}_T$
2	AAD 15BH	ATLS	jet + $\cancel{E}_T$
3	AAD 15CF	ATLS	$H^0 + \cancel{E}_T$
4	AAD 15CS	ATLS	$\gamma + \cancel{E}_T$
5	KHACHATRYAN 15AG	CMS	$t\bar{t} + \cancel{E}_T$
6	KHACHATRYAN 15AL	CMS	jet + $\cancel{E}_T$
7	KHACHATRYAN 15T	CMS	$\ell + \cancel{E}_T$
8	AAD 14AI	ATLS	$W + \cancel{E}_T$
9	AAD 14K	ATLS	$Z + \cancel{E}_T$
10	AAD 14O	ATLS	$Z + \cancel{E}_T$
11	AAD 13AD	ATLS	jet + $\cancel{E}_T$
12	AAD 13C	ATLS	$\gamma + \cancel{E}_T$
13	AALTONEN 12K	CDF	$t + \cancel{E}_T$
14	AALTONEN 12M	CDF	jet + $\cancel{E}_T$
15	CHATRCHYAN 12AP	CMS	jet + $\cancel{E}_T$
16	CHATRCHYAN 12T	CMS	$\gamma + \cancel{E}_T$

<sup>1</sup> AAD 15AS search for events with one or more bottom quark and missing  $E_T$ , and also events with a top quark pair and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See their Figs. 5 and 6 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}700$  GeV.

<sup>2</sup> AAD 15BH search for events with a jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See their Fig. 12 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1200$  GeV.

<sup>3</sup> AAD 15CF search for events with a  $H^0$  ( $\rightarrow \gamma\gamma$ ) and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See paper for limits on the strength of some contact interactions containing  $X^0$  and the Higgs fields.

<sup>4</sup> AAD 15CS search for events with a photon and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See their Fig. 13 (see also erratum) for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1000$  GeV.

<sup>5</sup> KHACHATRYAN 15AG search for events with a top quark pair and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 19.7 \text{ fb}^{-1}$ . See their Fig. 8 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}200$  GeV.

<sup>6</sup> KHACHATRYAN 15AL search for events with a jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 19.7 \text{ fb}^{-1}$ . See their Fig. 5 and Tables 4–6 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1000$  GeV.

<sup>7</sup> KHACHATRYAN 15T search for events with a lepton and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 19.7 \text{ fb}^{-1}$ . See their Fig. 17 for translated limits on  $X^0$ -proton cross section for  $m = 1\text{--}1000$  GeV.

- <sup>8</sup> AAD 14AI search for events with a  $W$  and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See their Fig. 4 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1500$  GeV.
- <sup>9</sup> AAD 14K search for events with a  $Z$  and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8$  TeV with  $L = 20.3 \text{ fb}^{-1}$ . See their Fig. 5 and 6 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}10^3$  GeV.
- <sup>10</sup> AAD 14O search for  $ZH^0$  production with  $H^0$  decaying to invisible final states. See their Fig. 4 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}60$  GeV in Higgs-portal  $X^0$  scenario.
- <sup>11</sup> AAD 13AD search for events with a jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.7 \text{ fb}^{-1}$ . See their Figs. 5 and 6 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1300$  GeV.
- <sup>12</sup> AAD 13C search for events with a photon and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.6 \text{ fb}^{-1}$ . See their Fig. 3 for translated limits on  $X^0$ -nucleon cross section for  $m = 1\text{--}1000$  GeV.
- <sup>13</sup> AALTONEN 12K search for events with a top quark and missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 7.7 \text{ fb}^{-1}$ . Upper limits on  $\sigma(tX^0)$  in the range 0.4–2 pb (95% CL) is given for  $m_{X^0} = 0\text{--}150$  GeV.
- <sup>14</sup> AALTONEN 12M search for events with a jet and missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 6.7 \text{ fb}^{-1}$ . Upper limits on the cross section in the range 2–10 pb (90% CL) is given for  $m_{X^0} = 1\text{--}300$  GeV. See their Fig. 2 for translated limits on  $X^0$ -nucleon cross section.
- <sup>15</sup> CHATRCHYAN 12AP search for events with a jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 5.0 \text{ fb}^{-1}$ . See their Fig. 4 for translated limits on  $X^0$ -nucleon cross section for  $m_{X^0} = 0.1\text{--}1000$  GeV.
- <sup>16</sup> CHATRCHYAN 12T search for events with a photon and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 5.0 \text{ fb}^{-1}$ . Upper limits on the cross section in the range 13–15 fb (90% CL) is given for  $m_{X^0} = 1\text{--}1000$  GeV. See their Fig. 2 for translated limits on  $X^0$ -nucleon cross section.

## CONCENTRATION OF STABLE PARTICLES IN MATTER

### Concentration of Heavy (Charge +1) Stable Particles in Matter

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$<4 \times 10^{-17}$	95	1 YAMAGATA	93	SPEC Deep sea water, $M=5\text{--}1600m_p$
$<6 \times 10^{-15}$	95	2 VERKERK	92	SPEC Water, $M=10^5$ to $3 \times 10^7$ GeV
$<7 \times 10^{-15}$	95	2 VERKERK	92	SPEC Water, $M=10^4$ , $6 \times 10^7$ GeV
$<9 \times 10^{-15}$	95	2 VERKERK	92	SPEC Water, $M=10^8$ GeV
$<3 \times 10^{-23}$	90	3 HEMMICK	90	SPEC Water, $M = 1000m_p$
$<2 \times 10^{-21}$	90	3 HEMMICK	90	SPEC Water, $M = 5000m_p$
$<3 \times 10^{-20}$	90	3 HEMMICK	90	SPEC Water, $M = 10000m_p$
$<1. \times 10^{-29}$		SMITH	82B	SPEC Water, $M=30\text{--}400m_p$
$<2. \times 10^{-28}$		SMITH	82B	SPEC Water, $M=12\text{--}1000m_p$
$<1. \times 10^{-14}$		SMITH	82B	SPEC Water, $M >1000 m_p$
$<(0.2\text{--}1.) \times 10^{-21}$		SMITH	79	SPEC Water, $M=6\text{--}350 m_p$

- <sup>1</sup>YAMAGATA 93 used deep sea water at 4000 m since the concentration is enhanced in deep sea due to gravity.  
<sup>2</sup>VERKERK 92 looked for heavy isotopes in sea water and put a bound on concentration of stable charged massive particle in sea water. The above bound can be translated into into a bound on charged dark matter particle ( $5 \times 10^6$  GeV), assuming the local density,  $\rho=0.3$  GeV/cm<sup>3</sup>, and the mean velocity  $\langle v \rangle=300$  km/s.  
<sup>3</sup>See HEMMICK 90 Fig. 7 for other masses 100–10000  $m_p$ .

### Concentration of Heavy Stable Particles Bound to Nuclei

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$<1.2 \times 10^{-11}$	95	<sup>1</sup> JAVORSEK	01	SPEC Au, $M= 3$ GeV
$<6.9 \times 10^{-10}$	95	<sup>1</sup> JAVORSEK	01	SPEC Au, $M= 144$ GeV
$<1 \times 10^{-11}$	95	<sup>2</sup> JAVORSEK	01B	SPEC Au, $M= 188$ GeV
$<1 \times 10^{-8}$	95	<sup>2</sup> JAVORSEK	01B	SPEC Au, $M= 1669$ GeV
$<6 \times 10^{-9}$	95	<sup>2</sup> JAVORSEK	01B	SPEC Fe, $M= 188$ GeV
$<1 \times 10^{-8}$	95	<sup>2</sup> JAVORSEK	01B	SPEC Fe, $M= 647$ GeV
$<4 \times 10^{-20}$	90	<sup>3</sup> HEMMICK	90	SPEC C, $M = 100m_p$
$<8 \times 10^{-20}$	90	<sup>3</sup> HEMMICK	90	SPEC C, $M = 1000m_p$
$<2 \times 10^{-16}$	90	<sup>3</sup> HEMMICK	90	SPEC C, $M = 10000m_p$
$<6 \times 10^{-13}$	90	<sup>3</sup> HEMMICK	90	SPEC Li, $M = 1000m_p$
$<1 \times 10^{-11}$	90	<sup>3</sup> HEMMICK	90	SPEC Be, $M = 1000m_p$
$<6 \times 10^{-14}$	90	<sup>3</sup> HEMMICK	90	SPEC B, $M = 1000m_p$
$<4 \times 10^{-17}$	90	<sup>3</sup> HEMMICK	90	SPEC O, $M = 1000m_p$
$<4 \times 10^{-15}$	90	<sup>3</sup> HEMMICK	90	SPEC F, $M = 1000m_p$
$<1.5 \times 10^{-13}/\text{nucleon}$	68	<sup>4</sup> NORMAN	89	SPEC $^{206}\text{Pb}X^-$
$<1.2 \times 10^{-12}/\text{nucleon}$	68	<sup>4</sup> NORMAN	87	SPEC $^{56,58}\text{Fe}X^-$

- <sup>1</sup>JAVORSEK 01 search for (neutral) SIMPs (strongly interacting massive particles) bound to Au nuclei. Here  $M$  is the effective SIMP mass.  
<sup>2</sup>JAVORSEK 01B search for (neutral) SIMPs (strongly interacting massive particles) bound to Au and Fe nuclei from various origins with exposures on the earth's surface, in a satellite, heavy ion collisions, etc. Here  $M$  is the mass of the anomalous nucleus. See also JAVORSEK 02.  
<sup>3</sup>See HEMMICK 90 Fig. 7 for other masses 100–10000  $m_p$ .  
<sup>4</sup>Bound valid up to  $m_{X^-} \sim 100$  TeV.

## GENERAL NEW PHYSICS SEARCHES

This subsection lists some of the search experiments which look for general signatures characteristic of new physics, independent of the framework of a specific model.

The observed events are compatible with Standard Model expectation, unless noted otherwise.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●			
	<sup>1</sup> AAD	15AT ATLS	$t + \cancel{E}_T$
	<sup>2</sup> KHACHATRY...15F	CMS	$t + \cancel{E}_T$

3	AALTONEN	14J	CDF	$W + 2 \text{ jets}$
4	AAD	13A	ATLS	$WW \rightarrow \ell\nu\ell'\nu$
5	AAD	13C	ATLS	$\gamma + \cancel{E}_T$
6	AALTONEN	13I	CDF	Delayed $\gamma + \cancel{E}_T$
7	CHATRCHYAN	13	CMS	$\ell^+\ell^- + \text{jets} + \cancel{E}_T$
8	AAD	12C	ATLS	$t\bar{t} + \cancel{E}_T$
9	AALTONEN	12M	CDF	jet + $\cancel{E}_T$
10	CHATRCHYAN	12AP	CMS	jet + $\cancel{E}_T$
11	CHATRCHYAN	12Q	CMS	$Z + \text{jets} + \cancel{E}_T$
12	CHATRCHYAN	12T	CMS	$\gamma + \cancel{E}_T$
13	AAD	11S	ATLS	jet + $\cancel{E}_T$
14	AALTONEN	11AF	CDF	$\ell^\pm\ell^\pm$
15	CHATRCHYAN	11C	CMS	$\ell^+\ell^- + \text{jets} + \cancel{E}_T$
16	CHATRCHYAN	11U	CMS	jet + $\cancel{E}_T$
17	AALTONEN	10AF	CDF	$\gamma\gamma + \ell, \cancel{E}_T$
18	AALTONEN	09AF	CDF	$\ell\gamma b \cancel{E}_T$
19	AALTONEN	09G	CDF	$\ell\ell\ell \cancel{E}_T$

- <sup>1</sup> AAD 15AT search for events with a top quark and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 20.3 \text{ fb}^{-1}$ .
- <sup>2</sup> KHACHATRYAN 15F search for events with a top quark and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 19.7 \text{ fb}^{-1}$ .
- <sup>3</sup> AALTONEN 14J examine events with a  $W$  and two jets in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 8.9 \text{ fb}^{-1}$ . Invariant mass distributions of the two jets are consistent with the Standard Model expectation.
- <sup>4</sup> AAD 13A search for resonant  $WW$  production in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.7 \text{ fb}^{-1}$ .
- <sup>5</sup> AAD 13C search for events with a photon and missing  $\cancel{E}_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.6 \text{ fb}^{-1}$ .
- <sup>6</sup> AALTONEN 13I search for events with a photon and missing  $E_T$ , where the photon is detected after the expected timing, in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 6.3 \text{ fb}^{-1}$ . The data are consistent with the Standard Model expectation.
- <sup>7</sup> CHATRCHYAN 13 search for events with an opposite-sign lepton pair, jets, and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.98 \text{ fb}^{-1}$ .
- <sup>8</sup> AAD 12C search for events with a  $t\bar{t}$  pair and missing  $\cancel{E}_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 1.04 \text{ fb}^{-1}$ .
- <sup>9</sup> AALTONEN 12M search for events with a jet and missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 6.7 \text{ fb}^{-1}$ .
- <sup>10</sup> CHATRCHYAN 12AP search for events with a jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 5.0 \text{ fb}^{-1}$ .
- <sup>11</sup> CHATRCHYAN 12Q search for events with a  $Z$ , jets, and missing  $\cancel{E}_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.98 \text{ fb}^{-1}$ .
- <sup>12</sup> CHATRCHYAN 12T search for events with a photon and missing  $\cancel{E}_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 5.0 \text{ fb}^{-1}$ .
- <sup>13</sup> AAD 11S search for events with one jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 33 \text{ pb}^{-1}$ .
- <sup>14</sup> AALTONEN 11AF search for high- $p_T$  like-sign dileptons in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 6.1 \text{ fb}^{-1}$ .
- <sup>15</sup> CHATRCHYAN 11C search for events with an opposite-sign lepton pair, jets, and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 34 \text{ pb}^{-1}$ .

- <sup>16</sup> CHATRCHYAN 11U search for events with one jet and missing  $E_T$  in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 36 \text{ pb}^{-1}$ .
- <sup>17</sup> AALTONEN 10AF search for  $\gamma\gamma$  events with  $e, \mu, \tau$ , or missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 1.1\text{--}2.0 \text{ fb}^{-1}$ .
- <sup>18</sup> AALTONEN 09AF search for  $\ell\gamma b$  events with missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 1.9 \text{ fb}^{-1}$ . The observed events are compatible with Standard Model expectation including  $t\bar{t}\gamma$  production.
- <sup>19</sup> AALTONEN 09G search for  $\mu\mu\mu$  and  $\mu\mu e$  events with missing  $E_T$  in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 976 \text{ pb}^{-1}$ .

## LIMITS ON JET-JET RESONANCES

### Heavy Particle Production Cross Section

Limits are for a particle decaying to two hadronic jets.

Units(pb)	CL%	Mass(GeV)	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
			<sup>1</sup> AAD	13D ATLS	7 TeV $pp \rightarrow 2$ jets
			<sup>2</sup> AALTONEN	13R CDF	1.96 TeV $p\bar{p} \rightarrow 4$ jets
			<sup>3</sup> CHATRCHYAN	13A CMS	7 TeV $pp \rightarrow 2$ jets
			<sup>4</sup> CHATRCHYAN	13A CMS	7 TeV $pp \rightarrow b\bar{b}X$
			<sup>5</sup> AAD	12S ATLS	7 TeV $pp \rightarrow 2$ jets
			<sup>6</sup> CHATRCHYAN	12BL CMS	7 TeV $pp \rightarrow t\bar{t}X$
			<sup>7</sup> AAD	11AG ATLS	7 TeV $pp \rightarrow 2$ jets
			<sup>8</sup> AALTONEN	11M CDF	1.96 TeV $p\bar{p} \rightarrow W + 2$ jets
			<sup>9</sup> ABAZOV	11i D0	1.96 TeV $p\bar{p} \rightarrow W + 2$ jets
			<sup>10</sup> AAD	10 ATLS	7 TeV $pp \rightarrow 2$ jets
			<sup>11</sup> KHACHATRY...	10 CMS	7 TeV $pp \rightarrow 2$ jets
			<sup>12</sup> ABE	99F CDF	1.8 TeV $p\bar{p} \rightarrow b\bar{b} + \text{anything}$
			<sup>13</sup> ABE	97G CDF	1.8 TeV $p\bar{p} \rightarrow 2$ jets
<2603	95	200	<sup>14</sup> ABE	93G CDF	1.8 TeV $p\bar{p} \rightarrow 2$ jets
< 44	95	400	<sup>14</sup> ABE	93G CDF	1.8 TeV $p\bar{p} \rightarrow 2$ jets
< 7	95	600	<sup>14</sup> ABE	93G CDF	1.8 TeV $p\bar{p} \rightarrow 2$ jets

<sup>1</sup> AAD 13D search for dijet resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.8 \text{ fb}^{-1}$ . The observed events are compatible with Standard Model expectation. See their Fig. 6 and Table 2 for limits on resonance cross section in the range  $m = 1.0\text{--}4.0$  TeV.

<sup>2</sup> AALTONEN 13R search for production of a pair of jet-jet resonances in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 6.6 \text{ fb}^{-1}$ . See their Fig. 5 and Tables I, II for cross section limits.

<sup>3</sup> CHATRCHYAN 13A search for  $qq, qg$ , and  $gg$  resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.8 \text{ fb}^{-1}$ . See their Fig. 3 and Table 1 for limits on resonance cross section in the range  $m = 1.0\text{--}4.3$  TeV.

<sup>4</sup> CHATRCHYAN 13A search for  $b\bar{b}$  resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.8 \text{ fb}^{-1}$ . See their Fig. 8 and Table 4 for limits on resonance cross section in the range  $m = 1.0\text{--}4.0$  TeV.

<sup>5</sup> AAD 12S search for dijet resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 1.0 \text{ fb}^{-1}$ . See their Fig. 3 and Table 2 for limits on resonance cross section in the range  $m = 0.9\text{--}4.0$  TeV.

<sup>6</sup> CHATRCHYAN 12BL search for  $t\bar{t}$  resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 4.4 \text{ fb}^{-1}$ . See their Fig. 4 for limits on resonance cross section in the range  $m = 0.5\text{--}3.0$  TeV.

- <sup>7</sup> AAD 11AG search for dijet resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 36 \text{ pb}^{-1}$ . Limits on number of events for  $m = 0.6\text{--}4$  TeV are given in their Table 3.
- <sup>8</sup> AALTONEN 11M find a peak in two jet invariant mass distribution around 140 GeV in  $W + 2$  jet events in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 4.3 \text{ fb}^{-1}$ .
- <sup>9</sup> ABAZOV 11I search for two-jet resonances in  $W + 2$  jet events in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96$  TeV with  $L = 4.3 \text{ fb}^{-1}$  and give limits  $\sigma < (2.6\text{--}1.3) \text{ pb}$  (95% CL) for  $m = 110\text{--}170$  GeV. The result is incompatible with AALTONEN 11M.
- <sup>10</sup> AAD 10 search for narrow dijet resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 315 \text{ nb}^{-1}$ . Limits on the cross section in the range  $10\text{--}10^3 \text{ pb}$  is given for  $m = 0.3\text{--}1.7$  TeV.
- <sup>11</sup> KHACHATRYAN 10 search for narrow dijet resonances in  $pp$  collisions at  $E_{\text{cm}} = 7$  TeV with  $L = 2.9 \text{ pb}^{-1}$ . Limits on the cross section in the range  $1\text{--}300 \text{ pb}$  is given for  $m = 0.5\text{--}2.6$  TeV separately in the final states  $qq$ ,  $qg$ , and  $gg$ .
- <sup>12</sup> ABE 99F search for narrow  $b\bar{b}$  resonances in  $p\bar{p}$  collisions at  $E_{\text{cm}}=1.8$  TeV. Limits on  $\sigma(p\bar{p} \rightarrow X + \text{anything}) \times B(X \rightarrow b\bar{b})$  in the range  $3\text{--}10^3 \text{ pb}$  (95%CL) are given for  $m_X=200\text{--}750$  GeV. See their Table I.
- <sup>13</sup> ABE 97G search for narrow dijet resonances in  $p\bar{p}$  collisions with  $106 \text{ pb}^{-1}$  of data at  $E_{\text{cm}} = 1.8$  TeV. Limits on  $\sigma(p\bar{p} \rightarrow X + \text{anything}) \cdot B(X \rightarrow jj)$  in the range  $10^4\text{--}10^{-1} \text{ pb}$  (95%CL) are given for dijet mass  $m=200\text{--}1150$  GeV with both jets having  $|\eta| < 2.0$  and the dijet system having  $|\cos\theta^*| < 0.67$ . See their Table I for the list of limits. Supersedes ABE 93G.
- <sup>14</sup> ABE 93G give cross section times branching ratio into light ( $d$ ,  $u$ ,  $s$ ,  $c$ ,  $b$ ) quarks for  $\Gamma = 0.02 M$ . Their Table II gives limits for  $M = 200\text{--}900$  GeV and  $\Gamma = (0.02\text{--}0.2) M$ .

## LIMITS ON NEUTRAL PARTICLE PRODUCTION

### Production Cross Section of Radiatively-Decaying Neutral Particle

VALUE (pb)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
<(0.043–0.17)	95	<sup>1</sup> ABBIENDI 00D	OPAL	$e^+e^- \rightarrow X^0 Y^0$ , $X^0 \rightarrow Y^0 \gamma$
<(0.05–0.8)	95	<sup>2</sup> ABBIENDI 00D	OPAL	$e^+e^- \rightarrow X^0 X^0$ , $X^0 \rightarrow Y^0 \gamma$
<(2.5–0.5)	95	<sup>3</sup> ACKERSTAFF 97B	OPAL	$e^+e^- \rightarrow X^0 Y^0$ , $X^0 \rightarrow Y^0 \gamma$
<(1.6–0.9)	95	<sup>4</sup> ACKERSTAFF 97B	OPAL	$e^+e^- \rightarrow X^0 X^0$ , $X^0 \rightarrow Y^0 \gamma$

<sup>1</sup> ABBIENDI 00D associated production limit is for  $m_{X^0} = 90\text{--}188$  GeV,  $m_{Y^0}=0$  at  $E_{\text{cm}}=189$  GeV. See also their Fig. 9.

<sup>2</sup> ABBIENDI 00D pair production limit is for  $m_{X^0} = 45\text{--}94$  GeV,  $m_{Y^0}=0$  at  $E_{\text{cm}}=189$  GeV. See also their Fig. 12.

<sup>3</sup> ACKERSTAFF 97B associated production limit is for  $m_{X^0} = 80\text{--}160$  GeV,  $m_{Y^0}=0$  from  $10.0 \text{ pb}^{-1}$  at  $E_{\text{cm}} = 161$  GeV. See their Fig. 3(a).

<sup>4</sup> ACKERSTAFF 97B pair production limit is for  $m_{X^0} = 40\text{--}80$  GeV,  $m_{Y^0}=0$  from  $10.0 \text{ pb}^{-1}$  at  $E_{\text{cm}} = 161$  GeV. See their Fig. 3(b).



**Heavy Particle Production Cross Section**

<u>VALUE (cm<sup>2</sup>/N)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		<sup>1</sup> LEES	15E BABR	$e^+e^-$ collisions
		<sup>2</sup> ADAMS	97B KTEV	$m = 1.2\text{--}5$ GeV
$< 10^{-36}\text{--}10^{-33}$	90	<sup>3</sup> GALLAS	95 TOF	$m = 0.5\text{--}20$ GeV
$< (4\text{--}0.3) \times 10^{-31}$	95	<sup>4</sup> AKESSON	91 CNTR	$m = 0\text{--}5$ GeV
$< 2 \times 10^{-36}$	90	<sup>5</sup> BADIER	86 BDMP	$\tau = (0.05\text{--}1.) \times 10^{-8}$ s
$< 2.5 \times 10^{-35}$		<sup>6</sup> GUSTAFSON	76 CNTR	$\tau > 10^{-7}$ s

<sup>1</sup> LEES 15E search for long-lived neutral particles produced in  $e^+e^-$  collisions in the Upsilon region, which decays into  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $e^\pm\mu^\mp$ ,  $\pi^+\pi^-$ ,  $K^+K^-$ , or  $\pi^\pm K^\mp$ . See their Fig. 2 for cross section limits.

<sup>2</sup> ADAMS 97B search for a hadron-like neutral particle produced in  $pN$  interactions, which decays into a  $\rho^0$  and a weakly interacting massive particle. Upper limits are given for the ratio to  $K_L$  production for the mass range 1.2–5 GeV and lifetime  $10^{-9}\text{--}10^{-4}$  s. See also our Light Gluino Section.

<sup>3</sup> GALLAS 95 limit is for a weakly interacting neutral particle produced in 800 GeV/c  $pN$  interactions decaying with a lifetime of  $10^{-4}\text{--}10^{-8}$  s. See their Figs. 8 and 9. Similar limits are obtained for a stable particle with interaction cross section  $10^{-29}\text{--}10^{-33}$  cm<sup>2</sup>. See Fig. 10.

<sup>4</sup> AKESSON 91 limit is from weakly interacting neutral long-lived particles produced in  $pN$  reaction at 450 GeV/c performed at CERN SPS. Bourquin-Gaillard formula is used as the production model. The above limit is for  $\tau > 10^{-7}$  s. For  $\tau > 10^{-9}$  s,  $\sigma < 10^{-30}$  cm<sup>2</sup>/nucleon is obtained.

<sup>5</sup> BADIER 86 looked for long-lived particles at 300 GeV  $\pi^-$  beam dump. The limit applies for nonstrongly interacting neutral or charged particles with mass  $>2$  GeV. The limit applies for particle modes,  $\mu^+\pi^-$ ,  $\mu^+\mu^-$ ,  $\pi^+\pi^-X$ ,  $\pi^+\pi^-\pi^\pm$  etc. See their figure 5 for the contours of limits in the mass- $\tau$  plane for each mode.

<sup>6</sup> GUSTAFSON 76 is a 300 GeV FNAL experiment looking for heavy ( $m > 2$  GeV) long-lived neutral hadrons in the M4 neutral beam. The above typical value is for  $m = 3$  GeV and assumes an interaction cross section of 1 mb. Values as a function of mass and interaction cross section are given in figure 2.

**Production of New Penetrating Non- $\nu$  Like States in Beam Dump**

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●

<sup>1</sup> LOSECCO 81 CALO 28 GeV protons

<sup>1</sup> No excess neutral-current events leads to  $\sigma(\text{production}) \times \sigma(\text{interaction}) \times \text{acceptance} < 2.26 \times 10^{-71}$  cm<sup>4</sup>/nucleon<sup>2</sup> (CL = 90%) for light neutrals. Acceptance depends on models (0.1 to  $4. \times 10^{-4}$ ).

**LIMITS ON CHARGED PARTICLES IN  $e^+e^-$** **Heavy Particle Production Cross Section in  $e^+e^-$** 

Ratio to  $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$  unless noted. See also entries in Free Quark Search and Magnetic Monopole Searches.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●

<sup>1</sup> ACKERSTAFF 98P OPAL  $Q=1,2/3$ ,  $m=45\text{--}89.5$  GeV

<sup>2</sup> ABREU 97D DLPH  $Q=1,2/3$ ,  $m=45\text{--}84$  GeV

<sup>3</sup> BARATE 97K ALEP  $Q=1$ ,  $m=45\text{--}85$  GeV

$<2 \times 10^{-5}$	95	<sup>4</sup> AKERS	95R	OPAL	$Q=1, m=5-45$ GeV
$<1 \times 10^{-5}$	95	<sup>4</sup> AKERS	95R	OPAL	$Q=2, m=5-45$ GeV
$<2 \times 10^{-3}$	90	<sup>5</sup> BUSKULIC	93C	ALEP	$Q=1, m=32-72$ GeV
$<(10^{-2}-1)$	95	<sup>6</sup> ADACHI	90C	TOPZ	$Q=1, m=1-16, 18-27$ GeV
$<7 \times 10^{-2}$	90	<sup>7</sup> ADACHI	90E	TOPZ	$Q=1, m=5-25$ GeV
$<1.6 \times 10^{-2}$	95	<sup>8</sup> KINOSHITA	82	PLAS	$Q=3-180, m < 14.5$ GeV
$<5.0 \times 10^{-2}$	90	<sup>9</sup> BARTEL	80	JADE	$Q=(3,4,5)/3, 2-12$ GeV

<sup>1</sup> ACKERSTAFF 98P search for pair production of long-lived charged particles at  $E_{cm}$  between 130 and 183 GeV and give limits  $\sigma < (0.05-0.2)$  pb (95%CL) for spin-0 and spin-1/2 particles with  $m=45-89.5$  GeV, charge 1 and 2/3. The limit is translated to the cross section at  $E_{cm}=183$  GeV with the  $s$  dependence described in the paper. See their Figs. 2-4.

<sup>2</sup> ABREU 97D search for pair production of long-lived particles and give limits  $\sigma < (0.4-2.3)$  pb (95%CL) for various center-of-mass energies  $E_{cm}=130-136, 161,$  and  $172$  GeV, assuming an almost flat production distribution in  $\cos\theta$ .

<sup>3</sup> BARATE 97K search for pair production of long-lived charged particles at  $E_{cm} = 130, 136, 161,$  and  $172$  GeV and give limits  $\sigma < (0.2-0.4)$  pb (95%CL) for spin-0 and spin-1/2 particles with  $m=45-85$  GeV. The limit is translated to the cross section at  $E_{cm}=172$  GeV with the  $E_{cm}$  dependence described in the paper. See their Figs. 2 and 3 for limits on  $J = 1/2$  and  $J = 0$  cases.

<sup>4</sup> AKERS 95R is a CERN-LEP experiment with  $W_{cm} \sim m_Z$ . The limit is for the production of a stable particle in multihadron events normalized to  $\sigma(e^+e^- \rightarrow \text{hadrons})$ . Constant phase space distribution is assumed. See their Fig. 3 for bounds for  $Q = \pm 2/3, \pm 4/3$ .

<sup>5</sup> BUSKULIC 93C is a CERN-LEP experiment with  $W_{cm} = m_Z$ . The limit is for a pair or single production of heavy particles with unusual ionization loss in TPC. See their Fig. 5 and Table 1.

<sup>6</sup> ADACHI 90C is a KEK-TRISTAN experiment with  $W_{cm} = 52-60$  GeV. The limit is for pair production of a scalar or spin-1/2 particle. See Figs. 3 and 4.

<sup>7</sup> ADACHI 90E is KEK-TRISTAN experiment with  $W_{cm} = 52-61.4$  GeV. The above limit is for inclusive production cross section normalized to  $\sigma(e^+e^- \rightarrow \mu^+\mu^-) \cdot \beta(3-\beta^2)/2$ , where  $\beta = (1 - 4m^2/W_{cm}^2)^{1/2}$ . See the paper for the assumption about the production mechanism.

<sup>8</sup> KINOSHITA 82 is SLAC PEP experiment at  $W_{cm} = 29$  GeV using lexan and  $^{39}\text{Cr}$  plastic sheets sensitive to highly ionizing particles.

<sup>9</sup> BARTEL 80 is DESY-PETRA experiment with  $W_{cm} = 27-35$  GeV. Above limit is for inclusive pair production and ranges between  $1. \times 10^{-1}$  and  $1. \times 10^{-2}$  depending on mass and production momentum distributions. (See their figures 9, 10, 11).

## Branching Fraction of $Z^0$ to a Pair of Stable Charged Heavy Fermions

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$<5 \times 10^{-6}$	95	<sup>1</sup> AKERS	95R	OPAL	$m=40.4-45.6$ GeV
$<1 \times 10^{-3}$	95	AKRAWY	900	OPAL	$m=29-40$ GeV

<sup>1</sup> AKERS 95R give the 95% CL limit  $\sigma(X\bar{X})/\sigma(\mu\mu) < 1.8 \times 10^{-4}$  for the pair production of singly- or doubly-charged stable particles. The limit applies for the mass range 40.4-45.6 GeV for  $X^\pm$  and  $< 45.6$  GeV for  $X^{\pm\pm}$ . See the paper for bounds for  $Q = \pm 2/3, \pm 4/3$ .

## LIMITS ON CHARGED PARTICLES IN HADRONIC REACTIONS

### MASS LIMITS for Long-Lived Charged Heavy Fermions

Limits are for spin 1/2 particles with no color and  $SU(2)_L$  charge. The electric charge  $Q$  of the particle (in the unit of  $e$ ) is therefore equal to its weak hypercharge. Pair production by Drell-Yan like  $\gamma$  and  $Z$  exchange is assumed to derive the limits.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>660	95	<sup>1</sup> AAD	15BJ ATLS	$ Q  = 2$
>200	95	<sup>2</sup> CHATRCHYAN 13AB	CMS	$ Q  = 1/3$
>480	95	<sup>2</sup> CHATRCHYAN 13AB	CMS	$ Q  = 2/3$
>574	95	<sup>2</sup> CHATRCHYAN 13AB	CMS	$ Q  = 1$
>685	95	<sup>2</sup> CHATRCHYAN 13AB	CMS	$ Q  = 2$
>140	95	<sup>3</sup> CHATRCHYAN 13AR	CMS	$ Q  = 1/3$
>310	95	<sup>3</sup> CHATRCHYAN 13AR	CMS	$ Q  = 2/3$

• • • We do not use the following data for averages, fits, limits, etc. • • •

<sup>1</sup> AAD 15BJ use  $20.3 \text{ fb}^{-1}$  of  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$ . See paper for limits for  $|Q| = 3, 4, 5, 6$ .

<sup>2</sup> CHATRCHYAN 13AB use  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  and  $18.8 \text{ fb}^{-1}$  at  $E_{\text{cm}} = 8 \text{ TeV}$ . See paper for limits for  $|Q| = 3, 4, \dots, 8$ .

<sup>3</sup> CHATRCHYAN 13AR use  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$ .

### Heavy Particle Production Cross Section

VALUE (nb)	CL%	DOCUMENT ID	TECN	COMMENT
<1.2 $\times 10^{-3}$	95	<sup>3</sup> AAD	11i ATLS	$ q =10e, m=0.2-1 \text{ TeV}$
<1.0 $\times 10^{-5}$	95	<sup>4,5</sup> AALTONEN	09Z CDF	$m > 100 \text{ GeV}$ , noncolored
<4.8 $\times 10^{-5}$	95	<sup>4,6</sup> AALTONEN	09Z CDF	$m > 100 \text{ GeV}$ , colored
<0.31–0.04 $\times 10^{-3}$	95	<sup>7</sup> ABAZOV	09M D0	pair production
<0.19	95	<sup>8</sup> AKTAS	04C H1	$m=3-10 \text{ GeV}$
<0.05	95	<sup>9</sup> ABE	92J CDF	$m=50-200 \text{ GeV}$
<30–130		<sup>10</sup> CARROLL	78 SPEC	$m=2-2.5 \text{ GeV}$
<100		<sup>11</sup> LEIPUNER	73 CNTR	$m=3-11 \text{ GeV}$

• • • We do not use the following data for averages, fits, limits, etc. • • •

<sup>1</sup> AAIJ 15BD search for production of long-lived particles in  $pp$  collisions at  $E_{\text{cm}} = 7$  and  $8 \text{ TeV}$ . See their Table 6 for cross section limits.

<sup>2</sup> AAD 13AH search for production of long-lived particles with  $|q|=(2-6)e$  in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $4.4 \text{ fb}^{-1}$ . See their Fig. 8 for cross section limits.

<sup>3</sup> AAD 11i search for production of highly ionizing massive particles in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 3.1 \text{ pb}^{-1}$ . See their Table 5 for similar limits for  $|q| = 6e$  and  $17e$ , Table 6 for limits on pair production cross section.

<sup>4</sup> AALTONEN 09Z search for long-lived charged particles in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 1.0 \text{ fb}^{-1}$ . The limits are on production cross section for a particle of mass above  $100 \text{ GeV}$  in the region  $|\eta| \lesssim 0.7, p_T > 40 \text{ GeV}$ , and  $0.4 < \beta < 1.0$ .

<sup>5</sup> Limit for weakly interacting charge-1 particle.

<sup>6</sup> Limit for up-quark like particle.

<sup>7</sup> ABAZOV 09M search for pair production of long-lived charged particles in  $p\bar{p}$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 1.1 \text{ fb}^{-1}$ . Limit on the cross section of  $(0.31-0.04) \text{ pb}$  (95% CL) is given for the mass range of  $60-300 \text{ GeV}$ , assuming the kinematics of stau pair production.

<sup>8</sup> AKTAS 04C search for production of long-lived particles in  $pp$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 1.1 \text{ fb}^{-1}$ . Limit on the cross section of  $(0.31-0.04) \text{ pb}$  (95% CL) is given for the mass range of  $60-300 \text{ GeV}$ , assuming the kinematics of stau pair production.

<sup>9</sup> ABE 92J search for production of long-lived particles in  $pp$  collisions at  $E_{\text{cm}} = 1.96 \text{ TeV}$  with  $L = 1.1 \text{ fb}^{-1}$ . Limit on the cross section of  $(0.31-0.04) \text{ pb}$  (95% CL) is given for the mass range of  $60-300 \text{ GeV}$ , assuming the kinematics of stau pair production.

- <sup>8</sup> AKTAS 04C look for charged particle photoproduction at HERA with mean c.m. energy of 200 GeV.
- <sup>9</sup> ABE 92J look for pair production of unit-charged particles which leave detector before decaying. Limit shown here is for  $m=50$  GeV. See their Fig. 5 for different charges and stronger limits for higher mass.
- <sup>10</sup> CARROLL 78 look for neutral,  $S = -2$  dihyperon resonance in  $pp \rightarrow 2K^+X$ . Cross section varies within above limits over mass range and  $p_{\text{lab}} = 5.1\text{--}5.9$  GeV/ $c$ .
- <sup>11</sup> LEIPUNER 73 is an NAL 300 GeV  $p$  experiment. Would have detected particles with lifetime greater than 200 ns.

## Heavy Particle Production Differential Cross Section

VALUE ( $\text{cm}^2\text{sr}^{-1}\text{GeV}^{-1}$ )	CL%	DOCUMENT ID	TECN	CHG	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
$<2.6 \times 10^{-36}$	90	1 BALDIN	76	CNTR	– $Q=1, m=2.1\text{--}9.4$ GeV
$<2.2 \times 10^{-33}$	90	2 ALBROW	75	SPEC	$\pm$ $Q=\pm 1, m=4\text{--}15$ GeV
$<1.1 \times 10^{-33}$	90	2 ALBROW	75	SPEC	$\pm$ $Q=\pm 2, m=6\text{--}27$ GeV
$<8. \times 10^{-35}$	90	3 JOVANOVI...	75	CNTR	$\pm$ $m=15\text{--}26$ GeV
$<1.5 \times 10^{-34}$	90	3 JOVANOVI...	75	CNTR	$\pm$ $Q=\pm 2, m=3\text{--}10$ GeV
$<6. \times 10^{-35}$	90	3 JOVANOVI...	75	CNTR	$\pm$ $Q=\pm 2, m=10\text{--}26$ GeV
$<1. \times 10^{-31}$	90	4 APPEL	74	CNTR	$\pm$ $m=3.2\text{--}7.2$ GeV
$<5.8 \times 10^{-34}$	90	5 ALPER	73	SPEC	$\pm$ $m=1.5\text{--}24$ GeV
$<1.2 \times 10^{-35}$	90	6 ANTIPOV	71B	CNTR	– $Q=-, m=2.2\text{--}2.8$
$<2.4 \times 10^{-35}$	90	7 ANTIPOV	71C	CNTR	– $Q=-, m=1.2\text{--}1.7,$ $2.1\text{--}4$
$<2.4 \times 10^{-35}$	90	BINON	69	CNTR	– $Q=-, m=1\text{--}1.8$ GeV
$<1.5 \times 10^{-36}$		8 DORFAN	65	CNTR	Be target $m=3\text{--}7$ GeV
$<3.0 \times 10^{-36}$		8 DORFAN	65	CNTR	Fe target $m=3\text{--}7$ GeV

- <sup>1</sup> BALDIN 76 is a 70 GeV Serpukhov experiment. Value is per Al nucleus at  $\theta = 0$ . For other charges in range  $-0.5$  to  $-3.0$ , CL = 90% limit is  $(2.6 \times 10^{-36})/|(charge)|$  for mass range  $(2.1\text{--}9.4 \text{ GeV}) \times |(charge)|$ . Assumes stable particle interacting with matter as do antiprotons.
- <sup>2</sup> ALBROW 75 is a CERN ISR experiment with  $E_{\text{cm}} = 53$  GeV.  $\theta = 40$  mr. See figure 5 for mass ranges up to 35 GeV.
- <sup>3</sup> JOVANOVIICH 75 is a CERN ISR 26+26 and 15+15 GeV  $pp$  experiment. Figure 4 covers ranges  $Q = 1/3$  to 2 and  $m = 3$  to 26 GeV. Value is per GeV momentum.
- <sup>4</sup> APPEL 74 is NAL 300 GeV  $pW$  experiment. Studies forward production of heavy (up to 24 GeV) charged particles with momenta 24–200 GeV ( $-$ charge) and 40–150 GeV ( $+$ charge). Above typical value is for 75 GeV and is per GeV momentum per nucleon.
- <sup>5</sup> ALPER 73 is CERN ISR 26+26 GeV  $pp$  experiment.  $p > 0.9$  GeV,  $0.2 < \beta < 0.65$ .
- <sup>6</sup> ANTIPOV 71B is from same 70 GeV  $p$  experiment as ANTIPOV 71C and BINON 69.
- <sup>7</sup> ANTIPOV 71C limit inferred from flux ratio. 70 GeV  $p$  experiment.
- <sup>8</sup> DORFAN 65 is a 30 GeV/ $c$   $p$  experiment at BNL. Units are per GeV momentum per nucleus.

## Long-Lived Heavy Particle Invariant Cross Section

VALUE ( $\text{cm}^2/\text{GeV}^2/N$ )	CL%	DOCUMENT ID	TECN	CHG	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
$<5\text{--}700 \times 10^{-35}$	90	1 BERNSTEIN	88	CNTR	
$<5\text{--}700 \times 10^{-37}$	90	1 BERNSTEIN	88	CNTR	
$<2.5 \times 10^{-36}$	90	2 THRON	85	CNTR	– $Q=1, m=4\text{--}12$ GeV
$<1. \times 10^{-35}$	90	2 THRON	85	CNTR	+ $Q=1, m=4\text{--}12$ GeV

$<6. \times 10^{-33}$	90	<sup>3</sup> ARMITAGE	79	SPEC	$m=1.87$ GeV
$<1.5 \times 10^{-33}$	90	<sup>3</sup> ARMITAGE	79	SPEC	$m=1.5-3.0$ GeV
		<sup>4</sup> BOZZOLI	79	CNTR $\pm$	$Q = (2/3, 1, 4/3, 2)$
$<1.1 \times 10^{-37}$	90	<sup>5</sup> CUTTS	78	CNTR	$m=4-10$ GeV
$<3.0 \times 10^{-37}$	90	<sup>6</sup> VIDAL	78	CNTR	$m=4.5-6$ GeV

<sup>1</sup>BERNSTEIN 88 limits apply at  $x = 0.2$  and  $p_T = 0$ . Mass and lifetime dependence of limits are shown in the regions:  $m = 1.5-7.5$  GeV and  $\tau = 10^{-8}-2 \times 10^{-6}$  s. First number is for hadrons; second is for weakly interacting particles.

<sup>2</sup>THRON 85 is FNAL 400 GeV proton experiment. Mass determined from measured velocity and momentum. Limits are for  $\tau > 3 \times 10^{-9}$  s.

<sup>3</sup>ARMITAGE 79 is CERN-ISR experiment at  $E_{cm} = 53$  GeV. Value is for  $x = 0.1$  and  $p_T = 0.15$ . Observed particles at  $m = 1.87$  GeV are found all consistent with being antideuterons.

<sup>4</sup>BOZZOLI 79 is CERN-SPS 200 GeV  $pN$  experiment. Looks for particle with  $\tau$  larger than  $10^{-8}$  s. See their figure 11-18 for production cross-section upper limits vs mass.

<sup>5</sup>CUTTS 78 is  $p$ Be experiment at FNAL sensitive to particles of  $\tau > 5 \times 10^{-8}$  s. Value is for  $-0.3 < x < 0$  and  $p_T = 0.175$ .

<sup>6</sup>VIDAL 78 is FNAL 400 GeV proton experiment. Value is for  $x = 0$  and  $p_T = 0$ . Puts lifetime limit of  $< 5 \times 10^{-8}$  s on particle in this mass range.

### Long-Lived Heavy Particle Production ( $\sigma(\text{Heavy Particle}) / \sigma(\pi)$ )

VALUE	EVTS	DOCUMENT ID	TECN	CHG	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$<10^{-8}$		<sup>1</sup> NAKAMURA	89	SPEC $\pm$	$Q = (-5/3, \pm 2)$
	0	<sup>2</sup> BUSSIÈRE	80	CNTR $\pm$	$Q = (2/3, 1, 4/3, 2)$

<sup>1</sup>NAKAMURA 89 is KEK experiment with 12 GeV protons on Pt target. The limit applies for mass  $\lesssim 1.6$  GeV and lifetime  $\gtrsim 10^{-7}$  s.

<sup>2</sup>BUSSIÈRE 80 is CERN-SPS experiment with 200-240 GeV protons on Be and Al target. See their figures 6 and 7 for cross-section ratio vs mass.

### Production and Capture of Long-Lived Massive Particles

VALUE ( $10^{-36}$ cm <sup>2</sup> )	DOCUMENT ID	TECN	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$<20$ to 800	<sup>1</sup> ALEKSEEV	76	ELEC $\tau=5$ ms to 1 day
$<200$ to 2000	<sup>1</sup> ALEKSEEV	76B	ELEC $\tau=100$ ms to 1 day
$<1.4$ to 9	<sup>2</sup> FRANKEL	75	CNTR $\tau=50$ ms to 10 hours
$<0.1$ to 9	<sup>3</sup> FRANKEL	74	CNTR $\tau=1$ to 1000 hours

<sup>1</sup>ALEKSEEV 76 and ALEKSEEV 76B are 61-70 GeV  $p$  Serpukhov experiment. Cross section is per Pb nucleus.

<sup>2</sup>FRANKEL 75 is extension of FRANKEL 74.

<sup>3</sup>FRANKEL 74 looks for particles produced in thick Al targets by 300-400 GeV/c protons.

### Long-Lived Particle Search at Hadron Collisions

Limits are for cross section times branching ratio.

VALUE (pb/nucleon)	CL%	DOCUMENT ID	TECN	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$<2$	90	<sup>1</sup> BADIÈR	86	BDMP $\tau = (0.05-1.) \times 10^{-8}$ s
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<sup>1</sup>BADIÈR 86 looked for long-lived particles at 300 GeV  $\pi^-$  beam dump. The limit applies for nonstrongly interacting neutral or charged particles with mass  $>2$  GeV. The limit applies for particle modes,  $\mu^+ \pi^-$ ,  $\mu^+ \mu^-$ ,  $\pi^+ \pi^- X$ ,  $\pi^+ \pi^- \pi^\pm$  etc. See their figure 5 for the contours of limits in the mass- $\tau$  plane for each mode.

**Long-Lived Heavy Particle Cross Section**

<i>VALUE</i> (pb/sr)	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<34	95	<sup>1</sup> RAM	94	SPEC 1015 < $m_{X^{++}}$ < 1085 MeV
<75	95	<sup>1</sup> RAM	94	SPEC 920 < $m_{X^{++}}$ < 1025 MeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

<sup>1</sup>RAM 94 search for a long-lived doubly-charged fermion  $X^{++}$  with mass between  $m_N$  and  $m_N + m_\pi$  and baryon number +1 in the reaction  $pp \rightarrow X^{++}n$ . No candidate is found. The limit is for the cross section at  $15^\circ$  scattering angle at 460 MeV incident energy and applies for  $\tau(X^{++}) \gg 0.1 \mu\text{s}$ .

**LIMITS ON CHARGED PARTICLES IN COSMIC RAYS****Heavy Particle Flux in Cosmic Rays**

<i>VALUE</i> ( $\text{cm}^{-2}\text{sr}^{-1}\text{s}^{-1}$ )	<i>CL%</i>	<i>EVTS</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>CHG</i>	<i>COMMENT</i>	
< 1	$\times 10^{-8}$	90	0	<sup>1</sup> AGNESE	15	CDM2	$Q = 1/6$
~ 6	$\times 10^{-9}$		2	<sup>2</sup> SAITO	90		$Q \simeq 14, m \simeq 370m_p$
< 1.4	$\times 10^{-12}$	90	0	<sup>3</sup> MINCER	85	CALO	$m \geq 1 \text{ TeV}$
				<sup>4</sup> SAKUYAMA	83B	PLAS	$m \sim 1 \text{ TeV}$
< 1.7	$\times 10^{-11}$	99	0	<sup>5</sup> BHAT	82	CC	
< 1.	$\times 10^{-9}$	90	0	<sup>6</sup> MARINI	82	CNTR	$\pm$ $Q=1, m \sim 4.5m_p$
2.	$\times 10^{-9}$		3	<sup>7</sup> YOCK	81	SPRK	$\pm$ $Q=1, m \sim 4.5m_p$
			3	<sup>7</sup> YOCK	81	SPRK	Fractionally charged
3.0	$\times 10^{-9}$		3	<sup>8</sup> YOCK	80	SPRK	$m \sim 4.5 m_p$
(4 $\pm$ 1)	$\times 10^{-11}$		3	GOODMAN	79	ELEC	$m \geq 5 \text{ GeV}$
< 1.3	$\times 10^{-9}$	90		<sup>9</sup> BHAT	78	CNTR	$\pm$ $m > 1 \text{ GeV}$
< 1.0	$\times 10^{-9}$		0	BRIATORE	76	ELEC	
< 7.	$\times 10^{-10}$	90	0	YOCK	75	ELEC	$\pm$ $Q > 7e$ or $< -7e$
> 6.	$\times 10^{-9}$		5	<sup>10</sup> YOCK	74	CNTR	$m > 6 \text{ GeV}$
< 3.0	$\times 10^{-8}$		0	DARDO	72	CNTR	
< 1.5	$\times 10^{-9}$		0	TONWAR	72	CNTR	$m > 10 \text{ GeV}$
< 3.0	$\times 10^{-10}$		0	BJORNBOE	68	CNTR	$m > 5 \text{ GeV}$
< 5.0	$\times 10^{-11}$	90	0	JONES	67	ELEC	$m=5-15 \text{ GeV}$

<sup>1</sup>See AGNESE 15 Fig. 6 for limits extending down to  $Q = 1/200$ .

<sup>2</sup>SAITO 90 candidates carry about 450 MeV/nucleon. Cannot be accounted for by conventional backgrounds. Consistent with strange quark matter hypothesis.

<sup>3</sup>MINCER 85 is high statistics study of calorimeter signals delayed by 20–200 ns. Calibration with AGS beam shows they can be accounted for by rare fluctuations in signals from low-energy hadrons in the shower. Claim that previous delayed signals including BJORNBOE 68, DARDO 72, BHAT 82, SAKUYAMA 83B below may be due to this fake effect.

<sup>4</sup>SAKUYAMA 83B analyzed 6000 extended air shower events. Increase of delayed particles and change of lateral distribution above  $10^{17}$  eV may indicate production of very heavy parent at top of atmosphere.

- <sup>5</sup> BHAT 82 observed 12 events with delay  $> 2. \times 10^{-8}$  s and with more than 40 particles. 1 eV has good hadron shower. However all events are delayed in only one of two detectors in cloud chamber, and could not be due to strongly interacting massive particle.
- <sup>6</sup> MARINI 82 applied PEP-counter for TOF. Above limit is for velocity = 0.54 of light. Limit is inconsistent with YOCK 80 YOCK 81 events if isotropic dependence on zenith angle is assumed.
- <sup>7</sup> YOCK 81 saw another 3 events with  $Q = \pm 1$  and  $m$  about  $4.5m_p$  as well as 2 events with  $m > 5.3m_p$ ,  $Q = \pm 0.75 \pm 0.05$  and  $m > 2.8m_p$ ,  $Q = \pm 0.70 \pm 0.05$  and 1 event with  $m = (9.3 \pm 3.)m_p$ ,  $Q = \pm 0.89 \pm 0.06$  as possible heavy candidates.
- <sup>8</sup> YOCK 80 events are with charge exactly or approximately equal to unity.
- <sup>9</sup> BHAT 78 is at Kolar gold fields. Limit is for  $\tau > 10^{-6}$  s.
- <sup>10</sup> YOCK 74 events could be tritons.

### Superheavy Particle (Quark Matter) Flux in Cosmic Rays

<u>VALUE</u> ( $\text{cm}^{-2}\text{sr}^{-1}\text{s}^{-1}$ )	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		<sup>1</sup> ADRIANI	15	PMLA $4 < m < 1.2 \times 10^5 m_p$
$< 5 \times 10^{-16}$	90	<sup>2</sup> AMBROSIO	00B	MCRO $m > 5 \times 10^{14}$ GeV
$< 1.8 \times 10^{-12}$	90	<sup>3</sup> ASTONE	93	CNTR $m \geq 1.5 \times 10^{-13}$ gram
$< 1.1 \times 10^{-14}$	90	<sup>4</sup> AHLEN	92	MCRO $10^{-10} < m < 0.1$ gram
$< 2.2 \times 10^{-14}$	90	<sup>5</sup> NAKAMURA	91	PLAS $m > 10^{11}$ GeV
$< 6.4 \times 10^{-16}$	90	<sup>6</sup> ORITO	91	PLAS $m > 10^{12}$ GeV
$< 2.0 \times 10^{-11}$	90	<sup>7</sup> LIU	88	BOLO $m > 1.5 \times 10^{-13}$ gram
$< 4.7 \times 10^{-12}$	90	<sup>8</sup> BARISH	87	CNTR $1.4 \times 10^8 < m < 10^{12}$ GeV
$< 3.2 \times 10^{-11}$	90	<sup>9</sup> NAKAMURA	85	CNTR $m > 1.5 \times 10^{-13}$ gram
$< 3.5 \times 10^{-11}$	90	<sup>10</sup> ULLMAN	81	CNTR Planck-mass $10^{19}$ GeV
$< 7. \times 10^{-11}$	90	<sup>10</sup> ULLMAN	81	CNTR $m \leq 10^{16}$ GeV

- <sup>1</sup> ADRIANI 15 search for relatively light quark matter with charge  $Z = 1-8$ . See their Figs. 2 and 3 for flux upper limits.
- <sup>2</sup> AMBROSIO 00B searched for quark matter (“nuclearites”) in the velocity range  $(10^{-5}-1) c$ . The listed limit is for  $2 \times 10^{-3} c$ .
- <sup>3</sup> ASTONE 93 searched for quark matter (“nuclearites”) in the velocity range  $(10^{-3}-1) c$ . Their Table 1 gives a compilation of searches for nuclearites.
- <sup>4</sup> AHLEN 92 searched for quark matter (“nuclearites”). The bound applies to velocity  $< 2.5 \times 10^{-3} c$ . See their Fig. 3 for other velocity/ $c$  and heavier mass range.
- <sup>5</sup> NAKAMURA 91 searched for quark matter in the velocity range  $(4 \times 10^{-5}-1) c$ .
- <sup>6</sup> ORITO 91 searched for quark matter. The limit is for the velocity range  $(10^{-4}-10^{-3}) c$ .
- <sup>7</sup> LIU 88 searched for quark matter (“nuclearites”) in the velocity range  $(2.5 \times 10^{-3}-1) c$ . A less stringent limit of  $5.8 \times 10^{-11}$  applies for  $(1-2.5) \times 10^{-3} c$ .
- <sup>8</sup> BARISH 87 searched for quark matter (“nuclearites”) in the velocity range  $(2.7 \times 10^{-4}-5 \times 10^{-3}) c$ .
- <sup>9</sup> NAKAMURA 85 at KEK searched for quark-matter. These might be lumps of strange quark matter with roughly equal numbers of  $u$ ,  $d$ ,  $s$  quarks. These lumps or nuclearites were assumed to have velocity of  $(10^{-4}-10^{-3}) c$ .
- <sup>10</sup> ULLMAN 81 is sensitive for heavy slow singly charge particle reaching earth with vertical velocity 100–350 km/s.

**Highly Ionizing Particle Flux**

<u>VALUE</u> ( $\text{m}^{-2}\text{yr}^{-1}$ )	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.4	95	0	KINOSHITA	81B PLAS	$Z/\beta$ 30–100

• • • We do not use the following data for averages, fits, limits, etc. • • •

**SEARCHES FOR BLACK HOLE PRODUCTION**

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • •	We do not use the following data for averages, fits, limits, etc. • • •		
1	AAD 15AN ATLS	8 TeV $pp \rightarrow$ multijets	
2	AAD 14A ATLS	8 TeV $pp \rightarrow \gamma + \text{jet}$	
3	AAD 14AL ATLS	8 TeV $pp \rightarrow \ell + \text{jet}$	
4	AAD 14C ATLS	8 TeV $pp \rightarrow \ell + (\ell \text{ or jets})$	
5	AAD 13D ATLS	7 TeV $pp \rightarrow 2 \text{ jets}$	
6	CHATRCHYAN 13A CMS	7 TeV $pp \rightarrow 2 \text{ jets}$	
7	CHATRCHYAN 13AD CMS	8 TeV $pp \rightarrow$ multijets	
8	AAD 12AK ATLS	7 TeV $pp \rightarrow \ell + (\ell \text{ or jets})$	
9	CHATRCHYAN 12W CMS	7 TeV $pp \rightarrow$ multijets	
10	AAD 11AG ATLS	7 TeV $pp \rightarrow 2 \text{ jets}$	

<sup>1</sup> AAD 15AN search for black hole or string ball formation followed by its decay to multijet final states, in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 20.3 \text{ fb}^{-1}$ . See their Figs. 6–8 for limits.

<sup>2</sup> AAD 14A search for quantum black hole formation followed by its decay to a  $\gamma$  and a jet, in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 20 \text{ fb}^{-1}$ . See their Fig. 3 for limits.

<sup>3</sup> AAD 14AL search for quantum black hole formation followed by its decay to a lepton and a jet, in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 20.3 \text{ fb}^{-1}$ . See their Fig. 2 for limits.

<sup>4</sup> AAD 14C search for microscopic (semiclassical) black hole formation followed by its decay to final states with a lepton and  $\geq 2$  (leptons or jets), in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 20.3 \text{ fb}^{-1}$ . See their Figures 8–11, Tables 7, 8 for limits.

<sup>5</sup> AAD 13D search for quantum black hole formation followed by its decay to two jets, in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.8 \text{ fb}^{-1}$ . See their Fig. 8 and Table 3 for limits.

<sup>6</sup> CHATRCHYAN 13A search for quantum black hole formation followed by its decay to two jets, in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 5 \text{ fb}^{-1}$ . See their Figs. 5 and 6 for limits.

<sup>7</sup> CHATRCHYAN 13AD search for microscopic (semiclassical) black hole formation followed by its evaporation to multiparticle final states, in multijet (including  $\gamma, \ell$ ) events in  $pp$  collisions at  $E_{\text{cm}} = 8 \text{ TeV}$  with  $L = 12 \text{ fb}^{-1}$ . See their Figs. 5–7 for limits.

<sup>8</sup> AAD 12AK search for microscopic (semiclassical) black hole formation followed by its decay to final states with a lepton and  $\geq 2$  (leptons or jets), in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 1.04 \text{ fb}^{-1}$ . See their Fig. 4 and 5 for limits.

<sup>9</sup> CHATRCHYAN 12W search for microscopic (semiclassical) black hole formation followed by its evaporation to multiparticle final states, in multijet (including  $\gamma, \ell$ ) events in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 4.7 \text{ fb}^{-1}$ . See their Figs. 5–8 for limits.

<sup>10</sup> AAD 11AG search for quantum black hole formation followed by its decay to two jets, in  $pp$  collisions at  $E_{\text{cm}} = 7 \text{ TeV}$  with  $L = 36 \text{ pb}^{-1}$ . See their Fig. 11 and Table 4 for limits.



## REFERENCES FOR Searches for WIMPs and Other Particles

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AAD	15AN	JHEP 1507 032	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15AS	EPJ C75 92	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15AT	EPJ C75 79	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15BH	EPJ C75 299	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15BJ	EPJ C75 362	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15CF	PRL 115 131801	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	15CS	PR D91 012008	G. Aad <i>et al.</i>	(ATLAS Collab.)
Also		PR D92 059903 (errata.)	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAIJ	15BD	EPJ C75 595	R. Aaij <i>et al.</i>	(LHCb Collab.)
AARTSEN	15C	EPJ C75 20	M.G. Aartsen <i>et al.</i>	(IceCube Collab.)
AARTSEN	15E	EPJ C75 492	M.G. Aartsen <i>et al.</i>	(IceCube Collab.)
ABRAMOWSKI	15	PRL 114 081301	A. Abramowski <i>et al.</i>	(H.E.S.S. Collab.)
ACKERMANN	15	PR D91 122002	M. Ackermann <i>et al.</i>	(Fermi-LAT Collab.)
ACKERMANN	15A	JCAP 1509 008	M. Ackermann <i>et al.</i>	(Fermi-LAT Collab.)
ACKERMANN	15B	PRL 115 231301	M. Ackermann <i>et al.</i>	(Fermi-LAT Collab.)
ADRIANI	15	PRL 115 111101	O. Adriani <i>et al.</i>	(PAMELA Collab.)
ADRIAN-MAR...	15	JCAP 1510 068	S. Adrian-Martinez <i>et al.</i>	(ANTARES Collab.)
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AGNESE	15	PRL 114 111302	R. Agnese <i>et al.</i>	(CDMS Collab.)
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APRILE	15	PRL 115 091302	E. Aprile <i>et al.</i>	(XENON Collab.)
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NAKAMURA	15	PTEP 2015 043F01	K. Nakamura <i>et al.</i>	(NEWAGE Collab.)
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AAD	14C	JHEP 1408 103	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	14K	PR D90 012004	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	14O	PRL 112 201802	G. Aad <i>et al.</i>	(ATLAS Collab.)
AALTONEN	14J	PR D89 092001	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ACKERMANN	14	PR D89 042001	M. Ackermann <i>et al.</i>	(Fermi-LAT Collab.)
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AAD	13D	JHEP 1301 029	G. Aad <i>et al.</i>	(ATLAS Collab.)
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ABRAMOWSKI	13	PRL 110 041301	A. Abramowski <i>et al.</i>	(H.E.S.S. Collab.)
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ABRAMOWSKI	12	APJ 750 123	A. Abramowski <i>et al.</i>	(H.E.S.S. Collab.)
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ALIU	12	PR D85 062001	E. Aliu <i>et al.</i>	(VERITAS Collab.)
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DAW	12	ASP 35 397	E. Daw <i>et al.</i>	(DRIFT-1Id Collab.)
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KIM	12	PRL 108 181301	S.C. Kim <i>et al.</i>	(KIMS Collab.)
AAD	11AG	NJP 13 053044	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11I	PL B698 353	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11S	PL B705 294	G. Aad <i>et al.</i>	(ATLAS Collab.)
AALSETH	11	PRL 106 131301	C.E. Aalseth <i>et al.</i>	(CoGeNT Collab.)
AALSETH	11A	PRL 107 141301	C.E. Aalseth <i>et al.</i>	(CoGeNT Collab.)
AALTONEN	11AF	PRL 107 181801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	11M	PRL 106 171801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	11I	PRL 107 011804	V. M. Abazov <i>et al.</i>	(D0 Collab.)
ABBASI	11C	PR D84 022004	R. Abbasi <i>et al.</i>	(IceCube Collab.)
ABRAMOWSKI	11	PRL 106 161301	A. Abramowski <i>et al.</i>	(H.E.S.S. Collab.)
ACKERMANN	11	PRL 107 241302	M. Ackermann <i>et al.</i>	(Fermi-LAT Collab.)
AHLEN	11	PL B695 124	S. Ahlen <i>et al.</i>	(DMTPC Collab.)
AHMED	11	PR D83 112002	Z. Ahmed <i>et al.</i>	(CDMS Collab.)
AHMED	11A	PR D84 011102	Z. Ahmed <i>et al.</i>	(CDMS and EDELWEISS Collabs.)
AHMED	11B	PRL 106 131302	Z. Ahmed <i>et al.</i>	(CDMS Collab.)
AJELLO	11	PR D84 032007	M. Ajello <i>et al.</i>	(Fermi-LAT Collab.)
ANGLE	11	PRL 107 051301	J. Angle <i>et al.</i>	(XENON10 Collab.)
Also		PRL 110 249901 (errat.)	J. Angle <i>et al.</i>	(XENON10 Collab.)
APRILE	11	PR D84 052003	E. Aprile <i>et al.</i>	(XENON100 Collab.)
APRILE	11A	PR D84 061101	E. Aprile <i>et al.</i>	(XENON100 Collab.)
APRILE	11B	PRL 107 131302	E. Aprile <i>et al.</i>	(XENON100 Collab.)
ARMENGAUD	11	PL B702 329	E. Armengaud <i>et al.</i>	(EDELWEISS II Collab.)
BEHNKE	11	PRL 106 021303	E. Behnke <i>et al.</i>	(COUPP Collab.)
CHATRCHYAN	11C	JHEP 1106 026	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11U	PRL 107 201804	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
GERINGER-SA...	11	PRL 107 241303	A. Geringer-Sameth, S.M. Koushiappas	
HORN	11	PL B705 471	M. Horn <i>et al.</i>	(ZEPLIN-III Collab.)
TANAKA	11	APJ 742 78	T. Tanaka <i>et al.</i>	(Super-Kamiokande Collab.)
AAD	10	PRL 105 161801	G. Aad <i>et al.</i>	(ATLAS Collab.)
AALTONEN	10AF	PR D82 052005	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABBASI	10	PR D81 057101	R. Abbasi <i>et al.</i>	(IceCube Collab.)
AHMED	10	SCI 327 1619	Z. Ahmed <i>et al.</i>	(CDMS II Collab.)
AKERIB	10	PR D82 122004	D.S. Akerib <i>et al.</i>	(CDMS-II Collab.)

AKIMOV	10	PL B692 180	D.Yu. Akimov <i>et al.</i>	(ZEPLIN-III Collab.)
APRILE	10	PRL 105 131302	E. Aprile <i>et al.</i>	(XENON100 Collab.)
ARMENGAUD	10	PL B687 294	E. Armengaud <i>et al.</i>	(EDELWEISS II Collab.)
FELIZARDO	10	PRL 105 211301	M. Felizardo <i>et al.</i>	(The SIMPLE Collab.)
KHACHATRYAN	10	PRL 105 211801	V. Khachatryan <i>et al.</i>	(CMS Collab.)
Also		PRL 106 029902	V. Khachatryan <i>et al.</i>	(CMS Collab.)
MIUCHI	10	PL B686 11	K. Miuchi <i>et al.</i>	(NEWAGE Collab.)
AALTONEN	09AF	PR D80 011102	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09G	PR D79 052004	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09Z	PRL 103 021802	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	09M	PRL 102 161802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBASI	09B	PRL 102 201302	R. Abbasi <i>et al.</i>	(IceCube Collab.)
AHMED	09	PRL 102 011301	Z. Ahmed <i>et al.</i>	(CDMS Collab.)
ANGLE	09	PR D80 115005	J. Angle <i>et al.</i>	(XENON10 Collab.)
ANGLOHER	09	ASP 31 270	G. Angloher <i>et al.</i>	(CRESST Collab.)
ARCHAMBAULT	09	PL B682 185	S. Archambault <i>et al.</i>	(PICASSO Collab.)
LEBEDENKO	09A	PRL 103 151302	V.N. Lebedenko <i>et al.</i>	(ZEPLIN-III Collab.)
LIN	09	PR D79 061101	S.T. Lin <i>et al.</i>	(TEXONO Collab.)
AALSETH	08	PRL 101 251301	C.E. Aalseth <i>et al.</i>	(CoGeNT Collab.)
Also		PRL 102 109903 (err.)	C.E. Aalseth <i>et al.</i>	(CoGeNT Collab.)
ANGLE	08A	PRL 101 091301	J. Angle <i>et al.</i>	(XENON10 Collab.)
BEDNYAKOV	08	PAN 71 111	V.A. Bednyakov, H.P. Klapdor-Kleingrothaus, I.V. Krivosheina	
		Translated from YAF 71 112.		
ALNER	07	PL B653 161	G.J. Alner <i>et al.</i>	(ZEPLIN-II Collab.)
LEE	07A	PRL 99 091301	H.S. Lee <i>et al.</i>	(KIMS Collab.)
MIUCHI	07	PL B654 58	K. Miuchi <i>et al.</i>	
AKERIB	06	PR D73 011102	D.S. Akerib <i>et al.</i>	(CDMS Collab.)
SHIMIZU	06A	PL B633 195	Y. Shimizu <i>et al.</i>	
AKERIB	05	PR D72 052009	D.S. Akerib <i>et al.</i>	(CDMS Collab.)
ALNER	05	PL B616 17	G.J. Alner <i>et al.</i>	(UK Dark Matter Collab.)
BARNABE-HEIDER	05	PL B624 186	M. Barnabe-Heider <i>et al.</i>	(PICASSO Collab.)
BENOIT	05	PL B616 25	A. Benoit <i>et al.</i>	(EDELWEISS Collab.)
GIRARD	05	PL B621 233	T.A. Girard <i>et al.</i>	(SIMPLE Collab.)
GIULIANI	05	PRL 95 101301	F. Giuliani	
GIULIANI	05A	PR D71 123503	F. Giuliani, T.A. Girard	
KLAPDOR-KLEINGROTHAUS	05	PL B609 226	H.V. Klapdor-Kleingrothaus, I.V. Krivosheina, C. Tomei	
AKTAS	04C	EPJ C36 413	A. Aktas <i>et al.</i>	(H1 Collab.)
GIULIANI	04	PL B588 151	F. Giuliani, T.A. Girard	
GIULIANI	04A	PRL 93 161301	F. Giuliani	
MIUCHI	03	ASP 19 135	K. Miuchi <i>et al.</i>	
TAKEDA	03	PL B572 145	A. Takeda <i>et al.</i>	
ANGLOHER	02	ASP 18 43	G. Angloher <i>et al.</i>	(CRESST Collab.)
BELLI	02	PR D66 043503	P. Belli <i>et al.</i>	
BERNABEI	02C	EPJ C23 61	R. Bernabei <i>et al.</i>	(DAMA Collab.)
GREEN	02	PR D66 083003	A.M. Green	
JAVORSEK	02	PR D65 072003	D. Javorsek II <i>et al.</i>	
BAUDIS	01	PR D63 022001	L. Baudis <i>et al.</i>	(Heidelberg-Moscow Collab.)
JAVORSEK	01	PR D64 012005	D. Javorsek II <i>et al.</i>	
JAVORSEK	01B	PRL 87 231804	D. Javorsek II <i>et al.</i>	
SMITH	01	PR D64 043502	D. Smith, N. Weiner	
ULLIO	01	JHEP 0107 044	P. Ullio, M. Kamionkowski, P. Vogel	
ABBIENDI	00D	EPJ C13 197	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
AMBROSIO	00B	EPJ C13 453	M. Ambrosio <i>et al.</i>	(MACRO Collab.)
BENOIT	00	PL B479 8	A. Benoit <i>et al.</i>	(EDELWEISS Collab.)
BERNABEI	00D	NJP 2 15	R. Bernabei <i>et al.</i>	(DAMA Collab.)
COLLAR	00	PRL 85 3083	J.I. Collar <i>et al.</i>	(SIMPLE Collab.)
ABE	99F	PRL 82 2038	F. Abe <i>et al.</i>	(CDF Collab.)
AMBROSIO	99	PR D60 082002	M. Ambrosio <i>et al.</i>	(Macro Collab.)
BERNABEI	99	PL B450 448	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BERNABEI	99D	PRL 83 4918	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BRHLIK	99	PL B464 303	M. Brhlik, L. Roszkowski	
DERBIN	99	PAN 62 1886	A.V. Derbin <i>et al.</i>	
		Translated from YAF 62 2034.		
ACKERSTAFF	98P	PL B433 195	K. Akerstaff <i>et al.</i>	(OPAL Collab.)
KLIMENKO	98	JETPL 67 875	A.A. Klimenko <i>et al.</i>	
		Translated from ZETFP 67 835.		
ABE	97G	PR D55 R5263	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	97D	PL B396 315	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACKERSTAFF	97B	PL B391 210	K. Akerstaff <i>et al.</i>	(OPAL Collab.)
ADAMS	97B	PRL 79 4083	J. Adams <i>et al.</i>	(FNAL KTeV Collab.)
BARATE	97K	PL B405 379	R. Barate <i>et al.</i>	(ALEPH Collab.)
SARSA	97	PR D56 1856	M.L. Sarsa <i>et al.</i>	(ZARA)

ALESSAND...	96	PL B384 316	A. Alessandrello <i>et al.</i>	(MILA, MILAI, SASSO)
BELLI	96	PL B387 222	P. Belli <i>et al.</i>	(DAMA Collab.)
Also		PL B389 783 (erratum)	P. Belli <i>et al.</i>	(DAMA Collab.)
BELLI	96C	NC 19C 537	P. Belli <i>et al.</i>	(DAMA Collab.)
BERNABEI	96	PL B389 757	R. Bernabei <i>et al.</i>	(DAMA Collab.)
COLLAR	96	PRL 76 331	J.I. Collar	(SCUC)
SARSA	96	PL B386 458	M.L. Sarsa <i>et al.</i>	(ZARA)
Also		PR D56 1856	M.L. Sarsa <i>et al.</i>	(ZARA)
SMITH	96	PL B379 299	P.F. Smith <i>et al.</i>	(RAL, SHEF, LOIC+)
SNOWDEN-...	96	PRL 76 332	D.P. Snowden-Ifft, E.S. Freeman, P.B. Price	(UCB)
AKERS	95R	ZPHY C67 203	R. Akers <i>et al.</i>	(OPAL Collab.)
GALLAS	95	PR D52 6	E. Gallas <i>et al.</i>	(MSU, FNAL, MIT, FLOR)
GARCIA	95	PR D51 1458	E. Garcia <i>et al.</i>	(ZARA, SCUC, PNL)
QUENBY	95	PL B351 70	J.J. Quenby <i>et al.</i>	(LOIC, RAL, SHEF+)
SNOWDEN-...	95	PRL 74 4133	D.P. Snowden-Ifft, E.S. Freeman, P.B. Price	(UCB)
Also		PRL 76 331	J.I. Collar	(SCUC)
Also		PRL 76 332	D.P. Snowden-Ifft, E.S. Freeman, P.B. Price	(UCB)
BECK	94	PL B336 141	M. Beck <i>et al.</i>	(MPIH, KIAE, SASSO)
RAM	94	PR D49 3120	S. Ram <i>et al.</i>	(TELA, TRIU)
ABE	93G	PRL 71 2542	F. Abe <i>et al.</i>	(CDF Collab.)
ASTONE	93	PR D47 4770	P. Astone <i>et al.</i>	(ROMA, ROMAI, CATA, FRAS)
BUSKULIC	93C	PL B303 198	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
YAMAGATA	93	PR D47 1231	T. Yamagata, Y. Takamori, H. Utsunomiya	(KONAN)
ABE	92J	PR D46 R1889	F. Abe <i>et al.</i>	(CDF Collab.)
AHLEN	92	PRL 69 1860	S.P. Ahlen <i>et al.</i>	(MACRO Collab.)
BACCI	92	PL B293 460	C. Bacci <i>et al.</i>	(Beijing-Roma-Saclay Collab.)
VERKERK	92	PRL 68 1116	P. Verkerk <i>et al.</i>	(ENSP, SACL, PAST)
AKESSON	91	ZPHY C52 219	T. Akesson <i>et al.</i>	(HELIOS Collab.)
NAKAMURA	91	PL B263 529	S. Nakamura <i>et al.</i>	
ORITO	91	PRL 66 1951	S. Orito <i>et al.</i>	(ICEPP, WASC, NIHO, ICRR)
REUSSER	91	PL B255 143	D. Reusser <i>et al.</i>	(NEUC, CIT, PSI)
ADACHI	90C	PL B244 352	I. Adachi <i>et al.</i>	(TOPAZ Collab.)
ADACHI	90E	PL B249 336	I. Adachi <i>et al.</i>	(TOPAZ Collab.)
AKRAWY	90O	PL B252 290	M.Z. Akrawy <i>et al.</i>	(OPAL Collab.)
HEMMICK	90	PR D41 2074	T.K. Hemmick <i>et al.</i>	(ROCH, MICH, OHIO+)
SAITO	90	PRL 65 2094	T. Saito <i>et al.</i>	(ICRR, KOBE)
NAKAMURA	89	PR D39 1261	T.T. Nakamura <i>et al.</i>	(KYOT, TMT)
NORMAN	89	PR D39 2499	E.B. Norman <i>et al.</i>	(LBL)
BERNSTEIN	88	PR D37 3103	R.M. Bernstein <i>et al.</i>	(STAN, WISC)
CALDWELL	88	PRL 61 510	D.O. Caldwell <i>et al.</i>	(UCSB, UCB, LBL)
LIU	88	PRL 61 271	G. Liu, B. Barish	
BARISH	87	PR D36 2641	B.C. Barish, G. Liu, C. Lane	(CIT)
NORMAN	87	PRL 58 1403	E.B. Norman, S.B. Gazes, D.A. Bennett	(LBL)
BADIER	86	ZPHY C31 21	J. Badier <i>et al.</i>	(NA3 Collab.)
MINCER	85	PR D32 541	A. Mincer <i>et al.</i>	(UMD, GMAS, NSF)
NAKAMURA	85	PL 161B 417	K. Nakamura <i>et al.</i>	(KEK, INUS)
THRON	85	PR D31 451	J.L. Thron <i>et al.</i>	(YALE, FNAL, IOWA)
SAKUYAMA	83B	LNC 37 17	H. Sakuyama, N. Suzuki	(MEIS)
Also		LNC 36 389	H. Sakuyama, K. Watanabe	(MEIS)
Also		NC 78A 147	H. Sakuyama, K. Watanabe	(MEIS)
Also		NC 6C 371	H. Sakuyama, K. Watanabe	(MEIS)
BHAT	82	PR D25 2820	P.N. Bhat <i>et al.</i>	(TATA)
KINOSHITA	82	PRL 48 77	K. Kinoshita, P.B. Price, D. Fryberger	(UCB+)
MARINI	82	PR D26 1777	A. Marini <i>et al.</i>	(FRAS, LBL, NWES, STAN+)
SMITH	82B	NP B206 333	P.F. Smith <i>et al.</i>	(RAL)
KINOSHITA	81B	PR D24 1707	K. Kinoshita, P.B. Price	(UCB)
LOSECCO	81	PL 102B 209	J.M. LoSecco <i>et al.</i>	(MICH, PENN, BNL)
ULLMAN	81	PRL 47 289	J.D. Ullman	(LEHM, BNL)
YOCK	81	PR D23 1207	P.C.M. Yock	(AUCK)
BARTEL	80	ZPHY C6 295	W. Bartel <i>et al.</i>	(JADE Collab.)
BUSSIERE	80	NP B174 1	A. Bussiere <i>et al.</i>	(BGNA, SACL, LAPP)
YOCK	80	PR D22 61	P.C.M. Yock	(AUCK)
ARMITAGE	79	NP B150 87	J.C.M. Armitage <i>et al.</i>	(CERN, DARE, FOM+)
BOZZOLI	79	NP B159 363	W. Bozzoli <i>et al.</i>	(BGNA, LAPP, SACL+)
GOODMAN	79	PR D19 2572	J.A. Goodman <i>et al.</i>	(UMD)
SMITH	79	NP B149 525	P.F. Smith, J.R.J. Bennett	(RHEL)
BHAT	78	PRAM 10 115	P.N. Bhat, P.V. Ramana Murthy	(TATA)
CARROLL	78	PRL 41 777	A.S. Carroll <i>et al.</i>	(BNL, PRIN)
CUTTS	78	PRL 41 363	D. Cutts <i>et al.</i>	(BROW, FNAL, ILL, BARI+)
VIDAL	78	PL 77B 344	R.A. Vidal <i>et al.</i>	(COLU, FNAL, STON+)

ALEKSEEV	76	SJNP 22 531	G.D. Alekseev <i>et al.</i>	(JINR)
		Translated from YAF 22	1021.	
ALEKSEEV	76B	SJNP 23 633	G.D. Alekseev <i>et al.</i>	(JINR)
		Translated from YAF 23	1190.	
BALDIN	76	SJNP 22 264	B.Y. Baldin <i>et al.</i>	(JINR)
		Translated from YAF 22	512.	
BRIATORE	76	NC 31A 553	L. Briatore <i>et al.</i>	(LCGT, FRAS, FREIB)
GUSTAFSON	76	PRL 37 474	H.R. Gustafson <i>et al.</i>	(MICH)
ALBROW	75	NP B97 189	M.G. Albrow <i>et al.</i>	(CERN, DARE, FOM+)
FRANKEL	75	PR D12 2561	S. Frankel <i>et al.</i>	(PENN, FNAL)
JOVANOVI...	75	PL 56B 105	J.V. Jovanovich <i>et al.</i>	(MANI, AACH, CERN+)
YOCK	75	NP B86 216	P.C.M. Yock	(AUCK, SLAC)
APPEL	74	PRL 32 428	J.A. Appel <i>et al.</i>	(COLU, FNAL)
FRANKEL	74	PR D9 1932	S. Frankel <i>et al.</i>	(PENN, FNAL)
YOCK	74	NP B76 175	P.C.M. Yock	(AUCK)
ALPER	73	PL 46B 265	B. Alper <i>et al.</i>	(CERN, LIVP, LUND, BOHR+)
LEIPUNER	73	PRL 31 1226	L.B. Leipuner <i>et al.</i>	(BNL, YALE)
DARDO	72	NC 9A 319	M. Dardo <i>et al.</i>	(TORI)
TONWAR	72	JP A5 569	S.C. Tonwar, S. Naranan, B.V. Sreekantan	(TATA)
ANTIPOV	71B	NP B31 235	Y.M. Antipov <i>et al.</i>	(SERP)
ANTIPOV	71C	PL 34B 164	Y.M. Antipov <i>et al.</i>	(SERP)
BINON	69	PL 30B 510	F.G. Binon <i>et al.</i>	(SERP)
BJORNBOE	68	NC B53 241	J. Bjornboe <i>et al.</i>	(BOHR, TATA, BERN+)
JONES	67	PR 164 1584	L.W. Jones	(MICH, WISC, LBL, UCLA, MINN+)
DORFAN	65	PRL 14 999	D.E. Dorfman <i>et al.</i>	(COLU)

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