## Neutrino Mixing

With the exception of the short-baseline anomalies such as LSND, current neutrino data can be described within the framework of a  $3\times 3$  mixing matrix between the flavor eigenstates  $\nu_e$ ,  $\nu_{\mu}$ , and  $\nu_{\tau}$  and the mass eigenstates  $\nu_1$ ,  $\nu_2$ , and  $\nu_3$ . (See Eq. (14.6) of the review "Neutrino Mass, Mixing, and Oscillations" by K. Nakamura and S.T. Petcov.) The Listings are divided into the following sections:

(A) Neutrino fluxes and event ratios: shows measurements which correspond to various oscillation tests for Accelerator, Reactor, Atmospheric, and Solar neutrino experiments. Typically ratios involve a measurement in a realm sensitive to oscillations compared to one for which no oscillation effect is expected.

(B) Three neutrino mixing parameters: shows measurements of  $\sin^2(\theta_{12})$ ,  $\sin^2(\theta_{23})$ ,  $\Delta m_{21}^2$ ,  $\Delta m_{32}^2$ ,  $\sin^2(\theta_{13})$  and  $\delta_{CP}$  which are all interpretations of data based on the three neutrino mixing scheme described in the review "Neutrino Mass, Mixing, and Oscillations." by K. Nakamura and S.T. Petcov. Many parameters have been calculated in the two-neutrino approximation.

(C) Other neutrino mixing results: shows measurements and limits for the probability of oscillation for experiments which might be relevant to the LSND oscillation claim. Included are experiments which are sensitive to  $\nu_{\mu} \rightarrow \nu_{e}, \bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}$ , eV sterile neutrinos, and CPT tests.

#### (A) Neutrino fluxes and event ratios

#### Events (observed/expected) from accelerator $\nu_{\mu}$ experiments.

Some neutrino oscillation experiments compare the flux in two or more detectors. This is usually quoted as the ratio of the event rate in the far detector to the expected rate based on an extrapolation from the near detector in the absence of oscillations.

VALUE	DOCUMENT ID		TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the following	owing data for ave	erages,	, fits, lim	nits, etc. • • •
$1.01 \pm 0.10$	<sup>1</sup> ABE	<b>1</b> 4B	T2K	$\nu_{e}$ rate in T2K near detect.
$0.71 \pm 0.08$	<sup>2</sup> AHN	06A	K2K	K2K to Super-K
$0.64 \pm 0.05$	<sup>3</sup> MICHAEL	06	MINS	All charged current events
$0.71\substack{+0.08\\-0.09}$	<sup>4</sup> ALIU	05	K2K	KEK to Super-K
$0.70 \substack{+ 0.10 \\ - 0.11}$	<sup>5</sup> AHN	03	K2K	KEK to Super-K

<sup>1</sup> The rate of  $\nu_e$  from  $\mu$  decay was measured to be 0.68  $\pm$  0.30 compared to the predicted flux. From *K* decay 1.10  $\pm$  0.14 compared to the predicted flux. <sup>2</sup> Based on the observation of 112 events when  $158.1^{+9.2}_{-8.6}$  were expected without os-

<sup>2</sup>Based on the observation of 112 events when  $158.1^{+9.2}_{-8.6}$  were expected without oscillations. Including not only the number of events but also the shape of the energy distribution, the evidence for oscillation is at the level of about 4.3  $\sigma$ . Supersedes ALIU 05.

<sup>3</sup> This ratio is based on the observation of 215 events compared to an expectation of 336  $\pm$  14 without oscillations. See also ADAMSON 08. <sup>4</sup> This ratio is based on the observation of 107 events at the far detector 250 km away

<sup>4</sup>This ratio is based on the observation of 107 events at the far detector 250 km away from KEK, and an expectation of  $151^{+12}_{-10}$ .

<sup>5</sup> This ratio is based on the observation of 56 events with an expectation of 80.1 + 6.2

#### Events (observed/expected) from reactor $\overline{\nu}_e$ experiments.

The quoted values are the ratios of the measured reactor  $\overline{\nu}_e$  event rate at the quoted distances, and the rate expected without oscillations. The expected rate is based on the experimental data for the most significant reactor fuels ( $^{235}$ U,  $^{239}$ Pu,  $^{241}$ Pu) and on calculations for  $^{238}$ U.

A recent re-evaluation of the spectral conversion of electron to  $\overline{\nu}_e$  in MUELLER 11 results in an upward shift of the reactor  $\overline{\nu}_e$  spectrum by 3% and, thus, might require revisions to the ratios listed in this table.

VALUE	DOCUMENT ID		TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the	e following data for	r avera	ges, fits	, limits, etc. • • •
$0.949 \pm 0.0028$	<sup>1</sup> AN	17A	DAYA	DayaBay, Ling Ao/Ao II reactors
	<sup>2</sup> AN	16	DAYA	DayaBay, Ling Ao/Ao II reactors
$0.944 \!\pm\! 0.007 \!\pm\! 0.003$	<sup>3</sup> AN	13	DAYA	DayaBay, Ling Ao/Ao II reactors
$0.944 \pm 0.016 \ \pm 0.040$	<sup>4</sup> ABE	12	DCHZ	Chooz reactors
$0.920 \pm 0.009 \ \pm 0.014$	<sup>5</sup> AHN	12	RENO	Yonggwang reactors
$0.940 \pm 0.011 \ \pm 0.004$	<sup>6</sup> AN	12	DAYA	DayaBay, LIng Ao/Ao II reactors
$1.08\ \pm 0.21\ \pm 0.16$	<sup>7</sup> DENIZ	10	TEXO	Kuo-Sheng reactor, 28 m
$0.658 \!\pm\! 0.044 \ \pm 0.047$	<sup>8</sup> ARAKI	05	KLND	Japanese react. $\sim 180$ km
$0.611 \!\pm\! 0.085 \ \pm 0.041$	<sup>9</sup> EGUCHI	03	KLND	Japanese react. $\sim 180$ km
$1.01 \ \pm 0.024 \ \pm 0.053$	<sup>10</sup> военм	01		Palo Verde react. 0.75–0.89 km

$1.01 \ \pm 0.028 \ \pm 0.027$		IOZ Chooz reactors 1 km
$0.987 \pm 0.006 \pm 0.037$	<sup>12</sup> GREENWOOD 96	Savannah River, 18.2 m
$0.988 \pm 0.004 \ \pm 0.05$	ACHKAR 95 CN	ITR Bugey reactor, 15 m
$0.994 \pm 0.010 \ \pm 0.05$	ACHKAR 95 CN	ITR Bugey reactor, 40 m
$0.915 \!\pm\! 0.132 \ \pm 0.05$		ITR Bugey reactor, 95 m
$0.987 \pm 0.014 \pm 0.027$		ITR Bugey reactor, 15 m
$0.985 \!\pm\! 0.018 \ \pm 0.034$	KUVSHINN 91 CN	ITR Rovno reactor
$1.05 \pm 0.02 \pm 0.05$	VUILLEUMIER82	Gösgen reactor
$0.955 \pm 0.035 \ \pm 0.110$	<sup>14</sup> KWON 81	$\overline{\nu}_e p \rightarrow e^+ n$
$0.89 \ \pm 0.15$	<sup>14</sup> BOEHM 80	$\overline{\nu}_e p \rightarrow e^+ n$

<sup>1</sup> AN 17A use 1230 days of data of the Daya Bay experiment. The reported deficit is based on the ratio of the  $\overline{\nu}_e$  signal observed by the far detectors versus the near detectors and correcting for differences in detector exposure. Supersedes AN 13.

 $^2$  AN 16 use 217 days of data (338k events) to determine the neutrino flux ratio relative to the prediction of Mueller-Huber and ILL-Vogel models (see AN 16 for details). The reported flux ratios were corrected for  $\theta_{13}$  oscillation effect. The flux measurement is consistent with results from previous short-baseline reactor experiments. The measured inverse beta decay yield is (1.55  $\pm$  0.04)  $\times$  10 $^{-18}$  cm $^2/({\rm GW}$  day) or  $\sigma_f$  = (5.92  $\pm$ 

 $0.14)\times 10^{-43}~{\rm cm}^2/{\rm fission}.$  About  $4\sigma$  excess of events was observed in the 4–6 MeV prompt energy region.

- <sup>3</sup> AN 13 use six identical detectors, with three placed near the reactor cores (flux-weighted baselines of 470 and 576 m) and the remaining three at the far hall (at the flux averaged distance of 1648 m from all six reactor cores) to determine the mixing angle  $\theta_{13}$  using the  $\overline{\nu}_e$  observed interaction rate ratios. This rate-only analysis excludes the no-oscillation hypothesis at 7.7 standard deviations. The value of  $\Delta m_{31}^2 = 2.32 \times 10^{-3} \text{ eV}^2$  was assumed in the analysis. This is an improved result (2.5 times increase in statistics) compared to AN 12. Superseded by AN 17A.
- <sup>4</sup> ABE 12 determine the  $\overline{\nu}_e$  interaction rate in a single detector, located 1050 m from the cores of two reactors. The rate normalization is fixed by the results of the Bugey4 reactor experiment, thus avoiding any dependence on possible very short baseline oscillations.
- <sup>5</sup> AHN 12 use two identical detectors, placed at flux weighted distances of 408.56 m and 1433.99m from six reactor cores, to determine the  $\overline{\nu}_e$  interaction rate ratio.
- <sup>6</sup> AN 12 use six identical detectors with three placed near the reactor cores (flux-weighted baselines of 470 m and 576 m) and the remaining three at the far hall (at the flux averaged distance of 1648 m from all six reactor cores) to determine the  $\overline{\nu}_e$  interaction rate ratios. Superseded by AN 13.
- <sup>7</sup> DENIZ 10 observe reactor  $\overline{\nu}_e e$  scattering with recoil kinetic energies 3–8 MeV using CsI(TI) detectors. The observed rate is consistent with the Standard Model prediction, leading to a constraint on  $\sin^2 \theta_{W} = 0.251 \pm 0.031(\text{stat}) \pm 0.024(\text{sys})$ .
- <sup>8</sup> Updated result of KamLAND, including the data used in EGUCHI 03. Note that the survival probabilities for different periods are not directly comparable because the effective baseline varies with power output of the reactor sources involved, and there were large variations in the reactor power production in Japan in 2003.
- $^9\,\text{EGUCHI}$  03 observe reactor neutrino disappearance at  $\sim\,$  180 km baseline to various Japanese nuclear power reactors.
- $^{10}$  BOEHM 01 search for neutrino oscillations at 0.75 and 0.89 km distance from the Palo Verde reactors.
- <sup>11</sup> APOLLONIO 99, APOLLONIO 98 search for neutrino oscillations at 1.1 km fixed distance from Chooz reactors. They use  $\overline{\nu}_e p \rightarrow e^+ n$  in Gd-loaded scintillator target. APOLLONIO 99 supersedes APOLLONIO 98. See also APOLLONIO 03 for detailed description.
- $^{12}\,\mathrm{GREENWOOD}$  96 search for neutrino oscillations at 18 m and 24 m from the reactor at Savannah River.

- <sup>13</sup>DECLAIS 94 result based on integral measurement of neutrons only. Result is ratio of measured cross section to that expected in standard V-A theory. Replaced by ACHKAR 95.
- $^{14}\,\rm KWON$  81 represents an analysis of a larger set of data from the same experiment as BOEHM 80.

#### — Atmospheric neutrinos ——

Neutrinos and antineutrinos produced in the atmosphere induce  $\mu$ -like and e-like events in underground detectors. The ratio of the numbers of the two kinds of events is defined as  $\mu/e$ . It has the advantage that systematic effects, such as flux uncertainty, tend to cancel, for both experimental and theoretical values of the ratio. The "ratio of the ratios" of experimental to theoretical  $\mu/e$ ,  $R(\mu/e)$ , or that of experimental to theoretical  $\mu/total$ ,  $R(\mu/total)$  with total =  $\mu+e$ , is reported below. If the actual value is not unity, the value obtained in a given experiment may depend on the experimental conditions. In addition, the measured "up-down asymmetry" for  $\mu$  (N<sub>up</sub>( $\mu$ )/N<sub>down</sub>( $\mu$ )) or e (N<sub>up</sub>(e)/N<sub>down</sub>(e)) is reported. The expected "up-down asymmetry" is nearly unity if there is no neutrino oscillation.

#### $R(\mu/e) = (Measured Ratio \mu/e) / (Expected Ratio \mu/e)$

VALUE	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following	g data for averages	s, fits,	limits, e	etc. • • •
$0.658\!\pm\!0.016\!\pm\!0.035$	<sup>1</sup> ASHIE	05	SKAM	sub-GeV
$0.702^{+0.032}_{-0.030}{\pm}0.101$	<sup>2</sup> ASHIE	05	SKAM	multi-GeV
$0.69\ \pm 0.10\ \pm 0.06$	<sup>3</sup> SANCHEZ <sup>4</sup> FUKUDA	03 96в		Calorimeter raw data Water Cherenkov
$1.00 \ \pm 0.15 \ \pm 0.08$	<sup>5</sup> DAUM	95	FREJ	Calorimeter
$0.60 \begin{array}{c} +0.06 \\ -0.05 \end{array} \pm 0.05$	<sup>6</sup> FUKUDA	94	KAMI	sub-GeV
$0.57 \begin{array}{c} +0.08 \\ -0.07 \end{array} \pm 0.07$	<sup>7</sup> FUKUDA	94	KAMI	multi-Gev
	<sup>8</sup> BECKER-SZ	<b>92</b> B	IMB	Water Cherenkov

- <sup>1</sup>ASHIE 05 results are based on an exposure of 92 kton yr during the complete Super-Kamiokande I running period. The analyzed data sample consists of fully-contained single-ring *e*-like events with 0.1 GeV/c <  $p_e$  and  $\mu$ -like events 0.2 GeV/c <  $p_{\mu}$ , both having a visible energy < 1.33 GeV. These criteria match the definition used by FUKUDA 94.
- <sup>2</sup> ASHIE 05 results are based on an exposure of 92 kton yr during the complete Super-Kamiokande I running period. The analyzed data sample consists of fully-contained single-ring events with visible energy > 1.33 GeV and partially-contained events. All partially-contained events are classified as  $\mu$ -like.
- <sup>3</sup>SANCHEZ 03 result is based on an exposure of 5.9 kton yr, and updates ALLISON 99 result. The analyzed data sample consists of fully-contained *e*-flavor and  $\mu$ -flavor events having lepton momentum > 0.3 GeV/c.

<sup>4</sup> FUKUDA 96B studied neutron background in the atmospheric neutrino sample observed in the Kamiokande detector. No evidence for the background contamination was found.

<sup>5</sup> DAUM 95 results are based on an exposure of 2.0 kton yr which includes the data used by BERGER 90B. This ratio is for the contained and semicontained events. DAUM 95 also report  $R(\mu/e) = 0.99 \pm 0.13 \pm 0.08$  for the total neutrino induced data sample which includes upward going stopping muons and horizontal muons in addition to the contained and semicontained events.

- <sup>6</sup> FUKUDA 94 result is based on an exposure of 7.7 kton yr and updates the HIRATA 92 result. The analyzed data sample consists of fully-contained *e*-like events with 0.1 <  $p_e < 1.33 \text{ GeV}/c$  and fully-contained  $\mu$ -like events with 0.2 <  $p_{\mu} < 1.5 \text{ GeV}/c$ .
- <sup>7</sup> FUKUDA 94 analyzed the data sample consisting of fully contained events with visible energy > 1.33 GeV and partially contained  $\mu$ -like events.
- <sup>8</sup> BECKER-SZENDY 92B reports the fraction of nonshowering events (mostly muons from atmospheric neutrinos) as  $0.36 \pm 0.02 \pm 0.02$ , as compared with expected fraction  $0.51 \pm 0.01 \pm 0.05$ . After cutting the energy range to the Kamiokande limits, BEIER 92 finds  $R(\mu/e)$  very close to the Kamiokande value.

#### $R(\nu_{\mu}) = (Measured Flux of \nu_{\mu}) / (Expected Flux of \nu_{\mu})$

	DOCUMENT ID		TECN	COMMENT			
VALUE	DOCUMENT ID		TECN	COMMENT			
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$							
$0.84 \pm 0.12$	<sup>1</sup> ADAMSON	06	MINS	MINOS atmospheric			
$0.72\!\pm\!0.026\!\pm\!0.13$	<sup>2</sup> AMBROSIO	01	MCRO	upward through-going			
$0.57 \pm 0.05 \pm 0.15$	<sup>3</sup> AMBROSIO	00	MCRO	upgoing partially contained			
$0.71 \pm 0.05 \pm 0.19$	<sup>4</sup> AMBROSIO	00	MCRO	downgoing partially contained			
	_			+ upgoing stopping			
$0.74\!\pm\!0.036\!\pm\!0.046$	<sup>5</sup> AMBROSIO	98	MCRO	Streamer tubes			
	<sup>6</sup> CASPER	91	IMB	Water Cherenkov			
	<sup>7</sup> AGLIETTA	89	NUSX				
$0.95 \pm 0.22$	<sup>8</sup> BOLIEV	81		Baksan			
$0.62 \pm 0.17$	CROUCH	78		Case Western/UCI			

 $^1$  ADAMSON 06 uses a measurement of 107 total neutrinos compared to an expected rate of 127  $\pm$  13 without oscillations.

- $^2$  AMBROSIO 01 result is based on the upward through-going muon tracks with  $E_{\mu} > 1$  GeV. The data came from three different detector configurations, but the statistics is largely dominated by the full detector run, from May 1994 to December 2000. The total live time, normalized to the full detector configuration, is 6.17 years. The first error is the statistical error, the second is the systematic error, dominated by the theoretical error in the predicted flux.
- <sup>3</sup> AMBROSIO 00 result is based on the upgoing partially contained event sample. It came from 4.1 live years of data taking with the full detector, from April 1994 to February 1999. The average energy of atmospheric muon neutrinos corresponding to this sample is 4 GeV. The first error is statistical, the second is the systematic error, dominated by the 25% theoretical error in the rate (20% in the flux and 15% in the cross section, added in quadrature). Within statistics, the observed deficit is uniform over the zenith angle.
- <sup>4</sup> AMBROSIO 00 result is based on the combined samples of downgoing partially contained events and upgoing stopping events. These two subsamples could not be distinguished due to the lack of timing information. The result came from 4.1 live years of data taking with the full detector, from April 1994 to February 1999. The average energy of atmospheric muon neutrinos corresponding to this sample is 4 GeV. The first error is statistical, the second is the systematic error, dominated by the 25% theoretical error in the rate (20% in the flux and 15% in the cross section, added in quadrature). Within statistics, the observed deficit is uniform over the zenith angle.
- <sup>5</sup> AMBROSIO 98 result is for all nadir angles and updates AHLEN 95 result. The lower cutoff on the muon energy is 1 GeV. In addition to the statistical and systematic errors, there is a Monte Carlo flux error (theoretical error) of  $\pm 0.13$ . With a neutrino oscillation hypothesis, the fit either to the flux or zenith distribution independently yields  $\sin^2 2\theta = 1.0$  and  $\Delta(m^2) \sim$  a few times  $10^{-3} \text{ eV}^2$ . However, the fit to the observed zenith distribution gives a maximum probability for  $\chi^2$  of only 5% for the best oscillation hypothesis.
- <sup>6</sup> CASPER 91 correlates showering/nonshowering signature of single-ring events with parent atmospheric-neutrino flavor. They find nonshowering ( $\approx \nu_{\mu}$  induced) fraction is  $0.41 \pm 0.03 \pm 0.02$ , as compared with expected  $0.51 \pm 0.05$  (syst).

- <sup>7</sup>AGLIETTA 89 finds no evidence for any anomaly in the neutrino flux. They define  $\rho$  = (measured number of  $\nu_e$ 's)/(measured number of  $\nu_\mu$ 's). They report  $\rho$ (measured)= $\rho$ (expected) = 0.96 $\stackrel{+0.32}{-0.28}$
- $^8\,{\rm From}$  this data BOLIEV 81 obtain the limit  $\Delta(m^2)~\leq~6\times 10^{-3}~{\rm eV}^2$  for maximal mixing,  $u_{\mu} \not\rightarrow 
  u_{\mu}$  type oscillation.

#### $R(\mu/total) = (Measured Ratio \mu/total) / (Expected Ratio \mu/total)$ VALUE

DOCUMENT ID TECN COMMENT

• • • We do not use the following data for averages, fits, limits, etc. • • •

 $1.1^{+0.07}_{-0.12}{\pm}0.11$ <sup>1</sup> CLARK 97 IMB multi-GeV

<sup>1</sup> CLARK 97 obtained this result by an analysis of fully contained and partially contained events in the IMB water-Cherenkov detector with visible energy > 0.95 GeV.

#### $N_{\rm up}(\mu)/N_{\rm down}(\mu)$

VALUE	DOCUMENT ID		TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the following	g data for averages	, fits,	limits, e	etc. • • •
$0.71 \pm 0.06$	<sup>1</sup> ADAMSON	12B	MINS	contained-vertex muons
$0.551^{+0.035}_{-0.033}{\pm}0.004$	<sup>2</sup> ASHIE	05	SKAM	multi-GeV

<sup>1</sup>ADAMSON 12B reports the atmospheric neutrino results obtained with MINOS far detector in 2,553 live days (an exposure of 37.9 kton yr). This result is obtained with a sample of high resolution contained-vertex muons. The quoted error is statistical only.

<sup>2</sup>ASHIE 05 results are based on an exposure of 92 kton yr during the complete Super-Kamiokande I running period. The analyzed data sample consists of fully-contained single-ring  $\mu$ -like events with visible energy > 1.33 GeV and partially-contained events. All partially-contained events are classified as µ-like. Upward-going events are those with  $-1 < \cos(\text{zenith angle}) < -0.2$  and downward-going events are those with 0.2 <cos(zenith angle) < 1. The  $\mu$ -like up-down ratio for the multi-GeV data deviates from 1 (the expectation for no atmospheric  $u_{\mu}$  oscillations) by more than 12 standard deviations.

#### $N_{\rm up}(e)/N_{\rm down}(e)$

VALUE	DOCUMENT ID		TECN COMMENT	
$\bullet$ $\bullet$ $\bullet$ We do not use the following	ng data for average	s, fits,	, limits, etc. • • •	
$0.961^{+0.086}_{-0.079}{\pm}0.016$	<sup>1</sup> ASHIE	05	SKAM multi-GeV	

<sup>1</sup>ASHIE 05 results are based on an exposure of 92 kton yr during the complete Super-Kamiokande I running period. The analyzed data sample consists of fully-contained single-ring e-like events with visible energy > 1.33 GeV. Upward-going events are those with  $-1 < \cos(\text{zenith angle}) < -0.2$  and downward-going events are those with 0.2  $<\cos( ext{zenith angle}) < 1$ . The e-like up-down ratio for the multi-GeV data is consistent with 1 (the expectation for no atmospheric  $\nu_e$  oscillations).

#### $R(up/down; \mu) = (Measured up/down; \mu) / (Expected up/down; \mu)$

VALUE	DOCUMENT ID		TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the	e following data for	averag	ges, fits,	limits, etc. • • •
$0.62\!\pm\!0.05\!\pm\!0.02$	<sup>1</sup> ADAMSON	12B	MINS	contained-vertex muons
$0.62^{+0.19}_{-0.14}{\pm}0.02$	<sup>2</sup> ADAMSON	06	MINS	atmospheric $\nu$ with far detector

- <sup>1</sup> ADAMSON 12B reports the atmospheric neutrino results obtained with MINOS far detector in 2,553 live days (an exposure of 37.9 kton·yr). This result is obtained with a sample of high resolution contained-vertex muons. The expected ratio is calculated with no neutrino oscillation.
- <sup>2</sup>ADAMSON 06 result is obtained with the MINOS far detector with an exposure of 4.54 kton yr. The expected ratio is calculated with no neutrino oscillation.

#### $N(\mu^+)/N(\mu^-)$

VALUE	DOCUMENT ID			COMMENT
$\bullet$ $\bullet$ $\bullet$ We do not use the following data for averages, fits,			limits, o	etc. • • •
$0.46 \substack{+ \ 0.05 \\ - \ 0.04}$	<sup>1,2</sup> ADAMSON	12B	MINS	contained-vertex muons
$0.63^{+0.09}_{-0.08}$	$^{1,3}$ ADAMSON	12B	MINS	u-induced rock-muons

- <sup>1</sup> ADAMSON 12B reports the atmospheric neutrino results obtained with MINOS far detector in 2,553 live days (an exposure of 37.9 kton·yr). The muon charge ratio  $N(\mu^+)/N(\mu^-)$  represents the  $\overline{\nu}_{\mu}/\nu_{\mu}$  ratio.
- <sup>2</sup> This result is obtained with a charge-separated sample of high resolution contained-vertex muons. The quoted error is statistical only.
- <sup>3</sup> This result is obtained with a charge-separated sample of high resolution neutrino-induced rock-muons. The quoted error is statistical only.

$R(\mu^+/\mu^-) = (Measured)$	$N(\mu^+)/N(\mu^-)) /$	′ (Expe	cted N( $\mu^+$ )/N( $\mu^-$ ))
VALUE	DOCUMENT ID	TECN	COMMENT

• • • We do not use t	he following data for	avera	ges, fits,	limits, etc. • • •
$0.93\!\pm\!0.09\!\pm\!0.09$	<sup>1,2</sup> ADAMSON	12B	MINS	contained-vertex muons
$1.29^{+0.19}_{-0.17}{\pm}0.16$	<sup>1,3</sup> ADAMSON	12B	MINS	u-induced rock-muons
$1.03\!\pm\!0.08\!\pm\!0.08$	<sup>1,4</sup> ADAMSON	12B	MINS	contained
$1.39^{+0.35+0.08}_{-0.46-0.14}$	<sup>5</sup> ADAMSON	07	MINS	Upward and horizontal $\mu$ with far detector
$0.96^{+0.38}_{-0.27}{\pm}0.15$	<sup>6</sup> ADAMSON	06	MINS	atmospheric $\nu$ with far detector

<sup>1</sup> ADAMSON 12B reports the atmospheric neutrino results obtained with MINOS far detector in 2,553 live days (an exposure of 37.9 kton·yr). The muon charge ratio N( $\mu^+$ )/N( $\mu^-$ ) represents the  $\overline{\nu}_{\mu}/\nu_{\mu}$  ratio. As far as the same oscillation parameters are used for  $\nu$ s and  $\overline{\nu}$ s, the expected  $\overline{\nu}_{\mu}/\nu_{\mu}$  ratio is almost entirely independent of any input oscillations.

- $^2$  This result is obtained with a charge-separated sample of high resolution contained-vertex  $_{2}$  muons.
- <sup>3</sup> This result is obtained with a charge-separated sample of high resolution neutrino-induced rock-muons.
- <sup>4</sup> The charge-separated samples of high resolution contained-vertex muons and neutrinoinduced rock-muons are combined to obtain this result which is consistent with unity.
- <sup>5</sup> ADAMSON 07 result is obtained with the MINOS far detector in 854.24 live days, based on neutrino-induced upward-going and horizontal muons. This result is consistent with *CPT* conservation.
- <sup>6</sup> ADAMSON 06 result is obtained with the MINOS far detector with an exposure of 4.54 kton yr, based on contained events. The expected ratio is calculated by assuming the same oscillation parameters for neutrinos and antineutrinos.

— Solar neutrinos —

Solar neutrinos are produced by thermonuclear fusion reactions in the Sun. Radiochemical experiments measure particular combinations of fluxes from various neutrino-producing reactions, whereas water-Cherenkov experiments mainly measure a flux of neutrinos from decay of <sup>8</sup>B. Solar neutrino fluxes are composed of all active neutrino species,  $\nu_e$ ,  $\nu_{\mu}$ , and  $\nu_{\tau}$ . In addition, some other mechanisms may cause antineutrino components in solar neutrino fluxes. Each measurement method is sensitive to a particular component or a combination of components of solar neutrino fluxes. For details, see Section 13.4 of Reviews, Tables, and Plots.

#### $\nu_e$ Capture Rates from Radiochemical Experiments

1 SNU (Solar Neutrino Unit) =  $10^{-36}$  captures per atom per second.

VALUE (SNU)	DOCUMENT ID		TECN	COMMENT
• • • We do not use the	e following data fo	r aver	ages, fits	s, limits, etc. • • •
$73.4 \begin{array}{c} +6.1 \\ -6.0 \end{array} \begin{array}{c} +3.7 \\ -4.1 \end{array}$	<sup>1</sup> KAETHER	10		GALX reanalysis
$67.6 \pm 4.0 \pm 3.2$	<sup>2</sup> KAETHER	10		GNO+GALX reanalysis combined
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	<sup>3</sup> ABDURASHI	. 09	SAGE	$^{71}\text{Ga} \rightarrow ~^{71}\text{Ge}$
62.9 $^{+5.5}_{-5.3}$ $\pm 2.5$	<sup>4</sup> ALTMANN	05	GNO	$^{71}\text{Ga} \rightarrow ~^{71}\text{Ge}$
$69.3 \pm 4.1 \pm 3.6$	<sup>5</sup> ALTMANN	05	GNO	GNO+GALX combined
77.5 $\pm 6.2 \ +4.3 \\ -4.7$	<sup>6</sup> HAMPEL	99	GALX	$71_{Ga} \rightarrow 71_{Ge}$
$2.56\!\pm\!0.16\!\pm\!0.16$	<sup>7</sup> CLEVELAND	98	HOME	$37_{Cl} \rightarrow 37_{Ar}$

<sup>1</sup> KAETHER 10 reports the reanalysis results of a complete GALLEX data (GALLEX I+II+III+IV, reported in HAMPEL 99) based on the event selection with a new pulse shape analysis, which provides a better background reduction than the rise time analysis adopted in HAMPEL 99.

<sup>2</sup> Combined result of GALLEX I+II+III+IV reanalysis and GNO I+II+III (ALTMANN 05).

- <sup>3</sup> ABDURASHITOV 09 reports a combined analysis of 168 extractions of the SAGE solar neutrino experiment during the period January 1990 through December 2007, and updates the ABDURASHITOV 02 result. The data are consistent with the assumption that the solar neutrino production rate is constant in time. Note that a  $\sim 15\%$  systematic uncertainty in the overall normalization may be added to the ABDURASHITOV 09 result, because calibration experiments for gallium solar neutrino measurements using intense  $^{51}$ Cr (twice by GALLEX and once by SAGE) and  $^{37}$ Ar (by SAGE) result in an average ratio of 0.87  $\pm$  0.05 of the observed to calculated rates.
- <sup>4</sup> ALTMANN 05 reports the complete result from the GNO solar neutrino experiment (GNO I+II+III), which is the successor project of GALLEX. Experimental technique of GNO is essentially the same as that of GALLEX. The run data cover the period 20 May 1998 through 9 April 2003.
- <sup>5</sup> Combined result of GALLEX I+II+III+IV (HAMPEL 99) and GNO I+II+III.
- <sup>6</sup> HAMPEL 99 report the combined result for GALLEX I+II+III+IV (65 runs in total), which update the HAMPEL 96 result. The GALLEX IV result (12 runs) is 118.4  $\pm$  17.8  $\pm$  6.6 SNU. (HAMPEL 99 discuss the consistency of partial results with the mean.) The GALLEX experimental program has been completed with these runs. The total run data cover the period 14 May 1991 through 23 January 1997. A total of 300 <sup>71</sup>Ge events were observed. Note that a  $\sim$  15% systematic uncertainty in the overall normalization may be added to the HAMPEL 99 result, because calibration experiments for gallium solar neutrino measurements using intense <sup>51</sup>Cr (twice by GALLEX and once by SAGE)

and <sup>37</sup>Ar (by SAGE) result in an average ratio of  $0.87 \pm 0.05$  of the observed to calculated rates. <sup>7</sup>CLEVELAND 98 is a detailed report of the <sup>37</sup>Cl experiment at the Homestake Mine.

CLEVELAND 98 is a detailed report of the <sup>57</sup>Cl experiment at the Homestake Mine. The average solar neutrino-induced <sup>37</sup>Ar production rate from 108 runs between 1970 and 1994 updates the DAVIS 89 result.

## φ<sub>ES</sub> (<sup>8</sup>B)

 $^8{\rm B}$  solar-neutrino flux measured via  $\nu\,e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross-section difference,  $\sigma(\nu_{\mu,\tau}\,e)\sim 0.16\sigma(\nu_e\,e)$ . If the  $^8{\rm B}$  solar-neutrino flux involves nonelectron flavor active neutrinos, their contribution to the flux is  $\sim 0.16$  times of  $\nu_e$ .

<u>VALUE (10<sup>6</sup> cm<math>^{-2}</math>s<math>^{-1}</math>)</u>	DOCUMENT ID		TECN	COMMENT
• • • We do not use t	he following data fo	r aver	ages, fits	s, limits, etc. • • •
$2.345\!\pm\!0.014\!\pm\!0.036$	<sup>1</sup> ABE	<b>16</b> C	SKAM	SK-I+II+III+IV average flux
$2.308 \!\pm\! 0.020 \!+\! 0.039 \\ - 0.040$	<sup>2</sup> ABE	16C	SKAM	SK-IV average flux
$2.250^{+0.030}_{-0.029}{\pm}0.038$	<sup>2</sup> ABE	16C	SKAM	SK-IV day flux
$2.364 \!\pm\! 0.029 \!\pm\! 0.040$	<sup>2</sup> ABE	<b>16</b> C	SKAM	SK-IV night flux
$2.404 \pm 0.039 \pm 0.053$	<sup>3</sup> ABE	<b>16</b> C	SKAM	SK-III average flux
$2.41 \ \pm 0.05 \ \begin{array}{c} +0.16 \\ -0.15 \end{array}$	<sup>4</sup> ABE	11	SKAM	SK-II average flux
$2.38\ \pm 0.02\ \pm 0.08$	<sup>5</sup> ABE	11	SKAM	SK-I average flux
$2.77\ \pm 0.26\ \pm 0.32$	<sup>6</sup> ABE	<b>11</b> B	KLND	average flux
$2.4 \pm 0.4 \pm 0.1$	<sup>7</sup> BELLINI	10A	BORX	average flux
$1.77 \begin{array}{r} +0.24 \\ -0.21 \end{array} \begin{array}{r} +0.09 \\ -0.10 \end{array}$	<sup>8</sup> AHARMIM	08	SNO	Phase III
$2.38 \ \pm 0.05 \ {}^{+0.16}_{-0.15}$	<sup>9</sup> CRAVENS	08	SKAM	average flux
$2.35\ \pm 0.02\ \pm 0.08$	<sup>10</sup> HOSAKA	06	SKAM	average flux
$2.35 \ \pm 0.22 \ \pm 0.15$	<sup>11</sup> AHARMIM	05A	SNO	Salty D <sub>2</sub> O; <sup>8</sup> B shape not con- strained
$2.34 \ \pm 0.23 \ +0.15 \\ -0.14$	<sup>11</sup> AHARMIM	05A	SNO	Salty D <sub>2</sub> O; <sup>8</sup> B shape con- strained
$2.39 \begin{array}{c} +0.24 \\ -0.23 \end{array} \pm 0.12$	<sup>12</sup> AHMAD	02	SNO	average flux
$2.39 \ \pm 0.34 \ {}^{+0.16}_{-0.14}$	<sup>13</sup> AHMAD	01	SNO	average flux
$2.80\ \pm 0.19\ \pm 0.33$	<sup>14</sup> FUKUDA	96	KAMI	average flux
$2.70 \pm 0.27$	<sup>14</sup> FUKUDA	96	KAMI	day flux
$2.87 \begin{array}{c} +0.27 \\ -0.26 \end{array}$	<sup>14</sup> FUKUDA	96	KAMI	night flux

<sup>1</sup>ABE 16C reports the combined results of the four phases of the Super-Kamiokande average flux measurements. Here the revised Super-Kamiokande-III result is used.

<sup>2</sup> ABE 16C reports the Super-Kamiokande-IV results for 1664 live days from September 2008 to February 2014. The analysis threshold is total electron energy of 4.0 MeV.

<sup>3</sup> ABE 16C revised the Super-Kamiokande-III average flux value reported in ABE 11. Super-Kamiokande-III results are for 548 live days from August 4, 2006 to August 18, 2008. The analysis threshold is 5.0 MeV, but the event sample in the 5.0–6.5 MeV total electron energy range has a total live time of 298 days.

<sup>4</sup>ABE 11 recalculated the Super-Kamiokande-II results using <sup>8</sup>B spectrum of WIN-\_ TER 06A.

<sup>5</sup> ABE 11 recalculated the Super-Kamiokande-I results using <sup>8</sup>B spectrum of WINTER 06A.

- <sup>6</sup>ABE 11B use a 123 kton day exposure of the KamLAND liquid scintillation detector to measure the  $^8$ B solar neutrino flux. They utilize u - e elastic scattering above a reconstructed-energy threshold of 5.5 MeV, corresponding to 5 MeV electron recoil energy. 299 electron recoil candidate events are reported, of which 157  $\pm$  23.6 are assigned to background.
- <sup>7</sup>BELLINI 10A reports the Borexino result with 3 MeV energy threshold for scattered electrons. The data correspond to 345.3 live days with a target mass of 100 t, between July 15, 2007 and August 23, 2009.
- <sup>8</sup>AHARMIM 08 reports the results from SNO Phase III measurement using an array of <sup>3</sup>He proportional counters to measure the rate of NC interactions in heavy water, over the period between November 27, 2004 and November 28, 2006, corresponding to 385.17 live days. A simultaneous fit was made for the number of NC events detected by the proportional counters and the numbers of NC, CC, and ES events detected by the PMTs, where the spectral distributions of the ES and CC events were not constrained to the  ${}^{8}B$ shape.
- 9 CRAVENS 08 reports the Super-Kamiokande-II results for 791 live days from December 2002 to October 2005. The photocathode coverage of the detector is 19% (reduced from 40% of that of Super-Kamiokande-I due to an accident in 2001). The analysis threshold for the average flux is 7 MeV.
- $^{10}$  HOSAKA 06 reports the final results for 1496 live days with Super-Kamiokande-I between May 31, 1996 and July 15, 2001, and replace FUKUDA 02 results. The analysis threshold is 5 MeV except for the first 280 live days (6.5 MeV).
- $^{11}$ AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, corresponding to 391.4 live days, and update AHMED 04A. The CC, ES, and NC events were statistically separated. In one method, the <sup>8</sup>B energy spectrum was not constrained. In the other method, the constraint of an undistorted <sup>8</sup>B energy spectrum was added for comparison with AHMAD 02 results  $^{12}$  AHMAD 02 reports the  $^{8}$ B solar-neutrino flux measured via  $\nu e$  elastic scattering above
- the kinetic energy threshold of 5 MeV. The data correspond to 306.4 live days with SNO between November 2, 1999 and May 28, 2001, and updates AHMAD 01 results.
- $^{13}$ AHMAD 01 reports the  $^8$ B solar-neutrino flux measured via  $\nu e$  elastic scattering above the kinetic energy threshold of 6.75 MeV. The data correspond to 241 live days with SNO between November 2, 1999 and January 15, 2001.
- $^{14}$  FUKUDA 96 results are for a total of 2079 live days with Kamiokande II and III from January 1987 through February 1995, covering the entire solar cycle 22, with threshold  ${\sf E}_e > 9.3\,{\sf MeV}$  (first 449 days),  $> 7.5\,{\sf MeV}$  (middle 794 days), and  $> 7.0\,{\sf MeV}$  (last 836 days). These results update the HIRATA 90 result for the average <sup>8</sup>B solar-neutrino flux and HIRATA 91 result for the day-night variation in the <sup>8</sup>B solar-neutrino flux. The total data sample was also analyzed for short-term variations: within experimental errors, no strong correlation of the solar-neutrino flux with the sunspot numbers was found.

 $\phi_{CC}$  (<sup>8</sup>B) <sup>8</sup>B solar-neutrino flux measured with charged-current reaction which is sensitive ex-

$VALUE (10^{6} \text{ cm}^{-2} \text{s}^{-1})$	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following	data for averages	s, fits,	limits, e	etc. • • •
$1.67 \substack{+0.05 + 0.07 \\ -0.04 - 0.08}$	<sup>1</sup> AHARMIM	08	SNO	Phase III
$1.68 \!\pm\! 0.06 \! \substack{+ 0.08 \\ - 0.09}$	<sup>2</sup> AHARMIM	05A	SNO	Salty D <sub>2</sub> O; <sup>8</sup> B shape
$1.72\!\pm\!0.05\!\pm\!0.11$	<sup>2</sup> AHARMIM	05A	SNO	not const. Salty D <sub>2</sub> O; <sup>8</sup> B shape constrained
$1.76 \substack{+0.06 \\ -0.05} \pm 0.09$	<sup>3</sup> AHMAD	02	SNO	average flux
$1.75 \pm 0.07 \substack{+0.12 \\ -0.11} \pm 0.05$	<sup>4</sup> AHMAD	01	SNO	average flux
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- <sup>1</sup> AHARMIM 08 reports the results from SNO Phase III measurement using an array of <sup>3</sup>He proportional counters to measure the rate of NC interactions in heavy water, over the period between November 27, 2004 and November 28, 2006, corresponding to 385.17 live days. A simultaneous fit was made for the number of NC events detected by the proportional counters and the numbers of NC, CC, and ES events detected by the PMTs, where the spectral distributions of the ES and CC events were not constrained to the <sup>8</sup>B shape.
- <sup>2</sup> AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, corresponding to 391.4 live days, and update AHMED 04A. The *CC*, *ES*, and *NC* events were statistically separated. In one method, the <sup>8</sup>B energy spectrum was not constrained. In the other method, the constraint of an undistorted <sup>8</sup>B energy spectrum was added for comparison with AHMAD 02 results.
   <sup>3</sup> AHMAD 02 reports the SNO result of the <sup>8</sup>B solar-neutrino flux measured with charged-
- <sup>3</sup> AHMAD 02 reports the SNO result of the <sup>8</sup>B solar-neutrino flux measured with chargedcurrent reaction on deuterium,  $\nu_e d \rightarrow ppe^-$ , above the kinetic energy threshold of 5 MeV. The data correspond to 306.4 live days with SNO between November 2, 1999 and May 28, 2001, and updates AHMAD 01 results. The complete description of the SNO Phase I data set is given in AHARMIM 07.
- <sup>4</sup> AHMAD 01 reports the first SNO result of the <sup>8</sup>B solar-neutrino flux measured with the charged-current reaction on deuterium,  $\nu_e d \rightarrow ppe^-$ , above the kinetic energy threshold of 6.75 MeV. The data correspond to 241 live days with SNO between November 2, 1999 and January 15, 2001.

## φ<sub>NC</sub> (<sup>8</sup>B)

 $^{8}\text{B}$  solar neutrino flux measured with neutral-current reaction, which is equally sensitive to  $\nu_{e},\,\nu_{\mu}$ , and  $\nu_{\tau}.$ 

VALUE $(10^{6} \text{ cm}^{-2} \text{s}^{-1})$	DOCUMENT ID		TECN	COMMENT
$\bullet$ $\bullet$ $\bullet$ We do not use the	s, limits, etc. • • •			
$5.25 \ \pm 0.16 \ +0.11 \\ -0.13$	<sup>1</sup> AHARMIM	13	SNO	All three phases combined
$5.140 \substack{+ 0.160 + 0.132 \\ - 0.158 - 0.117}$	<sup>2</sup> AHARMIM	10	SNO	Phase I+II, low threshold
$5.54 \begin{array}{c} +0.33 \\ -0.31 \end{array} \begin{array}{c} +0.36 \\ -0.34 \end{array}$	<sup>3</sup> AHARMIM	08	SNO	Phase III, prop. counter $+ PMT$
$\begin{array}{rrr} 4.94 \ \pm 0.21 \ \begin{array}{r} +0.38 \\ -0.34 \end{array}$	<sup>4</sup> AHARMIM	05A	SNO	Salty D <sub>2</sub> O; <sup>8</sup> B shape not const.
$\begin{array}{rrr} 4.81 \ \pm 0.19 \ \begin{array}{r} +0.28 \\ -0.27 \end{array}$	<sup>4</sup> AHARMIM	05A	SNO	Salty $D_2O$ ; <sup>8</sup> B shape constrained
$5.09 \begin{array}{c} +0.44 \\ -0.43 \end{array} \begin{array}{c} +0.46 \\ -0.43 \end{array}$	<sup>5</sup> AHMAD	02	SNO	average flux; <sup>8</sup> B shape const.
$6.42 \ \pm 1.57 \ \begin{array}{c} +0.55 \\ -0.58 \end{array}$	<sup>5</sup> AHMAD	02	SNO	average flux; <sup>8</sup> B shape not const.

<sup>1</sup> AHARMIM 13 obtained this result from a combined analysis of the data from all three phases, SNO-I, II, and III. The measurement of the <sup>8</sup>B flux mostly comes from the NC signal, however, CC contribution is included in the fit.

<sup>2</sup>AHARMIM 10 reports this result from a joint analysis of SNO Phase I+II data with the "effective electron kinetic energy" threshold of 3.5 MeV. This result is obtained with a "binned-histogram unconstrained fit" where binned probability distribution functions of the neutrino signal observables were used without any model constraints on the shape of the neutrino spectrum.

<sup>3</sup> AHARMIM 08 reports the results from SNO Phase III measurement using an array of <sup>3</sup>He proportional counters to measure the rate of NC interactions in heavy water, over the period between November 27, 2004 and November 28, 2006, corresponding to 385.17 live days. A simultaneous fit was made for the number of NC events detected by the proportional counters and the numbers of NC, CC, and ES events detected by the PMTs,

where the spectral distributions of the ES and CC events were not constrained to the  ${}^{8}\text{B}$  shape.

- <sup>4</sup> AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, corresponding to 391.4 live days, and update AHMED 04A. The *CC*, *ES*, and *NC* events were statistically separated. In one method, the <sup>8</sup>B energy spectrum was not constrained. In the other method, the constraint of an undistorted <sup>8</sup>B energy spectrum was added for comparison with AHMAD 02 results.
- <sup>5</sup>AHMAD 02 reports the first SNO result of the <sup>8</sup>B solar-neutrino flux measured with the neutral-current reaction on deuterium,  $\nu_{\ell} d \rightarrow n p \nu_{\ell}$ , above the neutral-current reaction threshold of 2.2 MeV. The data correspond to 306.4 live days with SNO between November 2, 1999 and May 28, 2001. The complete description of the SNO Phase I data set is given in AHARMIM 07.

### $\phi_{\nu_{\mu}+\nu_{\tau}}$ (<sup>8</sup>B)

Nonelectron-flavor active neutrino component ( $\nu_{\mu}$  and  $\nu_{\tau})$  in the  $^{8}{\rm B}$  solar-neutrino flux.

$VALUE (10^{6} \text{ cm}^{-2} \text{s}^{-1})$	DOCUMENT ID		TECN	COMMENT
$\bullet$ $\bullet$ $\bullet$ We do not use the	following data for	avera	ges, fits,	limits, etc. • • •
$3.26 \!\pm\! 0.25 \!+\! 0.40 \\ -0.35$	<sup>1</sup> AHARMIM	05A	SNO	From $\phi_{NC}$ , $\phi_{CC}$ , and $\phi_{ES}$ ; <sup>8</sup> B shape not const.
$3.09 \!\pm\! 0.22 \! \substack{+0.30 \\ -0.27}$	<sup>1</sup> AHARMIM	05A	SNO	From $\phi_{NC}$ , $\phi_{CC}$ , and $\phi_{ES}$ ; <sup>8</sup> B shape constrained
$3.41 \pm 0.45 {+0.48 \atop -0.45}$	<sup>2</sup> AHMAD	02	SNO	From $\phi_{NC}$ , $\phi_{CC}$ , and $\phi_{ES}$
$3.69 {\pm} 1.13$	<sup>3</sup> AHMAD	01		Derived from SNO+SuperKam, water Cherenkov

<sup>1</sup> AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, corresponding to 391.4 live days, and update AHMED 04A. The *CC*, *ES*, and *NC* events were statistically separated. In one method, the <sup>8</sup>B energy spectrum was not constrained. In the other method, the constraint of an undistorted <sup>8</sup>B energy spectrum was added for comparison with AHMAD 02 results.

 $^2$ AHMAD 02 deduced the nonelectron-flavor active neutrino component ( $u_{\mu}$  and  $u_{ au}$ )

in the <sup>8</sup>B solar-neutrino flux, by combining the charged-current result, the  $\nu e$  elastic-scattering result and the neutral-current result. The complete description of the SNO Phase I data set is given in AHARMIM 07.

<sup>3</sup>AHMAD 01 deduced the nonelectron-flavor active neutrino component ( $u_{\mu}$  and  $u_{ au}$ ) in

the <sup>8</sup>B solar-neutrino flux, by combining the SNO charged-current result (AHMAD 01) and the Super-Kamiokande  $\nu e$  elastic-scattering result (FUKUDA 01).

#### Total Flux of Active <sup>8</sup>B Solar Neutrinos

Total flux of active neutrinos ( $\nu_{\rm e},\,\nu_{\mu}$  , and  $\nu_{\tau}$  ).

$VALUE (10^6 \text{ cm}^{-2} \text{s}^{-1})$	DOCUMENT IL	)	TECN	COMMENT		
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$						
$5.25\ \pm 0.16\ +0.11\\-0.13$	<sup>1</sup> AHARMIM	13	SNO	All three phases combined		
$5.046\substack{+0.159+0.107\\-0.152-0.123}$	<sup>2</sup> AHARMIM	10	SNO	From $\phi_{NC}$ in Phase I+II, low threshold		
$5.54 \begin{array}{c} +0.33 \\ -0.31 \end{array} \begin{array}{c} +0.36 \\ -0.34 \end{array}$				$\phi_{NC}$ in Phase III		
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$\begin{array}{rrr} \text{4.94} \hspace{0.2cm} \pm 0.21 \hspace{0.2cm} \begin{array}{c} +0.38 \\ -0.34 \end{array}$	<sup>4</sup> AHARMIM	05A SNO	From $\phi_{NC}$ ; <sup>8</sup> B shape not const.
$\begin{array}{rrr} \textbf{4.81} \hspace{0.2cm} \pm 0.19 \hspace{0.2cm} \begin{array}{c} +0.28 \\ -0.27 \end{array}$	<sup>4</sup> AHARMIM	05A SNO	From $\phi_{NC}$ ; <sup>8</sup> B shape constrained
$5.09 \begin{array}{c} +0.44 \\ -0.43 \end{array} \begin{array}{c} +0.46 \\ -0.43 \end{array}$	<sup>5</sup> AHMAD	02 SNO	Direct measurement from $\phi_{\it NC}$
5.44 ±0.99	<sup>6</sup> AHMAD	01	Derived from SNO+SuperKam, water Cherenkov

<sup>1</sup> AHARMIM 13 obtained this result from a combined analysis of the data from all three phases, SNO-I, II, and III. The measurement of the <sup>8</sup>B flux mostly comes from the NC signal, however, CC contribution is included in the fit.

 $^2$  AHARMIM 10 reports this result from a joint analysis of SNO Phase I+II data with the "effective electron kinetic energy" threshold of 3.5 MeV. This result is obtained with the assumption of unitarity, which relates the NC, CC, and ES rates. The data were fit with the free parameters directly describing the total  $^8{\rm B}$  neutrino flux and the energy-dependent  $\nu_e$  survival probability.

<sup>3</sup> AHARMIM 08 reports the results from SNO Phase III measurement using an array of <sup>3</sup>He proportional counters to measure the rate of NC interactions in heavy water, over the period between November 27, 2004 and November 28, 2006, corresponding to 385.17 live days. A simultaneous fit was made for the number of NC events detected by the proportional counters and the numbers of NC, CC, and ES events detected by the PMTs, where the spectral distributions of the ES and CC events were not constrained to the <sup>8</sup>B shape.

- <sup>4</sup> AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, corresponding to 391.4 live days, and update AHMED 04A. The CC, ES, and NC events were statistically separated. In one method, the <sup>8</sup>B energy spectrum was not constrained. In the other method, the constraint of an undistorted <sup>8</sup>B energy spectrum was added for comparison with AHMAD 02 results.
- 5 AHMAD 02 determined the total flux of active <sup>8</sup>B solar neutrinos by directly measuring the neutral-current reaction,  $\nu_{\ell}d \rightarrow np\nu_{\ell}$ , which is equally sensitive to  $\nu_e$ ,  $\nu_{\mu}$ , and  $\nu_{\tau}$ . The complete description of the SNO Phase I data set is given in AHARMIM 07.
- <sup>6</sup> AHMAD 01 deduced the total flux of active <sup>8</sup>B solar neutrinos by combining the SNO charged-current result (AHMAD 01) and the Super-Kamiokande  $\nu e$  elastic-scattering result (FUKUDA 01).

#### **Day-Night Asymmetry (<sup>8</sup>B)** $A = (\phi_{\text{night}} - \phi_{\text{day}}) / \phi_{\text{average}}$

γ (Ψnight Ψda	y) / vaverage			
VALUE	DOCUMENT ID		TECN	COMMENT
$0.033 \pm 0.010 \pm 0.005$	<sup>1</sup> ABE	16C	SKAM	SK combined; Based on $\phi_{ES}$
$\bullet \bullet \bullet$ We do not use the fo	ollowing data for a	verage	es, fits, l	imits, etc. • • •
$0.036\!\pm\!0.016\!\pm\!0.006$	<sup>2</sup> ABE	<b>16</b> C	SKAM	SK-IV; Based on $\phi_{ES}$
$0.032\!\pm\!0.011\!\pm\!0.005$	<sup>3</sup> RENSHAW			Based on $\phi_{ES}$
$0.063\!\pm\!0.042\!\pm\!0.037$	<sup>4</sup> CRAVENS	08	SKAM	Based on $\phi_{ES}^{-\sim}$
$0.021 \!\pm\! 0.020 \!+\! 0.012 \\ -0.013$	<sup>5</sup> HOSAKA	06	SKAM	Based on $\phi_{ES}$
$0.017 \!\pm\! 0.016 \! \substack{+0.012 \\ -0.013}$	<sup>6</sup> HOSAKA	06	SKAM	Fitted in the LMA region
$-0.056\!\pm\!0.074\!\pm\!0.053$	<sup>7</sup> AHARMIM	05A	SNO	From salty SNO $\phi_{CC}$
$-0.037 \pm 0.063 \pm 0.032$	<sup>7</sup> AHARMIM	05A	SNO	From salty SNO $\phi_{CC}^{-}$ ; const. of no $\phi_{NC}^{-}$ asymmetry
$0.14 \ \pm 0.063 {+0.015 \atop -0.014}$	<sup>8</sup> AHMAD	<b>02</b> B	SNO	Derived from SNO $\phi_{CC}$
$0.07 \ \pm 0.049 {+0.013 \atop -0.012}$	<sup>9</sup> AHMAD	<b>02</b> B	SNO	Const. of no $\phi_{\it NC}$ asymmetry
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- <sup>1</sup>ABE 16C reports the combined day-night flux asymmetry results of the four phases of the Super-Kamiokande measurements. Amplitude fit method is used. See footnote to RENSHAW 14.
- <sup>2</sup> ABE 16C reports the Super-Kamiokande-IV results for 1664 live days from September 2008 to February 2014. The analysis threshold for day-night flux asymmetry is recoil electron energy of 4.49 MeV (total electron energy of 5.0 MeV). Amplitude fit method is used. See footnote to RENSHAW 14.
- $^3$  RENSHAW 14 obtains this result by using the "amplitude fit" introduced in SMY 04. The data from the Super-Kamiokande(SK)-I, -II, -III, and 1306 live days of the SK-IV measurements are used. The analysis threshold is recoil-electron kinetic energy of 4.5 MeV for SK-III, and SK-IV except for 250 live days in SK-III (6.0 MeV). The analysis threshold for SK-I and SK-IV except for 250 live days in SK-III (6.0 MeV). The analysis threshold for SK-I and SK-II is the same as in the previous reports. (Note that in the previous SK solar-neutrino results, the analysis threshold is quoted as recoil-electron total energy.) This day-night asymmetry result is consistent with neutrino oscillations for  $4\times10^{-5}~{\rm eV}^2~<\Delta m_{21}^2~<7\times10^{-5}~{\rm eV}^2$  and large mixing values of  $\theta_{12}$  at the 68% CL.
- <sup>4</sup> CRAVENS 08 reports the Super-Kamiokande-II results for 791 live days from December 2002 to October 2005. The photocathode coverage of the detector is 19% (reduced from 40% of that of Super-Kamiokande-I due to an accident in 2001). The analysis threshold for the day and night fluxes is 7.5 MeV except for the first 159 live days (8.0 MeV).
- <sup>5</sup> HOSAKA 06 reports the final results for 1496 live days with Super-Kamiokande-I between May 31, 1996 and July 15, 2001, and replace FUKUDA 02 results. The analysis threshold is 5 MeV except for the first 280 live days (6.5 MeV).
- <sup>6</sup> This result with reduced statistical uncertainty is obtained by assuming two-neutrino oscillations within the LMA (large mixing angle) region and by fitting the time variation of the solar neutrino flux measured via  $\nu_e$  elastic scattering to the variations expected from neutrino oscillations. For details, see SMY 04. There is an additional small systematic error of  $\pm 0.0004$  coming from uncertainty of oscillation parameters.
- <sup>7</sup> AHARMIM 05A measurements were made with dissolved NaCl (0.195% by weight) in heavy water over the period between July 26, 2001 and August 28, 2003, with 176.5 days of the live time recorded during the day and 214.9 days during the night. This result is obtained with the spectral distribution of the CC events not constrained to the <sup>8</sup>B shape.
- <sup>8</sup> AHMAD 02B results are based on the charged-current interactions recorded between November 2, 1999 and May 28, 2001, with the day and night live times of 128.5 and 177.9 days, respectively. The complete description of the SNO Phase I data set is given in AHARMIM 07.
- <sup>9</sup>AHMAD 02B results are derived from the charged-current interactions, neutral-current interactions, and  $\nu e$  elastic scattering, with the total flux of active neutrinos constrained to have no asymmetry. The data were recorded between November 2, 1999 and May 28, 2001, with the day and night live times of 128.5 and 177.9 days, respectively. The complete description of the SNO Phase I data set is given in AHARMIM 07.

## φ**ES** (<sup>7</sup>Be)

<sup>7</sup>Be solar-neutrino flux measured via  $\nu_e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross-section difference,  $\sigma(\nu_{\mu,\tau}e) \sim 0.2 \ \sigma(\nu_e e)$ . If the <sup>7</sup>Be solar-neutrino flux involves nonelectron flavor active neutrinos, their contribution to the flux is  $\sim 0.2$  times that of  $\nu_e$ .

$VALUE (10^9 \text{ cm}^{-2} \text{ s}^{-1})$	DOCUMENT ID		TECN	COMMENT
• • • We do not use the following	g data for average	s, fits,	limits, e	etc. • • •
$3.26 \pm 0.52$ $3.10 \pm 0.15$	<sup>1</sup> GANDO <sup>2</sup> BELLINI			average flux average flux

- <sup>1</sup>GANDO 15 uses 165.4 kton day exposure of the KamLAND liquid scintillator detector to measure the 862 keV <sup>7</sup>Be solar neutrino flux via  $\nu e$  elastic scattering
- <sup>2</sup> BELLINI 11A reports the <sup>7</sup>Be solar neutrino flux measured via  $\nu e$  elastic scattering. The data correspond to 740.7 live days between May 16, 2007 and May 8, 2010, and also correspond to 153.6 ton-year fiducial exposure. BELLINI 11A measured the 862 keV <sup>7</sup>Be solar neutrino flux, which is an 89.6% branch of the <sup>7</sup>Be solar neutrino flux, to be  $(2.78 \pm 0.13) \times 10^9$  cm<sup>-2</sup> s<sup>-1</sup>. Supercedes ARPESELLA 08A.

#### $\phi_{ES} (pep)$

pep solar-neutrino flux measured via  $\nu_e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross section difference,  $\sigma(\nu_{\mu,\tau}\ e)\sim~0.2\ \sigma(\nu_e\ e)$ . If the pep solar-neutrino flux involves non-electron flavor active neutrinos, their contribution to the flux is  $\sim~0.2$  times that of  $\nu_e$ .

VALUE $(10^8 \text{ cm}^{-2} \text{s}^{-1})$	DOCUMENT ID		TECN	COMMENT	
• • • We do not use the follow	ing data for averages	s, fits, l	limits, e	etc. • • •	
$1.0 \pm 0.2$	<sup>1</sup> BELLINI	12A	BORX	average flux	

<sup>1</sup> BELLINI 12A reports 1.44 MeV *pep* solar-neutrino flux measured via  $\nu_e$  elastic scattering. The data were collected between January 13, 2008 and May 9, 2010, corresponding to 20,4009 ton·day fiducial exposure. The listed flux value is calculated from the observed rate of *pep* solar neutrino interactions in Borexino (3.1 ± 0.6 ± 0.3 counts/(day·100 ton)) and the corresponding rate expected for no neutrino flavor oscillations (4.47 ± 0.05 counts/(day·100 ton)), using the SSM prediction for the *pep* solar neutrino flux of (1.441 ± 0.012) × 10<sup>8</sup> cm<sup>-2</sup>s<sup>-1</sup>.

#### $\phi_{ES}$ (CNO)

CNO solar-neutrino flux measured via  $\nu_e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross section difference,  $\sigma(\nu_{\mu,\tau} \ e) \sim 0.2 \ \sigma(\nu_e e)$ . If the CNO solar-neutrino flux involves non-electron flavor active neutrinos, their contribution to the flux is  $\sim 0.2$  times that of  $\nu_e$ .

VALUE ( $10^8 \text{ cm}^{-2} \text{s}^{-2}$	<sup>-1</sup> ) <u>CL%</u>	DOCUMENT IL	)	TECN	COMMENT
• • • We do no	t use the foll	owing data for a	verages,	fits, lim	nits, etc. • • •
<7.7	90	<sup>1</sup> BELLINI	12A	BORX	MSW-LMA solution assumed

<sup>1</sup> BELLINI 12A reports an upper limit of the CNO solar neutrino flux measured via  $\nu_e$  elastic scattering. The data were collected between January 13, 2008 and May 9, 2010, corresponding to 20,409 ton day fiducial exposure.

#### $\phi_{ES}(pp)$

pp solar-neutrino flux measured via  $\nu e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross section difference,  $\sigma(\nu_{\mu,\tau}\ e)\sim~0.3\ \sigma(\nu_e e)$ . If the pp solar-neutrino flux involves nonelectron flavor active neutrinos, their contribution to the flux is  $\sim~0.3$  times of  $\nu_e$ .

VALUE (
$$10^{10} \text{ cm}^{-2} \text{ s}^{-1}$$
)DOCUMENT IDTECNCOMMENT• • • We do not use the following data for averages, fits, limits, etc. • •

4.4 $\pm$ 0.5 <sup>1</sup> BELLINI 14A BORX average flux

<sup>1</sup> BELLINI 14A reports pp solar-neutrino flux measured via  $\nu e$  elastic scattering. The data were collected between January 2012 and May 2013, corresponding to 408 days of data. The pp neutrino interaction rate in Borexino is measured to be 144  $\pm$  13  $\pm$  10 counts/(day-100 ton) by fitting the measured energy spectrum of events in the 165–590

keV recoil electron kinetic energy window with the expected signal + background spectrum. The listed flux value  $\phi_{ES}(pp)$  is calculated from the observed rate and the number of  $(3.307 \pm 0.003) \times 10^{31}$  electrons for 100 tons of the Borexino scintillator, and the  $\nu_e e$  integrated cross section over the pp neutrino spectrum,  $\sigma(\nu_e e) = 11.38 \times 10^{-46}$  cm<sup>2</sup>.

#### $\phi_{CC}(pp)$

pp solar-neutrino flux measured with charged-current reaction which is sensitive exclusively to  $\nu_{e}$ .

<u>VALUE (10<sup>10</sup> cm<sup>-2</sup> s<sup>-1</sup>)</u>	DOCUMENT ID	TECN	COMMENT
$\bullet$ • • We do not use the follow	ving data for averages, fits,	limits,	etc. • • •
$3.38 \pm 0.47$	<sup>1</sup> ABDURASHI 09	FIT	Fit existing solar- $\nu$ data

<sup>1</sup> ABDURASHITOV 09 reports the *pp* solar-neutrino flux derived from the Ga solar neutrino capture rate by subtracting contributions from <sup>8</sup>B, <sup>7</sup>Be, *pep* and CNO solar neutrino fluxes determined by other solar neutrino experiments as well as neutrino oscillation parameters determined from available world neutrino oscillation data.

#### $\phi_{ES}$ (hep)

hep solar-neutrino flux measured via  $\nu e$  elastic scattering. This process is sensitive to all active neutrino flavors, but with reduced sensitivity to  $\nu_{\mu}$ ,  $\nu_{\tau}$  due to the cross-section difference,  $\sigma(\nu_{\mu,\tau} e) \sim 0.16 \sigma(\nu_e e)$ . If the hep solar-neutrino flux involves nonelectron flavor active neutrinos, their contribution to the flux is  $\sim 0.16$  times of  $\nu_e$ .

$VALUE (10^3 \text{ cm}^{-2} \text{s}^{-1})$	CL%	DOCUMENT ID		TECN	
$\bullet \bullet \bullet$ We do not use the	following	data for averages	, fits,	limits, etc.	•
<73	90	<sup>1</sup> HOSAKA	06	SKAM	

 $^1\,{\rm HOSAKA}$  06 result is obtained from the recoil electron energy window of 18–21 MeV, and updates FUKUDA 01 result.

## $\phi_{\overline{\nu}_e}$ (<sup>8</sup>B)

Searches are made for electron antineutrino flux from the Sun. Flux limits listed here are derived relative to the BS05(OP) Standard Solar Model <sup>8</sup>B solar neutrino flux  $(5.69 \times 10^6 \text{ cm}^{-2} \text{ s}^{-1})$ , with an assumption that solar  $\overline{\nu}_e$ s follow an unoscillated <sup>8</sup>B neutrino spectrum.

VALUE (%)	CL%	DOCUMENT ID		TECN	COMMENT		
$\bullet \bullet \bullet$ We do not use the	e following	data for average	s, fits,	limits, e	etc. • • •		
<0.013	90	BELLINI	11	BORX	$E_{\overline{ u}_{m{ m e}}} > 1.8 \; MeV$		
<1.9	90	<sup>1</sup> BALATA	06		$1.8 < E_{\overline{\nu}_e} < 20.0 \; MeV$		
<0.72	90	AHARMIM	04	SNO	$4.0 < E_{\overline{\nu}_{o}} < 14.8 \text{ MeV}$		
<0.022	90	EGUCHI	04	KLND	$8.3 < E_{\overline{ u}_e}^{c} < 14.8 \; { m MeV}$		
<0.7	90	GANDO	03	SKAM	$8.0 < E_{\overline{ u}_e} < 20.0 \text{ MeV}$		
<1.7	90	AGLIETTA			$7 < E_{\overline{\nu}_{ m e}} < 17  \mathrm{MeV}$		
$^1$ BALATA 06 obtained this result from the search for $\overline{ u}_e$ interactions with Counting Test							

Facility (the prototype of the Borexino detector).

#### (B) Three-neutrino mixing parameters

# See the related review(s): Three-Neutrino Mixing Parameters

## $\sin^2(\theta_{12})$

If an experin VALUE	nent reports $\sin^2(2\theta_{12})$ where $\underline{DOCUMENT ID}$	ve co	e value to $\sin^2(\theta_{12})$ . <u>COMMENT</u>	
$0.307^{+0.013}_{-0.012}$	<sup>1</sup> ABE	16C	FIT	KamLAND+global solar; $3 u$
• • • We do not u	use the following data for	, limits, etc. • • •		
$0.310\!\pm\!0.014$	<sup>2</sup> ABE	16C	FIT	SKAM+SNO; $3\nu$
$0.334 \substack{+ \ 0.027 \\ - \ 0.023}$	<sup>3</sup> ABE	16C	FIT	SK-I+II+III+IV; 3 <i>ν</i>
$0.327\substack{+0.026\\-0.031}$	<sup>4</sup> ABE	16C	FIT	SK-IV; 3 <i>ν</i>
$0.323 \!\pm\! 0.016$	<sup>5</sup> FORERO	14	FIT	$3\nu$
$0.304 \substack{+ 0.013 \\ - 0.012}$	<sup>6</sup> GONZALEZ	14	FIT	Either mass ordering; global fit
$0.299\substack{+0.014\\-0.014}$	<sup>7,8</sup> AHARMIM	13	FIT	global solar: $2\nu$
$0.307\substack{+0.016\\-0.013}$	<sup>8,9</sup> AHARMIM	13	FIT	global solar: $3\nu$
$0.304\substack{+0.022\\-0.018}$	<sup>8,10</sup> AHARMIM	13	FIT	$KamLAND + global \; solar: \; 3 u$
$0.304 \substack{+ 0.014 \\ - 0.013}$	<sup>11</sup> gando	13	FIT	${\sf KamLAND}+{\sf global \ solar}+{\sf SBL}+{\sf accelerator:}\ 3 u$
$0.304 \substack{+ 0.014 \\ - 0.013}$	<sup>12</sup> GANDO	13	FIT	$KamLAND + global \ solar: \ 3 u$
$0.325\substack{+0.039\\-0.039}$	<sup>13</sup> gando	13	FIT	KamLAND: $3\nu$
$0.30 \ {}^{+0.02}_{-0.01}$	<sup>14</sup> ABE	11	FIT	$KamLAND + global \; solar: \; 2  u$
$0.30 \begin{array}{c} +0.02 \\ -0.01 \end{array}$	<sup>15</sup> ABE	11	FIT	global solar: $2\nu$
$0.31 \ \begin{array}{c} + \ 0.03 \\ - \ 0.02 \end{array}$	<sup>16</sup> ABE	11	FIT	KamLAND + global solar: $3 u$
$0.31 \ \begin{array}{c} + \ 0.03 \\ - \ 0.03 \end{array}$	<sup>17</sup> ABE	11	FIT	global solar: $3 u$
$0.314 \substack{+\ 0.015 \\ -\ 0.012}$	<sup>18</sup> BELLINI	11A	FIT	KamLAND+global solar: 2 $ u$
$0.319^{+0.017}_{-0.015}$	<sup>19</sup> BELLINI	11A	FIT	global solar: $2\nu$
$0.311 \substack{+ 0.016 \\ - 0.016}$	<sup>20</sup> gando	11	FIT	$KamLAND + solar: \ 3 u$
$0.304 \substack{+ 0.046 \\ - 0.042}$	<sup>21</sup> gando	11	FIT	KamLAND: $3\nu$
$0.314 \substack{+ 0.018 \\ - 0.014}$	<sup>22,23</sup> AHARMIM	10	FIT	KamLAND + global solar: 2 $ u$
$0.314\substack{+0.017\\-0.020}$	<sup>22,24</sup> AHARMIM	10	FIT	global solar: $2 u$
$0.319\substack{+0.019\\-0.016}$	<sup>22,25</sup> AHARMIM	10	FIT	KamLAND+global solar: 3 $ u$
$0.319^{+0.023}_{-0.024}$	22,26 AHARMIM	10	FIT	global solar: $3 u$

$0.36 \begin{array}{c} +0.05 \\ -0.04 \end{array}$	<sup>27</sup> ABE	08A FIT	KamLAND
$0.32 \pm 0.03$	<sup>28</sup> ABE	08A FIT	KamLAND + global fit
$0.32 \pm 0.02$	<sup>29</sup> AHARMIM	08 FIT	KamLAND + global  solar
$0.31 \begin{array}{c} +0.04 \\ -0.04 \end{array}$	<sup>30</sup> HOSAKA	06 FIT	KamLAND + global  solar
$0.31 \begin{array}{c} +0.04 \\ -0.03 \end{array}$	<sup>31</sup> HOSAKA	06 FIT	$SKAM{+}SNO{+}KamLAND$
$0.31 \begin{array}{c} +0.03 \\ -0.04 \end{array}$	<sup>32</sup> HOSAKA	06 FIT	SKAM+SNO
$\substack{0.31 \ +0.02 \\ -0.03}$	<sup>33</sup> AHARMIM	05A FIT	KamLAND + global  solar
0.25-0.39	<sup>34</sup> AHARMIM	05A FIT	global solar
$0.29 \pm 0.03$	<sup>35</sup> ARAKI	05 FIT	KamLAND + global solar
$0.29 \begin{array}{c} +0.03 \\ -0.02 \end{array}$	<sup>36</sup> AHMED	04A FIT	KamLAND + global  solar
0.23–0.37	<sup>37</sup> AHMED	04A FIT	global solar
$0.31 \begin{array}{c} +0.04 \\ -0.04 \end{array}$	<sup>38</sup> SMY	04 FIT	KamLAND + global  solar
$0.29 \begin{array}{c} +0.04 \\ -0.04 \end{array}$	<sup>39</sup> SMY	04 FIT	global solar
$0.32 \begin{array}{c} +0.06 \\ -0.05 \end{array}$	<sup>40</sup> SMY	04 FIT	SKAM + SNO
0.19–0.33	<sup>41</sup> AHMAD	02B FIT	global solar
0.19-0.39	<sup>42</sup> FUKUDA	02 FIT	global solar
4			

<sup>1</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of  $\sin^2(\theta_{13}) = 0.0219 \pm 0.0014$  coming from reactor neutrino experiments, using all solar data and KamLAND data. *CPT* invariance is assumed.

- <sup>2</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of sin<sup>2</sup>( $\theta_{13}$ ) = 0.0219 ± 0.0014 coming from reactor neutrino experiments, using Super-Kamiokande (I+II+III+IV) and SNO data.
- <sup>3</sup>ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of  $\sin^2(\theta_{13}) = 0.0219 \pm 0.0014$  coming from reactor neutrino experiments, by combining the four phases of the Super-Kamiokande solar data.
- <sup>4</sup>ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of sin<sup>2</sup>( $\theta_{13}$ ) = 0.0219 ± 0.0014 coming from reactor neutrino experiments, using the Super-Kamiokande-IV data.

<sup>5</sup> FORERO 14 performs a global fit to neutrino oscillations using solar, reactor, longbaseline accelerator, and atmospheric neutrino data.

 $^6$  GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as  $0.304 \substack{+0.013 \\ -0.012}$  for normal and  $0.305 \substack{+0.012 \\ -0.013}$  for inverted mass ordering.

<sup>7</sup>AHARMIM 13 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data.

- <sup>8</sup>AHARMIM 13 global solar neutrino data include SNO's all-phases-combined analysis results on the total active <sup>8</sup>B neutrino flux and energy-dependent  $\nu_e$  survival probability parameters, measurements of CI (CLEVELAND 98), Ga (ABDURASHITOV 09 which contains combined analysis with GNO (ALTMANN 05 and Ph.D. thesis of F. Kaether)), and <sup>7</sup>Be (BELLINI 11A) rates, and <sup>8</sup>B solar-neutrino recoil electron measurements of SK-I (HOSAKA 06) zenith, SK-II (CRAVENS 08) and SK-III (ABE 11) day/night spectra, and Borexino (BELLINI 10A) spectra.
- <sup>9</sup>AHARMIM 13 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{32}^2$  fixed to 2.45 × 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino data.

- <sup>10</sup>AHARMIM 13 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{32}^2$  fixed to 2.45 × 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino and KamLAND (GANDO 11) data. CPT invariance is assumed.
- <sup>11</sup> GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND, global solar neutrino, short-baseline (SBL) reactor, and accelerator data, assuming CPT invariance. Supersedes GANDO 11.
- <sup>12</sup> GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND and global solar neutrino data, assuming CPT invariance. Supersedes GANDO 11.
- <sup>13</sup>GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND data. Supersedes GANDO 11.
- <sup>14</sup> ABE 11 obtained this result by a two-neutrino oscillation analysis using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, SAGE, and KamLAND data. CPT invariance is assumed.
- <sup>15</sup> ABE 11 obtained this result by a two-neutrino oscillation analysis using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, and SAGE data.
- $^{16}$  ABE 11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{32}$  fixed to  $2.4\times 10^{-3}~{\rm eV^2}$ , using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, SAGE, and KamLAND data. The normal neutrino mass ordering and CPT invariance are assumed.
- <sup>17</sup> ABE 11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{32}^2$  fixed to 2.4 × 10<sup>-3</sup> eV<sup>2</sup>, using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, and GALLEX/GNO data. The normal neutrino mass ordering is assumed.
- <sup>18</sup> BELLINI 11A obtained this result by a two-neutrino oscillation analysis using KamLAND, Homestake, SAGE, Gallex, GNO, Kamiokande, Super-Kamiokande, SNO, and Borexino (BELLINI 11A) data and the SSM flux prediction in SERENELLI 11 (Astrophysical Journal **743** 24 (2011)) with the exception that the <sup>8</sup>B flux was left free. CPT invariance is assumed.
- <sup>19</sup> BELLINI 11A obtained this result by a two-neutrino oscillation analysis using Homestake, SAGE, Gallex, GNO, Kamiokande, Super-Kamiokande, SNO, and Borexino (BELLINI 11A) data and the SSM flux prediction in SERENELLI 11 (Astrophysical Journal **743** 24 (2011)) with the exception that the <sup>8</sup>B flux was left free.
- $^{20}$  GANDO 11 obtain this result with three-neutrino fit using the KamLAND + solar data. Superseded by GANDO 13.
- $^{21}$  GANDO 11 obtain this result with three-neutrino fit using the KamLAND data only. Superseded by GANDO 13.
- <sup>22</sup> AHARMIM 10 global solar neutrino data include SNO's low-energy-threshold analysis survival probability day/night curves, SNO Phase III integral rates (AHARMIM 08), CI (CLEVELAND 98), SAGE (ABDURASHITOV 09), Gallex/GNO (HAMPEL 99, ALT-MANN 05), Borexino (ARPESELLA 08A), SK-I zenith (HOSAKA 06), and SK-II day/night spectra (CRAVENS 08).
- <sup>23</sup>AHARMIM 10 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data and KamLAND data (ABE 08A). *CPT* invariance is assumed.
- $^{24}$  AHARMIM 10 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data.
- <sup>25</sup> AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{31}^2$  fixed to  $2.3 \times 10^{-3}$  eV<sup>2</sup>, using global solar neutrino data and KamLAND data (ABE 08A). *CPT* invariance is assumed.
- <sup>26</sup>AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{31}^2$  fixed to 2.3 × 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino data.
- <sup>27</sup> ABE 08A obtained this result by a rate + shape + time combined geoneutrino and reactor two-neutrino fit for  $\Delta m_{21}^2$  and  $\tan^2 \theta_{12}$ , using KamLAND data only. Superseded by GANDO 11.

- $^{28}$  ABE 08A obtained this result by means of a two-neutrino fit using KamLAND, Homestake, SAGE, GALLEX, GNO, SK (zenith angle and E-spectrum), the SNO  $\chi^2$ -map, and solar flux data. CPT invariance is assumed. Superseded by GANDO 11.
- <sup>29</sup> The result given by AHARMIM 08 is  $\theta = (34.4 + 1.3)^{\circ}$ . This result is obtained by a two-neutrino oscillation analysis using solar neutrino data including those of Borexino (ARPESELLA 08A) and Super-Kamiokande-I (HOSAKA 06), and KamLAND data (ABE 08A). CPT invariance is assumed.
- <sup>30</sup> HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using SK  $\nu_e$  data, CC data from other solar neutrino experiments, and KamLAND data (ARAKI 05). CPT invariance is assumed.
- $^{31}$  HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using the data from Super-Kamiokande, SNO (AHMAD 02 and AHMAD 02B), and KamLAND (ARAKI 05) experiments. CPT invariance is assumed.
- $^{32}$ HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using the Super-Kamiokande and SNO (AHMAD 02 and AHMAD 02B) solar neutrino data.
- $^{33}$  The result given by AHARMIM 05A is heta = (33.9  $\pm$  1.6) $^{\circ}$ . This result is obtained by a two-neutrino oscillation analysis using SNO pure deuteron and salt phase data, SK  $\nu_{e}$  data, CI and Ga CC data, and KamLAND data (ARAKI 05). CPT invariance is assumed. AHARMIM 05A also quotes  $\theta = (33.9^{+2.4}_{-2.2})^{\circ}$  as the error enveloping the 68% CL two-dimensional region. This translates into  $\sin^2 2 \theta = 0.86^{+0.05}_{-0.06}$ .

- $^{34}$ AHARMIM 05A obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in figure 35a of AHARMIM 05A. AHARMIM 05A also quotes  $\tan^2\theta = 0.45 \stackrel{+0.09}{-0.08}$  as the error enveloping the 68% CL two-dimensional region. This translates into  $\sin^2 2 \theta = 0.86 \stackrel{+0.05}{-0.07}$
- $^{35}$  ARAKI 05 obtained this result by a two-neutrino oscillation analysis using KamLAND and solar neutrino data. *CPT* invariance is assumed. The  $1\sigma$  error shown here is translated from the number provided by the KamLAND collaboration,  $\tan^2\theta = 0.40 + 0.07 - 0.05$ . The corresponding number quoted in ARAKI 05 is  $\tan^2\theta = 0.40 {+0.10 \atop -0.07}$  ( $\sin^2 2 \theta = 0.82 \pm$ 0.07), which envelops the 68% CL two-dimensional region.
- <sup>36</sup> The result given by AHMED 04A is  $\theta = (32.5^{+1.7}_{-1.6})^{\circ}$ . This result is obtained by a two-neutrino oscillation analysis using solar neutrino and KamLAND data (EGUCHI 03). *CPT* invariance is assumed. AHMED 04A also quotes  $\theta = (32.5^{+2.4}_{-2.3})^{\circ}$  as the error enveloping the 68% CL two-dimensional region. This translates into  $\sin^2 2 \theta = 0.82 \pm 0.06$ .
- $^{37}$ AHMED 04A obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in Fig. 5(a) of AHMED 04A. The best-fit point is  $\Delta(m^2) = 6.5 \times 10^{-5} \text{ eV}^2$ ,  $\tan^2 \theta = 0.40 (\sin^2 2 \theta = 0.82)$ .
- $^{38}$  The result given by SMY 04 is tan $^2 heta = 0.44 \pm 0.08$ . This result is obtained by a twoneutrino oscillation analysis using solar neutrino and KamLAND data (IANNI 03). CPT invariance is assumed.
- $^{39}$ SMY 04 obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The  $1\sigma$  errors are read from Fig. 6(a) of SMY 04.
- $^{40}$ SMY 04 obtained this result by a two-neutrino oscillation analysis using the Super-Kamiokande and SNO (AHMAD 02 and AHMAD 02B) solar neutrino data. The  $1\sigma$ errors are read from Fig. 6(a) of SMY 04.
- $^{41}$ AHMAD 02B obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in Fig. 4(b) of AHMAD 02B. The best fit point is  $\Delta(m^2) = 5.0 \times 10^{-5} \text{ eV}^2$  and  $\tan \theta = 0.34$  ( $\sin^2 2 \theta = 0.76$ ).
- $^{42}$  FUKUDA 02 obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95%

CL two-dimensional region shown in Fig. 4 of FUKUDA 02. The best fit point is  $\Delta(m^2) = 6.9 \times 10^{-5} \text{ eV}^2$  and  $\tan^2\theta = 0.38$  ( $\sin^2 2 \theta = 0.80$ ).

#### $\Delta m_{21}^2$

<u>21</u> <u>VALUE (10<sup>-5</sup> eV<sup>2</sup>)</u>				COMMENT		
7.53±0.18	<sup>1</sup> GANDO	13	FIT	$KamLAND + global \ solar + SBL + accelerator: \ 3 u$		
• • • We do not use	e the following data fo	r aver	ages, fit			
$7.49 \substack{+0.19 \\ -0.18}$	<sup>2</sup> ABE	<b>16</b> C	FIT	KamLAND+global solar; 3 $ u$		
$4.8 \begin{array}{c} +1.3 \\ -0.6 \end{array}$	<sup>3</sup> ABE	<b>16</b> C	FIT	SKAM+SNO; $3\nu$		
$\begin{array}{c} 4.8 & +1.5 \\ -0.8 \end{array}$	<sup>4</sup> ABE	16C	FIT	SK-I+II+III+IV; 3 <i>v</i>		
$3.2 \ {}^{+2.8}_{-0.2}$	<sup>5</sup> ABE	16C	FIT	SK-IV; 3 <i>ν</i>		
$7.6 \begin{array}{c} +0.19 \\ -0.18 \end{array}$	<sup>6</sup> FORERO	14	FIT	3ν		
$7.50 \substack{+0.19 \\ -0.17}$	<sup>7</sup> GONZALEZ	14	FIT	Either mass ordering; global fit		
$5.13 \substack{+1.29 \\ -0.96}$	<sup>8,9</sup> AHARMIM	13	FIT	global solar: $2 u$		
$5.13 \substack{+1.49 \\ -0.98}$	<sup>9,10</sup> AHARMIM	13	FIT	global solar: $3 u$		
$7.46 \substack{+0.20 \\ -0.19}$	<sup>9,11</sup> AHARMIM	13	FIT	KamLAND + global solar: $3 u$		
$7.53 \substack{+0.19 \\ -0.18}$	<sup>12</sup> GANDO	13	FIT	$KamLAND+global\;solar:\;3 u$		
$7.54 \substack{+ \ 0.19 \\ - \ 0.18}$	<sup>13</sup> gando	13	FIT	KamLAND: $3\nu$		
$7.6 \pm 0.2$	<sup>14</sup> ABE	11	FIT	$KamLAND+global\;solar:\;2 u$		
$6.2 \begin{array}{c} +1.1 \\ -1.9 \end{array}$	<sup>15</sup> ABE	11	FIT	global solar: $2 u$		
7.7 ±0.3	<sup>16</sup> ABE	11	FIT	${\sf KamLAND}+{\sf global}$ solar: 3 $ u$		
$6.0 \ \begin{array}{c} +2.2 \\ -2.5 \end{array}$	<sup>17</sup> ABE	11	FIT	global solar: $3 u$		
$7.50 \substack{+0.16 \\ -0.24}$	<sup>18</sup> BELLINI	11A	FIT	KamLAND + global solar: 2 $ u$		
$5.2 \ \begin{array}{c} +1.5 \\ -0.9 \end{array}$	<sup>19</sup> BELLINI	11A	FIT	global solar: $2\nu$		
$7.50 \substack{+0.19 \\ -0.20}$	<sup>20</sup> GANDO	11	FIT	$KamLAND + solar: \ 3 u$		
$7.49 \pm 0.20$	<sup>21</sup> GANDO	11	FIT	KamLAND: $3\nu$		
$7.59 \substack{+0.20 \\ -0.21}$	<sup>22,23</sup> AHARMIM	10	FIT	KamLAND + global solar: $2 u$		
$5.89^{+2.13}_{-2.16}$	<sup>22,24</sup> AHARMIM	10	FIT	global solar: $2 u$		
$7.59 \pm 0.21$	<sup>22,25</sup> AHARMIM	10	FIT	$KamLAND+global\;solar:\;3 u$		
$6.31^{+2.49}_{-2.58}$	<sup>22,26</sup> AHARMIM	10	FIT	global solar: $3\nu$		
$7.58^{+0.14}_{-0.13}{\pm}0.15$	<sup>27</sup> ABE	08A	FIT	KamLAND		
$7.59 \pm 0.21$	<sup>28</sup> ABE	08A	FIT	KamLAND + global  solar		
$7.59\substack{+0.19 \\ -0.21}$	<sup>29</sup> AHARMIM	08	FIT	KamLAND + global  solar		
= 0.21 8.0 $\pm 0.3$	<sup>30</sup> HOSAKA	06	FIT	KamLAND + global  solar		

8.0 ±0.3	<sup>31</sup> HOSAKA	06	FIT	SKAM+SNO+KamLAND
$6.3 \begin{array}{c} +3.7 \\ -1.5 \end{array}$	<sup>32</sup> HOSAKA	06	FIT	SKAM+SNO
5–12	<sup>33</sup> HOSAKA	06	FIT	SKAM day/night in the LMA region
$8.0 \begin{array}{c} +0.4 \\ -0.3 \end{array}$	<sup>34</sup> AHARMIM	05A	FIT	KamLAND + global  solar LMA
3.3–14.4	<sup>35</sup> AHARMIM	05A	FIT	global solar
$7.9 \ {}^{+0.4}_{-0.3}$	<sup>36</sup> ARAKI	05	FIT	KamLAND + global  solar
$7.1 \ \begin{array}{c} +1.0 \\ -0.3 \end{array}$	<sup>37</sup> AHMED	04A	FIT	KamLAND + global  solar
3.2–13.7	<sup>38</sup> AHMED	04A	FIT	global solar
$7.1 \begin{array}{c} +0.6 \\ -0.5 \end{array}$	<sup>39</sup> SMY	04	FIT	KamLAND + global  solar
$6.0 \ \begin{array}{c} +1.7 \\ -1.6 \end{array}$	<sup>40</sup> SMY	04	FIT	global solar
$6.0 \ +2.5 \ -1.6$	<sup>41</sup> SMY	04	FIT	SKAM + SNO
2.8–12.0	<sup>42</sup> AHMAD	<b>0</b> 2B	FIT	global solar
3.2–19.1	<sup>43</sup> FUKUDA	02	FIT	global solar

<sup>1</sup> GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND, global solar neutrino, short-baseline (SBL) reactor, and accelerator data, assuming CPT invariance. Supersedes GANDO 11.

- <sup>2</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of  $\sin^2(\theta_{13}) = 0.0219 \pm 0.0014$  coming from reactor neutrino experiments, using all solar data and KamLAND data. *CPT* invariance is assumed. <sup>3</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint
- <sup>3</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of  $\sin^2(\theta_{13}) = 0.0219 \pm 0.0014$  coming from reactor neutrino experiments, using Super-Kamiokande (I+II+III+IV) and SNO data.
- <sup>4</sup> ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of  $\sin^2(\theta_{13}) = 0.0219 \pm 0.0014$  coming from reactor neutrino experiments, by combining the four phases of the Super-Kamiokande solar data.
- <sup>5</sup>ABE 16C obtained this result by a three-neutrino oscillation analysis, with a constraint of sin<sup>2</sup>( $\theta_{13}$ ) = 0.0219 ± 0.0014 coming from reactor neutrino experiments, using the Super-Kamiokande-IV data.
- $^6$  FORERO 14 performs a global fit to  $\Delta m^2_{21}$  using solar, reactor, long-baseline accelerator, and atmospheric neutrino data.
- <sup>7</sup> GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as  $(7.50^{+0.19}_{-0.17}) \times 10^{-5} \text{ eV}^2$  for normal and  $(7.50^{+0.18}_{-0.17}) \times 10^{-5} \text{ eV}^2$  for inverted mass ordering.
- $^{8}$ AHARMIM 13 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data.
- <sup>9</sup>AHARMIM 13 global solar neutrino data include SNO's all-phases-combined analysis results on the total active <sup>8</sup>B neutrino flux and energy-dependent  $\nu_e$  survival probability parameters, measurements of Cl (CLEVELAND 98), Ga (ABDURASHITOV 09 which contains combined analysis with GNO (ALTMANN 05 and Ph.D. thesis of F. Kaether)), and <sup>7</sup>Be (BELLINI 11A) rates, and <sup>8</sup>B solar-neutrino recoil electron measurements of SK-I (HOSAKA 06) zenith, SK-II (CRAVENS 08), and SK-III (ABE 11) day/night spectra, and Borexino (BELLINI 10A) spectra.
- $^{10}$  AHARMIM 13 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{31}$  fixed to 2.45  $\times$  10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino data.

- $^{11}$ AHARMIM 13 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{31}$  fixed to 2.45 imes 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino and KamLAND data (GANDO 11). CPT invariance is assumed.
- $^{12}$  GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND and global solar neutrino data, assuming CPT invariance. Supersedes GANDO 11.
- $^{13}$  GANDO 13 obtained this result by a three-neutrino oscillation analysis using KamLAND data. Supersedes GANDO 11.
- $^{14}$  ABE 11 obtained this result by a two-neutrino oscillation analysis using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, SAGE, and KamLAND data. CPT invariance is assumed.
- $^{15}$  ABE 11 obtained this result by a two-neutrino oscillation analysis using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, and SAGE data.
- $^{16}$  ABE 11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{32}$  fixed to  $2.4\times 10^{-3}~{\rm eV}^2$ , using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, SAGE, and KamLAND data. The normal neutrino mass ordering and CPT invariance are assumed.
- $^{17}\operatorname{ABE}$  11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{32}$  fixed to 2.4 × 10<sup>-3</sup> eV<sup>2</sup>, using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, and GALLEX/GNO data. The normal neutrino mass ordering is assumed.
- $^{18}$  BELLINI 11A obtained this result by a two-neutrino oscillation analysis using KamLAND, Homestake, SAGE, Gallex, GNO, Kamiokande, Super-Kamiokande, SNO, and Borexino (BELLINI 11A) data and the SSM flux prediction in SERENELLI 11 (Astrophysical Journal **743** 24 (2011)) with the exception that the <sup>8</sup>B flux was left free. CPT invariance is assumed.
- <sup>19</sup>BELLINI 11A obtained this result by a two-neutrino oscillation analysis using Homestake, SAGE, Gallex, GNO, Kamiokande, Super-Kamiokande, SNO, and Borexino (BELLINI 11A) data and the SSM flux prediction in SERENELLI 11 (Astrophysical Journal **743** 24 (2011)) with the exception that the <sup>8</sup>B flux was left free.
- $^{20}$  GANDO 11 obtain this result with three-neutrino fit using the KamLAND + solar data. Superseded by GANDO 13.
- $^{21}$ GANDO 11 obtain this result with three-neutrino fit using the KamLAND data only. Supersedes ABE 08A.
- $^{22}$ AHARMIM 10 global solar neutrino data include SNO's low-energy-threshold analysis survival probability day/night curves, SNO Phase III integral rates (AHARMIM 08), CI (CLEVELAND 98), SAGE (ABDURASHITOV 09), Gallex/GNO (HAMPEL 99, ALT-MANN 05), Borexino (ARPESELLA 08A), SK-I zenith (HOSAKA 06), and SK-II day/night spectra (CRAVENS 08).
- $^{23}$ AHARMIM 10 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data and KamLAND data (ABE 08A). CPT invariance is assumed.
- <sup>24</sup>AHARMIM 10 obtained this result by a two-neutrino oscillation analysis using global solar neutrino data.
- <sup>25</sup> AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{31}$  fixed to  $2.3 \times 10^{-3}$  eV<sup>2</sup>, using global solar neutrino data and KamLAND data (ABE 08A). CPT invariance is assumed.
- <sup>26</sup> AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{31}$  fixed to  $2.3 \times 10^{-3}$  eV<sup>2</sup>, using global solar neutrino data.
- $^{27}$  ABE 08A obtained this result by a rate + shape + time combined geoneutrino and reactor two-neutrino fit for  $\Delta m_{21}^2$  and  $\tan^2 \theta_{12}$ , using KamLAND data only. Superseded by GANDO 11.
- $^{28}$  ABE 08A obtained this result by means of a two-neutrino fit using KamLAND, Homestake, SAGE, GALLEX, GNO, SK (zenith angle and E-spectrum), the SNO  $\chi^2$ -map, and solar flux data. CPT invariance is assumed. Superseded by GANDO 11.

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- <sup>29</sup> AHARMIM 08 obtained this result by a two-neutrino oscillation analysis using all solar neutrino data including those of Borexino (ARPESELLA 08A) and Super-Kamiokande-I (HOSAKA 06), and KamLAND data (ABE 08A). CPT invariance is assumed.
- <sup>30</sup> HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using solar neutrino and KamLAND data (ARAKI 05). CPT invariance is assumed.
- <sup>31</sup> HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using the data from Super-Kamiokande, SNO (AHMAD 02 and AHMAD 02B), and KamLAND (ARAKI 05) experiments. CPT invariance is assumed.
- <sup>32</sup> HOSAKA 06 obtained this result by a two-neutrino oscillation analysis using the Super-Kamiokande and SNO (AHMAD 02 and AHMAD 02B) solar neutrino data.
- <sup>33</sup> HOSAKA 06 obtained this result from the consistency between the observed and expected day-night flux asymmetry amplitude. The listed 68% CL range is derived from the  $1\sigma$  boundary of the amplitude fit to the data. Oscillation parameters are constrained to be in the LMA region. The mixing angle is fixed at  $\tan^2\theta = 0.44$  because the fit depends only very weekly on it.
- <sup>34</sup> AHARMIM 05Å obtained this result by a two-neutrino oscillation analysis using solar neutrino and KamLAND data (ARAKI 05). *CPT* invariance is assumed. AHARMIM 05Å also quotes  $\Delta(m^2) = (8.0^{+0.6}_{-0.4}) \times 10^{-5} \text{ eV}^2$  as the error enveloping the 68% CL two-dimensional region.
- $^{35}$  AHARMIM 05Å obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in figure 35a of AHARMIM 05A. AHARMIM 05A also quotes  $\Delta(m^2) = (6.5^{+4.4}_{-2.3}) \times 10^{-5} \ {\rm eV}^2$  as the error enveloping the 68% CL two-dimensional region.
- <sup>36</sup> ARAKI 05 obtained this result by a two-neutrino oscillation analysis using KamLAND and solar neutrino data. *CPT* invariance is assumed. The  $1\sigma$  error shown here is provided by the KamLAND collaboration. The error quoted in ARAKI 05,  $\Delta(m^2) = (7.9^{+0.6}_{-0.5}) \times 10^{-5}$ 
  - $10^{-5}$ , envelops the 68% CL two-dimensional region.
- <sup>37</sup> AHMED 04A obtained this result by a two-neutrino oscillation analysis using solar neutrino and KamLAND data (EGUCHI 03). *CPT* invariance is assumed. AHMED 04A also quotes  $\Delta(m^2) = (7.1^{+1.2}_{-0.6}) \times 10^{-5} \text{ eV}^2$  as the error enveloping the 68% CL two-dimensional region.
- <sup>38</sup> AHMED 04A obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in Fig. 5(a) of AHMED 04A. The best-fit point is  $\Delta(m^2) = 6.5 \times 10^{-5} \text{ eV}^2$ ,  $\tan^2 \theta = 0.40 (\sin^2 2 \theta = 0.82)$ .
- <sup>39</sup> SMY 04 obtained this result by a two-neutrino oscillation analysis using solar neutrino and KamLAND data (IANNI 03). *CPT* invariance is assumed.
- <sup>40</sup> SMY 04 obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The  $1\sigma$  errors are read from Fig. 6(a) of SMY 04.
- $^{41}\,\rm SMY$  04 obtained this result by a two-neutrino oscillation analysis using the Super-Kamiokande and SNO (AHMAD 02 and AHMAD 02B) solar neutrino data. The  $1\sigma$  errors are read from Fig. 6(a) of SMY 04.
- <sup>42</sup> AHMAD 02B obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in Fig. 4(b) of AHMAD 02B. The best fit point is  $\Delta(m^2) = 5.0 \times 10^{-5} \text{ eV}^2$  and  $\tan \theta = 0.34 (\sin^2 2 \theta = 0.76)$ .
- <sup>43</sup> FUKUDA 02 obtained this result by a two-neutrino oscillation analysis using the data from all solar neutrino experiments. The listed range of the parameter envelops the 95% CL two-dimensional region shown in Fig. 4 of FUKUDA 02. The best fit point is  $\Delta(m^2) = 6.9 \times 10^{-5} \text{ eV}^2$  and  $\tan^2\theta = 0.38$  ( $\sin^2 2 \theta = 0.80$ ).

## $\sin^2(\theta_{23})$

The reported limits below correspond to the projection onto the  $\sin^2(\theta_{23})$  axis of the 90% CL contours in the  $\sin^2(\theta_{23}) - \Delta m_{32}^2$  plane presented by the authors. Unless otherwise specified, the limits are 90% CL and the reported uncertainties are 68% CL.

If an experiment reports  $\sin^2(2\theta_{23})$  we convert the value to  $\sin^2(\theta_{23})$ . <u>VALUE</u> <u>DOCUMENT ID</u> <u>TECN</u> <u>COMMENT</u>

VALUE	DOCUMENT IL	<u>TECN</u>	COMMENT
$0.421^{+0.033}_{-0.025}$ OU	JR FIT Error in	cludes scale fac	ctor of 1.3. Inverted ordering, quadrant I
$0.592^{+0.023}_{-0.030}$ OU	JR FIT Error in	cludes scale fac	ctor of 1.1. Inverted ordering, quadrant II
$0.417^{+0.025}_{-0.028}$ OU	JR FIT Error in	cludes scale fac	ctor of 1.2. Normal ordering, quadrant I
$0.597^{+0.024}_{-0.030}$ OU	JR FIT Error in	cludes scale fac	ctor of 1.2. Normal ordering, quadrant II
$0.55 \begin{array}{c} +0.05 \\ -0.09 \end{array}$	$^1$ ABE	17F T2K	normal mass ordering
$\substack{0.55 \\ -0.08}^{+0.05}$	<sup>1</sup> ABE	17F T2K	inverted mass ordering
$0.404\substack{+0.022\\-0.030}$	ADAMSON	17A NOVA	normal mass ordering; first quadrant
$0.624\substack{+0.022\\-0.030}$	ADAMSON	17A NOVA	normal mass ordering; second quadrant
$0.398\substack{+0.030\\-0.022}$	ADAMSON	17A NOVA	inverted mass ordering; first quadrant
$0.618\substack{+0.022\\-0.030}$	ADAMSON	17A NOVA	inverted mass ordering; second quadrant
$0.53 \begin{array}{c} +0.09 \\ -0.12 \end{array}$	<sup>2</sup> AARTSEN	15A ICCB	3 u osc; normal mass ordering
$\substack{0.51 + 0.09 \\ -0.11}$	<sup>2</sup> AARTSEN	15A ICCB	3 u osc; inverted mass ordering
$0.41 \ {}^{+0.23}_{-0.06}$	<sup>3</sup> ADAMSON	14 MINS	3 u osc., normal mass ordering
$0.41 \begin{array}{c} +0.26 \\ -0.07 \end{array}$	<sup>3</sup> ADAMSON	14 MINS	3 u osc.; inverted mass ordering
	use the following	g data for avera	ages, fits, limits, etc. • • •
$0.532\substack{+0.061\\-0.087}$	<sup>1</sup> ABE	17A T2K	normal mass ordering
$0.534\substack{+0.061\\-0.087}$	<sup>1</sup> ABE	17A T2K	inverted mass ordering
$0.51 \begin{array}{c} +0.08 \\ -0.07 \end{array}$	ABE	17c T2K	normal mass ordering with neutrinos
$0.42 \ \begin{array}{c} +0.25 \\ -0.07 \end{array}$	ABE	17c T2K	normal mass ordering with antineutrinos
$0.52 \begin{array}{c} +0.075 \\ -0.09 \end{array}$	ABE	17C T2K	normal mass ordering with neutrinos and antineutrinos
$0.45 \begin{array}{c} +0.19 \\ -0.07 \end{array}$	<sup>4</sup> ABE	16D T2K	$3\nu$ osc; normal mass ordering; $\overline{\nu}$ beam
0.38 to 0.65	<sup>5</sup> ADAMSON		normal mass ordering
0.37 to 0.64	<sup>5</sup> ADAMSON	16A NOVA	inverted mass ordering
$0.514 \substack{+ 0.055 \\ - 0.056}$	<sup>6</sup> ABE	14 T2K	3 u osc.; normal mass ordering
$0.511 \!\pm\! 0.055$	<sup>6</sup> ABE	14 T2K	3 u osc.; inverted mass ordering
$0.567\substack{+0.032\\-0.128}$	<sup>7</sup> FORERO	14 FIT	Normal mass ordering
$0.573^{\mathrm{+0.025}}_{\mathrm{-0.043}}$	<sup>7</sup> FORERO	14 FIT	Inverted mass ordering
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	0			
$0.452 \substack{+0.052 \\ -0.028}$	<sup>8</sup> GONZALEZ	14	FIT	Normal mass ordering; global fit
$0.579 \substack{+0.025 \\ -0.037}$	<sup>8</sup> GONZALEZ	14	FIT	Inverted mass ordering; global fit
0.24 to 0.76	<sup>9</sup> AARTSEN	<b>13</b> B	ICCB	DeepCore, $2\nu$ oscillation
$0.514 \pm 0.082$	<sup>10</sup> ABE	13G	T2K	3 u osc.; normal mass ordering
$0.388\substack{+0.051\\-0.053}$	<sup>11</sup> ADAMSON	<b>13</b> B	MINS	$Beam + Atmospheric; identical \ \nu \ \& \ \overline{\nu}$
0.3 to 0.7	<sup>12</sup> ABE	12A	T2K	off-axis beam
0.28 to 0.72	<sup>13</sup> ADAMSON	12	MINS	$\overline{ u}$ beam
0.25 to 0.75	<sup>14,15</sup> ADAMSON	12B	MINS	MINOS atmospheric
0.27 to 0.73	14,16 ADAMSON	12B	MINS	MINOS pure atmospheric $ u$
0.21 to 0.79	<sup>14,16</sup> ADAMSON	12B	MINS	MINOS pure atmospheric $\overline{ u}$
0.15 to 0.85	<sup>17</sup> ADRIAN-MAR.	.12	ANTR	atmospheric $\boldsymbol{\nu}$ with deep see telescope
0.39 to 0.61	<sup>18</sup> ABE	11C	SKAM	Super-Kamiokande
0.34 to 0.66	ADAMSON	11	MINS	2 u osc.; maximal mixing
$0.31 \ \begin{array}{c} +0.10 \\ -0.07 \end{array}$	<sup>19</sup> ADAMSON	11B	MINS	$\overline{ u}$ beam
0.41 to 0.59	<sup>20</sup> WENDELL	10	SKAM	$3 u$ osc. with solar terms; $ heta_{13}{=}0$
0.39 to 0.61	<sup>21</sup> WENDELL	10	SKAM	$3\nu$ osc.; normal mass ordering
0.37 to 0.63	<sup>22</sup> WENDELL	10	SKAM	3 u osc.; inverted mass ordering
0.31 to 0.69	ADAMSON	08A	MINS	MINOS
0.05 to 0.95	<sup>23</sup> ADAMSON	06	MINS	atmospheric $ u$ with far detector
0.18 to 0.82	<sup>24</sup> AHN	06A	K2K	KEK to Super-K
0.23 to 0.77	<sup>25</sup> MICHAEL	06	MINS	MINOS
0.18 to 0.82	<sup>26</sup> ALIU	05	K2K	KEK to Super-K
0.18 to 0.82	<sup>27</sup> ALLISON	05	SOU2	
0.36 to 0.64	<sup>28</sup> ASHIE	05	SKAM	•
0.28 to 0.72	<sup>29</sup> AMBROSIO	04		MACRO
0.34 to 0.66	<sup>30</sup> ASHIE	04	SKAM	1
0.08 to 0.92	<sup>31</sup> AHN	03	K2K	KEK to Super-K
0.13 to 0.87	<sup>32</sup> AMBROSIO	03		MACRO
0.26 to 0.74	<sup>33</sup> AMBROSIO	03	MCRO	MACRO
0.15 to 0.85	<sup>34</sup> SANCHEZ	03	SOU2	Soudan-2 Atmospheric
0.28 to 0.72	<sup>35</sup> AMBROSIO	01		upward $\mu$
0.29 to 0.71	<sup>36</sup> AMBROSIO	01		upward $\mu$
0.13 to 0.87	<sup>37</sup> FUKUDA	<b>99</b> C	SKAM	. ,
0.23 to 0.77	<sup>38</sup> FUKUDA		SKAM	
0.08 to 0.92	<sup>39</sup> FUKUDA	<b>99</b> D	SKAM	stop $\mu$ / through
0.29 to 0.71	<sup>40</sup> FUKUDA	98C	SKAM	•
0.08 to 0.92	41 HATAKEYAMA		KAMI	Kamiokande
0.24 to 0.76	42 HATAKEYAMA		KAMI	Kamiokande
0.20 to 0.80	<sup>43</sup> FUKUDA	94	KAMI	Kamiokande

<sup>1</sup> Errors are from the projections of the 68% contour on 2D plot of  $\Delta m^2$  versus sin<sup>2</sup>( $\theta_{23}$ ). ABE 17F supersedes ABE 17A.

 $^2$  AARTSEN 15A obtains this result by a three-neutrino oscillation analysis using 10–100 GeV muon neutrino sample from a total of 953 days of measurement with the low-energy subdetector DeepCore of the IceCube neutrino telescope.

<sup>3</sup> ADAMSON 14 uses a complete set of accelerator and atmospheric data. The analysis combines the  $\nu_{\mu}$  disappearance and  $\nu_{e}$  appearance data using three-neutrino oscillation fit. The fit results are obtained for normal and inverted mass ordering assumptions. The best fit is for lower  $\theta_{23}$  quadrant and inverted mass ordering.

 $^4\,{\rm ABE}$  16D reports oscillation results using  $\overline{\nu}_\mu$  disappearance in an off-axis beam.

- <sup>5</sup> ADAMSON 16A obtains  $\sin^2(\theta_{23})$  in the 68% C.L. range [0.38, 0.65] ([0.37, 0.64]), with two statistically degenerate best-fit values of 0.44 and 0.59 (0.44 and 0.59) for normal (inverted) mass ordering. Superseded by ADAMSON 17A.
- <sup>6</sup>ABE 14 results are based on  $\nu_{\mu}$  disappearance using three-neutrino oscillation fit. The confidence intervals are derived from one dimensional profiled likelihoods. Superseded by \_ABE 17A.
- <sup>7</sup> FORERO 14 performs a global fit to neutrino oscillations using solar, reactor, longbaseline accelerator, and atmospheric neutrino data.
- <sup>8</sup> GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as 68% CL intervals of 0.433–0.496 or 0.530–0.594 for normal and 0.514–0.612 for inverted mass ordering.
- <sup>9</sup> AARTSEN 13B obtained this result by a two-neutrino oscillation analysis using 20–100 GeV muon neutrino sample from a total of 318.9 days of live-time measurement with the low-energy subdetector DeepCore of the IceCube neutrino telescope.
- $^{10}$  The best fit value is sin $^2( heta_{23})=$  0.514  $\pm$  0.082. Superseded by ABE 14.
- <sup>11</sup> ADAMSON 13B obtained this result from  $\nu_{\mu}$  and  $\overline{\nu}_{\mu}$  disappearance using  $\nu_{\mu}$  (10.71 × 10<sup>20</sup> POT) and  $\overline{\nu}_{\mu}$  (3.36 × 10<sup>20</sup> POT) beams, and atmospheric (37.88kton-years) data from MINOS The fit assumed two-flavor neutrino hypothesis and identical  $\nu_{\mu}$  and  $\overline{\nu}_{\mu}$  oscillation parameters. Superseded by ADAMSON 14.
- <sup>12</sup> ABE 12A obtained this result by a two-neutrino oscillation analysis. The best-fit point is  $sin^2(2\theta_{23}) = 0.98$ .
- <sup>13</sup>ADAMSON 12 is a two-neutrino oscillation analysis using antineutrinos. The best fit value is  $\sin^2(2\theta_{23}) = 0.95 + 0.10 \pm 0.01$ .
- $^{14}$  ADAMSON 12B obtained this result by a two-neutrino oscillation analysis of the L/E distribution using 37.9 kton·yr atmospheric neutrino data with the MINOS far detector.
- <sup>15</sup> The best fit point is  $\Delta m^2 = 0.0019 \text{ eV}^2$  and  $\sin^2 2\theta = 0.99$ . The 90% single-parameter confidence interval at the best fit point is  $\sin^2 2\theta > 0.86$ .
- <sup>16</sup> The data are separated into pure samples of  $\nu$ s and  $\overline{\nu}$ s, and separate oscillation parameters for  $\nu$ s and  $\overline{\nu}$ s are fit to the data. The best fit point is  $(\Delta m^2, \sin^2 2\theta) = (0.0022 \text{ eV}^2, 0.99)$  and  $(\Delta \overline{m}^2, \sin^2 2\overline{\theta}) = (0.0016 \text{ eV}^2, 1.00)$ . The quoted result is taken from the 90% C.L. contour in the  $(\Delta m^2, \sin^2 2\theta)$  plane obtained by minimizing the four parameter log-likelihood function with respect to the other oscillation parameters.
- <sup>17</sup> ADRIAN-MARTINEZ 12 measured the oscillation parameters of atmospheric neutrinos with the ANTARES deep sea neutrino telescope using the data taken from 2007 to 2010 (863 days of total live time).
- <sup>18</sup> ABE 11C obtained this result by a two-neutrino oscillation analysis using the Super-Kamiokande-I+II+III atmospheric neutrino data. ABE 11C also reported results under a two-neutrino disappearance model with separate mixing parameters between  $\nu$  and  $\overline{\nu}$ , and obtained  $\sin^2 2\theta > 0.93$  for  $\nu$  and  $\sin^2 2\theta > 0.83$  for  $\overline{\nu}$  at 90% C.L.
- <sup>19</sup> ADAMSON 11B obtained this result by a two-neutrino oscillation analysis of antineutrinos in an antineutrino enhanced beam with  $1.71 \times 10^{20}$  protons on target. This results is consistent with the neutrino measurements of ADAMSON 11 at 2% C.L.
- consistent with the neutrino measurements of ADAMSON 11 at 2% C.L. <sup>20</sup>WENDELL 10 obtained this result ( $\sin^2\theta_{23} = 0.407-0.583$ ) by a three-neutrino oscillation analysis using the Super-Kamiokande-I+II+III atmospheric neutrino data, assuming  $\theta_{13} = 0$  but including the solar oscillation parameters  $\Delta m_{21}^2$  and  $\sin^2\theta_{12}$  in the fit.
- $^{21}$  WENDELL 10 obtained this result (sin<sup>2</sup> $\theta_{23} = 0.43-0.61$ ) by a three-neutrino oscillation analysis with one mass scale dominance ( $\Delta m^2_{21} = 0$ ) using the Super-Kamiokande-I+II+III atmospheric neutrino data, and updates the HOSAKA 06A result.
- $^{22}$  WENDELL 10 obtained this result (sin<sup>2</sup> $\theta_{23} = 0.44-0.63$ ) by a three-neutrino oscillation analysis with one mass scale dominance ( $\Delta m^2_{21} = 0$ ) using the Super-Kamiokande-I+II+III atmospheric neutrino data, and updates the HOSAKA 06A result.

- <sup>23</sup> ADAMSON 06 obtained this result by a two-neutrino oscillation analysis of the L/E distribution using 4.54 kton yr atmospheric neutrino data with the MINOS far detector. 24 Supersolution 4.101-05.
- <sup>24</sup> Supercedes ALIU 05.
- $^{25}$  MICHAEL 06 best fit is for maximal mixing. See also ADAMSON 08.
- <sup>26</sup> The best fit is for maximal mixing.
- <sup>27</sup> ALLISON 05 result is based upon atmospheric neutrino interactions including upwardstopping muons, with an exposure of 5.9 kton yr. From a two-flavor oscillation analysis the best-fit point is  $\Delta m^2 = 0.0017 \text{ eV}^2$  and  $\sin^2(2\theta) = 0.97$ .
- <sup>28</sup>ASHIE 05 obtained this result by a two-neutrino oscillation analysis using 92 kton yr atmospheric neutrino data from the complete Super-Kamiokande I running period.
- $^{29}$  AMBROSIO 04 obtained this result, without using the absolute normalization of the neutrino flux, by combining the angular distribution of upward through-going muon tracks with  $E_{\mu} > 1$  GeV,  $N_{low}$  and  $N_{high}$ , and the numbers of InDown + UpStop and InUp events. Here,  $N_{low}$  and  $N_{high}$  are the number of events with reconstructed neutrino energies < 30 GeV and > 130 GeV, respectively. InDown and InUp represent events with downward and upward-going tracks starting inside the detector due to neutrino interactions, while UpStop represents entering upward-going tracks which stop in the detector. The best fit is for maximal mixing.
- <sup>30</sup>ASHIE 04 obtained this result from the L(flight length)/E(estimated neutrino energy) distribution of  $\nu_{\mu}$  disappearance probability, using the Super-Kamiokande-I 1489 live-day atmospheric neutrino data.
- <sup>31</sup> There are several islands of allowed region from this K2K analysis, extending to high values of  $\Delta m^2$ . We only include the one that overlaps atmospheric neutrino analyses. The best fit is for maximal mixing.
- $^{32}$  AMBROSIO 03 obtained this result on the basis of the ratio R = N\_{low}/N\_{high}, where N<sub>low</sub> and N<sub>high</sub> are the number of upward through-going muon events with reconstructed neutrino energy < 30 GeV and > 130 GeV, respectively. The data came from the full detector run started in 1994. The method of FELDMAN 98 is used to obtain the limits.
- <sup>33</sup> AMBROSIO 03 obtained this result by using the ratio R and the angular distribution of the upward through-going muons. R is given in the previous note and the angular distribution is reported in AMBROSIO 01. The method of FELDMAN 98 is used to obtain the limits. The best fit is to maximal mixing.
- <sup>34</sup>SANCHEZ 03 is based on an exposure of 5.9 kton yr. The result is obtained using a likelihood analysis of the neutrino L/E distribution for a selection  $\mu$  flavor sample while the *e*-flavor sample provides flux normalization. The method of FELDMAN 98 is used to obtain the allowed region. The best fit is  $\sin^2(2\theta) = 0.97$ .
- <sup>35</sup> AMBROSIO 01 result is based on the angular distribution of upward through-going muon tracks with  $E_{\mu} > 1$  GeV. The data came from three different detector configurations, but the statistics is largely dominated by the full detector run, from May 1994 to December 2000. The total live time, normalized to the full detector configuration is 6.17 years. The best fit is obtained outside the physical region. The method of FELDMAN 98 is used to obtain the limits. The best fit is for maximal mixing.
- $^{36}$  AMBROSIO 01 result is based on the angular distribution and normalization of upward through-going muon tracks with  $\textit{E}_{\mu}~>1$  GeV. See the previous footnote.
- <sup>37</sup> FUKUDA 99C obtained this result from a total of 537 live days of upward through-going muon data in Super-Kamiokande between April 1996 to January 1998. With a threshold of  $E_{\mu} > 1.6$  GeV, the observed flux is  $(1.74 \pm 0.07 \pm 0.02) \times 10^{-13}$  cm<sup>-2</sup>s<sup>-1</sup>sr<sup>-1</sup>. The best fit is sin<sup>2</sup>(2 $\theta$ ) = 0.95.
- $^{38}$  FUKUDA 99D obtained this result from a simultaneous fitting to zenith angle distributions of upward-stopping and through-going muons. The flux of upward-stopping muons of minimum energy of 1.6 GeV measured between April 1996 and January 1998 is (0.39  $\pm$  0.04  $\pm$  0.02)  $\times 10^{-13}$  cm $^{-2}$ s $^{-1}$ sr $^{-1}$ . This is compared to the expected flux of (0.73  $\pm$  0.16 (theoretical error))  $\times 10^{-13}$  cm $^{-2}$ s $^{-1}$ sr $^{-1}$ . The best fit is to maximal mixing.

- <sup>39</sup> FUKUDA 99D obtained this result from the zenith dependence of the upwardstopping/through-going flux ratio. The best fit is to maximal mixing.
- <sup>40</sup> FUKUDA 98C obtained this result by an analysis of 33.0 kton yr atmospheric neutrino data. The best fit is for maximal mixing.
- <sup>41</sup> HATAKEYAMA 98 obtained this result from a total of 2456 live days of upward-going muon data in Kamiokande between December 1985 and May 1995. With a threshold of  $E_{\mu} > 1.6 \text{ GeV}$ , the observed flux of upward through-going muons is  $(1.94 \pm 0.10 + 0.07) \times 10^{-13} \text{ cm}^{-2} \text{s}^{-1} \text{sr}^{-1}$ . This is compared to the expected flux of ( $2.46 \pm 0.54$  (theoretical error))  $\times 10^{-13} \text{ cm}^{-2} \text{s}^{-1} \text{sr}^{-1}$ . The best fit is for maximal mixing.
- <sup>42</sup> HATAKEYAMA 98 obtained this result from a combined analysis of Kamiokande contained events (FUKUDA 94) and upward going muon events. The best fit is  $\sin^2(2\theta) = 0.95$ .
- <sup>43</sup> FUKUDA 94 obtained the result by a combined analysis of sub- and multi-GeV atmospheric neutrino events in Kamiokande. The best fit is for maximal mixing.

#### $\Delta m_{32}^2$

The sign of  $\Delta m_{32}^2$  is not known at this time. If given, values are shown separately for the normal and inverted mass ordering. Unless otherwise specified, the ranges below correspond to the projection onto the  $\Delta m_{32}^2$  axis of the 90% CL contours in the sin<sup>2</sup>(2 $\theta_{23}$ ) –  $\Delta m_{32}^2$  plane presented by the authors. If uncertainties are reported with the value, they correspond to one standard deviation uncertainty.

VALUE $(10^{-3} \text{ eV}^2)$	DOCUMENT ID		TECN	COMMENT
$-2.56 \pm 0.04$ OUR FIT	Assuming inverte	d mas	s hierard	chy
2.51 $\pm$ 0.05 OUR FIT	Error includes sc	ale fac	tor of 1.	1. Assuming normal mass hier-
archy	-			
$2.54 \pm 0.08$	<sup>1</sup> ABE	17F	T2K	Normal mass ordering with
$-2.51 \ \pm 0.08$	<sup>1</sup> ABE	17F	T2K	neutrinos and antineutrinos Inverted mass ordering with neutrinos and antineutrinos
$2.67 \pm 0.11$	ADAMSON	17A	NOVA	
$-2.72 \pm 0.11$	ADAMSON	17A	NOVA	$3\nu$ osc; inverted mass ordering
$2.45 \ \pm 0.06 \ \pm 0.06$	<sup>2</sup> AN	17A	DAYA	$3\nu$ osc; normal mass ordering
$-2.56~\pm 0.06~\pm 0.06$	<sup>2</sup> AN	17A	DAYA	$3\nu$ osc; inverted mass ordering
$2.56 \begin{array}{c} +0.21 \\ -0.23 \end{array} \begin{array}{c} +0.12 \\ -0.13 \end{array}$	<sup>3</sup> CHOI	16	RENO	3 u osc; normal mass ordering
$\begin{array}{rrrr} -2.69 & +0.23 & +0.13 \\ & -0.21 & -0.12 \end{array}$	<sup>3</sup> CHOI	16	RENO	$3\nu$ osc; inverted mass ordering
$2.72 \begin{array}{c} +0.19 \\ -0.20 \end{array}$	<sup>4</sup> AARTSEN	15A	ICCB	3 u osc; normal mass ordering
$-2.73 \begin{array}{c} +0.21 \\ -0.18 \end{array}$	<sup>4</sup> AARTSEN	15A	ICCB	3 u osc; inverted mass ordering
2.37 ±0.09	<sup>5</sup> ADAMSON	14	MINS	$3\nu$ osc., accel., atmospheric; normal mass ordering
$-2.41 \begin{array}{c} +0.09 \\ -0.12 \end{array}$	<sup>5</sup> ADAMSON	14	MINS	$3\nu$ osc., accel., atmsopheric; inverted mass ordering
$\bullet \bullet \bullet$ We do not use the f	ollowing data for a	averag	es, fits, l	imits, etc. • • •
$2.53 \begin{array}{c} +0.15 \\ -0.13 \end{array}$	ABE	17C	T2K	normal mass ordering with neutrinos
$2.55 \begin{array}{c} +0.33 \\ -0.27 \end{array}$	ABE	<b>17</b> C	T2K	normal mass ordering with antineutrinos
$2.55 \begin{array}{c} +0.08 \\ -0.08 \end{array}$	ABE	17C	T2K	Normal mass ordering with neutrinos and antineutrinos

$-2.63 \   {+0.08 \atop -0.08}$	ABE	17C	T2K	Inverted mass ordering with neutrinos and antineutrinos
$2.51 \begin{array}{c} +0.29 \\ -0.25 \end{array}$	<sup>6</sup> ABE	<b>16</b> D	T2K	$3\nu$ osc.; normal mass ordering; $\overline{ u}$ beam
$2.52 \begin{array}{c} +0.20 \\ -0.18 \end{array}$	<sup>7</sup> ADAMSON	16A	NOVA	$3\nu$ osc; normal mass ordering
$\begin{array}{rrrr} -2.56 & \pm 0.19 \\ 2.0{-}5.0 \\ 2.37 & \pm 0.11 \\ -2.47 & \pm 0.11 \\ 2.51 & \pm 0.10 \\ -2.56 & \pm 0.10 \end{array}$	<sup>7</sup> ADAMSON <sup>8</sup> AGAFONOVA <sup>9</sup> AN <sup>9</sup> AN <sup>10</sup> ABE <sup>10</sup> ABE	16A 15A 15 15 14 14	NOVA OPER DAYA DAYA T2K T2K	$3\nu$ osc; inverted mass ordering 90% CL, 5 events $3\nu$ osc.; normal mass ordering $3\nu$ osc.; inverted mass ordering $3\nu$ osc.; normal mass ordering $3\nu$ osc.; inverted mass ordering
$2.54 \begin{array}{c} +0.19 \\ -0.20 \end{array}$	<sup>11</sup> AN	14	DAYA	3 u osc.; normal mass ordering
$-2.64 \begin{array}{c} +0.20 \\ -0.19 \end{array}$	<sup>11</sup> AN	14	DAYA	$3\nu$ osc.; inverted mass ordering
$2.48 \begin{array}{c} +0.05 \\ -0.07 \end{array}$	<sup>12</sup> FORERO	14	FIT	$3\nu$ ; normal mass ordering
$-2.38 \begin{array}{c} +0.06 \\ -0.05 \end{array}$	<sup>12</sup> FORERO	14	FIT	$3\nu$ ; inverted mass ordering
2.457±0.047	<sup>13,14</sup> GONZALEZ	14	FIT	Normal mass ordering; global fit
$-2.449\substack{+0.047\\-0.048}$	<sup>13</sup> GONZALEZ	14	FIT	Inverted mass ordering; global fit
$2.3 \begin{array}{c} +0.6 \\ -0.5 \end{array}$	<sup>15</sup> AARTSEN	<b>13</b> B	ICCB	DeepCore, $2\nu$ oscillation
$2.44 \begin{array}{c} +0.17 \\ -0.15 \end{array}$	<sup>16</sup> ABE	13G	T2K	$3\nu$ osc.; normal mass ordering
$2.41 \begin{array}{c} +0.09 \\ -0.10 \end{array}$	<sup>17</sup> ADAMSON	<b>13</b> B	MINS	$2\nu$ osc.; beam + atmospheric; identical $\nu \& \overline{\nu}$
2.2–3.1	<sup>18</sup> ABE	12A	T2K	off-axis beam
$2.62 \begin{array}{c} +0.31 \\ -0.28 \end{array} \pm 0.09$	<sup>19</sup> ADAMSON	12	MINS	$\overline{ u}$ beam
1.35–2.55 1.4–5.6 0.9–2.5 1.8–5.0	20,21 ADAMSON 20,22 ADAMSON 20,22 ADAMSON 23 ADRIAN-MAR	12B 12B 12B 12	MINS MINS MINS ANTR	MINOS atmospheric MINOS pure atmospheric $\nu$ MINOS pure atmospheric $\overline{\nu}$ atm. $\nu$ with deep see tele- scope
1.3-4.0	<sup>24</sup> ABE	11C	SKAM	
$2.32 \begin{array}{c} +0.12 \\ -0.08 \end{array}$	ADAMSON	11	MINS	2 u oscillation; maximal mixing
$3.36 \begin{array}{c} +0.46 \\ -0.40 \end{array}$	<sup>25</sup> ADAMSON	<b>11</b> B	MINS	$\overline{ u}$ beam
< 3.37 1.9-2.6 -1.72.7 2.43 $\pm 0.13$ 0.07-50	<ul> <li><sup>26</sup> ADAMSON</li> <li><sup>27</sup> WENDELL</li> <li><sup>27</sup> WENDELL</li> <li>ADAMSON</li> <li><sup>28</sup> ADAMSON</li> </ul>	11C 10 10 08A 06		MINOS $3\nu$ osc.; normal mass ordering $3\nu$ osc.; inverted mass ordering MINOS atmospheric $\nu$ with far detec-
1.9-4.0 2.2-3.8 1.9-3.6 0.3-12 1.5-3.4 0.6-8.0	29,30 AHN 31 MICHAEL 29 ALIU 32 ALLISON 33 ASHIE 34 AMBROSIO	06A 06 05 05 05 04	K2K MINS K2K SOU2 SKAM MCRO	tor KEK to Super-K MINOS KEK to Super-K atmospheric neutrino MACRO

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1.9 to 3.0	<sup>35</sup> ASHIE	04	SKAM	L/E distribution
1.5–3.9	<sup>36</sup> AHN	03	K2K	KEK to Super-K
0.25-9.0	<sup>37</sup> AMBROSIO	03	MCRO	MACRO
0.6–7.0	<sup>38</sup> AMBROSIO	03	MCRO	MACRO
0.15–15	<sup>39</sup> SANCHEZ	03	SOU2	Soudan-2 Atmospheric
0.6–15	<sup>40</sup> AMBROSIO	01	MCRO	upward $\mu$
1.0-6.0	<sup>41</sup> AMBROSIO	01	MCRO	upward $\mu$
1.0-50	<sup>42</sup> FUKUDA	<b>99</b> C	SKAM	upward $\mu$
1.5-15.0	<sup>43</sup> FUKUDA			upward $\mu$
0.7–18	<sup>44</sup> FUKUDA	<b>99</b> D	SKAM	stop $\mu$ / through
0.5–6.0	<sup>45</sup> FUKUDA	98C	SKAM	Super-Kamiokande
0.55–50	<sup>46</sup> HATAKEYAMA	۹8	KAMI	Kamiokande
4–23	<sup>47</sup> HATAKEYAMA	۹8	KAMI	Kamiokande
5–25	<sup>48</sup> FUKUDA	94	KAMI	Kamiokande
-				

<sup>1</sup>Supersedes ABE 17C.

<sup>2</sup> AN 17A report results from combined rate and spectral shape analysis of 1230 days of data taken with the Daya Bay reactor experiment. The data set contains more than  $2.5 \times 10^6$  inverse beta-decay events with neutron capture on Gd. The fit to the data gives  $\Delta_{ee}^2 = (2.50 \pm 0.06 \pm 0.06) \times 10^{-3}$  eV. Supersedes AN 15.

<sup>3</sup> CHOI 16 reports result of the RENO experiment from a rate and shape analysis of 500 days of data. A simultaneous fit to  $\theta_{13}$  and  $\Delta m_{ee}^2$  yields  $\Delta m_{ee}^2 = (2.62 \substack{+0.21 + 0.12 \\ -0.23 - 0.13}) \times 10^{-3}$  eV. We convert the results to  $\Delta m_{32}^2$  using PDG 14 values of  $\sin^2(\theta_{12})$  and  $\Delta m_{21}^2$ .

<sup>4</sup> AARTSEN 15A obtains this result by a three-neutrino oscillation analysis using 10–100 GeV muon neutrino sample from a total of 953 days of measurements with the low-energy subdetector DeepCore of the IceCube neutrino telescope.

<sup>5</sup> ADAMSON 14 uses a complete set of accelerator and atmospheric data. The analysis combines The analysis combines the  $\nu_{\mu}$  disappearance and  $\nu_{e}$  appearance data using three-neutrino oscillation fit. The fit results are obtained for normal and inverted mass ordering assumptions.

 $^{6}\mathrm{ABE}$  16D reports oscillation results using  $\overline{\nu}_{\mu}$  disappearance in an off-axis beam.

<sup>7</sup> Superseded by ADAMSON 17A.

<sup>8</sup> AGAFONOVA 15A result is based on 5  $\nu_{\mu} \rightarrow \nu_{\tau}$  appearance candidates with an expected background of 0.25  $\pm$  0.05 events. The best fit is for  $\Delta m_{32}^2 = 3.3 \times 10^{-3} \text{ eV}^2$ .

- <sup>9</sup> AN 15 uses all eight identical detectors, with four placed near the reactor cores and the remaining four at the far hall to determine prompt energy spectra. The results correspond to the exposure of  $6.9 \times 10^5$  GW<sub>th</sub>-ton-days. They derive  $\Delta m_{ee}^2 = (2.42 \pm 0.11) \times 10^{-3}$  eV<sup>2</sup>. Assuming the normal (inverted) ordering, the fitted  $\Delta m_{32}^2 = (2.37 \pm 0.11) \times 10^{-3}$  ((2.47 ± 0.11) × 10<sup>-3</sup>) eV<sup>2</sup>. Superseded by AN 17A.
- ((2.47 ± 0.11) × 10<sup>-3</sup>) eV<sup>2</sup>. Superseded by AN 17A. <sup>10</sup> ABE 14 results are based on  $\nu_{\mu}$  disappearance using three-neutrino oscillation fit. The confidence intervals are derived from one dimensional profiled likelihoods. In ABE 14 the inverted mass ordering result is reported as  $\Delta m_{13}^2 = (2.48 \pm 0.10) \times 10^{-3} \text{ eV}^2$  which we converted to  $\Delta m_{32}^2$  by adding PDG 14 value of  $\Delta m_{21}^2 = (7.53 \pm 0.18) \times 10^{-5} \text{ eV}^2$ . Superseded by ABE 17C.
- <sup>11</sup> AN 14 uses six identical detectors, with three placed near the reactor cores (flux-weighted baselines of 512 and 561 m) and the remaining three at the far hall (at the flux averaged distance of 1579 m from all six reactor cores) to determine prompt energy spectra and derive  $\Delta m_{ee}^2 = (2.59^{+0.19}_{-0.20}) \times 10^{-3} \text{ eV}^2$ . Assuming the normal (inverted) ordering, the fitted  $\Delta m_{32}^2 = (2.54^{+0.19}_{-0.20}) \times 10^{-3} ((2.64^{+0.19}_{-0.20}) \times 10^{-3}) \text{ eV}^2$ . Superseded by AN 15.

- $^{12}$  FORERO 14 performs a global fit to  $\Delta m^2_{31}$  using solar, reactor, long-baseline accelerator, and atmospheric neutrino data.
- <sup>13</sup> GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as  $(2.460 \pm 0.046) \times 10^{-3} \text{ eV}^2$  for normal and  $(2.445^{+0.047}_{-0.045}) \times 10^{-3} \text{ eV}^2$  for inverted mass ordering.
- $^{14}\,{\rm The}$  value for normal mass ordering is actually a measurement of  $\Delta m^2_{31}$  which differs from  $\Delta m_{32}^2$  by a much smaller value of  $\Delta m_{12}^2$ .
- $^{15}$  AARTSEN 13B obtained this result by a two-neutrino oscillation analysis using 20–100 GeV muon neutrino sample from a total of 318.9 days of live-time measurement with the low-energy subdetector DeepCore of the IceCube neutrino telescope.
- $^{16}$  Based on the observation of 58  $\nu_{\mu}$  events with 205  $\pm$  17(syst) expected in the absence of neutrino oscillations. Superseded by ABE 14.
- $^{17}$  ADAMSON 13B obtained this result from  $\nu_{\mu}$  and  $\overline{\nu}_{\mu}$  disappearance using  $\nu_{\mu}$  (10.71  $\times$  $10^{20}$  POT) and  $\overline{
  u}_{\mu}$  (3.36 imes 10<sup>20</sup> POT) beams, and atmospheric (37.88 kton-years) data from MINOS. The fit assumed two-flavor neutrino hypothesis and identical  $u_{\mu}$  and  $\overline{
  u}_{\mu}$ oscillation parameters.
- $^{18}$  ABE 12A obtained this result by a two-neutrino oscillation analysis. The best-fit point is  $\Delta m^2_{32}=2.65\times 10^{-3}~eV^2.$
- $^{19}$  ADAMSON 12 is a two-neutrino oscillation analysis using antineutrinos.
- $^{20}$  ADAMSON 12B obtained this result by a two-neutrino oscillation analysis of the L/E distribution using 37.9 kton yr atmospheric neutrino data with the MINOS far detector.
- $^{21}$  The 90% single-parameter confidence interval at the best fit point is  $\Delta {
  m m}^2 =$  0.0019  $\pm$  $0.0004 \text{ eV}^2$
- <sup>22</sup> The data are separated into pure samples of  $\nu$ s and  $\overline{\nu}$ s, and separate oscillation parameters for  $\nu$ s and  $\overline{\nu}$ s are fit to the data. The best fit point is  $(\Delta m^2, \sin^2 2\theta) = (0.0022 \text{ eV}^2, \sin^2 2\theta)$ 0.99) and  $(\Delta \overline{m}^2, \sin^2 2\overline{\theta}) = (0.0016 \text{ eV}^2, 1.00)$ . The quoted result is taken from the 90% C.L. contour in the ( $\Delta m^2$ , sin<sup>2</sup>2 $\theta$ ) plane obtained by minimizing the four parameter log-likelihood function with respect to the other oscillation parameters.
- <sup>23</sup> ADRIAN-MARTINEZ 12 measured the oscillation parameters of atmospheric neutrinos with the ANTARES deep sea neutrino telescope using the data taken from 2007 to 2010 (863 days of total live time).
- <sup>24</sup> ABE 11C obtained this result by a two-neutrino oscillation analysis with separate mixing parameters between neutrinos and antineutrinos, using the Super-Kamiokande-I+II+III atmospheric neutrino data. The corresponding 90% CL neutrino oscillation parameter range obtained from this analysis is  $\Delta m^2 = 1.7 - 3.0 \times 10^{-3} \text{ eV}^2$ .
- <sup>25</sup> ADAMSON 11B obtained this result by a two-neutrino oscillation analysis of antineutrinos in an antineutrino enhanced beam with  $1.71 \times 10^{20}$  protons on target. This results is consistent with the neutrino measurements of ADAMSON 11 at 2% C.L.
- $^{26}$  ADAMSON 11C obtains this result based on a study of antineutrinos in a neutrino beam and assumes maximal mixing in the two-flavor approximation.
- $^{27}$  WENDELL 10 obtained this result by a three-neutrino oscillation analysis with one mass scale dominance ( $\Delta m^2_{21} = 0$ ) using the Super-Kamiokande-I+II+III atmospheric neutrino data, and updates the HOSAKA 06A result.
- $^{28}$  ADAMSON 06 obtained this result by a two-neutrino oscillation analysis of the L/E distribution using 4.54 kton yr atmospheric neutrino data with the MINOS far detector. <sup>29</sup> The best fit in the physical region is for  $\Delta m^2 = 2.8 \times 10^{-3} \text{ eV}^2$ .
- <sup>30</sup> Supercedes ALIU 05.
- <sup>31</sup> MICHAEL 06 best fit is  $2.74 \times 10^{-3}$  eV<sup>2</sup>. See also ADAMSON 08.
- $^{32}$  ALLISON 05 result is based on an atmospheric neutrino observation with an exposure of 5.9 kton yr. From a two-flavor oscillation analysis the best-fit point is  $\Delta m^2 = 0.0017$  $eV^2$  and  $sin^2 2 \theta = 0.97$ .

- <sup>33</sup>ASHIE 05 obtained this result by a two-neutrino oscillation analysis using 92 kton yr atmospheric neutrino data from the complete Super-Kamiokande I running period. The best fit is for  $\Delta m^2 = 2.1 \times 10^{-3} \text{ eV}^2$ .
- <sup>34</sup> AMBROSIO 04 obtained this result, without using the absolute normalization of the neutrino flux, by combining the angular distribution of upward through-going muon tracks with  $E_{\mu} > 1$  GeV,  $N_{low}$  and  $N_{high}$ , and the numbers of InDown + UpStop and InUp events. Here,  $N_{low}$  and  $N_{high}$  are the number of events with reconstructed neutrino energies < 30 GeV and > 130 GeV, respectively. InDown and InUp represent events with downward and upward-going tracks starting inside the detector due to neutrino interactions, while UpStop represents entering upward-going tracks which stop in the detector. The best fit is for  $\Delta m^2 = 2.3 \times 10^{-3} \text{ eV}^2$ .
- $^{35}$  ASHIE 04 obtained this result from the L(flight length)/E(estimated neutrino energy) distribution of  $\nu_{\mu}$  disappearance probability, using the Super-Kamiokande-I 1489 live-day
- atmospheric neutrino data. The best fit is for  $\Delta m^2 = 2.4 imes 10^{-3} ext{ eV}^2$ .
- <sup>36</sup> There are several islands of allowed region from this K2K analysis, extending to high values of  $\Delta m^2$ . We only include the one that overlaps atmospheric neutrino analyses. The best fit is for  $\Delta m^2 = 2.8 \times 10^{-3} \text{ eV}^2$ .
- <sup>37</sup> AMBROSIO 03 obtained this result on the basis of the ratio R = N<sub>low</sub>/N<sub>high</sub>, where N<sub>low</sub> and N<sub>high</sub> are the number of upward through-going muon events with reconstructed neutrino energy < 30 GeV and > 130 GeV, respectively. The data came from the full detector run started in 1994. The method of FELDMAN 98 is used to obtain the limits. The best fit is for  $\Delta m^2 = 2.5 \times 10^{-3} \text{ eV}^2$ .
- <sup>38</sup> AMBROSIO 03 obtained this result by using the ratio R and the angular distribution of the upward through-going muons. R is given in the previous note and the angular distribution is reported in AMBROSIO 01. The method of FELDMAN 98 is used to obtain the limits. The best fit is for  $\Delta m^2 = 2.5 \times 10^{-3} \text{ eV}^2$ .
- <sup>39</sup>SANCHEZ 03 is based on an exposure of 5.9 kton yr. The result is obtained using a likelihood analysis of the neutrino L/E distribution for a selection  $\mu$  flavor sample while the *e*-flavor sample provides flux normalization. The method of FELDMAN 98 is used to obtain the allowed region. The best fit is for  $\Delta m^2 = 5.2 \times 10^{-3} \text{ eV}^2$ .
- $^{40}$  AMBROSIO 01 result is based on the angular distribution of upward through-going muon tracks with  $E_{\mu} > 1$  GeV. The data came from three different detector configurations, but the statistics is largely dominated by the full detector run, from May 1994 to December 2000. The total live time, normalized to the full detector configuration is 6.17 years. The best fit is obtained outside the physical region. The method of FELDMAN 98 is used to obtain the limits.
- <sup>41</sup>AMBROSIO 01 result is based on the angular distribution and normalization of upward through-going muon tracks with  $E_{\mu} > 1$  GeV. See the previous footnote.
- $^{42}$  FUKUDA 99C obtained this result from a total of 537 live days of upward through-going muon data in Super-Kamiokande between April 1996 to January 1998. With a threshold of  $E_{\mu}~>1.6$  GeV, the observed flux is  $(1.74\pm0.07\pm0.02)\times10^{-13}~{\rm cm}^{-2}{\rm s}^{-1}{\rm sr}^{-1}$ .

The best fit is for  $\Delta m^2 = 5.9 \times 10^{-3} \text{ eV}^2$ .

<sup>43</sup> FUKUDA 99D obtained this result from a simultaneous fitting to zenith angle distributions of upward-stopping and through-going muons. The flux of upward-stopping muons of minimum energy of 1.6 GeV measured between April 1996 and January 1998 is  $(0.39 \pm 0.04 \pm 0.02) \times 10^{-13} \text{ cm}^{-2}\text{s}^{-1}\text{sr}^{-1}$ . This is compared to the expected flux of  $(0.73 \pm 0.16 \text{ (theoretical error)}) \times 10^{-13} \text{ cm}^{-2}\text{s}^{-1}\text{sr}^{-1}$ . The best fit is for  $\Delta m^2 = 3.9 \times 10^{-3} \text{ eV}^2$ .

<sup>44</sup> FUKUDA 99D obtained this result from the zenith dependence of the upwardstopping/through-going flux ratio. The best fit is for  $\Delta m^2 = 3.1 \times 10^{-3} \text{ eV}^2$ .

- <sup>45</sup> FUKUDA 98C obtained this result by an analysis of 33.0 kton yr atmospheric neutrino data. The best fit is for  $\Delta m^2 = 2.2 \times 10^{-3} \text{ eV}^2$ .
- <sup>46</sup> HATAKEYAMA 98 obtained this result from a total of 2456 live days of upward-going muon data in Kamiokande between December 1985 and May 1995. With a threshold of

 $E_{\mu} > 1.6 \text{ GeV}$ , the observed flux of upward through-going muons is  $(1.94 \pm 0.10 + 0.07) \times 10^{-13} = -2^{-1} = 1^{-1}$ . This is a standard through going muons is  $(1.94 \pm 0.10 + 0.07) \times 10^{-13}$ .

 $^{\mu}$  10<sup>-13</sup> cm<sup>-2</sup>s<sup>-1</sup>sr<sup>-1</sup>. This is compared to the expected flux of (2.46±0.54 (theoretical error)) × 10<sup>-13</sup> cm<sup>-2</sup>s<sup>-1</sup>sr<sup>-1</sup>. The best fit is for  $\Delta m^2 = 2.2 \times 10^{-3}$  eV<sup>2</sup>.

- <sup>47</sup> HATAKEYAMA 98 obtained this result from a combined analysis of Kamiokande contained events (FUKUDA 94) and upward going muon events. The best fit is for  $\Delta m^2 = 13 \times 10^{-3} \text{ eV}^2$ .
- <sup>48</sup> FUKUDA 94 obtained the result by a combined analysis of sub- and multi-GeV atmospheric neutrino events in Kamiokande. The best fit is for  $\Delta m^2 = 16 \times 10^{-3} \text{ eV}^2$ .

#### $\sin^2(\theta_{13})$

At present time direct measurements of  $\sin^2(\theta_{13})$  are derived from the reactor  $\overline{\nu}_e$  disappearance at distances corresponding to the  $\Delta m^2_{32}$  value, i.e. L  $\sim 1$ km. Alternatively, limits can also be obtained from the analysis of the solar neutrino data and accelerator-based  $\nu_{\mu} \rightarrow \nu_e$  experiments.

If an experiment reports  $\sin^2(2\theta_{13})$  we convert the value to  $\sin^2(\theta_{13})$ .

<u>VALUE</u> (units 10 <sup>-2</sup> ) <u>CL%</u>	DOCUMENT ID		TECN	COMMENT
2.12 $\pm$ 0.08 OUR AVERAGE				
$2.149 \pm 0.071 \pm 0.050$	<sup>1</sup> AN	17A	DAYA	DayaBay, LingAo/Ao II reactors
$2.25 \begin{array}{c} + & 0.87 \\ - & 0.86 \end{array}$	<sup>2</sup> ABE	<b>16</b> B	DCHZ	Chooz reactors
$1.81~\pm~0.29$	<sup>3</sup> AN	16A	DAYA	DayaBay, Ling Ao/Ao II reactors
$2.09~\pm~0.23~\pm0.16$	<sup>4</sup> CHOI	16		Yonggwang reactors
• • • We do not use the following		fits,	limits, et	.c. ● ● ●
$2.7$ $\pm$ 0.7	<sup>5</sup> ABE	17F	T2K	Normal mass ordering, T2K only
$2.15~\pm~0.13$	<sup>6</sup> AN	15	DAYA	DayaBay, Ling Ao/Ao II reactors
$2.6  \begin{array}{c} +  1.2 \\ -  1.1 \end{array}$	<sup>7</sup> ABE	14A	DCHZ	Chooz reactors
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	<sup>8</sup> ABE	14C	T2K	Inverted mass ordering
$3.6  \begin{array}{c} +  1.0 \\ -  0.9 \end{array}$	<sup>8</sup> ABE	14C	T2K	Normal mass ordering
$2.3  \begin{array}{c} + & 0.9 \\ - & 0.8 \end{array}$	<sup>9</sup> ABE	14H	DCHZ	Chooz reactors
$2.3 \pm 0.2$	<sup>10</sup> AN	14	DAYA	DayaBay, Ling Ao/Ao
$2.12~\pm~0.47$	$^{11}$ AN	<b>14</b> B	DAYA	II reactors DayaBay, Ling Ao/Ao II reactors
$2.34~\pm~0.20$	<sup>12</sup> FORERO	14	FIT	Normal mass ordering
$2.40~\pm~0.19$	<sup>12</sup> FORERO	14	FIT	Inverted mass ordering
$2.18~\pm~0.10$	<sup>13</sup> GONZALEZ	14	FIT	Normal mass ordering; global fit
$2.19 \ + \ 0.11 \\ - \ 0.10$	<sup>13</sup> GONZALEZ	14	FIT	Inverted mass ordering; global fit
$2.5 ~\pm~ 0.9 ~\pm 0.9$	<sup>14</sup> ABE	13C	DCHZ	Chooz reactors
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	<sup>15</sup> ABE	13E	T2K	Normal mass ordering
$2.8 \ + \ 1.6 \ - \ 1.2$	<sup>15</sup> ABE	13E	T2K	Inverted mass ordering
$1.6  \begin{array}{c} +  1.3 \\ -  0.9 \end{array}$	<sup>16</sup> ADAMSON	13A	MINS	Normal mass ordering
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3.0	$^{+}$ 1.8 $^{-}$ 1.6			<sup>16</sup> ADAMSON	13A	MINS	Inverted mass ordering
<13			90	AGAFONOVA	13	OPER	OPERA: $3\nu$
< 3.6			95	<sup>17</sup> AHARMIM	13	FIT	global solar: $3\nu$
2.3	$\pm$ 0.3	$\pm 0.1$		<sup>18</sup> AN	13	DAYA	DayaBay, LIng Ao/Ao II reactors
2.2	$\pm$ 1.1	$\pm 0.8$		<sup>19</sup> ABE	12	DCHZ	Chooz reactors
2.8	$\pm$ 0.8	$\pm 0.7$		<sup>20</sup> ABE	12B	DCHZ	Chooz reactors
2.9	$\pm$ 0.3	$\pm 0.5$		<sup>21</sup> AHN	12	RENO	Yonggwang reactors
2.4	± 0.4	$\pm 0.1$		<sup>22</sup> AN	12	DAYA	DayaBay, Ling Ao/Ao II reactors
2.5	$\begin{array}{r}+ 1.8\\- 1.6\end{array}$			<sup>23</sup> ABE	11	FIT	KamLAND + global solar
< 6.1			95	<sup>24</sup> ABE	11	FIT	Global solar
1.3	to 5.6		68	<sup>25</sup> ABE	11A	T2K	Normal mass ordering
1.5	to 5.6		68	<sup>26</sup> ABE	11A	T2K	Inverted mass ordering
0.3	to 2.3		68	<sup>27</sup> ADAMSON	11D	MINS	Normal mass ordering
0.8	to 3.9		68	28 ADAMSON	11D	MINS	Inverted mass ordering
8	± 3			<sup>29</sup> FOGLI	11	FIT	Global neutrino data
7.8	± 6.2			<sup>30</sup> GANDO	11	FIT	KamLAND + solar: $3\nu$
12.4	$\pm 13.3$			<sup>31</sup> GANDO	11	FIT	KamLAND: $3\nu$
3	+ 9 - 7		90	<sup>32</sup> ADAMSON	10A	MINS	Normal mass ordering
6	$^{+14}_{-6}$		90	<sup>33</sup> ADAMSON	10A	MINS	Inverted mass ordering
8	$^{+}$ 8 $^{-}$ 7			<sup>34,35</sup> AHARMIM	10	FIT	${\sf KamLAND}+{\sf global}$ solar: 3 $ u$
< 30			95	<sup>34,36</sup> AHARMIM	10	FIT	global solar: $3\nu$
< 15			90	<sup>37</sup> WENDELL	10	SKAM	$3\nu$ osc.; normal <i>m</i> or- dering
< 33			90	<sup>37</sup> WENDELL	10	SKAM	$3\nu$ osc.; inverted <i>m</i> ordering
11	$^{+11}_{-8}$			<sup>38</sup> ADAMSON	09	MINS	Normal mass ordering
18	$^{+15}_{-11}$			<sup>39</sup> ADAMSON	09	MINS	Inverted mass ordering
6	$\pm$ 4			<sup>40</sup> FOGLI	08	FIT	Global neutrino data
8	$\pm$ 7			<sup>41</sup> FOGLI	08	FIT	Solar + KamLAND
5	$\pm$ 5			<sup>42</sup> FOGLI	08	FIT	data Atmospheric + LBL + CHOOZ
< 36			90	<sup>43</sup> YAMAMOTO	06	K2K	Accelerator experiment
< 48			90	<sup>44</sup> AHN	04	K2K	Accelerator experiment
< 36			90	<sup>45</sup> BOEHM	01		Palo Verde react.
< 45			90	46 BOEHM	00		Palo Verde react.
< 15			90	<sup>47</sup> APOLLONIO	99	CHOZ	Reactor Experiment

<sup>1</sup> AN 17A report results from combined rate and spectral shape analysis of 1230 days of data taken with the Daya Bay reactor experiment. The data set contains more than  $2.5 \times 10^6$  inverse beta-decay events with neutron capture on Gd. A simultaneous fit to  $\theta_{13}$  and  $\Delta m^2_{ee}$  is performed. Supersedes AN 15.

 $^2$  ABE 16B uses 455.57 live days of data from a detector 1050 m away from two reactor cores of the Chooz nuclear power station, to determine the mixing parameter  $\sin^2(2\theta_{13})$ . This analysis uses 7.15 reactor-off days for constraining backgrounds. A rate and shape analysis is performed on combined neutron captures on H and Gd. Supersedes ABE 14H and ABE 13C.

- <sup>3</sup> AN 16A uses data from the eight antineutrino detectors (404 days) and six antineutrino detectors (217 days) runs to determine the mixing parameter  $\sin^2(2\theta_{13})$  using the neutron capture on H only. Supersedes AN 14B.
- <sup>4</sup> CHOI 16 reports results of the RENO experiment using about 500 days of data, performing a rate and shape analysis. Compared to AHN 12, a significant reduction of the systematic uncertainties is reported. A 3% excess of events near 5 MeV of the prompt energy is observed. Supersedes AHN 12.
- <sup>5</sup> Using T2K data only. For inverted mass ordering, all values of  $\theta_{13}$  are ruled out at 68% CL.
- $^{6}$ AN 15 uses all eight identical detectors, with four placed near the reactor cores and the remaining four at the far hall to determine the mixing angle  $\theta_{13}$  using the  $\overline{\nu}_{e}$  observed interaction rates with neutron capture on Gd and energy spectra. The result corresponds to the exposure of  $6.9 \times 10^{5}$  GW<sub>th</sub>-ton-days. Superseded by AN 17A.
- <sup>7</sup> ABE 14A uses 467.9 live days of one detector, 1050 m away from two reactor cores of the Chooz nuclear power station, to determine the mixing parameter  $\sin^2(2 \theta_{13})$ . The Bugey4 data (DECLAIS 94) is used to constrain the neutrino flux. The data set includes 7.24 reactor-off days. A "rate-modulation" analysis is performed. Supercedes ABE 12B.
- 7.24 reactor-off days. A "rate-modulation" analysis is performed. Supercedes ABE 12B. <sup>8</sup> ABE 14C result is for  $\nu_e$  appearance and assumes  $\Delta m_{32}^2 = 2.4 \times 10^{-3} \text{ eV}^2$ ,  $\sin^2(\theta_{23}) = 0.5$ , and  $\delta = 0$ .
- <sup>9</sup> ABE 14H uses 467.9 live days of one detector, 1050 m away from two reactor cores of the Chooz nuclear power station, to determine the mixing parameter  $\sin^2(2 \theta_{13})$ . The Bugey4 data (DECLAIS 94) is used to constrain the neutrino flux. The data set includes 7.24 reactor-off days. A rate and shape analysis is performed. Superceded by ABE 16B.
- <sup>10</sup> AN 14 uses six identical detectors, with three placed near the reactor cores (flux-weighted baselines of 512 and 561 m) and the remaining three at the far hall (at the flux averaged distance of 1579 m from all six reactor cores) to determine the mixing angle  $\theta_{13}$  using the  $\overline{\nu}_e$  observed interaction rates with neutron capture on Gd and energy spectra. Supersedes AN 13 and superseded by AN 15.
- <sup>11</sup> AN 14B uses six identical anti-neutrino detectors with flux-weighted baselines of  $\sim$  500 m and  $\sim$  1.6 km to six power reactors. This rate analysis uses a 217-day data set and neutron capture on protons (not Gd) only.  $\Delta m_{31}^2 = 2.32 \times 10^{-3} \text{ eV}^2$  is assumed. Superseded by AN 16A.
- <sup>12</sup> FORERO 14 performs a global fit to neutrino oscillations using solar, reactor, longbaseline accelerator, and atmospheric neutrino data.
- $^{13}$  GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as  $(2.18 \substack{+0.10 \\ -0.11}) \times 10^{-2} \ \text{eV}^2$  for normal and  $(2.19 \substack{+0.12 \\ -0.10}) \times 10^{-2} \ \text{eV}^2$  for inverted mass ordering.
- <sup>14</sup> ABE 13C uses delayed neutron capture on hydrogen instead of on Gd used previously. The physical volume is thus three times larger. The fit is based on the rate and shape analysis as in ABE 12B. The Bugey4 data (DECLAIS 94) is used to constrain the neutrino flux. Superseded by ABE 16B.
- $^{15}\mathrm{ABE}$  13E assumes maximal  $\theta_{23}$  mixing and CP phase  $\delta$  = 0.
- <sup>16</sup> ADAMSON 13A results obtained from  $\nu_e$  appearance, assuming  $\delta = 0$ , and sin<sup>2</sup>(2  $\theta_{23}$ ) = 0.957.
- <sup>17</sup> AHARMIM 13 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{32}^2$  fixed to 2.45 × 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino data. AHARMIM 13 global solar neutrino data include SNO's all-phases-combined analysis results on the total active <sup>8</sup>B neutrino flux and energy-dependent  $\nu_e$  survival probability parameters, measurements of Cl (CLEVELAND 98), Ga (ABDURASHITOV 09 which contains combined analysis with GNO (ALTMANN 05 and Ph.D. thesis of F. Kaether)), and <sup>7</sup>Be (BELLINI 11A) rates, and <sup>8</sup>B solar-neutrino recoil electron measurements of SK-I (HOSAKA 06) zenith, SK-II (CRAVENS 08) and SK-III (ABE 11) day/night spectra,

and Borexino (BELLINI 10A) spectra. AHARMIM 13 also reported a result combining global solar and KamLAND data, which is  $\sin^2(2 \theta_{13}) = (9.1 + 2.9) \times 10^{-2}$ .

- <sup>18</sup> AN 13 uses six identical detectors, with three placed near the reactor cores (flux-weighted baselines of 498 and 555 m) and the remaining three at the far hall (at the flux averaged distance of 1628 m from all six reactor cores) to determine the  $\overline{\nu}_e$  interaction rate ratios. Superseded by AN 14.
- <sup>19</sup> ABE 12 determines the  $\overline{\nu}_e$  interaction rate in a single detector, located 1050 m from the cores of two reactors. A rate and shape analysis is performed. The rate normalization is fixed by the results of the Bugey4 reactor experiment, thus avoiding any dependence on possible very short baseline oscillations. The value of  $\Delta m_{31}^2 = 2.4 \times 10^{-3} \text{ eV}^2$  is used in the analysis. Superseded by ABE 12B.
- <sup>20</sup>ABE 12B determines the neutrino mixing angle  $\theta_{13}$  using a single detector, located 1050 m from the cores of two reactors. This result is based on a spectral shape and rate analysis. The Bugey4 data (DECLAIS 94) is used to constrain the neutrino flux. Superseded by ABE 14A.
- <sup>21</sup> AHN 12 uses two identical detectors, placed at flux weighted distances of 408.56 m and 1433.99 m from six reactor cores, to determine the mixing angle  $\theta_{13}$ . This rate-only analysis excludes the no-oscillation hypothesis at 4.9 standard deviations. The value of  $\Delta m_{31}^2 = (2.32^{+0.12}_{-0.08}) \times 10^{-3} \text{ eV}^2$  was assumed in the analysis. Superseded by CHOI 16.
- $^{22}$  AN 12 uses six identical detectors with three placed near the reactor cores (flux-weighted baselines of 470 m and 576 m) and the remaining three at the far hall (at the flux averaged distance of 1648 m from all six reactor cores) to determine the mixing angle  $\theta_{13}$  using the  $\overline{\nu}_e$  observed interaction rate ratios. This rate-only analysis excludes the no-oscillation hypothesis at 5.2 standard deviations. The value of  $\Delta m^2_{31} = (2.32^{+0.12}_{-0.08}) \times 10^{-3} \ {\rm eV}^2$  was assumed in the analysis. Superseded by AN 13.
- $^{23}$  ABE 11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m^2_{32}$  fixed to  $2.4\times 10^{-3}~{\rm eV^2}$ , using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, GALLEX/GNO, SAGE, and KamLAND data. This result implies an upper bound of  $\sin^2\!\theta_{13} < 0.059~(95\%~{\rm CL})$  or  $\sin^2\!2\theta_{13} < 0.22~(95\%~{\rm CL})$ . The normal neutrino mass ordering and CPT invariance are assumed.
- <sup>24</sup> ABE 11 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{32}^2$  fixed to 2.4 × 10<sup>-3</sup> eV<sup>2</sup>, using solar neutrino data including Super-Kamiokande, SNO, Borexino (ARPESELLA 08A), Homestake, and GALLEX/GNO data. The normal neutrino mass ordering is assumed.
- <sup>25</sup> The quoted limit is for  $\Delta m_{32}^2 = 2.4 \times 10^{-3} \text{ eV}^2$ ,  $\theta_{23} = \pi/2$ ,  $\delta = 0$ , and the normal mass ordering. For other values of  $\delta$ , the 68% region spans from 0.03 to 0.25, and the 90% region from 0.02 to 0.32.
- <sup>26</sup> The quoted limit is for  $\Delta m_{32}^2 = 2.4 \times 10^{-3} \text{ eV}^2$ ,  $\theta_{23} = \pi/2$ ,  $\delta = 0$ , and the inverted mass ordering. For other values of  $\delta$ , the 68% region spans from 0.04 to 0.30, and the 90% region from 0.02 to 0.39.
- <sup>27</sup> The quoted limit is for  $\Delta m_{32}^2 = 2.32 \times 10^{-3} \text{ eV}^2$ ,  $\theta_{23} = \pi/2$ ,  $\delta = 0$ , and the normal mass ordering. For other values of  $\delta$ , the 68% region spans from 0.02 to 0.12, and the 90% region from 0 to 0.16.
- <sup>28</sup> The quoted limit is for  $\Delta m_{32}^2 = 2.32 \times 10^{-3} \text{ eV}^2$ ,  $\theta_{23} = \pi/2$ ,  $\delta = 0$ , and the inverted mass ordering. For other values of  $\delta$ , the 68% region spans from 0.02 to 0.16, and the 90% region from 0 to 0.21.
- $^{29}$  FOGLI 11 obtained this result from an analysis using the atmospheric, accelerator long baseline, CHOOZ, solar, and KamLAND data. Recently, MUELLER 11 suggested an average increase of about 3.5% in normalization of the reactor  $\overline{\nu}_e$  fluxess, and using these fluxes, the fitted result becomes 0.10  $\pm$  0.03.
- $^{30}$  GANDO 11 report  $\sin^2\theta_{13} = 0.020 \pm 0.016$ . This result was obtained with three-neutrino fit using the KamLAND + solar data.

- <sup>31</sup> GANDO 11 report  $\sin^2 \theta_{13} = 0.032 \pm 0.037$ . This result was obtained with three-neutrino fit using the KamLAND data only.
- $^{32}$  This result corresponds to the limit of <0.12 at 90% CL for  $\Delta m^2_{32} = 2.43 imes 10^{-3}$  eV<sup>2</sup>,  $\theta_{23} = \pi/2$ , and  $\delta = 0$ . For other values of  $\delta$ , the 90% CL region spans from 0 to 0.16.
- $^{33}$  This result corresponds to the limit of <0.20 at 90% CL for  $\Delta m^2_{32} = 2.43 imes 10^{-3}$  eV<sup>2</sup>,  $\theta_{23} = \pi/2$ , and  $\delta = 0$ . For other values of  $\delta$ , the 90% CL region spans from 0 to 0.21.
- $^{34}$ AHARMIM 10 global solar neutrino data include SNO's low-energy-threshold analysis survival probability day/night curves, SNO Phase III integral rates (AHARMIM 08), CI (CLEVELAND 98), SAGE (ABDURASHITOV 09), Gallex/GNO (HAMPEL 99, ALT-MANN 05), Borexino (ARPESELLA 08A), SK-I zenith (HOSAKA 06), and SK-II day/night spectra (CRAVENS 08).
- $^{35}$  AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{31}^2$  fixed to  $2.3 \times 10^{-3}$  eV<sup>2</sup>, using global solar neutrino data and KamLAND data (ABE 08A). CPT invariance is assumed. This result implies an upper bound of  $\sin^2\theta_{13} <$ 0.057 (95% CL) or  $\sin^2 2\theta_{13} < 0.22$  (95% CL).
- <sup>36</sup> AHARMIM 10 obtained this result by a three-neutrino oscillation analysis with the value of  $\Delta m_{31}^2$  fixed to 2.3 × 10<sup>-3</sup> eV<sup>2</sup>, using global solar neutrino data.
- $^{37}$  WENDELL 10 obtained this result by a three-neutrino oscillation analysis with one mass scale dominance ( $\Delta m_{21}^2 = 0$ ) using the Super-Kamiokande-I+II+III atmospheric neutrino data, and updates the HOSAKA 06A result.
- $^{38}$  The quoted limit is for  $\Delta m^2_{32}=2.43 imes 10^{-3}$  eV $^2$ ,  $heta_{23}=\pi/2$ , and  $\delta=0$ . For other
- values of  $\delta$ , the 68% CL region spans from 0.02 to 0.26. <sup>39</sup> The quoted limit is for  $\Delta m_{32}^2 = 2.43 \times 10^{-3} \text{ eV}^2$ ,  $\theta_{23} = \pi/2$ , and  $\delta = 0$ . For other values of  $\delta$ , the 68% CL region spans from 0.04 to 0.34.
- $^{40}$  FOGLI 08 obtained this result from a global analysis of all neutrino oscillation data, that is, solar + KamLAND + atmospheric + accelerator long baseline + CHOOZ.
- $^{41}$  FOGLI 08 obtained this result from an analysis using the solar and KamLAND neutrino oscillation data.
- <sup>42</sup> FOGLI 08 obtained this result from an analysis using the atmospheric, accelerator long baseline, and CHOOZ neutrino oscillation data.
- <sup>43</sup>YAMAMOTO 06 searched for  $u_{\mu} \rightarrow 
  u_{e}$  appearance. Assumes 2  $\sin^{2}(2\theta_{\mu e}) =$  $\sin^2(2 heta_{13})$ . The quoted limit is for  $\Delta m^2_{32} = 1.9 imes 10^{-3}$  eV<sup>2</sup>. That value of  $\Delta m^2_{32}$  is the one- $\sigma$  low value for AHN 06A. For the AHN 06A best fit value of  $2.8 \times 10^{-3}$  eV<sup>2</sup>, the sin<sup>2</sup>( $2\theta_{13}$ ) limit is < 0.26. Supersedes AHN 04.
- <sup>44</sup> AHN 04 searched for  $\nu_{\mu} \rightarrow \nu_{e}$  appearance. Assuming 2 sin<sup>2</sup>(2  $\theta_{\mu_{e}}$ ) = sin<sup>2</sup>(2  $\theta_{13}$ ), a limit on  $\sin^2(2 \theta_{\mu_e})$  is converted to a limit on  $\sin^2(2 \theta_{13})$ . The quoted limit is for  $\Delta m_{32}^2$ = 1.9 imes 10  $^{-3}$  eV  $^2$ . That value of  $\Delta m^2_{32}$  is the one- $\sigma$  low value for ALIU 05. For the
- ALIU 05 best fit value of  $2.8 \times 10^{-3} \text{ eV}^2$ , the  $\sin^2(2 \theta_{13})$  limit is < 0.30. <sup>45</sup> The quoted limit is for  $\Delta m_{32}^2 = 1.9 \times 10^{-3} \text{ eV}^2$ . That value of  $\Delta m_{32}^2$  is the 1- $\sigma$  low value for ALIU 05. For the ALIU 05 best fit value of  $2.8 \times 10^{-3}$  eV<sup>2</sup>, the sin<sup>2</sup>2  $\theta_{13}$  limit is < 0.19. In this range, the  $\theta_{13}$  limit is larger for lower values of  $\Delta m^2_{32}$ , and smaller for higher values of  $\Delta m_{32}^2$ .
- <sup>46</sup> The quoted limit is for  $\Delta m^2_{32} = 1.9 \times 10^{-3}$  eV<sup>2</sup>. That value of  $\Delta m^2_{32}$  is the 1- $\sigma$  low value for ALIU 05. For the ALIU 05 best fit value of  $2.8 \times 10^{-3}$  eV<sup>2</sup>, the sin<sup>2</sup>2  $\theta_{13}$
- limit is < 0.23. 47 The quoted limit is for  $\Delta m_{32}^2 = 2.43 \times 10^{-3} \text{ eV}^2$ . That value of  $\Delta m_{32}^2$  is the central value for ADAMSON 08. For the ADAMSON 08 1- $\sigma$  low value of 2.30  $\times$  10<sup>-3</sup> eV<sup>2</sup>, the sin<sup>2</sup>2  $\theta_{13}$  limit is < 0.16. See also APOLLONIO 03 for a detailed description of the experiment.

#### — CP violating phase -

### $\delta$ , *CP* violating phase

Measurements of  $\delta$  come from atmospheric and accelarator experiments looking at  $\nu_e$  appearance. We encode values between 0 and  $2\pi$ , though it is equivalent to use  $-\pi$  to  $\pi$ 

to $\pi$ . VALUE ( $\pi$ rad)	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do no	ot use t	the following data fo	r aver	ages, fits	s, limits, etc. ● ● ●
$1.45 \substack{+0.27 \\ -0.26}$		ABE	17F	T2K	Normal mass ordering
$1.54 \substack{+0.22 \\ -0.23}$		ABE	17F	T2K	Inverted mass ordering
$1.50^{+0.53}_{-0.57}$		<sup>1</sup> ADAMSON	<b>17</b> B	NOVA	Inverted mass ordering for angle second quadrant
$0.74 \substack{+0.57 \\ -0.93}$		<sup>1</sup> ADAMSON	<b>17</b> B	NOVA	Normal mass ordering for angle second quadrant
$1.48 \substack{+0.69 \\ -0.58}$		<sup>1</sup> ADAMSON	<b>17</b> B	NOVA	Normal mass ordering for angle first quadrant
0.0 to 0.1, 0.5	90	<sup>1,2</sup> ADAMSON	16	NOVA	Inverted mass ordering
to 2.0 0.0 to 2.0	90	<sup>2</sup> ADAMSON	16	NOVA	Normal mass ordering
0 to 0.15, 0.83 to 2	90	ABE	<b>15</b> D	T2K	Normal mass ordering
1.09 to 1.92	90	ABE	<b>15</b> D	T2K	Inverted mass ordering
0.05 to 1.2	90	<sup>3</sup> ADAMSON	14	MINS	Normal mass ordering
$1.34 \substack{+0.64 \\ -0.38}$		FORERO	14	FIT	Normal mass ordering
$1.48\substack{+0.34 \\ -0.32}$		FORERO	14	FIT	Inverted mass ordering
$1.70^{+0.22}_{-0.39}$		<sup>4</sup> GONZALEZ	14	FIT	Normal mass ordering; global fit
$1.41^{ig+0.35}_{-0.34}$		<sup>4</sup> GONZALEZ	14	FIT	Inverted mass ordering; global fit
0 to 1.5 or 1.9 to 2	90	<sup>5</sup> ADAMSON	13A	MINS	Normal mass ordering

 $^1\,{\rm Errors}$  are projections of 68% C.L. curve of  $\delta_{CP}$  vs.  ${\rm sin}^2\theta_{23}.$ 

 $^2$  ADAMSON 16 result is based on a data sample with 2.74  $\times$  10 $^{20}$  protons on target. The likelihood-based analysis observed 6  $\nu_e$  events with an expected background of 0.99  $\pm$  0.11 events.

<sup>3</sup>ADAMSON 14 result is based on three-flavor formalism and  $\theta_{23} > \pi/4$ . Likelihood as a function of  $\delta$  is also shown for the other three combinations of hierarchy and  $\theta_{23}$  quadrant; all values of  $\delta$  are allowed at 90% C.L.

<sup>4</sup> GONZALEZ-GARCIA 14 result comes from a frequentist global fit. The corresponding Bayesian global fit to the same data results are reported in BERGSTROM 15 as 68% CL intervals of 1.24–1.94 for normal and 1.15–1.77 for inverted mass ordering.

<sup>5</sup> ADAMSON 13A result is based on  $\nu_e$  appearance in MINOS and the calculated  $\sin^2(2\theta_{23}) = 0.957, \theta_{23} > \pi/4$ , and normal mass hierarchy. Likelihood as a function of  $\delta$  is also shown for the other three combinations of hierarchy and  $\theta_{23}$ quadrant; all values of  $\delta$  are allowed at 90% C.L.

#### (C) Other neutrino mixing results

The LSND collaboration reported in AGUILAR 01 a signal which is consistent with  $\overline{
u}_{\mu} \rightarrow \overline{
u}_e$  oscillations. In a three neutrino framework, this would be a measurement of  $\theta_{12}$  and  $\Delta m^2_{21}$ . This does not appear to be consistent with most of the other neutrino data. The MiniBooNE experiment, reported in AGUILAR-AREVALO 07, does a two-neutrino analysis which, assuming CP conservation, rules out AGUILAR 01. However, the MiniBooNE antineutrino data reported in AGUILAR-AREVALO 13A are consistent with the signal reported in AGUILAR 01. The following listings include results which might be relevant towards understanding these observations. They include searches for  $u_\mu o 
u_e, \ \overline{
u}_\mu o 
\overline{
u}_e$ , sterile neutrino oscillations, and CPT violation.

· · /		\ <i>\</i>	4 E/			
VALUE (e	V <sup>2</sup> )	CL%	DOCUMENT ID		TECN	COMMENT
• • • W	/e do not use the	following	data for averages	, fits,	limits, e	tc. ● ● ●
0.015	to 0.050	90	<sup>1</sup> AGUILAR-AR	.13A	MBNE	MiniBooNE
< 0.34		90	<sup>2</sup> MAHN	12	MBNE	MiniBooNE/SciBooNE
< 0.034		90	AGUILAR-AR	.07	MBNE	MiniBooNE
< 0.000	8	90	AHN	04	K2K	Water Cherenkov
<0.4		90	ASTIER	03	NOMD	CERN SPS
<2.4		90	AVVAKUMOV	02	NTEV	NUTEV FNAL
			<sup>3</sup> AGUILAR	01	LSND	$\nu \mu \rightarrow \nu_e \text{ osc.prob.}$
0.03	to 0.3	95	<sup>4</sup> ATHANASSO	.98	LSND	$\nu_{\mu} \rightarrow \nu_{e}$
<2.3		90	<sup>5</sup> LOVERRE	96		CHARM/CDHS
<0.9		90	VILAIN	94C	CHM2	CERN SPS
< 0.09		90	ANGELINI	86	HLBC	BEBC CERN PS
1_						

## $\Delta(m^2)$ for sin<sup>2</sup>(2 $\theta$ ) = 1 ( $\nu_{\mu} \rightarrow \nu_{e}$ )

<sup>1</sup>Based on  $\nu_{\mu} \rightarrow \nu_{e}$  appearance of 162.0  $\pm$  47.8 events; marginally compatible with two

neutrino oscillations. The best fit value is  $\Delta m^2=3.14~eV^2.$  ^2MAHN 12 is a combined spectral fit of MiniBooNE and SciBooNE neutrino data with the range of  $\Delta m^2$  up to 25 eV<sup>2</sup>. The best limit is 0.04 at 7 eV<sup>2</sup>.

<sup>3</sup>AGUILAR 01 is the final analysis of the LSND full data set. Search is made for the  $\nu_{\mu} \rightarrow \nu_{e}$  oscillations using  $\nu_{\mu}$  from  $\pi^{+}$  decay in flight by observing beam-on electron events from  $\nu_e C \rightarrow e^- X$ . Present analysis results in 8.1  $\pm$  12.2  $\pm$  1.7 excess events in the 60< $E_e$  < 200 MeV energy range, corresponding to oscillation probability of  $0.10\pm0.16\pm0.04\%$ . This is consistent, though less significant, with the previous result of ATHANASSOPOULOS 98, which it supersedes. The present analysis uses selection criteria developed for the decay at rest region, and is less effective in removing the background above 60 MeV than ATHANASSOPOULOS 98.

<sup>4</sup>ATHANASSOPOULOS 98 is a search for the  $\nu_{\mu} \rightarrow \nu_{e}$  oscillations using  $\nu_{\mu}$  from  $\pi^{+}$  decay in flight. The 40 observed beam-on electron events are consistent with  $\nu_{e}C \rightarrow$  $e^-$ X; the expected background is 21.9  $\pm$  2.1. Authors interpret this excess as evidence for an oscillation signal corresponding to oscillations with probability  $(0.26 \pm 0.10 \pm 0.05)\%$ . Although the significance is only  $2.3 \sigma$ , this measurement is an important and consistent cross check of ATHANASSOPOULOS 96 who reported evidence for  $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$  oscillations

from  $\mu^+$  decay at rest. See also ATHANASSOPOULOS 98B.

 $^5$ LOVERRE 96 uses the charged-current to neutral-current ratio from the combined CHARM (ALLABY 86) and CDHS (ABRAMOWICZ 86) data from 1986.

$\sin^2(2\theta)$ for "Large"	$\Delta(m^2)$	$( u_{\mu} \rightarrow  u_{e})$			
VALUE (units $10^{-3}$ )	CL%	DOCUMENT ID		TECN	COMMENT
• • • We do not use the	following	g data for averages	, fits,	limits, e	tc. ● ● ●
< 7.2	90				$\Delta(m^2) > 0.1 \; \mathrm{eV}^2$
0.8 to 3	90	<sup>1</sup> AGUILAR-AR	.13A	MBNE	MiniBooNE
< 11	90	<sup>2</sup> ANTONELLO	13	ICAR	$\nu_{\mu} \rightarrow \nu_{e}$
< 6.8	90	<sup>3</sup> ANTONELLO	13A	ICAR	$ u_{\mu} \rightarrow \nu_{e} $
<100	90		12		, MiniBooNE/SciBooNE
< 1.8	90	<sup>5</sup> AGUILAR-AR	.07	MBNE	MiniBooNE
<110	90	<sup>6</sup> AHN	04	K2K	Water Cherenkov
< 1.4	90	ASTIER	03	NOMD	CERN SPS
< 1.6	90	AVVAKUMOV	02	NTEV	NUTEV FNAL
		<sup>7</sup> AGUILAR	01	LSND	$ u_{\mu} \rightarrow \ \nu_{e} \ { m osc.prob}.$
0.5 to 30	95	<sup>8</sup> ATHANASSO	.98		$\nu_{\mu} \rightarrow \nu_{e}$
< 3.0	90	<sup>9</sup> LOVERRE	96		CHARM/CDHS
< 9.4	90	VILAIN	94C	CHM2	CERN SPS
< 5.6	90	<sup>10</sup> VILAIN	94C	CHM2	CERN SPS

<sup>1</sup>Based on  $\nu_{\mu} \rightarrow \nu_{e}$  appearance of 162.0  $\pm$  47.8 events; marginally compatible with two

, neutrino oscillations. The best fit value is sin^22 $\theta$  = 0.002. <sup>2</sup>ANTONELLO 13 use the ICARUS T600 detector at LNGS and  $\sim$  20 GeV beam of  $\nu_{\mu}$ from CERN 730 km away to search for an excess of  $\nu_e$  events. Two events are found

with 3.7  $\pm$  0.6 expected from conventional sources. This result excludes some parts of the parameter space expected by LSND. Superseded by ANTONELLO 13A.

 $^3\,\textsc{Based}$  on four events with a background of 6.4  $\pm$  0.9 from conventional sources with an average energy of 20 GeV and 730 km from the source of  $\nu_{\mu}$ .

<sup>4</sup> MAHN 12 is a combined fit of MiniBooNE and SciBooNE neutrino data. <sup>5</sup> The limit is  $\sin^2 2\theta < 0.9 \times 10^{-3}$  at  $\Delta m^2 = 2 \text{ eV}^2$ . That value of  $\Delta m^2$  corresponds to the smallest mixing angle consistent with the reported signal from LSND in AGUILAR 01. <sup>6</sup> The limit becomes  $\sin^2 2\theta < 0.15$  at  $\Delta m^2 = 2.8 \times 10^{-3} \text{ eV}^2$ , the bets-fit value of the  $\nu_{\mu}$  disappearance analysis in K2K.

 $^7\,{\rm AGUILAR}$  01 is the final analysis of the LSND full data set of the search for the  $\nu_\mu \rightarrow$  $\nu_e$  oscillations. See footnote in preceding table for further details.

 $^8$  ATHANASSOPOULOS 98 report (0.26  $\pm$  0.10  $\pm$  0.05)% for the oscillation probability; the value of  $\sin^2 2\theta$  for large  $\Delta m^2$  is deduced from this probability. See footnote in preceding table for further details, and see the paper for a plot showing allowed regions. If effect is due to oscillation, it is most likely to be intermediate  $\sin^2 2\theta$  and  $\Delta m^2$ . See also ATHANASSOPOULOS 98B. 9

LOVERRE 96 uses the charged-current to neutral-current ratio from the combined CHARM (ALLABY 86) and CDHS (ABRAMOWICZ 86) data from 1986.

 $^{10}$  VILAIN 94C limit derived by combining the  $\nu_{\mu}$  and  $\overline{\nu}_{\mu}$  data assuming CP conservation.

# $\Delta(m^2)$ for $\sin^2(2\theta) = 1$ $(\overline{ u}_{\mu} \rightarrow \overline{ u}_{e})$

VALUE (eV <sup>2</sup> )	<u>CL%</u>	DOCUMENT ID	TECN COMMENT
$\bullet \bullet \bullet$ We do not	use the following	g data for averages, fits,	limits, etc. • • •
0.023 to 0.060	90	<sup>1</sup> AGUILAR-AR13A	
<0.16	90	<sup>2</sup> CHENG 12	MBNE MiniBooNE/SciBooNE
0.03-0.09	90	<sup>3</sup> AGUILAR-AR10	MBNE E $_{ m  u}$ $>$ 475 MeV
0.03-0.07	90	<sup>4</sup> AGUILAR-AR10	MBNE $E_{\nu}^{\nu} > 200 \; MeV$
<0.06	90		MBNE MiniBooNE
<0.055	90	<sup>5</sup> ARMBRUSTER02	KAR2 Liquid Sci. calor.
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<2.6	90	AVVAKUMOV	02	NTEV	NUTEV FNAL
0.03–0.05		<sup>6</sup> AGUILAR	01	LSND	LAMPF
0.05-0.08	90	<sup>7</sup> ATHANASSO	96	LSND	LAMPF
0.048-0.090	80	<sup>8</sup> ATHANASSO	95		
<0.07	90	<sup>9</sup> HILL	95		
<0.9	90	VILAIN	94C	CHM2	CERN SPS
<0.14	90	<sup>10</sup> FREEDMAN	93	CNTR	LAMPF

<sup>1</sup>Based on  $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$  appearance of 78.4 ± 28.5 events. The best fit values are  $\Delta m^{2} = 0.043 \text{ eV}^{2}$  and  $\sin^{2}2\theta = 0.88$ .

 $^{2}$  CHENG 12 is a combined fit of MiniBooNE and SciBooNE antineutrino data.

 $^3$  This value is for a two neutrino oscillation analysis for excess antineutrino events with E $_{\nu}$  > 475 MeV. The best fit is at 0.07. The allowed region is consistent with LSND reported by AGUILAR 01. Supercedes AGUILAR-AREVALO 09B.

<sup>4</sup> This value is for a two neutrino oscillation analysis for excess antineutrino events with  $E_{\nu} > 200$  MeV with subtraction of the expected 12 events low energy excess seen in the

\_ neutrino component of the beam. The best fit value is 0.007 for  $\Delta(m^2) = 4.4 \ {
m eV}^2$ .

- <sup>5</sup> ARMBRUSTER 02 is the final analysis of the KARMEN 2 data for 17.7 m distance from the ISIS stopped pion and muon neutrino source. It is a search for  $\overline{\nu}_e$ , detected by the inverse  $\beta$ -decay reaction on protons and <sup>12</sup>C. 15 candidate events are observed, and 15.8  $\pm$  0.5 background events are expected, hence no oscillation signal is detected. The results exclude large regions of the parameter area favored by the LSND experiment.
- <sup>6</sup> AGUILAR 01 is the final analysis of the LSND full data set. It is a search for  $\overline{\nu}_e$  30 m from LAMPF beam stop. Neutrinos originate mainly for  $\pi^+$  decay at rest.  $\overline{\nu}_e$  are detected through  $\overline{\nu}_e p \rightarrow e^+ n (20 < E_{e^+} < 60 \text{ MeV})$  in delayed coincidence with  $np \rightarrow d\gamma$ . Authors observe 87.9  $\pm$  22.4  $\pm$  6.0 total excess events. The observation is attributed to  $\overline{\nu}_\mu \rightarrow \overline{\nu}_e$  oscillations with the oscillation probability of 0.264  $\pm$  0.067  $\pm$  0.045%, consistent with the previously published result. Taking into account all constraints, the most favored allowed region of oscillation parameters is a band of  $\Delta(m^2)$  from 0.2–2.0 eV<sup>2</sup>. Supersedes ATHANASSOPOULOS 95, ATHANASSOPOULOS 96, and \_ATHANASSOPOULOS 98.
- ATHANASSOPOULOS 98. 7 ATHANASSOPOULOS 96 is a search for  $\overline{\nu}_e$  30 m from LAMPF beam stop. Neutrinos originate mainly from  $\pi^+$  decay at rest.  $\overline{\nu}_e$  could come from either  $\overline{\nu}_\mu \rightarrow \overline{\nu}_e$  or  $\nu_e \rightarrow \overline{\nu}_e$ ; our entry assumes the first interpretation. They are detected through  $\overline{\nu}_e p \rightarrow e^+ n$  (20 MeV  $\langle E_{e^+} \rangle$  40 MeV) in delayed coincidence with  $np \rightarrow d\gamma$ . Authors observe 51 ± 20 ± 8 total excess events over an estimated background 12.5 ± 2.9. ATHANASSOPOULOS 96B is a shorter version of this paper.
- $^{8}$  ATHANASSOPOULOS 95 error corresponds to the  $1.6\sigma$  band in the plot. The expected background is 2.7  $\pm$  0.4 events. Corresponds to an oscillation probability of  $(0.34^{+}_{-0.18} \pm 0.07)\%$ . For a different interpretation, see HILL 95. Replaced by ATHANASSOPOULOS 96.
- <sup>9</sup> HILL 95 is a report by one member of the LSND Collaboration, reporting a different conclusion from the analysis of the data of this experiment (see ATHANASSOPOULOS 95). Contrary to the rest of the LSND Collaboration, Hill finds no evidence for the neutrino oscillation  $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$  and obtains only upper limits.
- <sup>10</sup> FREEDMAN 93 is a search at LAMPF for  $\overline{\nu}_e$  generated from any of the three neutrino types  $\nu_{\mu}$ ,  $\overline{\nu}_{\mu}$ , and  $\nu_e$  which come from the beam stop. The  $\overline{\nu}_e$ 's would be detected by the reaction  $\overline{\nu}_e p \rightarrow e^+ n$ . FREEDMAN 93 replaces DURKIN 88.

sin <sup>2</sup> (2 <i>θ</i> ) for "Large"	$\Delta(m^2)$	$(\overline{ u}_{\mu} \rightarrow \overline{ u}_{e})$	
VALUE (units $10^{-3}$ )	CL%	DOCUMENT ID	TECN COMMENT
$\bullet \bullet \bullet$ We do not use the	e following	g data for averages, fits,	limits, etc. • • •
<640	90	<sup>1</sup> ANTONELLO 13A	ICAR $\overline{\nu}_{e}$ appearance
<150	90	<sup>2</sup> CHENG 12	MBNE MiniBooNE/SciBooNE
0.4–9.0	99	<sup>3</sup> AGUILAR-AR10	MBNE $E_{\nu} > 475 \text{ MeV}$
0.4–9.0	99	<sup>4</sup> AGUILAR-AR10	MBNE $E_{\nu}$ > 200 MeV
< 3.3	90	<sup>5</sup> AGUILAR-AR09B	MBNE MiniBooNE
< 1.7	90	<sup>6</sup> ARMBRUSTER02	KAR2 Liquid Sci. calor.
< 1.1	90	AVVAKUMOV 02	NTEV NUTEV FNAL
$5.3 {\pm} 1.3 {\pm} 9.0$			LSND LAMPF
$6.2 \pm 2.4 \pm 1.0$		<sup>8</sup> ATHANASSO96	LSND LAMPF
3–12	80	<sup>9</sup> ATHANASSO95	
< 6	90	<sup>10</sup> HILL 95	

<sup>1</sup> ANTONELLO 13A obtained the limit by assuming  $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$  oscillation from the  $\sim 2\%$  of  $\overline{\nu}_{\mu}$  evnets contamination in the CNGS beam.

 $^{2}_{2}$ CHENG 12 is a combined fit of MiniBooNE and SciBooNE antineutrino data.

<sup>3</sup> This value is for a two neutrino oscillation analysis for excess antineutrino events with  $E_{\nu} > 475$  MeV. At 90% CL there is no solution at high  $\Delta(m^2)$ . The best fit is at maximal mixing. The allowed region is consistent with LSND reported by AGUILAR 01. Supercedes AGUILAR-AREVALO 09B.

<sup>4</sup> This value is for a two neutrino oscillation analysis for excess antineutrino events with  $E_{\nu} > 200$  MeV with subtraction of the expected 12 events low energy excess seen in the neutrino component of the beam. At 90% CL there is no solution at high  $\Delta(m^2)$ . The best fit value is 0.007 for  $\Delta(m^2) = 4.4 \text{ eV}^2$ .

 $^{5}$  This result is inconclusive with respect to small amplitude mixing suggested by LSND.

<sup>6</sup>ARMBRUSTER 02 is the final analysis of the KARMEN 2 data. See footnote in the \_preceding table for further details, and the paper for the exclusion plot.

- <sup>7</sup> AGUILAR 01 is the final analysis of the LSND full data set. The deduced oscillation probability is  $0.264 \pm 0.067 \pm 0.045\%$ ; the value of  $\sin^2 2\theta$  for large  $\Delta(m^2)$  is twice this probability (although these values are excluded by other constraints). See footnote in preceding table for further details, and the paper for a plot showing allowed regions. Supersedes ATHANASSOPOULOS 95, ATHANASSOPOULOS 96, and ATHANASSOPOULOS 98.
- <sup>8</sup>ATHANASSOPOULOS 96 reports  $(0.31 \pm 0.12 \pm 0.05)\%$  for the oscillation probability; the value of  $\sin^2 2\theta$  for large  $\Delta(m^2)$  should be twice this probability. See footnote in preceding table for further details, and see the paper for a plot showing allowed regions.
- $^{9}$ ATHANASSOPOULOS 95 error corresponds to the  $1.6\sigma$  band in the plot. The expected background is  $2.7 \pm 0.4$  events. Corresponds to an oscillation probability of  $(0.34^{+0.20}_{-0.18} \pm 0.07)\%$ . For a different interpretation, see HILL 95. Replaced by ATHANASSOPOULOS 96.
- <sup>10</sup> HILL 95 is a report by one member of the LSND Collaboration, reporting a different conclusion from the analysis of the data of this experiment (see ATHANASSOPOULOS 95). Contrary to the rest of the LSND Collaboration, Hill finds no evidence for the neutrino oscillation  $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$  and obtains only upper limits.

## $\Delta(m^2) ext{ for sin}^2(2 heta) = 1 \quad ( u_\mu(\overline{ u}_\mu) ightarrow \ u_e(\overline{ u}_e))$

VALUE (eV <sup>2</sup> )	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT	
<0.075	90	BORODOV	92	CNTR	BNL E776	
• • • We do not us	se the following	g data for average	s, fits,	, limits, e	etc. • • •	
<1.6	90	<sup>1</sup> ROMOSAN	97	CCFR	FNAL	
1 POMOSAN 07	uses wideband		ارس م			

<sup>1</sup> ROMOSAN 97 uses wideband beam with a 0.5 km decay region.

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VALUE (units $10^{-3}$ )	CL%	DOCUMENT ID	TECN	COMMENT
<1.8	90	<sup>1</sup> ROMOSAN 9	7 CCFR	FNAL
• • We do not u	se the follow	wing data for averages, f	its, limits, e	etc. • • •
<3.8	90	<sup>2</sup> MCFARLAND 9	5 CCFR	FNAL
<3	90	BORODOV 9	2 CNTR	BNL E776
<sup>2</sup> MCFARLAND	95 state t	and beam with a 0.5 km hat "This result is the excludes at 90%CL much on." See ATHANASSOF	e most stri	ngent to date for 250
$\Delta(m^2)$ for $\sin^2(m^2)$			TECN	COMMENT
<i>ALUE</i> (ev-)	<u></u>	DOCUMENT ID wing data for averages, f	ite limite	
< 0.01	90	<sup>1</sup> ACHKAR 9	5 CNIR	Bugey reactor
- ACHKAR 95 b	ound is for	L=15, 40, and 95 m.		
sin <sup>2</sup> (2 $ heta$ ) for "La	rge" Δ( <i>m</i>	<sup>2</sup> ) $(\overline{\nu}_e \neq \overline{\nu}_e)$		
• •	<u>CL%</u>		TECN	COMMENT
• • We do not u	se the follow	wing data for averages, f	its, limits, e	etc. ● ● ●
<0.02	90	<sup>1</sup> ACHKAR 9	5 CNTR	For $\Delta(m^2) = 0.6 \text{ eV}^2$
	1 · · c			( )
$\Delta(m^2)$ for sin <sup>2</sup> (	2 heta) = 1 (a	- Sterile neutrino lin $ u_{\mu} \rightarrow  u_{s})$		ce from the Bugey react
$\Delta(m^2)$ for sin <sup>2</sup> ( $ u_s$ means $ u_{ au}$	<b>2θ) = 1 (</b> ά or any ster	- Sterile neutrino lin	nits ——	
$\Delta(m^2)$ for sin <sup>2</sup> ( $ u_s$ means $ u_{ au}$ <i>/ALUE</i> (10 <sup>-5</sup> eV <sup>2</sup> )	<b>20) = 1 (i</b> or any ster	- Sterile neutrino lin $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u>	nits	1MENT
$\Delta(m^2)$ for sin <sup>2</sup> ( $\nu_s$ means $\nu_\tau$ <i>VALUE</i> (10 <sup>-5</sup> eV <sup>2</sup> ) ••• We do not us <3000 (or <550)	<b>20) = 1 (</b> or any ster <u>CL%</u> use the follow 90	- Sterile neutrino lim $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u> <u>7</u> wing data for averages, f <sup>1</sup> OYAMA 89	n <b>its</b> <u>FECN</u> <u>CON</u> Tits, limits, d (AMI Wat	<i>IMENT</i> etc. ● ● ● ter Cherenkov
$\Delta(m^2) \text{ for sin}^2(\nu_s \text{ means } \nu_\tau)$ $\frac{\Delta LUE (10^{-5} \text{ eV}^2)}{\bullet} \bullet \text{ We do not u}$ $< 3000 (or < 550)$ $< 4.2 \text{ or } > 54.$ $\frac{1}{OYAMA 89 \text{ giv}}$ argue that the for large mixing	$2\theta = 1$ (interpretation of the second seco	- Sterile neutrino lim $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u> 7 wing data for averages, f <sup>1</sup> OYAMA 89 M BIONTA 88 I of limits, depending on $\nu^{2}$ ) = (100–1000) × 10 <sup>-1</sup>	n <b>its</b> <u>FECN</u> <u>CON</u> Tits, limits, of (AMI Wat MB Flux assumptior	MENT etc. • • • ter Cherenkov $\times$ has $ u_{\mu}$ , $\overline{\nu}_{\mu}$ , $\nu_e$ , and $\overline{\nu}$ ns in their analysis. Th
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$\Delta(m^2) \text{ for sin}^2(\nu_s \text{ means } \nu_\tau \text{ Malue } (10^{-5} \text{ eV}^2)$ ••• We do not use of the second se	$2\theta = 1 (i)$ or any ster $CL\%$ $90$ $90$ $90$ $yes a range region \Delta(m) CL\% CL\% CL\% De follow$	- Sterile neutrino lim $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u> wing data for averages, ft <sup>1</sup> OYAMA BIONTA BIONTA BIONTA $\gamma^{2}$ ) = (100–1000) × 10 <sup>-1</sup> <u>DOCUMENT ID</u> wing data for averages, ft	nits <u>TECN</u> <u>CON</u> Tits, limits, of (AMI Wath MB Flux assumption -5 eV <sup>2</sup> is r <u>TECN</u> <u>CO</u> TECN <u>CO</u> TCCB IC	$\underline{MMENT}$ etc. • • • ter Cherenkov $<$ has $\nu_{\mu}$ , $\overline{\nu}_{\mu}$ , $\nu_{e}$ , and $\overline{\nu}_{\mu}$ has in their analysis. The not ruled out by any data $\underline{OMMENT}$ etc. • • •
$\Delta(m^{2}) \text{ for sin}^{2}(\nu_{s} \text{ means } \nu_{\tau} \text{ was superative}^{2})$ •• We do not use (10 <sup>-5</sup> eV <sup>2</sup> ) •• We do not use (3000 (or <550)) (4.2 or > 54.) •• OYAMA 89 gives argue that the for large mixing (1000 cm sector) (100	$2\theta = 1 (i)$ or any ster $CL\%$ use the follow 90 90 ves a range region $\Delta(mg)$ $E^{\mu} \rightarrow \nu_{g}$ $CL\%$ use the follow 90 95	- Sterile neutrino lim $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u> <u>7</u> wing data for averages, f <sup>1</sup> OYAMA 89 $\mu$ BIONTA 88 I of limits, depending on $\mu^{2}$ ) = (100–1000) × 10 <sup>-1</sup> <u>DOCUMENT ID</u> wing data for averages, f <sup>1</sup> AARTSEN 17B <sup>2</sup> ABDURASHI 17 <sup>3</sup> KO 17 <sup>4</sup> AARTSEN 16	r <u>ECN</u> <u>COM</u> Tits, limits, of (AMI Wat MB Flux assumption -5 eV <sup>2</sup> is r <u>TECN</u> <u>CC</u> TITS, limits, of ICCB IC T NEOS	$\underline{MMENT}$ etc. • • • ter Cherenkov $\times$ has $\nu_{\mu}$ , $\overline{\nu}_{\mu}$ , $\nu_{e}$ , and $\overline{\nu}_{\mu}$ has in their analysis. Th not ruled out by any data $\underline{DMMENT}$ etc. • • • eCube-DeepCore
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$\Delta(m^{2}) \text{ for sin}^{2}(\nu_{s} \text{ means } \nu_{\tau} \text{ walle } (10^{-5} \text{ eV}^{2})$ •• We do not use (3000 (or <550) < 4.2 or > 54. <sup>1</sup> OYAMA 89 give argue that the for large mixing (5) (1000 \text{ Gearch for } \nu_{\mu} \text{ or } (1000 \text{ Gearch for } \nu_	$2\theta = 1 (i)$ or any ster $CL\%$ $90$ 90 $90$ $yes a range region \Delta(m) \tau v_e \rightarrow v_s CL\% 90 95 90 90 90$	- Sterile neutrino lim $\nu_{\mu} \rightarrow \nu_{s}$ ) ile (noninteracting) $\nu$ . <u>DOCUMENT ID</u> <u>1</u> wing data for averages, f <sup>1</sup> OYAMA 89 $\mu$ BIONTA 89 $\mu$ BIONTA 88 I of limits, depending on $\mu^{2}$ ) = (100–1000) × 10 <sup>-1</sup> <u>DOCUMENT ID</u> wing data for averages, f <sup>1</sup> AARTSEN 17B <sup>2</sup> ABDURASHI 17 <sup>3</sup> KO 17 <sup>4</sup> AARTSEN 16 <sup>5</sup> ADAMSON 16B <sup>6</sup> ADAMSON 16C	nits <u>TECN</u> <u>CON</u> Tits, limits, of (AMI Wat MB Flux assumption -5 eV <sup>2</sup> is r <u>TECN</u> <u>CO</u> TITS, limits, of ICCB IC NEOS ICCB IC M	etc. • • • ter Cherenkov $\leftarrow$ has $\nu_{\mu}$ , $\overline{\nu}_{\mu}$ , $\nu_{e}$ , and $\overline{\nu}_{\mu}$ has in their analysis. The not ruled out by any da <u>DMMENT</u> etc. • • • eCube-DeepCore $\beta$ decay eCube
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- <sup>1</sup> AARTSEN 17B uses three years of upward-going atmospheric neutrino data in the energy range of 10-60 GeV to constrain their disappearance into light sterile neutrinos. The reported limit  $\sin^2\theta_{24} < 0.11$  at 90% C.L. is for  $\Delta m^2_{41} = 1.0 \text{ eV}^2$ . We convert the result to  $\sin^2 2\theta_{24}$  for the listing. AARTSEN 17B also reports  $\cos^2\theta_{24} \cdot \sin^2\theta_{34} < 0.15$  at 90% C.L. for  $\Delta m^2_{41} = 1.0 \text{ eV}^2$ .
- <sup>2</sup>ABDURASHITOV 17 use the Troitsk nu-mass experiment to search for sterile neutrinos with mass 0.1 - 2 keV. We convert the reported limit from  $U_{e4}^2 < 0.002$  to  $\sin^2 2\theta_{14} < 0.008$  assume  $U_{e4} \sim \sin \theta_{14}$ . The stated limit corresponds to the smallest  $U_{e4}^2$ . The exclusion curve begins at  $U_{e4}^2$  of 0.02 for m<sub>4</sub> = 0.1 keV.
- <sup>3</sup> KO 17 reports on short baseline reactor oscillation search ( $\overline{\nu}_e \rightarrow \overline{\nu}_s$ ), motivated be the so-called "reactor antineutrino anomaly". The experiment is conducted at 23.7 m from the core of unit 5 of the Hanbit Nuclear Power Complex in Korea. the reported limited on  $\sin^2(2\theta_{41})$  for sterile neutrinos was determined using the reactor antineutrino spectrum determined by the Daya Bay experiment for  $\Delta m_{14}^2$  around 0.55 eV<sup>2</sup> where the sensitivity is maximal. A fraction of the parameter space derived from the "reactor antineutrino anomaly" is excluded by this work. Compared to reactor models an event excess is observed at about 5 MeV, in agreement with other experiments.
- <sup>4</sup> AARTSEN 16 use one year of upward-going atmospheric muon neutrino data in the energy range of 320 GeV to 20 TeV to constrain their disappearance into light sterile neutrinos. Sterile neutrinos are expected to produce distinctive zenith distribution for these energies for 0.01  $\leq \Delta m^2 \leq 10 \text{ eV}^2$ . The stated limit is for  $\sin^2 2\theta_{24}$  at  $\Delta m^2$  around 0.3 eV<sup>2</sup>.
- <sup>5</sup> ADAMSON 16B combine the results of AN 16B, ADAMSON 16C, and Bugey-3 reactor experiments to constrain  $\nu_{\mu}$  to  $\nu_{e}$  mixing through oscillations into light sterile neutrinos. The stated limit for  $\sin^{2}2\theta_{\mu e}$  is at  $|\Delta m_{41}^{2}| = 1.2 \text{ eV}^{2}$ .
- <sup>6</sup> ADAMSON 16C use the NuMI beam and exposure of  $10.56 \times 10^{20}$  protons on target to search for the oscillation of  $\nu_{\mu}$  dominated beam into light sterile neutrinos with detectors at 1.04 and 735 km. The reported limit  $\sin^2(\theta_{24}) < 0.022$  at 95% C.L. is for  $|\Delta m_{41}^2| = 0.5 \text{ eV}^2$ . We convert the result to  $\sin^2(2\theta_{24})$  for the listing.
- <sup>7</sup> AN 16B utilize 621 days of data to place limits on the  $\overline{\nu}_e$  disappearance into a light sterile neutrino. The stated limit corresponds to the smallest  $\sin^2(2\theta_{14})$  at  $|\Delta m_{41}^2| \sim 3 \times 10^{-2} \text{ eV}^2$  (obtained from Figure 3 in AN 16B). The exclusion curve begins at  $|\Delta m_{41}^2| \sim 1.5 \times 10^{-4} \text{ eV}^2$  and extends to  $\sim 0.25 \text{ eV}^2$ . The analysis assumes  $\sin^2(2\theta_{12}) = 0.846 \pm 0.021$ ,  $\Delta m_{21}^2 = (7.53 \pm 0.18) \times 10^{-5} \text{ eV}^2$ , and  $|\Delta m_{32}^2| = (2.44 \pm 0.06) \times 10^{-3} \text{ eV}^2$ .
- <sup>8</sup>AMBROSIO 01 tested the pure 2-flavor  $\nu_{\mu} \rightarrow \nu_{s}$  hypothesis using matter effects which change the shape of the zenith-angle distribution of upward through-going muons. With maximum mixing and  $\Delta m^{2}$  around 0.0024 eV<sup>2</sup>, the  $\nu_{\mu} \rightarrow \nu_{s}$  oscillation is disfavored with 99% confidence level with respect to the  $\nu_{\mu} \rightarrow \nu_{\tau}$  hypothesis.
- <sup>9</sup> FUKUDA 00 tested the pure 2-flavor  $\nu_{\mu} \rightarrow \nu_{s}$  hypothesis using three complementary atmospheric-neutrino data samples. With this hypothesis, zenith-angle distributions are expected to show characteristic behavior due to neutral currents and matter effects. In the  $\Delta m^{2}$  and  $\sin^{2}2\theta$  region preferred by the Super-Kamiokande data, the  $\nu_{\mu} \rightarrow \nu_{s}$  hypothesis isrejected at the 99% confidence level, while the  $\nu_{\mu} \rightarrow \nu_{\tau}$  hypothesis consistently fits all of the data sample.

------ CPT tests

$\left< \Delta m_{21}^2 - \Delta \overline{m}_{21}^2 \right>$				
$VALUE (10^{-4} \text{ eV}^2)$	<u>CL%</u>	DOCUMENT ID	TECN	COMMENT
• • • We do not use the	e following	data for averages, fits	, limits, e	etc. • • •
<1.1	99.7	<sup>1</sup> DEGOUVEA 05	FIT	solar vs. reactor
<sup>1</sup> DEGOUVEA 05 obta solar neutrino data.	ined this b	bound at the 3 $\sigma$ CL from	m the Ka	mLAND (ARAKI 05) and
$\left< \Delta m_{32}^2 - \Delta \overline{m}_{32}^2 \right>$				
VALUE $(10^{-3} \text{ eV}^2)$	<u>CL%</u>	DOCUMENT ID	TECN	COMMENT
• • • We do not use the	e following	data for averages, fits	, limits, e	etc. • • •
$0.6^{+2.4}_{-0.8}$	90	<sup>1</sup> ADAMSON 12B	MINS	MINOS atmospheric

<sup>1</sup> The quoted result is the single-parameter 90% C.L. interval determined from the 90% C.L. contour in the  $(\Delta m^2, \Delta \overline{m}^2)$  plane, which is obtained by minimizing the four parameter log-likelihood function with respect to the other oscillation parameters.

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ABE	17A	PRL 118 151801	K. Abe <i>et al.</i>	(	Collab.)
ABE	17C	PR D96 011102	K. Abe <i>et al.</i>		Collab.)
ABE	17F	PR D96 092006	K. Abe <i>et al.</i>		Collab.)
ADAMSON	17A	PRL 118 151802	P. Adamson <i>et al.</i>		Collab.)
ADAMSON	17B	PRL 118 231801	P. Adamson <i>et al.</i>		Collab.)
AN	17A	PR D95 072006	F.P. An <i>et al.</i>	(Daya Bay	
KO	17	PRL 118 121802	Y.J. Ko <i>et al.</i>		Collab.)
AARTSEN	16	PRL 117 071801	M.G. Aartsen <i>et al.</i>	(IceCube	
ABE	16B	JHEP 1601 163	Y. Abe <i>et al.</i>	(Double Chooz	
ABE	16C	PR D94 052010	K. Abe <i>et al.</i>	(Super-Kamiokande	Collab.)
ABE	16D	PRL 116 181801	K. Abe <i>et al.</i>		Collab.)
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ADAMSON	16A	PR D93 051104	P. Adamson <i>et al.</i>	(NOvA	Collab.)
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AN	16	PRL 116 061801	F.P. An <i>et al.</i>	(Daya Bay	Collab.)
AN	16A	PR D93 072011	F.P. An <i>et al.</i>	(Daya Bay	Collab.)
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Also		PR D91 072010	K. Abe <i>et al.</i>	ÌΤ2K	Collab.)
ABE	14A	PL B735 51	Y. Abe <i>et al.</i>	(Double Chooz	Collab.)
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AN	14	PRL 112 061801	F.P. An et al.	(Daya Bay	
AN	14B	PR D90 071101	F.P. An et al.	(Daya Bay	
BELLINI	14A	NAT 512 383	G. Bellini <i>et al.</i>	(Borexino	,
FORERO	14	PR D90 093006	D.V. Forero, M. Tortola,		
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PDG	14	CP C38 070001	K. Olive <i>et al.</i>		Collab.)
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				(oupor ramonande	

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ABE	13C	PL B723 66	Y. Abe <i>et al.</i>	(Double Chooz Collab.)	)
ABE	13E	PR D88 032002	K. Abe <i>et al.</i>	(T2K Collab.)	)
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ADAMSON	13A	PRL 110 171801	P. Adamson <i>et al.</i>	(MINOS Collab.)	)
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			-		
AGUILAR-AR	13A	PRL 110 161801	A.A. Aguilar-Arevalo <i>et al.</i>	(MiniBooNE Collab.)	)
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ABE	12	PRL 108 131801	Y. Abe <i>et al.</i>	(Double Chooz Collab.)	)
ABE	12A	PR D85 031103	K. Abe <i>et al.</i>	` (T2K Collab.)	1
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CHENG	12	PR D86 052009	G. Cheng et al. (M	iniBooNE/SciBooNE Collab.)	)
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MAHN	12	PR D85 032007		iniBooNE/SciBooNE_Collab.)	
ABE	11	PR D83 052010	K. Abe <i>et al.</i>	(Super-Kamiokande Collab.)	)
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AGOILAN-AN		INC 103 101001	A.A. Aguillai-Alevalo el al.		۱.
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AHARMIM	10	PR C81 055504	B. Aharmim <i>et al.</i>	(SNO Collab.)	
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BELLINI	10A	PR D82 033006	G. Bellini <i>et al.</i>	(SNO Collab.) (Borexino Collab.)	)
BELLINI DENIZ	10A 10	PR D82 033006 PR D81 072001	G. Bellini <i>et al.</i> M. Deniz <i>et al.</i>	(SNO Collab.)	)
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BELLINI DENIZ KAETHER	10A 10 10	PR D82 033006 PR D81 072001 PL B685 47	G. Bellini <i>et al.</i> M. Deniz <i>et al.</i> F. Kaether <i>et al.</i>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.)	))
BELLINI DENIZ KAETHER WENDELL	10A 10 10 10	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004	G. Bellini <i>et al.</i> M. Deniz <i>et al.</i> F. Kaether <i>et al.</i> R. Wendell <i>et al.</i>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.)	)))
BELLINI DENIZ KAETHER WENDELL ABDURASHI	10A 10 10 10 09	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807	G. Bellini <i>et al.</i> M. Deniz <i>et al.</i> F. Kaether <i>et al.</i> R. Wendell <i>et al.</i> J.N. Abdurashitov <i>et al.</i>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.)	) ) )
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BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON	10A 10 10 10 09 09	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802	<ul> <li>G. Bellini <i>et al.</i></li> <li>M. Deniz <i>et al.</i></li> <li>F. Kaether <i>et al.</i></li> <li>R. Wendell <i>et al.</i></li> <li>J.N. Abdurashitov <i>et al.</i></li> <li>P. Adamson <i>et al.</i></li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.)	)))))))
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BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE	10A 10 10 10 09 09	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803	<ul> <li>G. Bellini <i>et al.</i></li> <li>M. Deniz <i>et al.</i></li> <li>F. Kaether <i>et al.</i></li> <li>R. Wendell <i>et al.</i></li> <li>J.N. Abdurashitov <i>et al.</i></li> <li>P. Adamson <i>et al.</i></li> <li>A.A. Aguilar-Arevalo <i>et al.</i></li> <li>S. Abe <i>et al.</i></li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.)	) )) )) ))
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR	10A 10 10 10 09 09 09B	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801	<ul> <li>G. Bellini <i>et al.</i></li> <li>M. Deniz <i>et al.</i></li> <li>F. Kaether <i>et al.</i></li> <li>R. Wendell <i>et al.</i></li> <li>J.N. Abdurashitov <i>et al.</i></li> <li>P. Adamson <i>et al.</i></li> <li>A.A. Aguilar-Arevalo <i>et al.</i></li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.)	) )) )) ))
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also	10A 10 10 09 09 09B 08A	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>S. Abe et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.)	) ) ) ) ) )
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON	10A 10 10 09 09 09B 08A 08	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.)	)))))))))))))))))))))))))))))))))))))))
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON	10A 10 10 09 09 09B 08A 08 08A	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.)	)))))))))))))))))))))))))))))))))))))))
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BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM	10A 10 10 09 09 09B 08A 08 08A	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 111301	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A. A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.)	
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM Also	10A 10 10 09 09B 08A 08 08A 08	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 111301 PR C87 015502	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.)	) ) ) ) ) ) ) )
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM	10A 10 10 09 09 09B 08A 08 08A	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 111301	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A. A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.)	) ) ) ) ) ) ) )
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE AIso ADAMSON ADAMSON AHARMIM AIso ARPESELLA	10A 10 10 09 09B 08A 08 08A 08	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 111301 PR C87 015502 PRL 101 091302	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A. A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>C. Arpesella et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (Borexino Collab.)	
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON ADAMSON AHARMIM Also ARPESELLA CRAVENS	10A 10 10 09 09B 08A 08 08A 08 08A 08	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 091302 PRL 101 091302 PR D78 032002	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>C. Arpesella et al.</li> <li>J.P. Cravens et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (Super-Kamiokande Collab.) (SAGE Collab.) (MINOS Collab.) (MiniBooNE Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.)	
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM Also ARPESELLA CRAVENS FOGLI	10A 10 10 09 09B 08A 08 08A 08 08A 08 08A 08	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 111301 PR C87 015502 PRL 101 091302 PR D78 032002 PRL 101 141801	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>C. Arpesella et al.</li> <li>J.P. Cravens et al.</li> <li>G.L. Fogli</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.)	
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BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM Also ARPESELLA CRAVENS FOGLI	10A 10 10 09 09B 08A 08 08A 08 08A 08 08A 08 08A 08 07	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 091302 PRL 101 091302 PRL 101 141801 PR D75 092003	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>C. Arpesella et al.</li> <li>J.P. Cravens et al.</li> <li>G.L. Fogli</li> <li>P. Adamson et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (KamLAND Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (Super-Kamiokande Collab.)	
BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM Also ARPESELLA CRAVENS FOGLI ADAMSON AGUILAR-AR	10A 10 10 09 09B 08A 08 08A 08 08A 08 08A 08 07 07	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 131802 PRL 101 1302 PRL 101 1302 PRL 101 091302 PR D78 032002 PRL 101 141801 PR D75 092003 PRL 98 231801	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>C. Arpesella et al.</li> <li>J.P. Cravens et al.</li> <li>G.L. Fogli</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (KamLAND Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (Super-Kamiokande Collab.) (MINOS Collab.) (MINOS Collab.) (Super-Kamiokande Collab.)	
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BELLINI DENIZ KAETHER WENDELL ABDURASHI ADAMSON AGUILAR-AR ABE Also ADAMSON ADAMSON AHARMIM AIso ARPESELLA CRAVENS FOGLI ADAMSON AGUILAR-AR AHARMIM ADAMSON AHN BALATA HOSAKA HOSAKA MICHAEL WINTER YAMAMOTO AHARMIM	10A 10 10 09 09B 08A 08 08A 08 08A 08 08A 08 07 07 07 06 06A 06 06A 06 06A 06 06A 06 05A 05	PR D82 033006 PR D81 072001 PL B685 47 PR D81 092004 PR C80 015807 PRL 103 261802 PRL 103 111801 PRL 100 221803 PRL 101 119904E PR D77 072002 PRL 101 131802 PRL 101 131802 PRL 101 091302 PRL 101 091302 PR D78 032002 PRL 101 141801 PR D75 092003 PRL 98 231801 PR C75 045502 PR D74 072003 EV D73 012002 PR D74 072003 EV D74 072002 EV D74 07207 EV D74 07207 E	<ul> <li>G. Bellini et al.</li> <li>M. Deniz et al.</li> <li>F. Kaether et al.</li> <li>R. Wendell et al.</li> <li>J.N. Abdurashitov et al.</li> <li>P. Adamson et al.</li> <li>A.A. Aguilar-Arevalo et al.</li> <li>S. Abe et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>B. Aharmim et al.</li> <li>G.L. Fogli</li> <li>P. Adamson et al.</li> <li>B. Aharmim et al.</li> <li>G.L. Fogli</li> <li>P. Adamson et al.</li> <li>M.H. Ahn et al.</li> <li>J. Hosaka et al.</li> <li>J. Hosaka et al.</li> <li>J. Hosaka et al.</li> <li>M. W.T. Winter et al.</li> <li>S. Yamamoto et al.</li> <li>B. Aharmim et al.</li> </ul>	(SNO Collab.) (Borexino Collab.) (TEXONO Collab.) (SAGE Collab.) (SAGE Collab.) (MINOS Collab.) (MINOS Collab.) (KamLAND Collab.) (KamLAND Collab.) (KamLAND Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (MINOS Collab.) (MINOS Collab.) (MINOS Collab.) (MINOS Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SNO Collab.) (SUper-Kamiokande Collab.) (MINOS Collab.)	

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ALTMANN	05	PL B616 174	M. Altmann <i>et al.</i>	(GNO Collab.)
ARAKI	05	PRL 94 081801	T. Araki <i>et al.</i>	(KamLAND Collab.)
ASHIE	05	PR D71 112005	Y. Ashie <i>et al.</i>	(Super-Kamiokande Collab.)
DEGOUVEA	05	PR D71 093002	A. de Gouvea, C. Pena-Garay	
AHARMIM	03	PR D70 093014	B. Aharmim <i>et al.</i>	
	04 04A			(SNO Collab.)
AHMED	-	PRL 92 181301	S.N. Ahmed <i>et al.</i>	(SNO Collab.)
AHN	04	PRL 93 051801	M.H. Ahn <i>et al.</i>	(K2K Collab.)
AMBROSIO	04	EPJ C36 323	M. Ambrosio <i>et al.</i>	(MACRO Collab.)
ASHIE	04	PRL 93 101801	Y. Ashie <i>et al.</i>	(Super-Kamiokande Collab.)
EGUCHI	04	PRL 92 071301	K. Eguchi <i>et al.</i>	(KamLAND Collab.)
SMY	04	PR D69 011104	M.B. Smy <i>et al.</i>	(Super-Kamiokande Collab.)
AHN	03	PRL 90 041801	M.H. Ahn <i>et al.</i>	(K2K Collab.)
AMBROSIO	03	PL B566 35	M. Ambrosio <i>et al.</i>	(MACRO Collab.)
APOLLONIO	03	EPJ C27 331	M. Apollonio <i>et al.</i>	(CHOOZ Collab.)
ASTIER	03	PL B570 19	P. Astier <i>et al.</i>	(NOMAD Collab.)
EGUCHI	03	PRL 90 021802	K. Eguchi <i>et al.</i>	(KamLAND Collab.)
GANDO	03	PRL 90 171302	Y. Gando <i>et al.</i>	(Super-Kamiokande Collab.)
IANNI	03	JP G29 2107	A. lanni	(INFN Gran Sasso)
SANCHEZ	03	PR D68 113004	M. Sanchez <i>et al.</i>	(Soudan 2 Collab.)
ABDURASHI				
ADDURASIII	02	JETP 95 181 Translated from ZETF 12	J.N. Abdurashitov <i>et al.</i>	(SAGE Collab.)
AHMAD	02		Q.R. Ahmad <i>et al.</i>	(SNO Callab)
		PRL 89 011301	•	(SNO Collab.)
AHMAD	02B	PRL 89 011302	Q.R. Ahmad <i>et al.</i>	(SNO Collab.)
ARMBRUSTER		PR D65 112001	B. Armbruster <i>et al.</i>	(KARMEN 2 Collab.)
AVVAKUMOV	02	PRL 89 011804	S. Avvakumov <i>et al.</i>	(NuTeV Collab.)
FUKUDA	02	PL B539 179	S. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
AGUILAR	01	PR D64 112007	A. Aguilar <i>et al.</i>	(LSND Collab.)
AHMAD	01	PRL 87 071301	Q.R. Ahmad <i>et al.</i>	(SNO Collab.)
AMBROSIO	01	PL B517 59	M. Ambrosio <i>et al.</i>	(MÀCRO Collab.)
BOEHM	01	PR D64 112001	F. Boehm <i>et al.</i>	( , , , , , , , , , , , , , , , , , , ,
FUKUDA	01	PRL 86 5651	S. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
AMBROSIO	00	PL B478 5	M. Ambrosio <i>et al.</i>	(MACRO Collab.)
BOEHM	00	PRL 84 3764	F. Boehm <i>et al.</i>	(MACITO COMBD.)
				(Super Kamiekanda, Callah.)
FUKUDA	00	PRL 85 3999	S. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
ALLISON	99	PL B449 137	W.W.M. Allison <i>et al.</i>	(Soudan 2 Collab.)
APOLLONIO	99	PL B466 415	M. Apollonio <i>et al.</i>	(CHOOZ Collab.)
Also		PL B472 434 (errat.)	M. Apollonio <i>et al.</i>	(CHOOZ Collab.)
FUKUDA	99C	PRL 82 2644	Y. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
FUKUDA	99D	PL B467 185	Y. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
HAMPEL	99	PL B447 127	W. Hampel <i>et al.</i>	(GALLEX Collab.)
AMBROSIO	98	PL B434 451	M. Ambrosio et al.	(MACRO Collab.)
APOLLONIO	98	PL B420 397	M. Apollonio <i>et al.</i>	(CHOOZ Collab.)
ATHANASSO		PRL 81 1774	C. Athanassopoulos et al.	(LSND Collab.)
ATHANASSO		PR C58 2489	C. Athanassopoulos et al.	(LSND Collab.)
CLEVELAND	98	APJ 496 505	B.T. Cleveland <i>et al.</i>	(Homestake Collab.)
FELDMAN	98	PR D57 3873	G.J. Feldman, R.D. Cousins	(Homestake Collab.)
FUKUDA	98C	PRL 81 1562	Y. Fukuda <i>et al.</i>	(Super-Kamiokande Collab.)
HATAKEYAMA				
		PRL 81 2016	S. Hatakeyama <i>et al.</i>	(Kamiokande Collab.)
CLARK	97	PRL 79 345	R. Clark <i>et al.</i>	(IMB Collab.)
ROMOSAN	97	PRL 78 2912	A. Romosan <i>et al.</i>	(CCFR Collab.)
AGLIETTA	96	JETPL 63 791	M. Aglietta <i>et al.</i>	(LSD Collab.)
	00	Translated from ZETFP		
ATHANASSO		PR C54 2685	C. Athanassopoulos <i>et al.</i>	(LSND Collab.)
ATHANASSO		PRL 77 3082	C. Athanassopoulos <i>et al.</i>	(LSND Collab.)
FUKUDA	96	PRL 77 1683	Y. Fukuda <i>et al.</i>	(Kamiokande Collab.)
FUKUDA	96B	PL B388 397	Y. Fukuda <i>et al.</i>	(Kamiokande Collab.)
GREENWOOD	96	PR D53 6054	Z.D. Greenwood <i>et al.</i>	(UCI, SVR, SCUC)
HAMPEL	96	PL B388 384	W. Hampel <i>et al.</i>	(GALLEX Collab.)
LOVERRE	96	PL B370 156	P.F. Loverre	, , , , , , , , , , , , , , , , , , ,
ACHKAR	95	NP B434 503		G, SACLD, CPPM, CDEF+)
AHLEN	95	PL B357 481	S.P. Ahlen <i>et al.</i>	(MACRO Collab.)
ATHANASSO		PRL 75 2650	C. Athanassopoulos <i>et al.</i>	(LSND Collab.)
DAUM	. 95 95		K. Daum <i>et al.</i>	(FREJUS Collab.)
		ZPHY C66 417	J.E. Hill	
	95 95	PRL 75 2654	-	(PENN)
	45	PRL 75 3993	K.S. McFarland <i>et al.</i>	(CCFR Collab.)
MCFARLAND				
MCFARLAND DECLAIS	94	PL B338 383	Y. Declais <i>et al.</i>	
MCFARLAND DECLAIS FUKUDA	94 94	PL B335 237	Y. Fukuda <i>et al.</i>	(Kamiokande Collab.)
MCFARLAND DECLAIS FUKUDA VILAIN	94 94 94C	PL B335 237 ZPHY C64 539	Y. Fukuda <i>et al.</i> P. Vilain <i>et al.</i>	(CHARM II Collab.)
MCFARLAND DECLAIS FUKUDA VILAIN FREEDMAN	94 94 94C 93	PL B335 237 ZPHY C64 539 PR D47 811	Y. Fukuda <i>et al.</i> P. Vilain <i>et al.</i> S.J. Freedman <i>et al.</i>	(CHARM II Collab.) (LAMPF E645 Collab.)
MCFARLAND DECLAIS FUKUDA VILAIN	94 94 94C	PL B335 237 ZPHY C64 539	Y. Fukuda <i>et al.</i> P. Vilain <i>et al.</i>	(CHARM II Collab.)

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