

SEARCHES NOT IN OTHER SECTIONS

Magnetic Monopole Searches

Isolated supermassive monopole candidate events have not been confirmed. The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$< 1.4 \times 10^{-16} \text{ cm}^{-2} \text{sr}^{-1} \text{s}^{-1} \quad \text{for } 1.1 \times 10^{-4} < \beta < 1$$

Supersymmetric Particle Searches

All supersymmetric mass bounds here are model dependent.

The limits assume:

1) $\tilde{\chi}_1^0$ is the lightest supersymmetric particle; 2) R -parity is conserved;

See the Particle Listings for a Note giving details of supersymmetry.

$\tilde{\chi}_i^0$ — neutralinos (mixtures of $\tilde{\gamma}$, \tilde{Z}^0 , and \tilde{H}_i^0)

Mass $m_{\tilde{\chi}_1^0} > 0 \text{ GeV}$, CL = 95%

[general MSSM, non-universal gaugino masses]

Mass $m_{\tilde{\chi}_1^0} > 46 \text{ GeV}$, CL = 95%

[all $\tan\beta$, all m_0 , all $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$]

Mass $m_{\tilde{\chi}_2^0} > 670 \text{ GeV}$, CL = 95%

[$3/4\ell + \cancel{E}_T$, Tn2n3B, $m_{\tilde{\chi}_1^0} < 200 \text{ GeV}$]

Mass $m_{\tilde{\chi}_3^0} > 670 \text{ GeV}$, CL = 95%

[$3/4\ell + \cancel{E}_T$, Tn2n3B, $m_{\tilde{\chi}_1^0} < 200 \text{ GeV}$]

Mass $m_{\tilde{\chi}_4^0} > 116 \text{ GeV}$, CL = 95%

[$1 < \tan\beta < 40$, all m_0 , all $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$]

$\tilde{\chi}_i^\pm$ — charginos (mixtures of \tilde{W}^\pm and \tilde{H}_i^\pm)

Mass $m_{\tilde{\chi}_1^\pm} > 94 \text{ GeV}$, CL = 95%

[$\tan\beta < 40$, $m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} > 3 \text{ GeV}$, all m_0]

Mass $m_{\tilde{\chi}_1^\pm} > 500 \text{ GeV}$, CL = 95%

[$2\ell^\pm + \cancel{E}_T$, Tchi1chi1B, $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$]

$\tilde{\chi}^\pm$ — long-lived chargino

Mass $m_{\tilde{\chi}^\pm} > 620 \text{ GeV}$, CL = 95% [stable $\tilde{\chi}^\pm$]

$\tilde{\nu}$ — sneutrinoMass $m > 41$ GeV, CL = 95% [model independent]Mass $m > 94$ GeV, CL = 95%[CMSSM, $1 \leq \tan\beta \leq 40$, $m_{\tilde{e}_R} - m_{\tilde{\chi}_1^0} > 10$ GeV]Mass $m > 2300$ GeV, CL = 95%[RPV, $\tilde{\nu}_\tau \rightarrow e\mu$, $\lambda'_{311} = 0.11$] \tilde{e} — scalar electron (selectron)Mass $m(\tilde{e}_L) > 107$ GeV, CL = 95% [all $m_{\tilde{e}_L} - m_{\tilde{\chi}_1^0}$]Mass $m > 410$ GeV, CL = 95%[RPV, $\geq 4\ell^\pm$, $\tilde{\ell} \rightarrow l\tilde{\chi}_1^0$, $\tilde{\chi}_1^0 \rightarrow \ell^\pm \ell^\mp \nu$] $\tilde{\mu}$ — scalar muon (smuon)Mass $m > 94$ GeV, CL = 95%[CMSSM, $1 \leq \tan\beta \leq 40$, $m_{\tilde{\mu}_R} - m_{\tilde{\chi}_1^0} > 10$ GeV]Mass $m > 410$ GeV, CL = 95%[RPV, $\geq 4\ell^\pm$, $\tilde{\ell} \rightarrow l\tilde{\chi}_1^0$, $\tilde{\chi}_1^0 \rightarrow \ell^\pm \ell^\mp \nu$] $\tilde{\tau}$ — scalar tau (stau)Mass $m > 81.9$ GeV, CL = 95%[$m_{\tilde{\tau}_R} - m_{\tilde{\chi}_1^0} > 15$ GeV, all θ_τ , $B(\tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0) = 100\%$]Mass $m > 286$ GeV, CL = 95% [long-lived $\tilde{\tau}$] \tilde{q} — squarks of the first two quark generationsMass $m > 1450$ GeV, CL = 95%[CMSSM, $\tan\beta = 30$, $A_0 = -2\max(m_0, m_{1/2})$, $\mu > 0$]Mass $m > 1550$ GeV, CL = 95%

[mass degenerate squarks]

Mass $m > 1050$ GeV, CL = 95%

[single light squark bounds]

 \tilde{q} — long-lived squarkMass $m > 1000$, CL = 95%[\tilde{t} , charge-suppressed interaction model]Mass $m > 845$, CL = 95% [\tilde{b} , stable, Regge model] \tilde{b} — scalar bottom (sbottom)Mass $m > 1230$ GeV, CL = 95%[jets+ \cancel{E}_T , Tsb01, $m_{\tilde{\chi}_1^0} = 0$ GeV]

\tilde{t} — scalar top (stop)

Mass $m > 1120$ GeV, CL = 95%
 [$1\ell + \text{jets} + \cancel{E}_T$, Tstop1, $m_{\tilde{\chi}_1^0} = 0$ GeV]

\tilde{g} — gluino

Mass $m > 1860$ GeV, CL = 95%
 [≥ 1 jets + \cancel{E}_T , Tglu1A, $m_{\tilde{\chi}_1^0} = 0$ GeV]

Technicolor

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full *Review* (the data listings).

Quark and Lepton Compositeness, Searches for

Scale Limits Λ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \bar{\psi}_L \gamma_\mu \psi_L \bar{\psi}_L \gamma^\mu \psi_L$$

(with $g^2/4\pi$ set equal to 1), then we define $\Lambda \equiv \Lambda_{LL}^\pm$. For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

| | |
|------------------------------------|------------------------|
| $\Lambda_{LL}^+(eeee)$ | > 8.3 TeV, CL = 95% |
| $\Lambda_{LL}^-(eeee)$ | > 10.3 TeV, CL = 95% |
| $\Lambda_{LL}^+(ee\mu\mu)$ | > 8.5 TeV, CL = 95% |
| $\Lambda_{LL}^-(ee\mu\mu)$ | > 9.5 TeV, CL = 95% |
| $\Lambda_{LL}^+(ee\tau\tau)$ | > 7.9 TeV, CL = 95% |
| $\Lambda_{LL}^-(ee\tau\tau)$ | > 7.2 TeV, CL = 95% |
| $\Lambda_{LL}^+(\ell\ell\ell\ell)$ | > 9.1 TeV, CL = 95% |
| $\Lambda_{LL}^-(\ell\ell\ell\ell)$ | > 10.3 TeV, CL = 95% |
| $\Lambda_{LL}^+(eeqq)$ | > 24 TeV, CL = 95% |
| $\Lambda_{LL}^-(eeqq)$ | > 37 TeV, CL = 95% |

| | |
|-----------------------------|---------------------------------------|
| $\Lambda_{LL}^+(eeuu)$ | > 23.3 TeV, CL = 95% |
| $\Lambda_{LL}^-(eeuu)$ | > 12.5 TeV, CL = 95% |
| $\Lambda_{LL}^+(eedd)$ | > 11.1 TeV, CL = 95% |
| $\Lambda_{LL}^-(eedd)$ | > 26.4 TeV, CL = 95% |
| $\Lambda_{LL}^+(eecc)$ | > 9.4 TeV, CL = 95% |
| $\Lambda_{LL}^-(eecc)$ | > 5.6 TeV, CL = 95% |
| $\Lambda_{LL}^+(eebb)$ | > 9.4 TeV, CL = 95% |
| $\Lambda_{LL}^-(eebb)$ | > 10.2 TeV, CL = 95% |
| $\Lambda_{LL}^+(\mu\mu qq)$ | > 20 TeV, CL = 95% |
| $\Lambda_{LL}^-(\mu\mu qq)$ | > 30 TeV, CL = 95% |
| $\Lambda(\ell\nu\ell\nu)$ | > 3.10 TeV, CL = 90% |
| $\Lambda(e\nu qq)$ | > 2.81 TeV, CL = 95% |
| $\Lambda_{LL}^+(qqqq)$ | > 13.1 none 17.4–29.5 TeV, CL = 95% |
| $\Lambda_{LL}^-(qqqq)$ | > 21.8 TeV, CL = 95% |
| $\Lambda_{LL}^+(\nu\nu qq)$ | > 5.0 TeV, CL = 95% |
| $\Lambda_{LL}^-(\nu\nu qq)$ | > 5.4 TeV, CL = 95% |

Excited Leptons

The limits from $\ell^{*+}\ell^{*-}$ do not depend on λ (where λ is the $\ell\ell^*$ transition coupling). The λ -dependent limits assume chiral coupling.

$e^{*\pm}$ — excited electron

Mass $m > 103.2$ GeV, CL = 95% (from e^*e^*)

Mass $m > 3.000 \times 10^3$ GeV, CL = 95% (from ee^*)

Mass $m > 356$ GeV, CL = 95% (if $\lambda_\gamma = 1$)

$\mu^{*\pm}$ — excited muon

Mass $m > 103.2$ GeV, CL = 95% (from $\mu^*\mu^*$)

Mass $m > 3.000 \times 10^3$ GeV, CL = 95% (from $\mu\mu^*$)

$\tau^{*\pm}$ — excited tau

Mass $m > 103.2$ GeV, CL = 95% (from $\tau^*\tau^*$)

Mass $m > 2.500 \times 10^3$ GeV, CL = 95% (from $\tau\tau^*$)

ν^* — excited neutrino

Mass $m > 1.600 \times 10^3$ GeV, CL = 95% (from $\nu^*\nu^*$)

Mass $m > 213$ GeV, CL = 95% (from ν^*X)

q^* — excited quark

Mass $m > 338$ GeV, CL = 95% (from $q^* q^*$)

Mass $m > 6.000 \times 10^3$ GeV, CL = 95% (from $q^* X$)

Color Sextet and Octet Particles

Color Sextet Quarks (q_6)

Mass $m > 84$ GeV, CL = 95% (Stable q_6)

Color Octet Charged Leptons (ℓ_8)

Mass $m > 86$ GeV, CL = 95% (Stable ℓ_8)

Color Octet Neutrinos (ν_8)

Mass $m > 110$ GeV, CL = 90% ($\nu_8 \rightarrow \nu g$)

Extra Dimensions

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

$R < 30 \mu\text{m}$, CL = 95% (direct tests of Newton's law)

$R < 10.9 \mu\text{m}$, CL = 95% ($pp \rightarrow jG$)

$R < 0.16\text{--}916$ nm (astrophysics; limits depend on technique and assumptions)

Constraints on the fundamental gravity scale

$M_{TT} > 8.4$ TeV, CL = 95% ($pp \rightarrow$ dijet, angular distribution)

$M_c > 4.16$ TeV, CL = 95% ($pp \rightarrow \ell\bar{\ell}$)

Constraints on the Kaluza-Klein graviton in warped extra dimensions

$M_G > 4.1$ TeV, CL = 95% ($pp \rightarrow \gamma\gamma$)

Constraints on the Kaluza-Klein gluon in warped extra dimensions

$M_{g_{KK}} > 2.5$ TeV, CL = 95% ($g_{KK} \rightarrow t\bar{t}$)