## Light Quarks ( $u, d, s$ )

OMITTED FROM SUMMARY TABLE

## u-QUARK MASS

The $u$-, $d$-, and $s$-quark masses are estimates of so-called "current-quark masses," in a mass- independent subtraction scheme such as $\overline{\mathrm{MS}}$. The ratios $m_{u} / m_{d}$ and $m_{s} / m_{d}$ are extracted from pion and kaon masses using chiral symmetry. The estimates of $d$ and $u$ masses are not without controversy and remain under active investigation. Within the literature there are even suggestions that the $u$ quark could be essentially massless. The s-quark mass is estimated from $\operatorname{SU}(3)$ splittings in hadron masses.

We have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2$ GeV . Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35. The values of "Our Evaluation" were determined in part via Figures 1 and 2.

| $\overline{\text { MS MASS }(\mathrm{MeV})}$ |  | DOCUMENT ID |  | TECN |
| :---: | :---: | :---: | :---: | :---: |
| $2.16{ }_{-0.26}^{+0.49}$ | OUR EVALUATION | See the ideogram | bel |  |
| $2.6 \pm 0.4$ |  | 1 DOMINGUEZ | 19 | THEO |
| $2.130 \pm 0.041$ |  | 2 BAZAVOV | 18 | LATT |
| $2.27 \pm 0.06$ | $\pm 0.06$ | 3 FODOR | 16 | LATT |
| $2.36 \pm 0.24$ |  | ${ }^{4}$ CARRASCO | 14 | LATT |
| $2.57 \pm 0.26$ | $\pm 0.07$ | ${ }^{5}$ AOKI | 12 | LATT |
| $2.24 \pm 0.10$ | $\pm 0.34$ | ${ }^{6}$ BLUM | 10 | LATT |
| $2.01 \pm 0.14$ |  | 7 MCNEILE | 10 | LATT |

-     - We do not use the following data for averages, fits, limits, etc.

| 2.15 | $\pm 0.03$ | $\pm 0.10$ | 8 | DURR |
| :--- | :--- | :--- | :--- | :--- |
| 1.9 | $\pm 0.2$ | 9 | 11 | LATT |
| 2.01 | $\pm 0.14$ | 7 | DAVIES | 10 |
| LATT |  |  |  |  |
| 2.9 | $\pm 0.2$ | 10 DOMINGUEZ | 09 | LATT |
| 2.9 | $\pm 0.8$ | 11 DEANDREA | 08 | THEO |
| 3.02 | $\pm 0.33$ | 12 BLUM | 07 | LATT |
| 2.7 | $\pm 0.4$ | 13 JAMIN | 06 | THEO |
| 1.9 | $\pm 0.2$ | 14 MASON | 06 | LATT |
| 2.8 | $\pm 0.2$ | 15 NARISON | 06 | THEO |
| 1.7 | $\pm 0.3$ | 16 AUBIN | $04 A$ | LATT |

${ }^{1}$ DOMINGUEZ 19 determine the quark mass from a QCD finite energy sum rule for the divergence of the axial current.
${ }^{2}$ BAZAVOV 18 determine the quark masses using a lattice computation with staggered fermions and four active quark flavors.
${ }^{3}$ FODOR 16 is a lattice simulation with $N_{f}=2+1$ dynamical flavors and includes partially quenched QED effects.
${ }^{4}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{5}$ AOKI 12 is a lattice computation using $1+1+1$ dynamical quark flavors.
${ }^{6}$ BLUM 10 determines light quark masses using a QCD plus QED lattice computation of the electromagnetic mass splittings of the low-lying hadrons. The lattice simulations use $2+1$ dynamical quark flavors.
${ }^{7}$ DAVIES 10 and MCNEILE 10 determine $\bar{m}_{c}(\mu) / \bar{m}_{s}(\mu)=11.85 \pm 0.16$ using a lattice computation with $N_{f}=2+1$ dynamical fermions of the pseudoscalar meson masses. Mass $m_{u}$ is obtained from this using the value of $m_{c}$ from ALLISON 08 or MCNEILE 10 and the BAZAVOV 10 values for the light quark mass ratios, $m_{s} / \bar{m}$ and $m_{u} / m_{d}$.
${ }^{8}$ DURR 11 determine quark mass from a lattice computation of the meson spectrum using $N_{f}=2+1$ dynamical flavors. The lattice simulations were done at the physical quark mass, so that extrapolation in the quark mass was not needed. The individual $m_{u}, m_{d}$ values are obtained using the lattice determination of the average mass $m_{u d}$ and of the ratio $m_{s} / m_{\mathrm{ud}}$ and the value of $Q=\left(m_{s}^{2}-m_{\mathrm{ud}}^{2}\right) /\left(m_{d}^{2}-m_{u}^{2}\right)$ as determined from $\eta \rightarrow 3 \pi$ decays.
${ }^{9}$ BAZAVOV 10 is a lattice computation using $2+1$ dynamical quark flavors.
${ }^{10}$ DOMINGUEZ 09 use QCD finite energy sum rules for the two-point function of the divergence of the axial vector current computed to order $\alpha_{s}^{4}$.
${ }^{11}$ DEANDREA 08 determine $m_{u}-m_{d}$ from $\eta \rightarrow 3 \pi^{0}$, and combine with the PDG 06 lattice average value of $m_{u}+m_{d}=7.6 \pm 1.6$ to determine $m_{u}$ and $m_{d}$.
${ }^{12}$ BLUM 07 determine quark masses from the pseudoscalar meson masses using a QED plus QCD lattice computation with two dynamical quark flavors.
13 JAMIN 06 determine $m_{u}(2 \mathrm{GeV})$ by combining the value of $m_{s}$ obtained from the spectral function for the scalar $K \pi$ form factor with other determinations of the quark mass ratios.
14 MASON 06 extract light quark masses from a lattice simulation using staggered fermions with an improved action, and three dynamical light quark flavors with degenerate $u$ and $d$ quarks. Perturbative corrections were included at NNLO order. The quark masses $m_{u}$ and $m_{d}$ were determined from their $\left(m_{u}+m_{d}\right) / 2$ measurement and AUBIN 04A $m_{u} / m_{d}$ value.
${ }^{15}$ NARISON 06 uses sum rules for $e^{+} e^{-} \rightarrow$ hadrons to order $\alpha_{s}^{3}$ to determine $m_{s}$ combined with other determinations of the quark mass ratios.
${ }^{16}$ AUBIN 04 A employ a partially quenched lattice calculation of the pseudoscalar meson masses.


## d-QUARK MASS

See the comment for the $u$ quark above.
We have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2$ GeV . Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 . The values of "Our Evaluation" were determined in part via Figures 1 and 2.

| $\overline{\text { MS MASS ( } \mathrm{MeV} \text { ) }}$ |  | DOCUMENT ID |  | TECN |
| :---: | :---: | :---: | :---: | :---: |
| $4.67{ }_{-0.17}^{+0.48}$ | OUR EVALUATION | See the ideogr | be |  |
| $5.3 \pm 0.4$ |  | 1 DOMINGUEZ | 19 | THEO |
| $4.675 \pm 0.056$ |  | 2 BAZAVOV | 18 | LATT |
| $4.67 \pm 0.06$ | $\pm 0.06$ | 3 FODOR | 16 | LATT |
| $5.03 \pm 0.26$ |  | ${ }^{4}$ CARRASCO | 14 | LATT |
| $3.68 \pm 0.29$ | $\pm 0.10$ | ${ }^{5}$ AOKI | 12 | LATT |
| $4.65 \pm 0.15$ | $\pm 0.32$ | ${ }^{6}$ BLUM | 10 | LATT |
| $4.77 \pm 0.15$ |  | 7 MCNEILE | 10 | LATT |

-     - We do not use the following data for averages, fits, limits, etc.
$4.79 \pm 0.07 \pm 0.12$
$4.6 \pm 0.3$
8 DURR
$4.79 \pm 0.16$
9 BAZAVOV 10 LATT
$5.3 \pm 0.4$
7 DAVIES 10 LATT
$4.7 \pm 0.8$
10 DOMINGUEZ 09 THEO
11 DEANDREA 08 THEO

| 5.49 | $\pm 0.39$ | 12 BLUM | 07 | LATT |
| :--- | :--- | :--- | :--- | :--- |
| 4.8 | $\pm 0.5$ | 13 JAMIN | 06 | THEO |
| 4.4 | $\pm 0.3$ | 14 MASON | 06 | LATT |
| 5.1 | $\pm 0.4$ | 15 NARISON | 06 | THEO |
| 3.9 | $\pm 0.5$ | 16 AUBIN | $04 A$ | LATT |

${ }^{1}$ DOMINGUEZ 19 determine the quark mass from a QCD finite energy sum rule for the divergence of the axial current.
${ }^{2}$ BAZAVOV 18 determine the quark masses using a lattice computation with staggered fermions and four active quark flavors.
${ }^{3}$ FODOR 16 is a lattice simulation with $N_{f}=2+1$ dynamical flavors and includes partially quenched QED effects.
${ }^{4}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{5}$ AOKI 12 is a lattice computation using $1+1+1$ dynamical quark flavors.
${ }^{6}$ BLUM 10 determines light quark masses using a QCD plus QED lattice computation of the electromagnetic mass splittings of the low-lying hadrons. The lattice simulations use $2+1$ dynamical quark flavors.
7 DAVIES 10 and MCNEILE 10 determine $\bar{m}_{c}(\mu) / \bar{m}_{S}(\mu)=11.85 \pm 0.16$ using a lattice computation with $N_{f}=2+1$ dynamical fermions of the pseudoscalar meson masses. Mass $m_{d}$ is obtained from this using the value of $m_{c}$ from ALLISON 08 or MCNEILE 10 and the BAZAVOV 10 values for the light quark mass ratios, $m_{s} / \bar{m}$ and $m_{u} / m_{d}$.
${ }^{8}$ DURR 11 determine quark mass from a lattice computation of the meson spectrum using $N_{f}=2+1$ dynamical flavors. The lattice simulations were done at the physical quark mass, so that extrapolation in the quark mass was not needed. The individual $m_{u}, m_{d}$ values are obtained using the lattice determination of the average mass $m_{u d}$ and of the ratio $m_{s} / m_{\mathrm{ud}}$ and the value of $Q=\left(m_{s}^{2}-m_{\mathrm{ud}}^{2}\right) /\left(m_{d}^{2}-m_{u}^{2}\right)$ as determined from $\eta \rightarrow 3 \pi$ decays.
${ }^{9}$ BAZAVOV 10 is a lattice computation using $2+1$ dynamical quark flavors.
10 DOMINGUEZ 09 use QCD finite energy sum rules for the two-point function of the divergence of the axial vector current computed to order $\alpha_{s}^{4}$.
11 DEANDREA 08 determine $m_{u}-m_{d}$ from $\eta \rightarrow 3 \pi^{0}$, and combine with the PDG 06 lattice average value of $m_{u}+m_{d}=7.6 \pm 1.6$ to determine $m_{u}$ and $m_{d}$.
12 BLUM 07 determine quark masses from the pseudoscalar meson masses using a QED plus QCD lattice computation with two dynamical quark flavors.
13 JAMIN 06 determine $m_{d}(2 \mathrm{GeV})$ by combining the value of $m_{s}$ obtained from the spectral function for the scalar $K \pi$ form factor with other determinations of the quark mass ratios.
14 MASON 06 extract light quark masses from a lattice simulation using staggered fermions with an improved action, and three dynamical light quark flavors with degenerate $u$ and $d$ quarks. Perturbative corrections were included at NNLO order. The quark masses $m_{u}$ and $m_{d}$ were determined from their $\left(m_{u}+m_{d}\right) / 2$ measurement and AUBIN 04A $m_{u} / m_{d}$ value.
15 NARISON 06 uses sum rules for $e^{+} e^{-} \rightarrow$ hadrons to order $\alpha_{s}^{3}$ to determine $m_{s}$ combined with other determinations of the quark mass ratios.
16 AUBIN 04A perform three flavor dynamical lattice calculation of pseudoscalar meson masses, with continuum estimate of electromagnetic effects in the kaon masses, and one-loop perturbative renormalization constant.


$$
\bar{m}=\left(m_{u}+m_{d}\right) / \mathbf{2}
$$

See the comments for the $u$ quark above.
We have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2$ GeV . Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 . The values of "Our Evaluation" were determined in part via Figures 1 and 2.

| $\overline{\text { MS MASS ( }} \mathrm{MeV}$ ) |  |  | DOCUMENT ID |  | TECN |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 3.45 | +0.55 -0.15 | OUR EVALUATION | See the ideogram below. |  |  |
| 3.9 | $\pm 0.3$ |  | 1 DOMINGUEZ | 19 | THEO |
| 4.7 | $\begin{array}{r} +0.8 \\ -0.7 \end{array}$ |  | 2 YUAN | 17 | THEO |
| 3.70 | $\pm 0.17$ |  | 3 CARRASCO | 14 | LATT |
| 3.45 | $\pm 0.12$ |  | ${ }^{4}$ ARTHUR | 13 | LATT |
| 3.469 | $\pm 0.047$ | $\pm 0.048$ - 5 | 5 DURR | 11 | LATT |
| 3.6 | $\pm 0.2$ |  | 6 BLOSSIER | 10 | LATT |
| 3.39 | $\pm 0.06$ |  | 7 MCNEILE | 10 | LATT |

-     - We do not use the following data for averages, fits, limits, etc.

| 3.59 | $\pm 0.21$ | 8 | AOKI | 11 A |
| :---: | :--- | :--- | :--- | :--- | LATT

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| $3.55{ }_{-0.28}^{+0.65}$ | 13 ISHIKAWA | 08 | LATT |
| :---: | :---: | :---: | :---: |
| $4.026 \pm 0.048$ | 14 NAKAMURA | 08 | LATT |
| $4.25 \pm 0.35$ | 15 BLUM | 07 | LATT |
| $4.08 \pm 0.25 \pm 0.42$ | 16 GOCKELER | 06 | LATT |
| $4.7 \pm 0.2 \pm 0.3$ | 17 GOCKELER | 06A | LATT |
| $3.2 \pm 0.3$ | 18 MASON | 06 | LATT |
| $3.95 \pm 0.3$ | 19 NARISON | 06 | THEO |
| $2.8 \pm 0.3$ | 20 AUBIN | 04 | LATT |
| $4.29 \pm 0.14 \pm 0.65$ | 21 AOKI | 03 | LATT |
| $3.223 \pm 0.3$ | 22 AOKI | 03B | LATT |
| $4.4 \pm 0.1 \quad \pm 0.4$ | 23 BECIREVIC | 03 | LATT |
| $4.1 \pm 0.3 \pm 1.0$ | 24 CHIU | 03 | LATT |

${ }^{1}$ DOMINGUEZ 19 determine the quark mass from a QCD finite energy sum rule for the divergence of the axial current.
${ }^{2}$ YUAN 17 determine $\bar{m}$ using QCD sum rules in the isospin $I=0$ scalar channel. At the end of the "Numerical Results" section of YUAN 17 the authors discuss the significance of their larger value of the light quark mass compared to previous determinations.
${ }^{3}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{4}$ ARTHUR 13 is a lattice computation using $2+1$ dynamical domain wall fermions. Masses at $\mu=3 \mathrm{GeV}$ have been converted to $\mu=2 \mathrm{GeV}$ using conversion factors given in their paper.
${ }^{5}$ DURR 11 determine quark mass from a lattice computation of the meson spectrum using $N_{f}=2+1$ dynamical flavors. The lattice simulations were done at the physical quark mass, so that extrapolation in the quark mass was not needed.
${ }^{6}$ BLOSSIER 10 determines quark masses from a computation of the hadron spectrum using $N_{f}=2$ dynamical twisted-mass Wilson fermions.
${ }^{7}$ DAVIES 10 and MCNEILE 10 determine $\bar{m}_{c}(\mu) / \bar{m}_{s}(\mu)=11.85 \pm 0.16$ using a lattice computation with $N_{f}=2+1$ dynamical fermions of the pseudoscalar meson masses. Mass $\bar{m}$ is obtained from this using the value of $m_{c}$ from ALLISON 08 or MCNEILE 10 and the BAZAVOV 10 values for the light quark mass ratio, $m_{s} / \bar{m}$.
${ }^{8}$ AOKI 11A determine quark masses from a lattice computation of the hadron spectrum using $N_{f}=2+1$ dynamical flavors of domain wall fermions.
${ }^{9}$ DOMINGUEZ 09 use QCD finite energy sum rules for the two-point function of the divergence of the axial vector current computed to order $\alpha_{S}^{4}$.
10 ALLTON 08 use a lattice computation of the $\pi, K$, and $\Omega$ masses with $2+1$ dynamical flavors of domain wall quarks, and non-perturbative renormalization.
11 BLOSSIER 08 use a lattice computation of pseudoscalar meson masses and decay constants with 2 dynamical flavors and non-perturbative renormalization.
12 DOMINGUEZ-CLARIMON 08B obtain an inequality from sum rules for the scalar twopoint correlator.
13 ISHIKAWA 08 use a lattice computation of the light meson spectrum with $2+1$ dynamical flavors of $\mathcal{O}(a)$ improved Wilson quarks, and one-loop perturbative renormalization.
14 NAKAMURA 08 do a lattice computation using quenched domain wall fermions and non-perturbative renormalization.
${ }^{15}$ BLUM 07 determine quark masses from the pseudoscalar meson masses using a QED plus QCD lattice computation with two dynamical quark flavors.
16 GOCKELER 06 use an unquenched lattice computation of the axial Ward Identity with $N_{f}=2$ dynamical light quark flavors, and non-perturbative renormalization, to obtain $\bar{m}(2 \mathrm{GeV})=4.08 \pm 0.25 \pm 0.19 \pm 0.23 \mathrm{MeV}$, where the first error is statistical, the second
and third are systematic due to the fit range and force scale uncertainties, respectively. We have combined the systematic errors linearly.
17 GOCKELER 06A use an unquenched lattice computation of the pseudoscalar meson masses with $N_{f}=2$ dynamical light quark flavors, and non-perturbative renormalization.
18 MASON 06 extract light quark masses from a lattice simulation using staggered fermions with an improved action, and three dynamical light quark flavors with degenerate $u$ and $d$ quarks. Perturbative corrections were included at NNLO order.
19 NARISON 06 uses sum rules for $e^{+} e^{-} \rightarrow$ hadrons to order $\alpha_{s}^{3}$ to determine $m_{s}$ combined with other determinations of the quark mass ratios.
${ }^{20}$ AUBIN 04 perform three flavor dynamical lattice calculation of pseudoscalar meson masses, with one-loop perturbative renormalization constant.
21 AOKI 03 uses quenched lattice simulation of the meson and baryon masses with degenerate light quarks. The extrapolations are done using quenched chiral perturbation theory.
22 The errors given in AOKI 03B were ${ }_{-0}^{+0.046}$. We changed them to $\pm 0.3$ for calculating the overall best values. AOKI 03B uses lattice simulation of the meson and baryon masses with two dynamical light quarks. Simulations are performed using the $\mathcal{O}(a)$ improved Wilson action.
23 BECIREVIC 03 perform quenched lattice computation using the vector and axial Ward identities. Uses $\mathcal{O}(a)$ improved Wilson action and nonperturbative renormalization.
${ }^{24}$ CHIU 03 determines quark masses from the pion and kaon masses using a lattice simulation with a chiral fermion action in quenched approximation.


## $\boldsymbol{m}_{\boldsymbol{u}} / \boldsymbol{m}_{\boldsymbol{d}}$ MASS RATIO

VALUE DOCUMENT ID TECN COMMENT


-     - We do not use the following data for averages, fits, limits, etc.

| 0.550 | $\pm 0.031$ |
| :--- | :--- |
| 0.43 | $\pm 0.08$ |
| 0.410 | $\pm 0.036$ |
| 0.553 | $\pm 0.043$ |


| 7 BLUM | 07 | LATT |
| :---: | :--- | :--- |
| 8 AUBIN | $04 A$ | LATT |
| 9 NELSON | 03 | LATT |
| 10 LEUTWYLER | 96 | THEO |

Compilation

${ }^{1}$ FODOR 16 is a lattice simulation with $N_{f}=2+1$ dynamical flavors and includes partially quenched QED effects.
${ }^{2}$ BASAK 15 is a lattice computation using $2+1$ dynamical quark flavors.
${ }^{3}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{4}$ AOKI 12 is a lattice computation using $1+1+1$ dynamical quark flavors.
${ }^{5}$ BAZAVOV 10 is a lattice computation using $2+1$ dynamical quark flavors.
${ }^{6}$ BLUM 10 is a lattice computation using $2+1$ dynamical quark flavors.
${ }^{7}$ BLUM 07 determine quark masses from the pseudoscalar meson masses using a QED plus QCD lattice computation with two dynamical quark flavors.
${ }^{8}$ AUBIN 04A perform three flavor dynamical lattice calculation of pseudoscalar meson masses, with continuum estimate of electromagnetic effects in the kaon masses.
${ }^{9}$ NELSON 03 computes coefficients in the order $p^{4}$ chiral Lagrangian using a lattice calculation with three dynamical flavors. The ratio $m_{u} / m_{d}$ is obtained by combining this with the chiral perturbation theory computation of the meson masses to order $p^{4}$.
10 LEUTWYLER 96 uses a combined fit to $\eta \rightarrow 3 \pi$ and $\psi^{\prime} \rightarrow J / \psi(\pi, \eta)$ decay rates, and the electromagnetic mass differences of the $\pi$ and $K$.

## $s$-QUARK MASS

See the comment for the $u$ quark above.
We have normalized the $\overline{\mathrm{MS}}$ masses at a renormalization scale of $\mu=2$ GeV . Results quoted in the literature at $\mu=1 \mathrm{GeV}$ have been rescaled by dividing by 1.35 .


-     - We do not use the following data for averages, fits, limits, etc. - • -

| $93.6 \pm 0.8$ | 9 CHAKRABOR.. 15 |  | LATT |
| :---: | :---: | :---: | :---: |
| $96.2 \pm 2.7$ | 10 AOKI | 11A | LATT |
| $95 \pm 6$ | 11 BLOSSIER | 10 | LATT |
| $97.6 \pm 2.9 \pm 5.5$ | 12 BLUM | 10 | LATT |
| $92.4 \pm 1.5$ | 13 DAVIES | 10 | LATT |
| $92.2 \pm 1.3$ | 13 MCNEILE | 10 | LATT |
| $107.3 \pm 11.7$ | 14 ALLTON | 08 | LATT |
| $105 \pm 3 \pm 9$ | 15 BLOSSIER | 08 | LATT |
| $102 \pm 8$ | 16 DOMINGUEZ | 08A | THEO |
| 90.1 +17.2 | 17 ISHIKAWA | 08 | LATT |
| $105.6 \pm 1.2$ | 18 NAKAMURA | 08 | LATT |
| $119.5 \pm 9.3$ | 19 BLUM | 07 | LATT |
| $105 \pm 6 \pm 7$ | 20 CHETYRKIN | 06 | THEO |
| $111 \pm 6 \pm 10$ | 21 GOCKELER | 06 | LATT |
| $119 \pm 5 \pm 8$ | 22 GOCKELER | 06A | LATT |
| $92 \pm 9$ | 23 JAMIN | 06 | THEO |
| $87 \pm 6$ | 24 MASON | 06 | LATT |
| $104 \pm 15$ | 25 NARISON | 06 | THEO |
| $\geq 71 \pm 4, \leq 151 \pm 14$ | 26 NARISON | 06 | THEO |
| $96 \begin{array}{r}+5 \\ -3\end{array}$ | 27 BAIKOV | 05 | THEO |


| 81 | $\pm 22$ | 28 GAMIZ | 05 | THEO |
| :---: | :---: | :---: | :---: | :---: |
| 125 | $\pm 28$ | 29 GORBUNOV | 05 | THEO |
| 93 | $\pm 32$ | 30 NARISON | 05 | THEO |
| 76 | $\pm 8$ | 31 AUBIN | 04 | LATT |
| 116 | $\pm 6 \pm 0.65$ | 32 AOKI | 03 | LATT |
| 84.5 | $\begin{aligned} & +12 \\ & -\quad 1.7 \end{aligned}$ | 33 AOKI | 03B | LATT |
| 106 | $\pm 2 \pm 8$ | 34 BECIREVIC | 03 | LATT |
| 92 | $\pm 9 \pm 16$ | 35 CHIU | 03 | LATT |
| 117 | $\pm 17$ | 36 GAMIZ | 03 | THEO |
| 103 | $\pm 17$ | 37 GAMIZ | 03 | THEO |

${ }^{1}$ BAZAVOV 18 determine the quark masses using a lattice computation with staggered fermions and four active quark flavors.
${ }^{2}$ LYTLE 18 combined with CHAKRABORTY 2015 determine $\bar{m}_{s}(3 \mathrm{GeV})=84.78 \pm 0.65$ MeV from a lattice simulation with $n_{f}=2+1+1$ flavors. They also determine the quoted value $\bar{m}_{s}(2 \mathrm{GeV})$ for $n_{f}=4$ dynamical flavors.
${ }^{3}$ ANANTHANARAYAN 16 determine $\bar{m}_{s}(2 \mathrm{GeV})=106.70 \pm 9.36 \mathrm{MeV}$ and $74.47 \pm 7.77$ MeV from fits to ALEPH and OPAL $\tau$ decay data, respectively. We have used the weighted average of the two.
${ }^{4}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{5}$ ARTHUR 13 is a lattice computation using $2+1$ dynamical domain wall fermions. Masses at $\mu=3 \mathrm{GeV}$ have been converted to $\mu=2 \mathrm{GeV}$ using conversion factors given in their paper.
$6_{\text {BODENSTEIN }} 13$ determines $m_{s}$ from QCD finite energy sum rules, and the perturbative computation of the pseudoscalar correlator to five-loop order.
${ }^{7}$ FRITZSCH 12 determine $m_{s}$ using a lattice computation with $N_{f}=2$ dynamical flavors.
8 DURR 11 determine quark mass from a lattice computation of the meson spectrum using $N_{f}=2+1$ dynamical flavors. The lattice simulations were done at the physical quark mass, so that extrapolation in the quark mass was not needed.
${ }^{9}$ CHAKRABORTY 15 is a lattice QCD computation that determines $m_{c}$ and $m_{c} / m_{s}$ using pseudoscalar mesons masses tuned on gluon field configurations with $2+1+1 \mathrm{dy}$ namical flavors of HISQ quarks with $u / d$ masses down to the physical value.
10 AOKI 11A determine quark masses from a lattice computation of the hadron spectrum using $N_{f}=2+1$ dynamical flavors of domain wall fermions.
11 BLOSSIER 10 determines quark masses from a computation of the hadron spectrum using $N_{f}=2$ dynamical twisted-mass Wilson fermions.
12 BLUM 10 determines light quark masses using a QCD plus QED lattice computation of the electromagnetic mass splittings of the low-lying hadrons. The lattice simulations use $2+1$ dynamical quark flavors.
13 DAVIES 10 and MCNEILE 10 determine $\bar{m}_{c}(\mu) / \bar{m}_{s}(\mu)=11.85 \pm 0.16$ using a lattice computation with $N_{f}=2+1$ dynamical fermions of the pseudoscalar meson masses. Mass $m_{s}$ is obtained from this using the value of $m_{c}$ from ALLISON 08 or MCNEILE 10.
14 ALLTON 08 use a lattice computation of the $\pi, K$, and $\Omega$ masses with $2+1$ dynamical flavors of domain wall quarks, and non-perturbative renormalization.
15 BLOSSIER 08 use a lattice computation of pseudoscalar meson masses and decay constants with 2 dynamical flavors and non-perturbative renormalization.
16 DOMINGUEZ 08A make determination from QCD finite energy sum rules for the pseudoscalar two-point function computed to order $\alpha_{s}^{4}$.
17 ISHIKAWA 08 use a lattice computation of the light meson spectrum with $2+1$ dynamical flavors of $\mathcal{O}(a)$ improved Wilson quarks, and one-loop perturbative renormalization.

18 NAKAMURA 08 do a lattice computation using quenched domain wall fermions and non-perturbative renormalization.
${ }^{19}$ BLUM 07 determine quark masses from the pseudoscalar meson masses using a QED plus QCD lattice computation with two dynamical quark flavors.
${ }^{20}$ CHETYRKIN 06 use QCD sum rules in the pseudoscalar channel to order $\alpha_{s}^{4}$.
21 GOCKELER 06 use an unquenched lattice computation of the axial Ward Identity with $N_{f}=2$ dynamical light quark flavors, and non-perturbative renormalization, to obtain $\bar{m}_{s}(2 \mathrm{GeV})=111 \pm 6 \pm 4 \pm 6 \mathrm{MeV}$, where the first error is statistical, the second and third are systematic due to the fit range and force scale uncertainties, respectively. We have combined the systematic errors linearly.
22 GOCKELER 06A use an unquenched lattice computation of the pseudoscalar meson masses with $N_{f}=2$ dynamical light quark flavors, and non-perturbative renormalization.
23 JAMIN 06 determine $\bar{m}_{S}(2 \mathrm{GeV})$ from the spectral function for the scalar $K \pi$ form factor.
24 MASON 06 extract light quark masses from a lattice simulation using staggered fermions with an improved action, and three dynamical light quark flavors with degenerate $u$ and $d$ quarks. Perturbative corrections were included at NNLO order.
25 NARISON 06 uses sum rules for $e^{+} e^{-} \rightarrow$ hadrons to order $\alpha_{s}^{3}$.
26 NARISON 06 obtains the quoted range from positivity of the spectral functions.
27 BAIKOV 05 determines $\bar{m}_{s}\left(M_{\tau}\right)=100_{-3}^{+5}+17$ from sum rules using the strange spectral function in $\tau$ decay. The computations were done to order $\alpha_{s}^{3}$, with an estimate of the $\alpha_{s}^{4}$ terms. We have converted the result to $\mu=2 \mathrm{GeV}$.
28 GAMIZ 05 determines $\bar{m}_{s}(2 \mathrm{GeV})$ from sum rules using the strange spectral function in $\tau$ decay. The computations were done to order $\alpha_{s}^{2}$, with an estimate of the $\alpha_{s}^{3}$ terms.
29 GORBUNOV 05 use hadronic tau decays to N3LO, including power corrections.
30 NARISON 05 determines $\bar{m}_{s}(2 \mathrm{GeV})$ from sum rules using the strange spectral function in $\tau$ decay. The computations were done to order $\alpha_{s}^{3}$.
${ }^{31}$ AUBIN 04 perform three flavor dynamical lattice calculation of pseudoscalar meson masses, with one-loop perturbative renormalization constant.
32 AOKI 03 uses quenched lattice simulation of the meson and baryon masses with degenerate light quarks. The extrapolations are done using quenched chiral perturbation theory. Determines $\mathrm{m}_{s}=113.8 \pm 2.3_{-2.9}^{+5.8}$ using $K$ mass as input and $\mathrm{m}_{s}=142.3 \pm 5.8_{-}^{+22}$ using $\phi$ mass as input. We have performed a weighted average of these values.
33 AOKI 03B uses lattice simulation of the meson and baryon masses with two dynamical light quarks. Simulations are performed using the $\mathcal{O}(a)$ improved Wilson action.
34 BECIREVIC 03 perform quenched lattice computation using the vector and axial Ward identities. Uses $\mathcal{O}(a)$ improved Wilson action and nonperturbative renormalization. They also quote $\bar{m} / \mathrm{m}_{s}=24.3 \pm 0.2 \pm 0.6$.
${ }^{35}$ CHIU 03 determines quark masses from the pion and kaon masses using a lattice simulation with a chiral fermion action in quenched approximation.
${ }^{36}$ GAMIZ 03 determines $m_{s}$ from SU(3) breaking in the $\tau$ hadronic width. The value of $V_{u s}$ is chosen to satisfy CKM unitarity.
37 GAMIZ 03 determines $m_{s}$ from SU(3) breaking in the $\tau$ hadronic width. The value of $V_{u s}$ is taken from the PDG.


## OTHER LIGHT QUARK MASS RATIOS

## $\boldsymbol{m}_{\boldsymbol{s}} / \boldsymbol{m}_{\boldsymbol{d}}$ MASS RATIO

$\qquad$

## 17-22 OUR EVALUATION

20.0
$18.9 \pm 0.8$
21
18
18 to 23

| DOCUMENT ID |  | TECN | COMMENT |
| :---: | :---: | :---: | :---: |
| 1 GAO | 97 | THEO |  |
| 2 LEUTWYLER | 96 | THEO | Compilation |
| 3 DONOGHUE | 92 | THEO |  |
| 4 GERARD | 90 | THEO |  |
| 5 LEUTWYLER | 90B | THEO |  |

${ }^{1}$ GAO 97 uses electromagnetic mass splittings of light mesons.
${ }^{2}$ LEUTWYLER 96 uses a combined fit to $\eta \rightarrow 3 \pi$ and $\psi^{\prime} \rightarrow J / \psi(\pi, \eta)$ decay rates, and the electromagnetic mass differences of the $\pi$ and $K$.
3 DONOGHUE 92 result is from a combined analysis of meson masses, $\eta \rightarrow 3 \pi$ using second-order chiral perturbation theory including nonanalytic terms, and $(\psi(2 S) \rightarrow$ $J / \psi(1 S) \pi) /(\psi(2 S) \rightarrow J / \psi(1 S) \eta)$.
${ }^{4}$ GERARD 90 uses large $N$ and $\eta-\eta^{\prime}$ mixing.
${ }^{5}$ LEUTWYLER 90B determines quark mass ratios using second-order chiral perturbation theory for the meson and baryon masses, including nonanalytic corrections. Also uses Weinberg sum rules to determine $L_{7}$.

## $\boldsymbol{m}_{\boldsymbol{s}} / \overline{\boldsymbol{m}}$ MASS RATIO

$\bar{m} \equiv\left(m_{u}+m_{d}\right) / 2$
VALUE DOCUMENTID TECN
$27.3{ }_{-1.3}^{+0.7}$ OUR EVALUATION See the ideogram below.

| $27.35 \pm 0.05{ }_{-0.07}^{+0.10}$ | 1 BAZAVOV | 14 A | LATT |
| :--- | :--- | :--- | :--- |
| $26.66 \pm 0.32$ | $2^{2}$ CARRASCO | 14 | LATT |
| $27.36 \pm 0.54$ | $3^{4}$ ARTHUR | 13 | LATT |
| $27.53 \pm 0.20 \pm 0.08$ | 4 DURR | 11 | LATT |

-     - We do not use the following data for averages, fits, limits, etc. - -

| 26.8 | $\pm 1.4$ | ${ }^{5}$ AOKI | 11 A | LATT |
| :--- | :--- | :--- | :--- | :--- |
| 27.3 | $\pm 0.9$ | ${ }^{\text {BLOSSIER }}$ | 10 | LATT |
| 28.8 | $\pm 1.65$ | 7 ALLTON | 08 | LATT |
| 27.3 | $\pm 0.3 \pm 1.2$ | 8 BLOSSIER | 08 | LATT |
| 23.5 | $\pm 1.5$ | 9 | OLLER | 07 A | THEO


$1^{1}$ BAZAVOV 14 A is a lattice computation using 4 dynamical flavors of HISQ fermions.
${ }^{2}$ CARRASCO 14 is a lattice QCD computation of light quark masses using $2+1+1$ dynamical quarks, with $m_{u}=m_{d} \neq m_{s} \neq m_{c}$. The $u$ and $d$ quark masses are obtained separately by using the $K$ meson mass splittings and lattice results for the electromagnetic contributions.
${ }^{3}$ ARTHUR 13 is a lattice computation using $2+1$ dynamical domain wall fermions.
${ }^{4}$ DURR 11 determine quark mass from a lattice computation of the meson spectrum using $N_{f}=2+1$ dynamical flavors. The lattice simulations were done at the physical quark mass, so that extrapolation in the quark mass was not needed.
${ }^{5}$ AOKI 11A determine quark masses from a lattice computation of the hadron spectrum using $N_{f}=2+1$ dynamical flavors of domain wall fermions.
${ }^{6}$ BLOSSIER 10 determines quark masses from a computation of the hadron spectrum using $N_{f}=2$ dynamical twisted-mass Wilson fermions.
${ }^{7}$ ALLTON 08 use a lattice computation of the $\pi, K$, and $\Omega$ masses with $2+1$ dynamical flavors of domain wall quarks, and non-perturbative renormalization.
${ }^{8}$ BLOSSIER 08 use a lattice computation of pseudoscalar meson masses and decay constants with 2 dynamical flavors and non-perturbative renormalization.
${ }^{9}$ OLLER 07A use unitarized chiral perturbation theory to order $p^{4}$.
${ }^{10}$ Three flavor dynamical lattice calculation of pseudoscalar meson masses.

## Q MASS RATIO

$$
\begin{aligned}
& Q \equiv \sqrt{\left(m_{s}^{2}-\bar{m}^{2}\right) /\left(m^{2} d-m_{u}^{2}\right)} ; \quad \bar{m} \equiv\left(m_{u}+m_{d}\right) / 2 \\
& \text { VALUE } \\
& \text { DOCUMENT ID } \\
& \hline
\end{aligned}
$$

-     - We do not use the following data for averages, fits, limits, etc. - • -

| $22.1 \pm 0.7$ | 1 COLANGELO | 18 | THEO |
| :--- | :--- | :--- | :--- |
| $22.0 \pm 0.7$ | 2 COLANGELO | 17 | THEO |
| $21.6 \pm 1.1$ | 3 GUO | 17 | THEO |
| $23.4 \pm 0.4 \pm 0.5$ | 4 FODOR | 16 | LATT |
| $21.4 \pm 0.4$ | 5 GUO | 15 F | THEO |
| $22.8 \pm 0.4$ | 6 MARTEMYA... 05 | THEO |  |
| $22.7 \pm 0.8$ | 7 ANISOVICH | 96 | THEO |

${ }^{1}$ COLANGELO 18 obtain $Q$ from a dispersive analysis of $\eta \rightarrow 3 \pi$ decay.
${ }^{2}$ COLANGELO 17 obtain $Q$ from a dispersive analysis of KLOE collaboration data on $\eta \rightarrow \pi^{+} \pi^{-} \pi^{0}$ decays and chiral perturbation theory input.
${ }^{3}$ GUO 17 determine $Q$ from a dispersive model fit to KLOE and WASA-at-COSY data on $\eta \rightarrow \pi^{+} \pi^{-} \pi^{0}$ decay and matching to chiral perturbation theory.
${ }^{4}$ FODOR 16 is a lattice simulation with $N_{f}=2+1$ dynamical flavors and includes partially quenched QED effects.
${ }^{5}$ GUO 15F determine $Q$ from a Khuri-Treiman analysis of $\eta \rightarrow 3 \pi$ decays.
${ }^{6}$ MARTEMYANOV 05 determine $Q$ from $\eta \rightarrow 3 \pi$ decay.
${ }^{7}$ ANISOVICH 96 find $Q$ from $\eta \rightarrow \pi^{+} \pi^{-} \pi^{0}$ decay using dispersion relations and chiral perturbation theory.

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