

# SEARCHES

## not in other sections

### Magnetic Monopole Searches

The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$< 1.4 \times 10^{-16} \text{ cm}^{-2} \text{sr}^{-1} \text{s}^{-1} \quad \text{for } 1.1 \times 10^{-4} < \beta < 1$$

### Supersymmetric Particle Searches

All supersymmetric mass bounds here are model dependent.

The limits assume:

1)  $\tilde{\chi}_1^0$  is the lightest supersymmetric particle; 2)  $R$ -parity is conserved, unless stated otherwise;

See the Particle Listings for a Note giving details of supersymmetry.

$\tilde{\chi}_i^0$  — neutralinos (mixtures of  $\tilde{\gamma}$ ,  $\tilde{Z}^0$ , and  $\tilde{H}_i^0$ )

Mass  $m_{\tilde{\chi}_1^0} > 0 \text{ GeV}$ , CL = 95%

[general MSSM, non-universal gaugino masses]

Mass  $m_{\tilde{\chi}_1^0} > 46 \text{ GeV}$ , CL = 95%

[all  $\tan\beta$ , all  $m_0$ , all  $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ]

Mass  $m_{\tilde{\chi}_2^0} > 62.4 \text{ GeV}$ , CL = 95%

[ $1 < \tan\beta < 40$ , all  $m_0$ , all  $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ]

Mass  $m_{\tilde{\chi}_3^0} > 99.9 \text{ GeV}$ , CL = 95%

[ $1 < \tan\beta < 40$ , all  $m_0$ , all  $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ]

Mass  $m_{\tilde{\chi}_4^0} > 116 \text{ GeV}$ , CL = 95%

[ $1 < \tan\beta < 40$ , all  $m_0$ , all  $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ]

Mass  $m_{\tilde{\chi}}$  none 200–670 GeV, CL = 95% [R-Parity Violating]

[wino production,  $\tilde{\chi} \rightarrow b + \ell/\nu + t/b$  via  $\lambda'_{i33}$  coupling]

$\tilde{\chi}_i^\pm$  — charginos (mixtures of  $\tilde{W}^\pm$  and  $\tilde{H}_i^\pm$ )

Mass  $m_{\tilde{\chi}_1^\pm} > 94 \text{ GeV}$ , CL = 95%

[ $\tan\beta < 40$ ,  $m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} > 3 \text{ GeV}$ , all  $m_0$ ]

Mass  $m_{\tilde{\chi}_1^\pm} > 1000$  GeV, CL = 95%

$[2\ell + \cancel{E}_T, \text{Tchi1chi1C}, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m_{\tilde{\chi}_1^\pm} > 1600$  GeV, CL = 95% [R-Parity Violating]

$[\text{Tchi1n2l}, \tilde{\chi}_1^0 \rightarrow \ell^\pm \ell^\mp \nu, \lambda_{12k} \neq 0, m_{\tilde{\chi}_1^0} = 1200 \text{ GeV}]$

$\tilde{\chi}^\pm$  — long-lived chargino

Mass  $m_{\tilde{\chi}^\pm} > 1050$  GeV, CL = 95%

$[\tilde{\chi}^\pm \rightarrow \tilde{\chi}_1^0 \pi^\pm, \text{wino LSP, stable}]$

Mass  $m_{\tilde{\chi}^\pm} > 1050$  GeV, CL = 95%

$[\tilde{\chi}^\pm \rightarrow \tilde{\chi}_1^0 \pi^\pm, \text{wino LSP}, \tau = 20 \text{ ns}]$

$\tilde{\nu}$  — sneutrino

Mass  $m > 41$  GeV, CL = 95% [model independent]

Mass  $m > 94$  GeV, CL = 95%

$[\text{CMSSM}, 1 \leq \tan\beta \leq 40, m_{\tilde{e}_R} - m_{\tilde{\chi}_1^0} > 10 \text{ GeV}]$

Mass  $m > 4200$  GeV, CL = 95% [R-Parity Violating]

$[1e + 1\mu, \nu_\tau \rightarrow e\mu, \lambda = \lambda' = 0.1]$

$\tilde{e}$  — scalar electron (selectron)

Mass  $m > 107$  GeV, CL = 95% [all  $m_{\tilde{e}_L} - m_{\tilde{\chi}_1^0}$ ]

Mass  $m > 700$  GeV, CL = 95%

$[2\ell + \cancel{E}_T, m_{\tilde{\ell}_R} = m_{\tilde{\ell}_L} \text{ and } \tilde{\ell} = \tilde{e}, \tilde{\mu}, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 250$  GeV, CL = 95%

$[\ell^\pm \ell^\mp + \cancel{E}_T, \tilde{e}_R, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 410$  GeV, CL = 95% [R-Parity Violating]

$[\geq 4\ell^\pm, \tilde{\ell} \rightarrow l \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \ell^\pm \ell^\mp \nu]$

Mass  $m > 1200$  GeV, CL = 95% [R-Parity Violating]

$[\geq 4\ell, \lambda_{12k} \neq 0, m_{\tilde{\chi}_1^0} = 900 \text{ GeV (m-degenerate } \tilde{\ell}_L, \tilde{\nu})]$

$\tilde{\mu}$  — scalar muon (smuon)

Mass  $m > 700$  GeV, CL = 95%

$[2\ell + \cancel{E}_T, m_{\tilde{\ell}_R} = m_{\tilde{\ell}_L} \text{ and } \tilde{\ell} = \tilde{e}, \tilde{\mu}, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 210$ , CL = 95%

$[\ell^\pm \ell^\mp + \cancel{E}_T, \tilde{\mu}_R, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 94$  GeV, CL = 95%

$[\text{CMSSM}, 1 \leq \tan\beta \leq 40, m_{\tilde{\mu}_R} - m_{\tilde{\chi}_1^0} > 10 \text{ GeV}]$

Mass  $m > 410$  GeV, CL = 95% [R-Parity Violating]

$[\geq 4\ell^\pm, \tilde{\ell} \rightarrow l \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \ell^\pm \ell^\mp \nu]$

Mass  $m > 1200$  GeV, CL = 95% [R-Parity Violating]

$[\geq 4\ell, \lambda_{12k} \neq 0, m_{\tilde{\chi}_1^0} = 900 \text{ GeV (m-degenerate } \tilde{\ell}_L, \tilde{\nu})]$

$\tilde{\tau}$  — scalar tau (stau)

Mass  $m > 81.9$  GeV, CL = 95%

$[m_{\tilde{\tau}_R} - m_{\tilde{\chi}_1^0} > 15 \text{ GeV, all } \theta_\tau, \text{B}(\tilde{\tau} \rightarrow \tau \tilde{\chi}_1^0) = 100\%]$

Mass  $m > 400$  GeV, CL = 95%

$[2 \text{ hadronic } \tau + \cancel{E}_T, \tilde{\tau}_{R,L} \rightarrow \tau \tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 1 \text{ GeV}]$

Mass  $m > 90$  GeV, CL = 95% [R-Parity Violating]

$[\tilde{\tau}_R, \text{indirect}, \Delta m > 5 \text{ GeV}]$

Mass  $m > 1200$ , CL = 95% [R-Parity Violating]

$[\geq 4\ell, \lambda_{12k} \neq 0, m_{\tilde{\chi}_1^0} = 900 \text{ GeV (m-degenerate } \tilde{\ell}_L, \tilde{\nu})]$

Mass  $m > 286$  GeV, CL = 95% [long-lived  $\tilde{\tau}$ ]

$\tilde{q}$  — squarks of the first two quark generations

Mass  $m > 1.220 \times 10^3$  GeV, CL = 95%

$[\text{jets} + \cancel{E}_T, \text{Tsqk1, 1 non-degenerate } \tilde{q}, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 1.600 \times 10^3$  GeV, CL = 95% [R-Parity Violating]

$[\tilde{q} \rightarrow q \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \ell \ell \nu, \lambda_{121}, \lambda_{122} \neq 0, m_{\tilde{g}} = 2400 \text{ GeV}]$

$\tilde{q}$  — long-lived squark

Mass  $m > 1340$ , CL = 95% [ $\tilde{t}$  R-hadrons]

Mass  $m > 1250$ , CL = 95% [ $\tilde{b}$  R-hadrons]

$\tilde{b}$  — scalar bottom (sbottom)

Mass  $m > 1.270 \times 10^3$  GeV, CL = 95%

$[b\text{-jets} + \cancel{E}_T, \text{Tsbol1, } m_{\tilde{\chi}_1^0} = 0 \text{ GeV}]$

Mass  $m > 307$  GeV, CL = 95% [R-Parity Violating]

$[\tilde{b} \rightarrow t d \text{ or } t s, \lambda_{332}'' \text{ or } \lambda_{331}'' \text{ coupling}]$

$\tilde{t}$  — scalar top (stop)

Mass  $m > 1.310 \times 10^3$  GeV, CL = 95%

$[\text{jets} + \cancel{E}_T, \text{Tstop1, } m_{\tilde{\chi}_1^0} < 300 \text{ GeV}]$

Mass  $m > 1100$  GeV, CL = 95% [R-Parity Violating]

$[\tilde{t} \rightarrow b e, \text{Tstop2RPV, prompt}]$

Mass  $m > 460$  GeV, CL = 95%

[R-Parity Violating, long-lived  $\tilde{t}$ ,  $\tilde{t} \rightarrow d \bar{\ell}$ ,  $0.01 \text{ cm} < c\tau < 1000 \text{ cm}$ ]

$\tilde{g}$  — gluino

Mass  $m > 2.300 \times 10^3$  GeV, CL = 95%

$[\text{jets} + \cancel{E}_T, \text{Tglu1A, } m_{\tilde{\chi}_1^0} < 200 \text{ GeV}]$

Mass  $m > 2.260 \times 10^3$  GeV, CL = 95% [R-Parity Violating]

$[\geq 4\ell, \lambda_{12k} \neq 0, m_{\tilde{\chi}_1^0} > 1000 \text{ GeV}]$

## Technicolor

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full *Review* (the data listings).

## Quark and Lepton Compositeness, Searches for

### Scale Limits $\Lambda$ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \bar{\psi}_L \gamma_\mu \psi_L \bar{\psi}_L \gamma^\mu \psi_L$$

(with  $g^2/4\pi$  set equal to 1), then we define  $\Lambda \equiv \Lambda_{LL}^\pm$ . For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

$\Lambda_{LL}^+(eeee)$	$> 8.3 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(eeee)$	$> 10.3 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(ee\mu\mu)$	$> 8.5 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(ee\mu\mu)$	$> 9.5 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(ee\tau\tau)$	$> 7.9 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(ee\tau\tau)$	$> 7.2 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(\ell\ell\ell\ell)$	$> 9.1 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(\ell\ell\ell\ell)$	$> 10.3 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(eeqq)$	$> 24 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(eeqq)$	$> 37 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(eeuu)$	$> 23.3 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(eeuu)$	$> 12.5 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(eedd)$	$> 11.1 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(eedd)$	$> 26.4 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(eccc)$	$> 9.4 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^-(eccc)$	$> 5.6 \text{ TeV, CL} = 95\%$
$\Lambda_{LL}^+(eebb)$	$> 9.4 \text{ TeV, CL} = 95\%$

$$\begin{aligned}
 \Lambda_{LL}^-(e e b b) &> 10.2 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^+(\mu \mu q q) &> 23.3 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^-(\mu \mu q q) &> 40.0 \text{ TeV, CL} = 95\% \\
 \Lambda(\ell \nu \ell \nu) &> 3.10 \text{ TeV, CL} = 90\% \\
 \Lambda(e \nu q q) &> 2.81 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^+(q q q q) &> 13.1 \text{ none } 17.4\text{--}29.5 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^-(q q q q) &> 21.8 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^+(\nu \nu q q) &> 5.0 \text{ TeV, CL} = 95\% \\
 \Lambda_{LL}^-(\nu \nu q q) &> 5.4 \text{ TeV, CL} = 95\%
 \end{aligned}$$

### Excited Leptons

The limits from  $\ell^{*+} \ell^{*-}$  do not depend on  $\lambda$  (where  $\lambda$  is the  $\ell \ell^*$  transition coupling). The  $\lambda$ -dependent limits assume chiral coupling.

$e^{*\pm}$  — excited electron

$$\begin{aligned}
 \text{Mass } m &> 103.2 \text{ GeV, CL} = 95\% \quad (\text{from } e^* e^*) \\
 \text{Mass } m &> 5.600 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } e e^*) \\
 \text{Mass } m &> 356 \text{ GeV, CL} = 95\% \quad (\text{if } \lambda_\gamma = 1)
 \end{aligned}$$

$\mu^{*\pm}$  — excited muon

$$\begin{aligned}
 \text{Mass } m &> 103.2 \text{ GeV, CL} = 95\% \quad (\text{from } \mu^* \mu^*) \\
 \text{Mass } m &> 5.700 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } \mu \mu^*)
 \end{aligned}$$

$\tau^{*\pm}$  — excited tau

$$\begin{aligned}
 \text{Mass } m &> 103.2 \text{ GeV, CL} = 95\% \quad (\text{from } \tau^* \tau^*) \\
 \text{Mass } m &> 4.600 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } \tau \tau^*)
 \end{aligned}$$

$\nu^*$  — excited neutrino

$$\begin{aligned}
 \text{Mass } m &> 1.600 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } \nu^* \nu^*) \\
 \text{Mass } m &> 213 \text{ GeV, CL} = 95\% \quad (\text{from } \nu^* X)
 \end{aligned}$$

$q^*$  — excited quark

$$\begin{aligned}
 \text{Mass } m &> 338 \text{ GeV, CL} = 95\% \quad (\text{from } q^* q^*) \\
 \text{Mass } m &> 6700 \text{ GeV, CL} = 95\% \quad (\text{from } q^* X)
 \end{aligned}$$

### Color Sextet and Octet Particles

Color Sextet Quarks ( $q_6$ )

$$\text{Mass } m > 84 \text{ GeV, CL} = 95\% \quad (\text{Stable } q_6)$$

Color Octet Charged Leptons ( $\ell_8$ )

$$\text{Mass } m > 86 \text{ GeV, CL} = 95\% \quad (\text{Stable } \ell_8)$$

Color Octet Neutrinos ( $\nu_8$ )

Mass  $m > 110$  GeV, CL = 90% ( $\nu_8 \rightarrow \nu g$ )

## Extra Dimensions

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

### Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

(direct tests of Newton's law)

$R < 3.8 \mu\text{m}$ , CL = 95% ( $pp \rightarrow j G$ )

$R < 0.16\text{--}916$  nm (astrophysics; limits depend on technique and assumptions)

### Constraints on the fundamental gravity scale

$M_{TT} > 9.02$  TeV, CL = 95% ( $pp \rightarrow$  dijet, angular distribution)

$M_c > 4.16$  TeV, CL = 95% ( $pp \rightarrow \ell \bar{\ell}$ )

### Constraints on the Kaluza-Klein graviton in warped extra dimensions

$M_G > 4.78$  TeV, CL = 95% ( $pp \rightarrow e^+ e^-, \mu^+ \mu^-$ )

### Constraints on the Kaluza-Klein gluon in warped extra dimensions

$M_{g_{KK}} > 3.8$  TeV, CL = 95% ( $g_{KK} \rightarrow t \bar{t}$ )

## WIMP and Dark Matter Searches

No confirmed evidence found for galactic

WIMPs from the GeV to the TeV mass scales and down to  $1 \times 10^{-10}$

pb spin independent cross section at  $M = 100$  GeV.