

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+) \text{ Status: } ****$$

The parity has not actually been measured, but + is of course expected.

Ξ⁰ MASS

The fit uses the Ξ⁰, Ξ⁻, and Ξ⁺ mass and mass difference measurements.

<u>VALUE (MeV)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
1314.83±0.20 OUR FIT				
1314.82±0.20 OUR AVERAGE				
1314.82±0.06±0.20	3120	FANTI	00 NA48	p Be, 450 GeV
1315.2 ±0.92	49	WILQUET	72 HLBC	
1313.4 ±1.8	1	PALMER	68 HBC	

$m_{\Xi^-} - m_{\Xi^0}$

The fit uses the Ξ⁰, Ξ⁻, and Ξ⁺ mass and mass difference measurements.

<u>VALUE (MeV)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
6.48±0.24 OUR FIT				
6.3 ±0.7 OUR AVERAGE				
6.9 ±2.2	29	LONDON	66 HBC	
6.1 ±0.9	88	PJERROU	65B HBC	
6.8 ±1.6	23	JAUNEAU	63 FBC	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
6.1 ±1.6	45	CARMONY	64B HBC	See PJERROU 65B

Ξ⁰ MEAN LIFE

<u>VALUE (10⁻¹⁰ s)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
2.90±0.09 OUR AVERAGE				
2.83±0.16	6300	¹ ZECH	77 SPEC	Neutral hyperon beam
2.88 ^{+0.21} _{-0.19}	652	BALTAY	74 HBC	1.75 GeV/c K ⁻ p
2.90 ^{+0.32} _{-0.27}	157	² MAYEUR	72 HLBC	2.1 GeV/c K ⁻
3.07 ^{+0.22} _{-0.20}	340	DAUBER	69 HBC	
3.0 ±0.5	80	PJERROU	65B HBC	
2.5 ^{+0.4} _{-0.3}	101	HUBBARD	64 HBC	
3.9 ^{+1.4} _{-0.8}	24	JAUNEAU	63 FBC	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
3.5 ^{+1.0} _{-0.8}	45	CARMONY	64B HBC	See PJERROU 65B

¹The ZECH 77 result is $\tau_{\Xi^0} = [2.77 - (\tau_{\Lambda} - 2.69)] \times 10^{-10}$ s, in which we use $\tau_{\Lambda} = 2.63 \times 10^{-10}$ s.

²The MAYEUR 72 value is modified by the erratum.

Ξ^0 MAGNETIC MOMENT

See the "Note on Baryon Magnetic Moments" in the Λ Listings.

VALUE (μ_N)	EVTS	DOCUMENT ID	TECN
-1.250 ± 0.014 OUR AVERAGE			
-1.253 ± 0.014	270k	COX	81 SPEC
-1.20 ± 0.06	42k	BUNCE	79 SPEC

Ξ^0 DECAY MODES

Mode	Fraction (Γ_i/Γ)	Confidence level
Γ_1 $\Lambda\pi^0$	(99.523 ± 0.013) %	
Γ_2 $\Lambda\gamma$	(1.17 ± 0.07) × 10 ⁻³	
Γ_3 $\Sigma^0\gamma$	(3.33 ± 0.10) × 10 ⁻³	
Γ_4 $\Sigma^+ e^- \bar{\nu}_e$	(2.7 ± 0.4) × 10 ⁻⁴	
Γ_5 $\Sigma^+ \mu^- \bar{\nu}_\mu$	(4.9 $\begin{smallmatrix} +2.1 \\ -1.6 \end{smallmatrix}$) × 10 ⁻⁶	

$\Delta S = \Delta Q$ (SQ) violating modes or $\Delta S = 2$ forbidden (S2) modes

Γ_6 $\Sigma^- e^+ \nu_e$	SQ	< 9	× 10 ⁻⁴	90%
Γ_7 $\Sigma^- \mu^+ \nu_\mu$	SQ	< 9	× 10 ⁻⁴	90%
Γ_8 $p\pi^-$	S2	< 8	× 10 ⁻⁶	90%
Γ_9 $p e^- \bar{\nu}_e$	S2	< 1.3	× 10 ⁻³	
Γ_{10} $p \mu^- \bar{\nu}_\mu$	S2	< 1.3	× 10 ⁻³	

CONSTRAINED FIT INFORMATION

An overall fit to 3 branching ratios uses 7 measurements and one constraint to determine 4 parameters. The overall fit has a $\chi^2 = 4.4$ for 4 degrees of freedom.

The following *off-diagonal* array elements are the correlation coefficients $\langle \delta x_i \delta x_j \rangle / (\delta x_i \delta x_j)$, in percent, from the fit to the branching fractions, $x_i \equiv \Gamma_i / \Gamma_{\text{total}}$. The fit constrains the x_i whose labels appear in this array to sum to one.

x_2	-54		
x_3	-78	0	
x_4	-30	0	0
	x_1	x_2	x_3

Ξ^0 BRANCHING RATIOS $\Gamma(\Lambda\gamma)/\Gamma(\Lambda\pi^0)$ Γ_2/Γ_1

<u>VALUE (units 10^{-3})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
1.17±0.07 OUR FIT				
1.17±0.07 OUR AVERAGE				
1.17±0.05±0.06	672	³ LAI	04A NA48	p Be, 450 GeV
1.91±0.34±0.19	31	⁴ FANTI	00 NA48	p Be, 450 GeV
1.06±0.12±0.11	116	JAMES	90 SPEC	FNAL hyperons

³ LAI 04A used our 2002 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Lambda\gamma)/\Gamma_{\text{total}} = (1.16 \pm 0.05 \pm 0.06) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

⁴ FANTI 00 used our 1998 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Lambda\gamma)/\Gamma_{\text{total}} = (1.90 \pm 0.34 \pm 0.19) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

 $\Gamma(\Sigma^0\gamma)/\Gamma(\Lambda\pi^0)$ Γ_3/Γ_1

<u>VALUE (units 10^{-3})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
3.35±0.10 OUR FIT				
3.35±0.10 OUR AVERAGE				
3.34±0.05±0.09	4045	ALAVI-HARATI01C	KTEV	p nucleus, 800 GeV
3.16±0.76±0.32	17	⁵ FANTI	00 NA48	p Be, 450 GeV
3.56±0.42±0.10	85	TEIGE	89 SPEC	FNAL hyperons

⁵ FANTI 00 used our 1998 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Sigma^0\gamma)/\Gamma_{\text{total}} = (3.14 \pm 0.76 \pm 0.32) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

 $\Gamma(\Sigma^+ e^- \bar{\nu}_e)/\Gamma_{\text{total}}$ Γ_4/Γ

<u>VALUE (units 10^{-4})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
2.7 ±0.4 OUR FIT				
2.71±0.22±0.31	176	AFFOLDER	99 KTEV	p nucleus 800 GeV

 $\Gamma(\Sigma^+ \mu^- \bar{\nu}_\mu)/\Gamma(\Sigma^+ e^- \bar{\nu}_e)$ Γ_5/Γ_4

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
0.018^{+0.007}_{-0.005} ±0.002	9	ABOUZAID	05 KTEV	p nucleus 800 GeV

 $\Gamma(\Sigma^+ \mu^- \bar{\nu}_\mu)/\Gamma(\Lambda\pi^0)$ Γ_5/Γ_1

<u>VALUE (units 10^{-3})</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<1.1	90	0	YEH	74 HBC	Effective denom.=2100
<1.5			DAUBER	69 HBC	
<7			HUBBARD	66 HBC	

 $\Gamma(\Sigma^- e^+ \nu_e)/\Gamma(\Lambda\pi^0)$ Γ_6/Γ_1

<u>VALUE (units 10^{-3})</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.9	90	0	YEH	74 HBC	Effective denom.=2500
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<1.5			DAUBER	69 HBC	
<6			HUBBARD	66 HBC	

$\Gamma(\Sigma^- \mu^+ \nu_\mu) / \Gamma(\Lambda \pi^0)$ Γ_7 / Γ_1

Test of $\Delta S = \Delta Q$ rule.

VALUE (units 10^{-3})	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<0.9	90	0	YEH	74 HBC	Effective denom.=2500
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<1.5			DAUBER	69 HBC	
<6			HUBBARD	66 HBC	

$\Gamma(p \pi^-) / \Gamma(\Lambda \pi^0)$ Γ_8 / Γ_1

$\Delta S=2$. Forbidden in first-order weak interaction.

VALUE (units 10^{-6})	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
< 8.2	90		WHITE	05 HYCP	p Cu, 800 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •					
< 36	90		GEWENIGER	75 SPEC	
<1800	90	0	YEH	74 HBC	Effective denom.=1300
< 900			DAUBER	69 HBC	
<5000			HUBBARD	66 HBC	

$\Gamma(p e^- \bar{\nu}_e) / \Gamma(\Lambda \pi^0)$ Γ_9 / Γ_1

$\Delta S=2$. Forbidden in first-order weak interaction.

VALUE (units 10^{-3})	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<1.3			DAUBER	69 HBC	
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<3.4	90	0	YEH	74 HBC	Effective denom.=670
<6			HUBBARD	66 HBC	

$\Gamma(p \mu^- \bar{\nu}_\mu) / \Gamma(\Lambda \pi^0)$ Γ_{10} / Γ_1

$\Delta S=2$. Forbidden in first-order weak interaction.

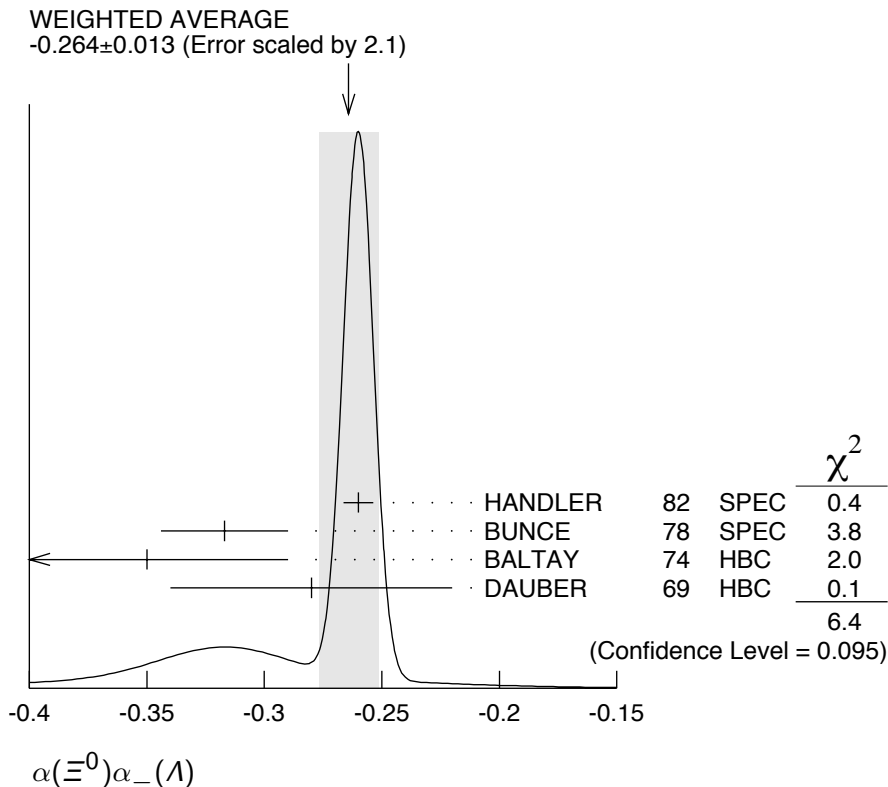
VALUE (units 10^{-3})	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<1.3			DAUBER	69 HBC	
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<3.5	90	0	YEH	74 HBC	Effective denom.=664
<6			HUBBARD	66 HBC	

Ξ^0 DECAY PARAMETERS

See the "Note on Baryon Decay Parameters" in the neutron Listings.

$\alpha(\Xi^0) \alpha_-(\Lambda)$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-0.264 ± 0.013 OUR AVERAGE				Error includes scale factor of 2.1. See the ideogram below.
$-0.260 \pm 0.004 \pm 0.005$	300k	HANDLER	82 SPEC	FNAL hyperons
-0.317 ± 0.027	6075	BUNCE	78 SPEC	FNAL hyperons
-0.35 ± 0.06	505	BALTAY	74 HBC	$K^- p$ 1.75 GeV/c
-0.28 ± 0.06	739	DAUBER	69 HBC	$K^- p$ 1.7-2.6 GeV/c



α FOR $\Xi^0 \rightarrow \Lambda\pi^0$

The above average, $\alpha(\Xi^0)\alpha_-(\Lambda) = -0.264 \pm 0.013$, where the error includes a scale factor of 2.1, divided by our current average $\alpha_-(\Lambda) = 0.642 \pm 0.013$, gives the following value for $\alpha(\Xi^0)$.

<u>VALUE</u>	<u>DOCUMENT ID</u>
-0.411 ± 0.022 OUR EVALUATION	Error includes scale factor of 2.1.

ϕ ANGLE FOR $\Xi^0 \rightarrow \Lambda\pi^0$

($\tan\phi = \beta/\gamma$)

<u>VALUE ($^\circ$)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
21 ± 12 OUR AVERAGE				
16 ± 17	652	BALTAY	74	HBC 1.75 GeV/c $K^- p$
38 ± 19	739	⁶ DAUBER	69	HBC
-8 ± 30	146	⁷ BERGE	66	HBC

⁶ DAUBER 69 uses $\alpha_\Lambda = 0.647 \pm 0.020$.

⁷ The errors have been multiplied by 1.2 due to approximations used for the Ξ polarization; see DAUBER 69 for a discussion.

RADIATIVE HYPERON DECAYS

Written September 2003 by J.D. Jackson (LBNL).

The weak radiative decays of spin-1/2 hyperons, $B_i \rightarrow B_f\gamma$, yield information about matrix elements (form factors) similar to that gained from weak hadronic decays. For a polarized

spin-1/2 hyperon decaying radiatively via a $\Delta Q = 0$, $\Delta S = 1$ transition, the angular distribution of the direction $\hat{\mathbf{p}}$ of the final spin-1/2 baryon in the hyperon rest frame is

$$\frac{dN}{d\Omega} = \frac{N}{4\pi} (1 + \alpha_\gamma \mathbf{P}_i \cdot \hat{\mathbf{p}}) . \quad (1)$$

Here \mathbf{P}_i is the polarization of the decaying hyperon, and α_γ is the asymmetry parameter. In terms of the form factors $F_1(q^2)$, $F_2(q^2)$, and $G(q^2)$ of the effective hadronic weak electromagnetic vertex,

$$F_1(q^2)\gamma_\lambda + iF_2(q^2)\sigma_{\lambda\mu}q^\mu + G(q^2)\gamma_\lambda\gamma_5 ,$$

α_γ is

$$\alpha_\gamma = \frac{2 \operatorname{Re}[G(0)F_M^*(0)]}{|G(0)|^2 + |F_M(0)|^2} , \quad (2)$$

where $F_M = (m_i - m_f)[F_2 - F_1/(m_i + m_f)]$. If the decaying hyperon is unpolarized, the decay baryon has a longitudinal polarization given by $P_f = -\alpha_\gamma$ [1].

The angular distribution for the weak hadronic decay, $B_i \rightarrow B_f\pi$, has the same form as Eq. (1), but of course with a different asymmetry parameter, α_π . Now, however, if the decaying hyperon is unpolarized, the decay baryon has a longitudinal polarization given by $P_f = +\alpha_\pi$ [2,3]. The difference of sign is because the spins of the pion and photon are different.

$\Xi^0 \rightarrow \Lambda\gamma$ decay—The radiative decay $\Xi^0 \rightarrow \Lambda\gamma$ of an unpolarized Ξ^0 uses the hadronic decay $\Lambda \rightarrow p\pi^-$ as the analyzer. As noted above, the longitudinal polarization of the Λ will be $P_\Lambda = -\alpha_{\Xi\Lambda\gamma}$. Let α_- be the $\Lambda \rightarrow p\pi^-$ asymmetry parameter and $\theta_{\Lambda p}$ be the angle, as seen in the Λ rest frame, between the Λ line of flight and the proton momentum. Then

the hadronic version of Eq. (1) applied to the $\Lambda \rightarrow p\pi^-$ decay gives

$$\frac{dN}{d\cos\theta_{\Lambda p}} = \frac{N}{2} (1 - \alpha_{\Xi\Lambda\gamma} \alpha_- \cos\theta_{\Lambda p}) \quad (3)$$

for the angular distribution of the proton in the Λ frame. The only published measurement of $\alpha_{\Xi\Lambda\gamma}$ [4] got the sign wrong, as explained in an erratum 12 years later [5]. The corrected result is $\alpha_{\Xi\Lambda\gamma} = -0.43 \pm 0.44$.

$\Xi^0 \rightarrow \Sigma^0\gamma$ decay—The asymmetry parameter here, $\alpha_{\Xi\Sigma\gamma}$, is measured by following the decay chain $\Xi^0 \rightarrow \Sigma^0\gamma$, $\Sigma^0 \rightarrow \Lambda\gamma$, $\Lambda \rightarrow p\pi^-$. Again, for an unpolarized Ξ^0 , the longitudinal polarization of the Σ^0 will be $P_\Sigma = -\alpha_{\Xi\Sigma\gamma}$. In the $\Sigma^0 \rightarrow \Lambda\gamma$ decay, a parity-conserving magnetic-dipole transition, the polarization of the Σ^0 is transferred to the Λ , as may be seen as follows. Let $\theta_{\Sigma\Lambda}$ be the angle seen in the Σ^0 rest frame between the Σ^0 line of flight and the Λ momentum. For Σ^0 helicity $+1/2$, the probability amplitudes for positive and negative spin states of the Σ^0 *along the Λ momentum* are $\cos(\theta_{\Sigma\Lambda}/2)$ and $\sin(\theta_{\Sigma\Lambda}/2)$. Then the amplitude for a negative helicity photon and a negative helicity Λ is $\cos(\theta_{\Sigma\Lambda}/2)$, while the amplitude for positive helicities for the photon and Λ is $\sin(\theta_{\Sigma\Lambda}/2)$. For Σ^0 helicity $-1/2$, the amplitudes are interchanged. If the Σ^0 has longitudinal polarization P_Σ , the probabilities for Λ helicities $\pm 1/2$ are therefore

$$p(\pm 1/2) = \frac{1}{2}(1 \mp P_\Sigma) \cos^2(\theta_{\Sigma\Lambda}/2) + \frac{1}{2}(1 \pm P_\Sigma) \sin^2(\theta_{\Sigma\Lambda}/2), \quad (4)$$

and the longitudinal polarization of the Λ is

$$P_\Lambda = -P_\Sigma \cos\theta_{\Sigma\Lambda} = +\alpha_{\Xi\Sigma\gamma} \cos\theta_{\Sigma\Lambda}. \quad (5)$$

Using Eq. (1) for the $\Lambda \rightarrow p\pi^-$ decay again, we get for the joint angular distribution of the $\Sigma^0 \rightarrow \Lambda\gamma$, $\Lambda \rightarrow p\pi^-$ chain,

$$\frac{d^2N}{d\cos\theta_{\Sigma\Lambda}d\cos\theta_{\Lambda p}} = \frac{N}{4} (1 + \alpha_{\Xi\Sigma\gamma} \cos\theta_{\Sigma\Lambda} \alpha_- \cos\theta_{\Lambda p}) . \quad (6)$$

The KTeV collaboration recently measured $\alpha_{\Xi\Sigma\gamma}$ to be -0.63 ± 0.09 [6]. The only other measurement has been withdrawn [7].

References

1. R.E. Behrends, Phys. Rev. **111**, 1691 (1958); see Eq. (7) or (8).
2. In ancient times, the signs of the asymmetry term in the angular distributions of radiative and hadronic decays of polarized hyperons were sometimes opposite. For roughly 40 years, however, the overwhelming convention has been to make them the same. The aim, not always achieved, is to remove ambiguities.
3. For the definition of α_π , see the note on “Baryon Decay Parameters,” in the Neutron Listings in this *Review*.
4. C. James *et al.*, Phys. Rev. Lett. **64**, 843 (1990).
5. C. James *et al.*, Phys. Rev. Lett. **89**, 169901 (2002) (erratum). The various sign conventions spelled out here are discussed.
6. A. Alavi-Harati *et al.*, Phys. Rev. Lett. **86**, 3239 (2001).
7. S. Teige *et al.*, Phys. Rev. Lett. **63**, 2717 (1989); erratum, Phys. Rev. Lett. **89**, 169902 (2002).

α FOR $\Xi^0 \rightarrow \Lambda\gamma$

See the note above on “Radiative Hyperon Decays.”

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
-0.73 ± 0.17 OUR AVERAGE				
$-0.78 \pm 0.18 \pm 0.06$	672	LAI	04A NA48	p Be, 450 GeV
-0.43 ± 0.44	87	⁸ JAMES	90 SPEC	FNAL hyperons

⁸The sign has been changed; see the erratum (JAMES 02, under JAMES 90).

α FOR $\Xi^0 \rightarrow \Sigma^0 \gamma$

See the note above on "Radiative Hyperon Decays."

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$-0.63 \pm 0.08 \pm 0.05$	4045	ALAVI-HARATI01C	KTEV	p nucleus, 800 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
$+0.20 \pm 0.32 \pm 0.05$	85	⁹ TEIGE	89 SPEC	FNAL hyperons

⁹This result has been withdrawn, due to an error. See the erratum (TEIGE 02, under TEIGE 89).

$g_1(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$+1.32^{+0.21}_{-0.17} \pm 0.05$	487	¹⁰ ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV

¹⁰ALAVI-HARATI 01I assumes here that the second-class current is zero and that the weak-magnetism term takes its exact SU(3) value.

$g_2(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$-1.7^{+2.1}_{-2.0} \pm 0.5$	487	¹¹ ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV

¹¹ALAVI-HARATI 01I thus assumes that $g_2 = 0$ in calculating g_1/f_1 , above.

$f_2(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$2.0 \pm 1.2 \pm 0.5$	487	ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV

Ξ^0 REFERENCES

ABOUZAID	05	PRL 95 081801	E. Abouzaid <i>et al.</i>	(FNAL KTeV Collab.)
WHITE	05	PRL 94 101804	C.G. White <i>et al.</i>	(FNAL HyperCP Collab.)
LAI	04A	PL B584 251	A. Lai <i>et al.</i>	(CERN NA48 Collab.)
JAMES	02	PRL 89 169901 (erratum)	C. James <i>et al.</i>	(MINN, MICH, WISC, RUTG)
TEIGE	02	PRL 89 169902 (erratum)	S. Teige <i>et al.</i>	(RUTG, MICH, MINN)
ALAVI-HARATI	01C	PRL 86 3239	A. Alavi-Harati <i>et al.</i>	(FNAL KTeV Collab.)
ALAVI-HARATI	01I	PRL 87 132001	A. Alavi-Harati <i>et al.</i>	(FNAL KTeV Collab.)
FANTI	00	EPJ C12 69	V. Fanti <i>et al.</i>	(CERN NA48 Collab.)
AFFOLDER	99	PRL 82 3751	A. Affolder <i>et al.</i>	(FNAL KTeV Collab.)
JAMES	90	PRL 64 843	C. James <i>et al.</i>	(MINN, MICH, WISC, RUTG)
Also		PRL 89 169901 (erratum)	C. James <i>et al.</i>	(MINN, MICH, WISC, RUTG)
TEIGE	89	PRL 63 2717	S. Teige <i>et al.</i>	(RUTG, MICH, MINN)
Also		PRL 89 169902 (erratum)	S. Teige <i>et al.</i>	(RUTG, MICH, MINN)
HANDLER	82	PR D25 639	R. Handler <i>et al.</i>	(WISC, MICH, MINN+)
COX	81	PRL 46 877	P.T. Cox <i>et al.</i>	(MICH, WISC, RUTG, MINN+)
BUNCE	79	PL 86B 386	G.R.M. Bunce <i>et al.</i>	(BNL, MICH, RUTG+)
BUNCE	78	PR D18 633	G.R.M. Bunce <i>et al.</i>	(WISC, MICH, RUTG)
ZECH	77	NP B124 413	G. Zech <i>et al.</i>	(SIEG, CERN, DORT, HEIDH)
GEWENIGER	75	PL 57B 193	C. Geweniger <i>et al.</i>	(CERN, HEIDH)
BALTAY	74	PR D9 49	C. Baltay <i>et al.</i>	(COLU, BING) J
YEH	74	PR D10 3545	N. Yeh <i>et al.</i>	(BING, COLU)
MAYEUR	72	NP B47 333	C. Mayeur <i>et al.</i>	(BRUX, CERN, TUFTS, LOUC)
Also		NP B53 268 (erratum)	C. Mayeur	
WILQUET	72	PL 42B 372	G. Wilquet <i>et al.</i>	(BRUX, CERN, TUFTS+)
DAUBER	69	PR 179 1262	P.M. Dauber <i>et al.</i>	(LRL)
PALMER	68	PL 26B 323	R.B. Palmer <i>et al.</i>	(BNL, SYRA)
BERGE	66	PR 147 945	J.P. Berge <i>et al.</i>	(LRL)
HUBBARD	66	Thesis UCRL 11510	J.R. Hubbard	(LRL)
LONDON	66	PR 143 1034	G.W. London <i>et al.</i>	(BNL, SYRA)

PJERROU	65B	PRL 14 275	G.M. Pjerrou <i>et al.</i>	(UCLA)
Also		Thesis	G.M. Pjerrou	(UCLA)
CARMONY	64B	PRL 12 482	D.D. Carmony <i>et al.</i>	(UCLA)
HUBBARD	64	PR 135B 183	J.R. Hubbard <i>et al.</i>	(LRL)
JAUNEAU	63	PL 4 49	L. Jauneau <i>et al.</i>	(EPOL, CERN, LOUC+)
Also		Siena Conf. 1 1	L. Jauneau <i>et al.</i>	(EPOL, CERN, LOUC+)
