

Heavy Bosons Other Than Higgs Bosons, Searches for

We list here various limits on charged and neutral heavy vector bosons (other than W 's and Z 's), heavy scalar bosons (other than Higgs bosons), vector or scalar leptoquarks, and axigluons.

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MASS LIMITS for W' (Heavy Charged Vector Boson Other Than W) in Hadron Collider Experiments

Couplings of W' to quarks and leptons are taken to be identical with those of W . The following limits are obtained from $p\bar{p} \rightarrow W'X$ with W' decaying to the mode indicated in the comments. New decay channels (e.g., $W' \rightarrow WZ$) are assumed to be suppressed. UA1 and UA2 experiments assume that the $t\bar{b}$ channel is not open.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>800	95	ABAZOV	04C D0	$W' \rightarrow q\bar{q}$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
none 200–610	95	¹ ABAZOV	06N D0	$W' \rightarrow tb$
225–536	95	² ACOSTA	03B CDF	$W' \rightarrow tb$
none 200–480	95	³ AFFOLDER	02C CDF	$W' \rightarrow WZ$
>786	95	⁴ AFFOLDER	01I CDF	$W' \rightarrow e\nu, \mu\nu$
>660	95	⁵ ABE	00 CDF	$W' \rightarrow \mu\nu$
none 300–420	95	⁶ ABE	97G CDF	$W' \rightarrow q\bar{q}$
>720	95	⁷ ABACHI	96C D0	$W' \rightarrow e\nu$
>610	95	⁸ ABACHI	95E D0	$W' \rightarrow e\nu, \tau\nu$
>652	95	⁹ ABE	95M CDF	$W' \rightarrow e\nu$
>251	90	¹⁰ ALITTI	93 UA2	$W' \rightarrow q\bar{q}$
none 260–600	95	¹¹ RIZZO	93 RVUE	$W' \rightarrow q\bar{q}$
>220	90	¹² ALBAJAR	89 UA1	$W' \rightarrow e\nu$
>209	90	¹³ ANSARI	87D UA2	$W' \rightarrow e\nu$

¹ The ABAZOV 06N quoted limit is for W' with SM-like coupling which interferes with the SM W boson. For W' with right-handed coupling, $M_{W'}$ between 200 and 630 (670) GeV is excluded for $M_{\nu_R} \ll M_{W'} (M_{\nu_R} > M_{W'})$.

² The ACOSTA 03B quoted limit is for $M_{W'} \gg M_{\nu_R}$. For $M_{W'} < M_{\nu_R}$, $M_{W'}$ between 225 and 566 GeV is excluded.

³ The quoted limit is obtained assuming $W' WZ$ coupling strength is the same as the ordinary WWZ coupling strength in the Standard Model. See their Fig. 2 for the limits on the production cross sections as a function of the W' width.

⁴ AFFOLDER 01I combine a new bound on $W' \rightarrow e\nu$ of 754 GeV with the bound of ABE 00 on $W' \rightarrow \mu\nu$ to obtain quoted bound.

⁵ ABE 00 assume that the neutrino from W' decay is stable and has a mass significantly less than $m_{W'}$.

⁶ ABE 97G search for new particle decaying to dijets.

⁷ For bounds on W_R with nonzero right-handed mass, see Fig. 5 from ABACHI 96C.

⁸ ABACHI 95E assume that the decay $W' \rightarrow WZ$ is suppressed and that the neutrino from W' decay is stable and has a mass significantly less $m_{W'}$.

⁹ ABE 95M assume that the decay $W' \rightarrow WZ$ is suppressed and the (right-handed) neutrino is light, noninteracting, and stable. If $m_{\nu} = 60$ GeV, for example, the effect on the mass limit is negligible.

¹⁰ ALITTI 93 search for resonances in the two-jet invariant mass. The limit assumes $\Gamma(W')/m_{W'} = \Gamma(W)/m_W$ and $B(W' \rightarrow jj) = 2/3$. This corresponds to W_R with $m_{\nu_R} > m_{W_R}$ (no leptonic decay) and $W_R \rightarrow t\bar{b}$ allowed. See their Fig. 4 for limits in the $m_{W'} - B(q\bar{q})$ plane.

¹¹ RIZZO 93 analyses CDF limit on possible two-jet resonances. The limit is sensitive to the inclusion of the assumed K factor.

¹² ALBAJAR 89 cross section limit at 630 GeV is $\sigma(W') B(e\nu) < 4.1$ pb (90% CL).

¹³ See Fig. 5 of ANSARI 87D for the excluded region in the $m_{W'} - [g_{W'q}]^2 B(W' \rightarrow e\bar{\nu})$ plane. Note that the quantity $(g_{W'q})^2 B(W' \rightarrow e\bar{\nu})$ is normalized to unity for the standard W couplings.

W_R (Right-Handed W Boson) MASS LIMITS

Assuming a light right-handed neutrino, except for BEALL 82, LANGACKER 89B, and COLANGELO 91. $g_R = g_L$ assumed. [Limits in the section MASS LIMITS for W' below are also valid for W_R if $m_{\nu_R} \ll m_{W_R}$.] Some limits assume manifest left-right symmetry, i.e., the equality of left- and right Cabibbo-Kobayashi-Maskawa matrices. For a comprehensive review, see LANGACKER 89B. Limits on the W_L - W_R mixing angle ζ are found in the next section. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
> 715	90	14 CZAKON	99	RVUE Electroweak
• • • We do not use the following data for averages, fits, limits, etc. • • •				
[> 3300]	95	15 CYBURT	05	COSM Nucleosynthesis; light ν_R
> 310	90	16 THOMAS	01	CNTR β^+ decay
> 137	95	17 ACKERSTAFF	99D	OPAL τ decay
> 1400	68	18 BARENBOIM	98	RVUE Electroweak, Z - Z' mixing
> 549	68	19 BARENBOIM	97	RVUE μ decay
> 220	95	20 STAHL	97	RVUE τ decay
> 220	90	21 ALLET	96	CNTR β^+ decay
> 281	90	22 KUZNETSOV	95	CNTR Polarized neutron decay
> 282	90	23 KUZNETSOV	94B	CNTR Polarized neutron decay
> 439	90	24 BHATTACH...	93	RVUE Z - Z' mixing
> 250	90	25 SEVERIJNS	93	CNTR β^+ decay
		26 IMAZATO	92	CNTR K^+ decay
> 475	90	27 POLAK	92B	RVUE μ decay
> 240	90	28 AQUINO	91	RVUE Neutron decay
> 496	90	28 AQUINO	91	RVUE Neutron and muon decay
> 700		29 COLANGELO	91	$m_{K_L^0} - m_{K_S^0}$
> 477	90	30 POLAK	91	RVUE μ decay
[none 540–23000]		31 BARBIERI	89B	ASTR SN 1987A; light ν_R
> 300	90	32 LANGACKER	89B	RVUE General
> 160	90	33 BALKE	88	CNTR $\mu \rightarrow e\nu\bar{\nu}$
> 406	90	34 JODIDIO	86	ELEC Any ζ
> 482	90	34 JODIDIO	86	ELEC $\zeta = 0$
> 800		MOHAPATRA	86	SU(2) _L × SU(2) _R × U(1)

> 400	95	35 STOKER	85	ELEC	Any ζ
> 475	95	35 STOKER	85	ELEC	$\zeta < 0.041$
		36 BERGSMA	83	CHRM	$\nu_\mu e \rightarrow \mu \nu_e$
> 380	90	37 CARR	83	ELEC	μ^+ decay
>1600		38 BEALL	82	THEO	$m_{K_L^0} - m_{K_S^0}$

¹⁴ CZAKON 99 perform a simultaneous fit to charged and neutral sectors.

¹⁵ CYBURT 05 limit follows by requiring that three light ν_R 's decouple when $T_{dec} > 140$ MeV. For different T_{dec} , the bound becomes $m_{W_R} > 3.3$ TeV ($T_{dec} / 140$ MeV) $^{3/4}$.

¹⁶ THOMAS 01 limit is from measurement of β^+ polarization in decay of polarized ^{12}N . The listed limit assumes no mixing.

¹⁷ ACKERSTAFF 99D limit is from τ decay parameters. Limit increase to 145 GeV for zero mixing.

¹⁸ BARENBOIM 98 assumes minimal left-right model with Higgs of $SU(2)_R$ in $SU(2)_L$ doublet. For Higgs in $SU(2)_L$ triplet, $m_{W_R} > 1100$ GeV. Bound calculated from effect of corresponding Z_{LR} on electroweak data through Z - Z_{LR} mixing.

¹⁹ The quoted limit is from μ decay parameters. BARENBOIM 97 also evaluate limit from K_L - K_S mass difference.

²⁰ STAHL 97 limit is from fit to τ -decay parameters.

²¹ ALLET 96 measured polarization-asymmetry correlation in $^{12}\text{N}\beta^+$ decay. The listed limit assumes zero L - R mixing.

²² KUZNETSOV 95 limit is from measurements of the asymmetry $\langle \vec{p}_\nu \cdot \sigma_n \rangle$ in the β decay of polarized neutrons. Zero mixing assumed. See also KUZNETSOV 94B.

²³ KUZNETSOV 94B limit is from measurements of the asymmetry $\langle \vec{p}_\nu \cdot \sigma_n \rangle$ in the β decay of polarized neutrons. Zero mixing assumed.

²⁴ BHATTACHARYYA 93 uses Z - Z' mixing limit from LEP '90 data, assuming a specific Higgs sector of $SU(2)_L \times SU(2)_R \times U(1)$ gauge model. The limit is for $m_t = 200$ GeV and slightly improves for smaller m_t .

²⁵ SEVERIJNS 93 measured polarization-asymmetry correlation in $^{107}\text{In}\beta^+$ decay. The listed limit assumes zero L - R mixing. Value quoted here is from SEVERIJNS 94 erratum.

²⁶ IMAZATO 92 measure positron asymmetry in $K^+ \rightarrow \mu^+ \nu_\mu$ decay and obtain $\xi P_\mu > 0.990$ (90% CL). If W_R couples to $u\bar{s}$ with full weak strength ($|V_{us}^R| = 1$), the result corresponds to $m_{W_R} > 653$ GeV. See their Fig. 4 for m_{W_R} limits for general $|V_{us}^R|^2 = 1 - |V_{ud}^R|^2$.

²⁷ POLAK 92B limit is from fit to muon decay parameters and is essentially determined by JODIDIO 86 data assuming $\zeta = 0$. Supersedes POLAK 91.

²⁸ AQUINO 91 limits obtained from neutron lifetime and asymmetries together with unitarity of the CKM matrix. Manifest left-right symmetry assumed. Stronger of the two limits also includes muon decay results.

²⁹ COLANGELO 91 limit uses hadronic matrix elements evaluated by QCD sum rule and is less restrictive than BEALL 82 limit which uses vacuum saturation approximation. Manifest left-right symmetry assumed.

³⁰ POLAK 91 limit is from fit to muon decay parameters and is essentially determined by JODIDIO 86 data assuming $\zeta = 0$. Superseded by POLAK 92B.

³¹ BARBIERI 89B limit holds for $m_{\nu_R} \leq 10$ MeV.

³² LANGACKER 89B limit is for any ν_R mass (either Dirac or Majorana) and for a general class of right-handed quark mixing matrices.

³³ BALKE 88 limit is for $m_{\nu_{eR}} = 0$ and $m_{\nu_{\mu R}} \leq 50$ MeV. Limits come from precise measurements of the muon decay asymmetry as a function of the positron energy.

³⁴ JODIDIO 86 is the same TRIUMF experiment as STOKER 85 (and CARR 83); however, it uses a different technique. The results given here are combined results of the

- two techniques. The technique here involves precise measurement of the end-point e^+ spectrum in the decay of the highly polarized μ^+ .
- ³⁵ STOKER 85 is same TRIUMF experiment as CARR 83. Here they measure the decay e^+ spectrum asymmetry above 46 MeV/c using a muon-spin-rotation technique. Assumed a light right-handed neutrino. Quoted limits are from combining with CARR 83.
- ³⁶ BERGSMA 83 set limit $m_{W_2}/m_{W_1} > 1.9$ at CL = 90%.
- ³⁷ CARR 83 is TRIUMF experiment with a highly polarized μ^+ beam. Looked for deviation from $V-A$ at the high momentum end of the decay e^+ energy spectrum. Limit from previous world-average muon polarization parameter is $m_{W_R} > 240$ GeV. Assumes a light right-handed neutrino.
- ³⁸ BEALL 82 limit is obtained assuming that W_R contribution to $K_L^0 - K_S^0$ mass difference is smaller than the standard one, neglecting the top quark contributions. Manifest left-right symmetry assumed.
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Limit on W_L - W_R Mixing Angle ζ

Lighter mass eigenstate $W_1 = W_L \cos\zeta - W_R \sin\zeta$. Light ν_R assumed unless noted. Values in brackets are from cosmological and astrophysical considerations.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
< 0.12	95	³⁹ ACKERSTAFF	99D	OPAL τ decay
< 0.013	90	⁴⁰ CZAKON	99	RVUE Electroweak
< 0.0333		⁴¹ BARENBOIM	97	RVUE μ decay
< 0.04	90	⁴² MISHRA	92	CCFR νN scattering
-0.0006 to 0.0028	90	⁴³ AQUINO	91	RVUE
[none 0.00001–0.02]		⁴⁴ BARBIERI	89B	ASTR SN 1987A
< 0.040	90	⁴⁵ JODIDIO	86	ELEC μ decay
-0.056 to 0.040	90	⁴⁵ JODIDIO	86	ELEC μ decay

³⁹ ACKERSTAFF 99D limit is from τ decay parameters.

⁴⁰ CZAKON 99 perform a simultaneous fit to charged and neutral sectors.

⁴¹ The quoted limit is from μ decay parameters. BARENBOIM 97 also evaluate limit from K_L - K_S mass difference.

⁴² MISHRA 92 limit is from the absence of extra large-x, large-y $\bar{\nu}_\mu N \rightarrow \bar{\nu}_\mu X$ events at Tevatron, assuming left-handed ν and right-handed $\bar{\nu}$ in the neutrino beam. The result gives $\zeta^2(1-2m_{W_1}^2/m_{W_2}^2) < 0.0015$. The limit is independent of ν_R mass.

⁴³ AQUINO 91 limits obtained from neutron lifetime and asymmetries together with unitarity of the CKM matrix. Manifest left-right asymmetry is assumed.

⁴⁴ BARBIERI 89B limit holds for $m_{\nu_R} \leq 10$ MeV.

⁴⁵ First JODIDIO 86 result assumes $m_{W_R} = \infty$, second is for unconstrained m_{W_R} .

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MASS LIMITS for Z' (Heavy Neutral Vector Boson Other Than Z)

Limits for Z'_{SM}

Z'_{SM} is assumed to have couplings with quarks and leptons which are identical to those of Z , and decays only to known fermions.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>1305	95	46 ABDALLAH	06C DLPH	$e^+ e^-$
> 850		47 ABULENCIA	06L CDF	$p\bar{p}; Z'_{SM} \rightarrow e^+ e^-$
>1500	95	48 CHEUNG	01B RVUE	Electroweak
• • • We do not use the following data for averages, fits, limits, etc. • • •				
> 825	95	49 ABULENCIA	05A CDF	$p\bar{p}; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$
> 399	95	50 ACOSTA	05R CDF	$\bar{p}p; Z'_{SM} \rightarrow \tau^+ \tau^-$
none 400–640	95	ABAZOV	04C D0	$p\bar{p}; Z'_{SM} \rightarrow q\bar{q}$
>1018	95	51 ABBIENDI	04G OPAL	$e^+ e^-$
> 670	95	52 ABAZOV	01B D0	$p\bar{p}, Z'_{SM} \rightarrow e^+ e^-$
> 710	95	53 ABREU	00S DLPH	$e^+ e^-$
> 898	95	54 BARATE	00I ALEP	$e^+ e^-$
> 809	95	55 ERLER	99 RVUE	Electroweak
> 690	95	56 ABE	97S CDF	$p\bar{p}; Z'_{SM} \rightarrow e^+ e^-, \mu^+ \mu^-$
> 490	95	ABACHI	96D D0	$p\bar{p}; Z'_{SM} \rightarrow e^+ e^-$
> 398	95	57 VILAIN	94B CHM2	$\nu_\mu e \rightarrow \nu_\mu e$ and $\bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$
> 237	90	58 ALITTI	93 UA2	$p\bar{p}; Z'_{SM} \rightarrow q\bar{q}$
none 260–600	95	59 RIZZO	93 RVUE	$p\bar{p}; Z'_{SM} \rightarrow q\bar{q}$
> 426	90	60 ABE	90F VNS	$e^+ e^-$

46 ABDALLAH 06C use data $\sqrt{s} = 130$ –207 GeV.

47 ABULENCIA 06L search for resonances decaying to $e^+ e^-$ in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

48 CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.

49 ABULENCIA 05A search for resonances decaying to electron or muon pairs in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

50 ACOSTA 05R search for resonances decaying to tau lepton pairs in $\bar{p}p$ collisions at $\sqrt{s} = 1.96$ TeV.

51 ABBIENDI 04G give 95%CL limit on $Z - Z'$ mixing $-0.00422 < \theta < 0.00091$. $\sqrt{s} = 91$ to 207 GeV.

52 ABAZOV 01B search for resonances in $p\bar{p} \rightarrow e^+ e^-$ at $\sqrt{s}=1.8$ TeV. They find $\sigma \cdot B(Z' \rightarrow ee) < 0.06$ pb for $M_{Z'} > 500$ GeV.

53 ABREU 00S uses LEP data at $\sqrt{s}=90$ to 189 GeV.

54 BARATE 00I search for deviations in cross section and asymmetries in $e^+ e^- \rightarrow$ fermions at $\sqrt{s}=90$ to 183 GeV. Assume $\theta=0$. Bounds in the mass-mixing plane are shown in their Figure 18.

55 ERLER 99 give 90%CL limit on the $Z-Z'$ mixing $-0.0041 < \theta < 0.0003$. $\rho_0=1$ is assumed.

56 ABE 97S find $\sigma(Z') \times B(e^+ e^-, \mu^+ \mu^-) < 40$ fb for $m_{Z'} > 600$ GeV at $\sqrt{s}=1.8$ TeV.

57 VILAIN 94B assume $m_t = 150$ GeV.

58 ALITTI 93 search for resonances in the two-jet invariant mass. The limit assumes $B(Z' \rightarrow q\bar{q})=0.7$. See their Fig. 5 for limits in the $m_{Z'} - B(q\bar{q})$ plane.

59 RIZZO 93 analyses CDF limit on possible two-jet resonances.

60 ABE 90F use data for R , $R_{\ell\ell}$, and $A_{\ell\ell}$. They fix $m_W = 80.49 \pm 0.43 \pm 0.24$ GeV and $m_Z = 91.13 \pm 0.03$ GeV.

Limits for Z_{LR}

Z_{LR} is the extra neutral boson in left-right symmetric models. $g_L = g_R$ is assumed unless noted. Values in parentheses assume stronger constraint on the Higgs sector, usually motivated by specific left-right symmetric models (see the Note on the W').

Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino. Direct search bounds assume decays to Standard Model fermions only, unless noted.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>518	95	61 ABBIENDI	04G OPAL	$e^+ e^-$
>860	95	62 CHEUNG	01B RVUE	Electroweak
>630	95	63 ABE	97S CDF	$p\bar{p}; Z'_{LR} \rightarrow e^+ e^-$, $\mu^+ \mu^-$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
>455	95	64 ABDALLAH	06C DLPH	$e^+ e^-$
>380	95	65 ABREU	00S DLPH	$e^+ e^-$
>436	95	66 BARATE	00I ALEP	$e^+ e^-$
>550	95	67 CHAY	00 RVUE	Electroweak
		68 ERLER	00 RVUE	Cs
		69 CASALBUONI	99 RVUE	Cs
(> 1205)	90	70 CZAKON	99 RVUE	Electroweak
>564	95	71 ERLER	99 RVUE	Electroweak
(> 1673)	95	72 ERLER	99 RVUE	Electroweak
(> 1700)	68	73 BARENBOIM	98 RVUE	Electroweak
>244	95	74 CONRAD	98 RVUE	$\nu_\mu N$ scattering
>253	95	75 VILAIN	94B CHM2	$\nu_\mu e \rightarrow \nu_\mu e$ and $\bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$
none 200–600	95	76 RIZZO	93 RVUE	$p\bar{p}; Z_{LR} \rightarrow q\bar{q}$
[> 2000]		WALKER	91 COSM	Nucleosynthesis; light ν_R
none 200–500		77 GRIFOLS	90 ASTR	SN 1987A; light ν_R
none 350–2400		78 BARBIERI	89B ASTR	SN 1987A; light ν_R

⁶¹ ABBIENDI 04G give 95% CL limit on $Z - Z'$ mixing $-0.00098 < \theta < 0.00190$. See their Fig. 20 for the limit contour in the mass-mixing plane. $\sqrt{s} = 91$ to 207 GeV.

⁶² CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.

⁶³ ABE 97S find $\sigma(Z') \times B(e^+ e^-, \mu^+ \mu^-) < 40$ fb for $m_{Z'} > 600$ GeV at $\sqrt{s} = 1.8$ TeV.

⁶⁴ ABDALLAH 06C give 95% CL limit $|\theta| < 0.0028$. See their Fig. 14 for limit contours in the mass-mixing plane.

⁶⁵ ABREU 00S give 95% CL limit on $Z - Z'$ mixing $|\theta| < 0.0018$. See their Fig. 6 for the limit contour in the mass-mixing plane. $\sqrt{s} = 90$ to 189 GeV.

⁶⁶ BARATE 00I search for deviations in cross section and asymmetries in $e^+ e^- \rightarrow$ fermions at $\sqrt{s} = 90$ to 183 GeV. Assume $\theta = 0$. Bounds in the mass-mixing plane are shown in their Figure 18.

⁶⁷ CHAY 00 also find $-0.0003 < \theta < 0.0019$. For g_R free, $m_{Z'} > 430$ GeV.

⁶⁸ ERLER 00 discuss the possibility that a discrepancy between the observed and predicted values of $Q_W(\text{Cs})$ is due to the exchange of Z' . The data are better described in a certain class of the Z' models including Z_{LR} and Z_χ .

⁶⁹ CASALBUONI 99 discuss the discrepancy between the observed and predicted values of $Q_W(\text{Cs})$. It is shown that the data are better described in a class of models including the Z_{LR} model.

- ⁷⁰ CZAKON 99 perform a simultaneous fit to charged and neutral sectors. Assumes manifest left-right symmetric model. Finds $|\theta| < 0.0042$.
- ⁷¹ ERLER 99 give 90% CL limit on the Z - Z' mixing $-0.0009 < \theta < 0.0017$.
- ⁷² ERLER 99 assumes 2 Higgs doublets, transforming as 10 of SO(10), embedded in E_6 .
- ⁷³ BARENBOIM 98 also gives 68% CL limits on the Z - Z' mixing $-0.0005 < \theta < 0.0033$. Assumes Higgs sector of minimal left-right model.
- ⁷⁴ CONRAD 98 limit is from measurements at CCFR, assuming no Z - Z' mixing.
- ⁷⁵ VILAIN 94B assume $m_t = 150$ GeV and $\theta=0$. See Fig. 2 for limit contours in the mass-mixing plane.
- ⁷⁶ RIZZO 93 analyses CDF limit on possible two-jet resonances.
- ⁷⁷ GRIFOLS 90 limit holds for $m_{\nu_R} \lesssim 1$ MeV. A specific Higgs sector is assumed. See also GRIFOLS 90D, RIZZO 91.
- ⁷⁸ BARBIERI 89B limit holds for $m_{\nu_R} \leq 10$ MeV. Bounds depend on assumed supernova core temperature.

Limits for Z_χ

Z_χ is the extra neutral boson in $SU(5) \rightarrow SU(5) \times U(1)_\chi$. $g_\chi = e/\cos\theta_W$ is assumed unless otherwise stated. We list limits with the assumption $\rho=1$ but with no further constraints on the Higgs sector. Values in parentheses assume stronger constraint on the Higgs sector motivated by superstring models. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
> 740		79 ABULENCIA	06L CDF	$p\bar{p}; Z_\chi \rightarrow e^+ e^-$
> 781	95	80 ABBIENDI	04G OPAL	$e^+ e^-$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
> 545	95	81 ABDALLAH	06C DLPH	$e^+ e^-$
> 690	95	82 ABULENCIA	05A CDF	$p\bar{p}; Z'_\chi \rightarrow e^+ e^-, \mu^+ \mu^-$
> 2100		83 BARGER	03B COSM	Nucleosynthesis; light ν_R
> 680	95	84 CHEUNG	01B RVUE	Electroweak
> 440	95	85 ABREU	00S DLPH	$e^+ e^-$
> 533	95	86 BARATE	00I ALEP	$e^+ e^-$
> 554	95	87 CHO	00 RVUE	Electroweak
		88 ERLER	00 RVUE	Cs
		89 ROSNER	00 RVUE	Cs
> 545	95	90 ERLER	99 RVUE	Electroweak
(> 1368)	95	91 ERLER	99 RVUE	Electroweak
> 215	95	92 CONRAD	98 RVUE	$\nu_\mu N$ scattering
> 595	95	93 ABE	97S CDF	$p\bar{p}; Z'_\chi \rightarrow e^+ e^-, \mu^+ \mu^-$
> 190	95	94 ARIMA	97 VNS	Bhabha scattering
> 262	95	95 VILAIN	94B CHM2	$\nu_\mu e \rightarrow \nu_\mu e; \bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$
[>1470]		96 FARAGGI	91 COSM	Nucleosynthesis; light ν_R
> 231	90	97 ABE	90F VNS	$e^+ e^-$
[> 1140]		98 GONZALEZ-G..90D	COSM	Nucleosynthesis; light ν_R
[> 2100]		99 GRIFOLS	90 ASTR	SN 1987A; light ν_R

- 79 ABULENCIA 06L search for resonances decaying to $e^+ e^-$ in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.
- 80 ABBIENDI 04G give 95% CL limit on $Z - Z'$ mixing $-0.00099 < \theta < 0.00194$. See their Fig. 20 for the limit contour in the mass-mixing plane. $\sqrt{s} = 91$ to 207 GeV.
- 81 ABDALLAH 06C give 95% CL limit $|\theta| < 0.0031$. See their Fig. 14 for limit contours in the mass-mixing plane.
- 82 ABULENCIA 05A search for resonances decaying to electron or muon pairs in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.
- 83 BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino $\delta N_\nu < 1$. The quark-hadron transition temperature $T_c = 150$ MeV is assumed. The limit with $T_c = 400$ MeV is > 4300 GeV.
- 84 CHEUNG 01B limit is derived from bounds on contact interactions in a global electroweak analysis.
- 85 ABREU 00S give 95% CL limit on $Z-Z'$ mixing $|\theta| < 0.0017$. See their Fig. 6 for the limit contour in the mass-mixing plane. $\sqrt{s} = 90$ to 189 GeV.
- 86 BARATE 00I search for deviations in cross section and asymmetries in $e^+ e^- \rightarrow$ fermions at $\sqrt{s} = 90$ to 183 GeV. Assume $\theta = 0$. Bounds in the mass-mixing plane are shown in their Figure 18.
- 87 CHO 00 use various electroweak data to constrain Z' models assuming $m_H = 100$ GeV. See Fig. 3 for limits in the mass-mixing plane.
- 88 ERLER 00 discuss the possibility that a discrepancy between the observed and predicted values of $Q_W(\text{Cs})$ is due to the exchange of Z' . The data are better described in a certain class of the Z' models including Z_{LR} and Z_χ .
- 89 ROSNER 00 discusses the possibility that a discrepancy between the observed and predicted values of $Q_W(\text{Cs})$ is due to the exchange of Z' . The data are better described in a certain class of the Z' models including Z_χ .
- 90 ERLER 99 give 90% CL limit on the $Z-Z'$ mixing $-0.0020 < \theta < 0.0015$.
- 91 ERLER 99 assumes 2 Higgs doublets, transforming as 10 of SO(10), embedded in E_6 .
- 92 CONRAD 98 limit is from measurements at CCFR, assuming no $Z-Z'$ mixing.
- 93 ABE 97S find $\sigma(Z') \times B(e^+ e^-, \mu^+ \mu^-) < 40$ fb for $m_{Z'} > 600$ GeV at $\sqrt{s} = 1.8$ TeV.
- 94 $Z-Z'$ mixing is assumed to be zero. $\sqrt{s} = 57.77$ GeV.
- 95 VILAIN 94B assume $m_t = 150$ GeV and $\theta = 0$. See Fig. 2 for limit contours in the mass-mixing plane.
- 96 FARAGGI 91 limit assumes the nucleosynthesis bound on the effective number of neutrinos $\Delta N_\nu < 0.5$ and is valid for $m_{\nu_R} < 1$ MeV.
- 97 ABE 90F use data for R , $R_{\ell\ell}$, and $A_{\ell\ell}$. ABE 90F fix $m_W = 80.49 \pm 0.43 \pm 0.24$ GeV and $m_Z = 91.13 \pm 0.03$ GeV.
- 98 Assumes the nucleosynthesis bound on the effective number of light neutrinos ($\delta N_\nu < 1$) and that ν_R is light ($\lesssim 1$ MeV).
- 99 GRIFOLS 90 limit holds for $m_{\nu_R} \lesssim 1$ MeV. See also GRIFOLS 90D, RIZZO 91.

Limits for Z_ψ

Z_ψ is the extra neutral boson in $E_6 \rightarrow \text{SO}(10) \times U(1)_\psi$. $g_\psi = e/\cos\theta_W$ is assumed unless otherwise stated. We list limits with the assumption $\rho = 1$ but with no further constraints on the Higgs sector. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>475	95	100 ABDALLAH 06C DLPH	$e^+ e^-$	
>725		101 ABULENCIA 06L CDF	$p\bar{p}; Z_\psi \rightarrow e^+ e^-$	

• • • We do not use the following data for averages, fits, limits, etc. • • •

>675	95	102 ABULENCIA	05A CDF	$p\bar{p}; Z'_\psi \rightarrow e^+e^-, \mu^+\mu^-$
>366	95	103 ABBIENDI	04G OPAL	e^+e^-
>600		104 BARGER	03B COSM	Nucleosynthesis; light ν_R
>350	95	105 ABREU	00S DLPH	e^+e^-
>294	95	106 BARATE	00I ALEP	e^+e^-
>137	95	107 CHO	00 RVUE	Electroweak
>146	95	108 ERLER	99 RVUE	Electroweak
> 54	95	109 CONRAD	98 RVUE	$\nu_\mu N$ scattering
>590	95	110 ABE	97S CDF	$p\bar{p}; Z'_\psi \rightarrow e^+e^-, \mu^+\mu^-$
>135	95	111 VILAIN	94B CHM2	$\nu_\mu e \rightarrow \nu_\mu e; \bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$
>105	90	112 ABE	90F VNS	e^+e^-
[> 160]		113 GONZALEZ-G..90D	COSM	Nucleosynthesis; light ν_R
[> 2000]		114 GRIFOLS	90D ASTR	SN 1987A; light ν_R

100 ABDALLAH 06C give 95% CL limit $|\theta| < 0.0027$. See their Fig. 14 for limit contours in the mass-mixing plane.

101 ABULENCIA 06L search for resonances decaying to e^+e^- in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

102 ABULENCIA 05A search for resonances decaying to electron or muon pairs in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

103 ABBIENDI 04G give 95% CL limit on $Z - Z'$ mixing $-0.00129 < \theta < 0.00258$. See their Fig. 20 for the limit contour in the mass-mixing plane. $\sqrt{s} = 91$ to 207 GeV.

104 BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino $\delta N_\nu < 1$. The quark-hadron transition temperature $T_c=150$ MeV is assumed. The limit with $T_c=400$ MeV is >1100 GeV.

105 ABREU 00S give 95% CL limit on $Z-Z'$ mixing $|\theta| < 0.0018$. See their Fig. 6 for the limit contour in the mass-mixing plane. $\sqrt{s}=90$ to 189 GeV.

106 BARATE 00I search for deviations in cross section and asymmetries in $e^+e^- \rightarrow$ fermions at $\sqrt{s}=90$ to 183 GeV. Assume $\theta=0$. Bounds in the mass-mixing plane are shown in their Figure 18.

107 CHO 00 use various electroweak data to constrain Z' models assuming $m_H=100$ GeV. See Fig. 3 for limits in the mass-mixing plane.

108 ERLER 99 give 90% CL limit on the $Z-Z'$ mixing $-0.0013 < \theta < 0.0024$.

109 CONRAD 98 limit is from measurements at CCFR, assuming no $Z-Z'$ mixing.

110 ABE 97S find $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) < 40$ fb for $m_{Z'} > 600$ GeV at $\sqrt{s}=1.8$ TeV.

111 VILAIN 94B assume $m_t = 150$ GeV and $\theta=0$. See Fig. 2 for limit contours in the mass-mixing plane.

112 ABE 90F use data for R , $R_{\ell\ell}$, and $A_{\ell\ell}$. ABE 90F fix $m_W = 80.49 \pm 0.43 \pm 0.24$ GeV and $m_Z = 91.13 \pm 0.03$ GeV.

113 Assumes the nucleosynthesis bound on the effective number of light neutrinos ($\delta N_\nu < 1$) and that ν_R is light ($\lesssim 1$ MeV).

114 GRIFOLS 90D limit holds for $m_{\nu_R} \lesssim 1$ MeV. See also RIZZO 91.

Limits for Z_η

Z_η is the extra neutral boson in E_6 models, corresponding to $Q_\eta = \sqrt{3/8} Q_\chi - \sqrt{5/8} Q_\psi$. $g_\eta = e/\cos\theta_W$ is assumed unless otherwise stated. We list limits with the assumption $\rho=1$ but with no further constraints on the Higgs sector. Values in parentheses assume stronger constraint on the Higgs sector motivated by superstring

models. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
> 745		115 ABULENCIA	06L	CDF $p\bar{p}; Z_\eta \rightarrow e^+e^-$
> 515	95	116 ABBIENDI	04G	OPAL e^+e^-
> 619	95	117 CHO	00	RVUE Electroweak
• • • We do not use the following data for averages, fits, limits, etc. • • •				
> 360	95	118 ABDALLAH	06C	DLPH e^+e^-
> 720	95	119 ABULENCIA	05A	CDF $p\bar{p}; Z'_\eta \rightarrow e^+e^-, \mu^+\mu^-$
> 1600		120 BARGER	03B	COSM Nucleosynthesis; light ν_R
> 310	95	121 ABREU	00S	DLPH e^+e^-
> 329	95	122 BARATE	00I	ALEP e^+e^-
> 365	95	123 ERLER	99	RVUE Electroweak
> 87	95	124 CONRAD	98	RVUE $\nu_\mu N$ scattering
> 620	95	125 ABE	97S	CDF $p\bar{p}; Z'_\eta \rightarrow e^+e^-, \mu^+\mu^-$
> 100	95	126 VILAIN	94B	CHM2 $\nu_\mu e \rightarrow \nu_\mu e; \bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$
> 125	90	127 ABE	90F	VNS e^+e^-
[> 820]		128 GONZALEZ-G.	90D	COSM Nucleosynthesis; light ν_R
[> 3300]		129 GRIFOLS	90	ASTR SN 1987A; light ν_R
[> 1040]		128 LOPEZ	90	COSM Nucleosynthesis; light ν_R

115 ABULENCIA 06L search for resonances decaying to e^+e^- in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

116 ABBIENDI 04G give 95% CL limit on $Z - Z'$ mixing $-0.00447 < \theta < 0.00331$. See their Fig. 20 for the limit contour in the mass-mixing plane. $\sqrt{s} = 91$ to 207 GeV.

117 CHO 00 use various electroweak data to constrain Z' models assuming $m_H=100$ GeV. See Fig. 3 for limits in the mass-mixing plane.

118 ABDALLAH 06C give 95% CL limit $|\theta| < 0.0092$. See their Fig. 14 for limit contours in the mass-mixing plane.

119 ABULENCIA 05A search for resonances decaying to electron or muon pairs in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV.

120 BARGER 03B limit is from the nucleosynthesis bound on the effective number of light neutrino $\delta N_\nu < 1$. The quark-hadron transition temperature $T_c=150$ MeV is assumed. The limit with $T_c=400$ MeV is > 3300 GeV.

121 ABREU 00S give 95% CL limit on $Z-Z'$ mixing $|\theta| < 0.0024$. See their Fig. 6 for the limit contour in the mass-mixing plane. $\sqrt{s}=90$ to 189 GeV.

122 BARATE 00I search for deviations in cross section and asymmetries in $e^+e^- \rightarrow$ fermions at $\sqrt{s}=90$ to 183 GeV. Assume $\theta=0$. Bounds in the mass-mixing plane are shown in their Figure 18.

123 ERLER 99 give 90% CL limit on the $Z-Z'$ mixing $-0.0062 < \theta < 0.0011$.

124 CONRAD 98 limit is from measurements at CCFR, assuming no $Z-Z'$ mixing.

125 ABE 97S find $\sigma(Z') \times B(e^+e^-, \mu^+\mu^-) < 40$ fb for $m_{Z'} > 600$ GeV at $\sqrt{s}=1.8$ TeV.

126 VILAIN 94B assume $m_t = 150$ GeV and $\theta=0$. See Fig. 2 for limit contours in the mass-mixing plane.

127 ABE 90F use data for R , $R_{\ell\ell}$, and $A_{\ell\ell}$. ABE 90F fix $m_W = 80.49 \pm 0.43 \pm 0.24$ GeV and $m_Z = 91.13 \pm 0.03$ GeV.

128 These authors claim that the nucleosynthesis bound on the effective number of light neutrinos ($\delta N_\nu < 1$) constrains Z' masses if ν_R is light ($\lesssim 1$ MeV).

129 GRIFOLS 90 limit holds for $m_{\nu_R} \lesssim 1$ MeV. See also GRIFOLS 90D, RIZZO 91.

Limits for other Z'

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •			
130 ABULENCIA	06M CDF	$Z' \rightarrow e\mu$	
131 ABAZOV	04A D0	$Z' \rightarrow t\bar{t}$	
132 BARGER	03B COSM	Nucleosynthesis; light ν_R	
133 CHO	00 RVUE	E_6 -motivated	
134 CHO	98 RVUE	E_6 -motivated	
135 ABE	97G CDF	$Z' \rightarrow \bar{q}q$	
130 ABULENCIA	06M	search for new particle with lepton flavor violating decay at $\sqrt{s} = 1.96$ TeV. See their Fig. 4 for an exclusion plot on a mass-coupling plane.	
131	Search for narrow resonance decaying to $t\bar{t}$. See their Fig. 2 for limit on σB .		
132	BARGER 03B	use the nucleosynthesis bound on the effective number of light neutrino δN_ν . See their Figs. 4–5 for limits in general E_6 motivated models.	
133	CHO 00	use various electroweak data to constrain Z' models assuming $m_H=100$ GeV. See Fig. 2 for limits in general E_6 -motivated models.	
134	CHO 98	study constraints on four-Fermi contact interactions obtained from low-energy electroweak experiments, assuming no Z - Z' mixing.	
135	Search for Z' decaying to dijets at $\sqrt{s}=1.8$ TeV. For Z' with electromagnetic strength coupling, no bound is obtained.		

Indirect Constraints on Kaluza-Klein Gauge Bosons

Bounds on a Kaluza-Klein excitation of the Z boson or photon in $d=1$ extra dimension. These bounds can also be interpreted as a lower bound on $1/R$, the size of the extra dimension. Unless otherwise stated, bounds assume all fermions live on a single brane and all gauge fields occupy the $4+d$ -dimensional bulk. See also the section on “Extra Dimensions” in the “Searches” Listings in this Review.

VALUE (TeV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
> 4.7		136 MUECK	02 RVUE	Electroweak
> 3.3	95	137 CORNET	00 RVUE	$e\nu qq'$
>5000		138 DELGADO	00 RVUE	ϵ_K
> 2.6	95	139 DELGADO	00 RVUE	Electroweak
> 3.3	95	140 RIZZO	00 RVUE	Electroweak
> 2.9	95	141 MARCIANO	99 RVUE	Electroweak
> 2.5	95	142 MASIP	99 RVUE	Electroweak
> 1.6	90	143 NATH	99 RVUE	Electroweak
> 3.4	95	144 STRUMIA	99 RVUE	Electroweak
136	MUECK 02 limit is 2σ and is from global electroweak fit ignoring correlations among observables. Higgs is assumed to be confined on the brane and its mass is fixed. For scenarios of bulk Higgs, of brane-SU(2) $_L$, bulk-U(1) $_Y$, and of bulk-SU(2) $_L$, brane-U(1) $_Y$, the corresponding limits are > 4.6 TeV, > 4.3 TeV and > 3.0 TeV, respectively.			
137	Bound is derived from limits on $e\nu qq'$ contact interaction, using data from HERA and the Tevatron.			
138	Bound holds only if first two generations of quarks lives on separate branes. If quark mixing is not complex, then bound lowers to 400 TeV from Δm_K .			
139	See Figs. 1 and 2 of DELGADO 00 for several model variations. Special boundary conditions can be found which permit KK states down to 950 GeV and that agree with the measurement of $Q_W(C_s)$. Quoted bound assumes all Higgs bosons confined to brane; placing one Higgs doublet in the bulk lowers bound to 2.3 TeV.			

- 140 Bound is derived from global electroweak analysis assuming the Higgs field is trapped on the matter brane. If the Higgs propagates in the bulk, the bound increases to 3.8 TeV.
 141 Bound is derived from global electroweak analysis but considering only presence of the KK W bosons.
 142 Global electroweak analysis used to obtain bound independent of position of Higgs on brane or in bulk.
 143 Bounds from effect of KK states on G_F , α , M_W , and M_Z . Hard cutoff at string scale determined using gauge coupling unification. Limits for $d=2,3,4$ rise to 3.5, 5.7, and 7.8 TeV.
 144 Bound obtained for Higgs confined to the matter brane with $m_H=500$ GeV. For Higgs in the bulk, the bound increases to 3.5 TeV.
-

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MASS LIMITS for Leptoquarks from Pair Production

These limits rely only on the color or electroweak charge of the leptoquark.

VALUE (GeV)	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
>251	95		145 ABAZOV	06A D0	Second generation
>226	95		146 ABULENCIA	06T CDF	Second generation
>256	95		147 ABAZOV	05H D0	First generation
>236	95		148 ACOSTA	05P CDF	First generation
>148	95		149 AFFOLDER	00K CDF	Third generation
• • • We do not use the following data for averages, fits, limits, etc. • • •					
>136	95		150 ABAZOV	06L D0	All generations
>117	95		151 ACOSTA	05I CDF	First generation
> 99	95		152 ABBIENDI	03R OPAL	First generation
>100	95		152 ABBIENDI	03R OPAL	Second generation
> 98	95		152 ABBIENDI	03R OPAL	Third generation
> 98	95		153 ABAZOV	02 D0	All generations
>225	95		154 ABAZOV	01D D0	First generation
> 85.8	95		155 ABBIENDI	00M OPAL	Superseded by ABBI-ENDI 03R
> 85.5	95		155 ABBIENDI	00M OPAL	Superseded by ABBI-ENDI 03R
> 82.7	95		155 ABBIENDI	00M OPAL	Superseded by ABBI-ENDI 03R
>200	95		156 ABBOTT	00C D0	Second generation
>123	95		157 AFFOLDER	00K CDF	Second generation
>160	95		158 ABBOTT	99J D0	Second generation
>225	95		159 ABBOTT	98E D0	First generation
> 94	95		160 ABBOTT	98J D0	Third generation
>202	95		161 ABE	98S CDF	Second generation
>242	95		162 GROSS-PILCH.98		First generation
> 99	95		163 ABE	97F CDF	Third generation
>213	95		164 ABE	97X CDF	First generation
> 45.5	95	2	165,166 ABREU	93J DLPH	First + second generation
> 44.4	95		167 ADRIANI	93M L3	First generation
> 44.5	95		167 ADRIANI	93M L3	Second generation
> 45	95		167 DECAMP	92 ALEP	Third generation
none 8.9–22.6	95		168 KIM	90 AMY	First generation
none 10.2–23.2	95		168 KIM	90 AMY	Second generation
none 5–20.8	95		169 BARTEL	87B JADE	
none 7–20.5	95		170 BEHREND	86B CELL	

- 145 ABAZOV 06A search for scalar leptoquarks using $\mu\mu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV and at 1.96 TeV. The limit above assumes $B(\mu q) = 1$. For $B(\mu q) = 0.5$, the limit becomes 204 GeV.
- 146 ABULENCIA 06T search for scalar leptoquarks using $\mu\mu jj$, $\nu\nu jj$, and $e\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.96$ TeV. The quoted limit assumes $B(\mu q) = 1$. For $B(\mu q) = 0.5$ or 0.1, the bound becomes 208 GeV or 143 GeV, respectively. See their Fig. 4 for the exclusion limit as a function of $B(\mu q)$.
- 147 ABAZOV 05H search for scalar leptoquarks using $ee jj$ and $e\nu jj$ events in $\bar{p}p$ collisions at $E_{cm} = 1.8$ TeV and 1.96 TeV. The limit above assumes $B(eq) = 1$. For $B(eq) = 0.5$ the bound becomes 234 GeV.
- 148 ACOSTA 05P search for scalar leptoquarks using $ee jj$, $e\nu jj$ events in $\bar{p}p$ collisions at $E_{cm} = 1.96$ TeV. The limit above assumes $B(eq) = 1$. For $B(eq) = 0.5$ and 0.1, the bound becomes 205 GeV and 145 GeV, respectively.
- 149 AFFOLDER 00K search for scalar leptoquark using $\nu\nu bb$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The quoted limit assumes $B(\nu b) = 1$. Bounds for vector leptoquarks are also given.
- 150 ABAZOV 06L search for scalar leptoquarks using $\nu\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV and at 1.96 TeV. The limit above assumes $B(\nu q) = 1$.
- 151 ACOSTA 05I search for scalar leptoquarks using $\nu\nu jj$ events in $\bar{p}p$ collisions at $E_{cm} = 1.96$ TeV. The limit above assumes $B(\nu q) = 1$.
- 152 ABBIENDI 03R search for scalar/vector leptoquarks in $e^+ e^-$ collisions at $\sqrt{s} = 189-209$ GeV. The quoted limits are for charge $-4/3$ isospin 0 scalar-leptoquark with $B(\ell q) = 1$. See their table 12 for other cases.
- 153 ABAZOV 02 search for scalar leptoquarks using $\nu\nu jj$ events in $\bar{p}p$ collisions at $E_{cm} = 1.8$ TeV. The bound holds for all leptoquark generations. Vector leptoquarks are likewise constrained to lie above 200 GeV.
- 154 ABAZOV 01D search for scalar leptoquarks using $e\nu jj$, $ee jj$, and $\nu\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The limit above assumes $B(eq) = 1$. For $B(eq) = 0.5$ and 0, the bound becomes 204 and 79 GeV, respectively. Bounds for vector leptoquarks are also given. Supersedes ABBOTT 98E.
- 155 ABBIENDI 00M search for scalar/vector leptoquarks in $e^+ e^-$ collisions at $\sqrt{s} = 183$ GeV. The quoted limits are for charge $-4/3$ isospin 0 scalar-leptoquarks with $B(\ell q) = 1$. See their Table 8 and Figs. 6–9 for other cases.
- 156 ABBOTT 00C search for scalar leptoquarks using $\mu\mu jj$, $\mu\nu jj$, and $\nu\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The limit above assumes $B(\mu q) = 1$. For $B(\mu q) = 0.5$ and 0, the bound becomes 180 and 79 GeV respectively. Bounds for vector leptoquarks are also given.
- 157 AFFOLDER 00K search for scalar leptoquark using $\nu\nu cc$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The quoted limit assumes $B(\nu c) = 1$. Bounds for vector leptoquarks are also given.
- 158 ABBOTT 99J search for leptoquarks using $\mu\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The quoted limit is for a scalar leptoquark with $B(\mu q) = B(\nu q) = 0.5$. Limits on vector leptoquarks range from 240 to 290 GeV.
- 159 ABBOTT 98E search for scalar leptoquarks using $e\nu jj$, $ee jj$, and $\nu\nu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The limit above assumes $B(eq) = 1$. For $B(eq) = 0.5$ and 0, the bound becomes 204 and 79 GeV, respectively.
- 160 ABBOTT 98J search for charge $-1/3$ third generation scalar and vector leptoquarks in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The quoted limit is for scalar leptoquark with $B(\nu b) = 1$.
- 161 ABE 98S search for scalar leptoquarks using $\mu\mu jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The limit is for $B(\mu q) = 1$. For $B(\mu q) = B(\nu q) = 0.5$, the limit is > 160 GeV.
- 162 GROSS-PILCHER 98 is the combined limit of the CDF and D \emptyset Collaborations as determined by a joint CDF/D \emptyset working group and reported in this FNAL Technical Memo. Original data published in ABE 97X and ABBOTT 98E.
- 163 ABE 97F search for third generation scalar and vector leptoquarks in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The quoted limit is for scalar leptoquark with $B(\tau b) = 1$.
- 164 ABE 97X search for scalar leptoquarks using $ee jj$ events in $p\bar{p}$ collisions at $E_{cm} = 1.8$ TeV. The limit is for $B(eq) = 1$.

- 165 Limit is for charge $-1/3$ isospin-0 leptoquark with $B(\ell q) = 2/3$.
- 166 First and second generation leptoquarks are assumed to be degenerate. The limit is slightly lower for each generation.
- 167 Limits are for charge $-1/3$, isospin-0 scalar leptoquarks decaying to $\ell^- q$ or νq with any branching ratio. See paper for limits for other charge-isospin assignments of leptoquarks.
- 168 KIM 90 assume pair production of charge $2/3$ scalar-leptoquark via photon exchange. The decay of the first (second) generation leptoquark is assumed to be any mixture of $d e^+$ and $u \bar{\nu}$ ($s \mu^+$ and $c \bar{\nu}$). See paper for limits for specific branching ratios.
- 169 BARTEL 87B limit is valid when a pair of charge $2/3$ spinless leptoquarks X is produced with point coupling, and when they decay under the constraint $B(X \rightarrow c \bar{\nu}_\mu) + B(X \rightarrow s \mu^+) = 1$.
- 170 BEHREND 86B assumed that a charge $2/3$ spinless leptoquark, χ , decays either into $s \mu^+$ or $c \bar{\nu}$: $B(\chi \rightarrow s \mu^+) + B(\chi \rightarrow c \bar{\nu}) = 1$.
-

MASS LIMITS for Leptoquarks from Single Production

These limits depend on the q - ℓ -leptoquark coupling g_{LQ} . It is often assumed that $g_{LQ}^2/4\pi=1/137$. Limits shown are for a scalar, weak isoscalar, charge $-1/3$ leptoquark.

<u>VALUE (GeV)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •				
>295	95	171 AKTAS 172 CHEKANOV	05B H1 05A ZEUS	First generation Lepton-flavor violation
>298	95	173 CHEKANOV	03B ZEUS	First generation
>197	95	174 ABBIENDI 175 CHEKANOV	02B OPAL 02 ZEUS	First generation Repl. by CHEKANOV 05A
>290	95	176 ADLOFF	01C H1	First generation
>204	95	177 BREITWEG 178 BREITWEG	01 ZEUS 00E ZEUS	First generation First generation
>161	95	179 ABREU	99G DLPH	First generation
>200	95	180 ADLOFF 181 DERRICK	99 H1 97 ZEUS	First generation Lepton-flavor violation
> 73	95	182 ABREU	93J DLPH	Second generation
>168	95	183 DERRICK	93 ZEUS	First generation

171 AKTAS 05B limit is for a scalar, weak isoscalar, charge $-1/3$ leptoquark coupled with e_R . See their Fig. 3 for limits on states with different quantum numbers.

172 CHEKANOV 05 search for various leptoquarks with lepton-flavor violating couplings. See their Figs.6–10 and Tables 1–8 for detailed limits.

173 CHEKANOV 03B limit is for a scalar, weak isoscalar, charge $-1/3$ leptoquark coupled with e_R . See their Figs. 11–12 and Table 5 for limits on states with different quantum numbers.

174 For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 4 and Fig. 5.

175 CHEKANOV 02 search for various leptoquarks with lepton-flavor violating couplings. See their Figs. 6–7 and Tables 5–6 for detailed limits.

176 For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 3.

177 See their Fig. 14 for limits in the mass-coupling plane.

178 BREITWEG 00E search for $F=0$ leptoquarks in $e^+ p$ collisions. For limits in mass-coupling plane, see their Fig. 11.

- 179 ABREU 99G limit obtained from process $e\gamma \rightarrow LQ+q$. For limits on vector and scalar states with different quantum numbers and the limits in the coupling-mass plane, see their Fig. 4 and Table 2.
- 180 For limits on states with different quantum numbers and the limits in the mass-coupling plane, see their Fig. 13 and Fig. 14. ADLOFF 99 also search for leptoquarks with lepton-flavor violating couplings. ADLOFF 99 supersedes AID 96B.
- 181 DERRICK 97 search for various leptoquarks with lepton-flavor violating couplings. See their Figs. 5–8 and Table 1 for detailed limits.
- 182 Limit from single production in Z decay. The limit is for a leptoquark coupling of electromagnetic strength and assumes $B(\ell q) = 2/3$. The limit is 77 GeV if first and second leptoquarks are degenerate.
- 183 DERRICK 93 search for single leptoquark production in ep collisions with the decay $e q$ and νq . The limit is for leptoquark coupling of electromagnetic strength and assumes $B(e q) = B(\nu q) = 1/2$. The limit for $B(e q) = 1$ is 176 GeV. For limits on states with different quantum numbers, see their Table 3.
-

Indirect Limits for Leptoquarks

<u>VALUE (TeV)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •				
> 1.7	96	184 CHEKANOV	05A ZEUS	Lepton-flavor violation
> 46	90	185 ADLOFF	03 H1	First generation
		186 CHANG	03 BELL	Pati-Salam type
		187 CHEKANOV	02 ZEUS	Repl. by CHEKANOV 05A
> 1.7	95	188 CHEUNG	01B RVUE	First generation
> 0.39	95	189 ACCIARRI	00P L3	$e^+ e^- \rightarrow q q$
> 1.5	95	190 ADLOFF	00 H1	First generation
> 0.2	95	191 BARATE	00I ALEP	$e^+ e^-$
		192 BARGER	00 RVUE	Cs
		193 GABRIELLI	00 RVUE	Lepton flavor violation
> 0.74	95	194 ZARNECKI	00 RVUE	S_1 leptoquark
		195 ABBIENDI	99 OPAL	
> 19.3	95	196 ABE	98V CDF	$B_s \rightarrow e^\pm \mu^\mp$, Pati-Salam type
		197 ACCIARRI	98J L3	$e^+ e^- \rightarrow q \bar{q}$
		198 ACKERSTAFF	98V OPAL	$e^+ e^- \rightarrow q \bar{q}$, $e^+ e^- \rightarrow b \bar{b}$
> 0.76	95	199 DEANDREA	97 RVUE	\tilde{R}_2 leptoquark
		200 DERRICK	97 ZEUS	Lepton-flavor violation
		201 GROSSMAN	97 RVUE	$B \rightarrow \tau^+ \tau^- (\chi)$
		202 JADACH	97 RVUE	$e^+ e^- \rightarrow q \bar{q}$
>1200		203 KUZNETSOV	95B RVUE	Pati-Salam type
		204 MIZUKOSHI	95 RVUE	Third generation scalar lepto-quark
> 0.3	95	205 BHATTACH...	94 RVUE	Spin-0 leptoquark coupled to $\bar{e}_R t_L$
		206 DAVIDSON	94 RVUE	
> 18		207 KUZNETSOV	94 RVUE	Pati-Salam type
> 0.43	95	208 LEURER	94 RVUE	First generation spin-1 lepto-quark
> 0.44	95	208 LEURER	94B RVUE	First generation spin-0 lepto-quark
		209 MAHANTA	94 RVUE	P and T violation
> 1		210 SHANKER	82 RVUE	Nonchiral spin-0 leptoquark
> 125		210 SHANKER	82 RVUE	Nonchiral spin-1 leptoquark

- 184 CHEKANOV 05 search for various leptoquarks with lepton-flavor violating couplings. See their Figs.6–10 and Tables 1–8 for detailed limits.
- 185 ADLOFF 03 limit is for the weak isotriplet spin-0 leptoquark at strong coupling $\lambda=\sqrt{4\pi}$. For the limits of leptoquarks with different quantum numbers, see their Table 3. Limits are derived from bounds on $e^\pm q$ contact interactions.
- 186 The bound is derived from $B(B^0 \rightarrow e^\pm \mu^\mp) < 1.7 \times 10^{-7}$.
- 187 CHEKANOV 02 search for lepton-flavor violation in $e p$ collisions. See their Tables 1–4 for limits on lepton-flavor violating and four-fermion interactions induced by various leptoquarks.
- 188 CHEUNG 01B quoted limit is for a scalar, weak isoscalar, charge $-1/3$ leptoquark with a coupling of electromagnetic strength. The limit is derived from bounds on contact interactions in a global electroweak analysis. For the limits of leptoquarks with different quantum numbers, see Table 5.
- 189 ACCIARRI 00P limit is for the weak isoscalar spin-0 leptoquark with the coupling of electromagnetic strength. For the limits of leptoquarks with different quantum numbers, see their Table 4.
- 190 ADLOFF 00 limit is for the weak isotriplet spin-0 leptoquark at strong coupling, $\lambda=\sqrt{4\pi}$. For the limits of leptoquarks with different quantum numbers, see their Table 2. ADLOFF 00 limits are from the Q^2 spectrum measurement of $e^+ p \rightarrow e^+ X$.
- 191 BARATE 00I search for deviations in cross section and jet-charge asymmetry in $e^+ e^- \rightarrow \bar{q}q$ due to t -channel exchange of a leptoquark at $\sqrt{s}=130$ to 183 GeV. Limits for other scalar and vector leptoquarks are also given in their Table 22.
- 192 BARGER 00 explain the deviation of atomic parity violation in cesium atoms from prediction is explained by scalar leptoquark exchange.
- 193 GABRIELLI 00 calculate various process with lepton flavor violation in leptoquark models.
- 194 ZARNECKI 00 limit is derived from data of HERA, LEP, and Tevatron and from various low-energy data including atomic parity violation. Leptoquark coupling with electromagnetic strength is assumed.
- 195 ABBIENDI 99 limits are from $e^+ e^- \rightarrow q\bar{q}$ cross section at 130–136, 161–172, 183 GeV. See their Fig. 8 and Fig. 9 for limits in mass-coupling plane.
- 196 ABE 98V quoted limit is from $B(B_s \rightarrow e^\pm \mu^\mp) < 8.2 \times 10^{-6}$. ABE 98V also obtain a similar limit on $M_{LQ} > 20.4$ TeV from $B(B_d \rightarrow e^\pm \mu^\mp) < 4.5 \times 10^{-6}$. Both bounds assume the non-canonical association of the b quark with electrons or muons under SU(4).
- 197 ACCIARRI 98J limit is from $e^+ e^- \rightarrow q\bar{q}$ cross section at $\sqrt{s}=130$ –172 GeV which can be affected by the t - and u -channel exchanges of leptoquarks. See their Fig. 4 and Fig. 5 for limits in the mass-coupling plane.
- 198 ACKERSTAFF 98V limits are from $e^+ e^- \rightarrow q\bar{q}$ and $e^+ e^- \rightarrow b\bar{b}$ cross sections at $\sqrt{s}=130$ –172 GeV, which can be affected by the t - and u -channel exchanges of leptoquarks. See their Fig. 21 and Fig. 22 for limits of leptoquarks in mass-coupling plane.
- 199 DEANDREA 97 limit is for \tilde{R}_2 leptoquark obtained from atomic parity violation (APV). The coupling of leptoquark is assumed to be electromagnetic strength. See Table 2 for limits of the four-fermion interactions induced by various scalar leptoquark exchange. DEANDREA 97 combines APV limit and limits from Tevatron and HERA. See Fig. 1–4 for combined limits of leptoquark in mass-coupling plane.
- 200 DERRICK 97 search for lepton-flavor violation in $e p$ collision. See their Tables 2–5 for limits on lepton-flavor violating four-fermion interactions induced by various leptoquarks.
- 201 GROSSMAN 97 estimate the upper bounds on the branching fraction $B \rightarrow \tau^+ \tau^- (X)$ from the absence of the B decay with large missing energy. These bounds can be used to constrain leptoquark induced four-fermion interactions.
- 202 JADACH 97 limit is from $e^+ e^- \rightarrow q\bar{q}$ cross section at $\sqrt{s}=172.3$ GeV which can be affected by the t - and u -channel exchanges of leptoquarks. See their Fig. 1 for limits on vector leptoquarks in mass-coupling plane.
- 203 KUZNETSOV 95B use π , K , B , τ decays and μe conversion and give a list of bounds on the leptoquark mass and the fermion mixing matrix in the Pati-Salam model. The quoted limit is from $K_L \rightarrow \mu e$ decay assuming zero mixing.

- 204 MIZUKOSHI 95 calculate the one-loop radiative correction to the Z -physics parameters in various scalar leptoquark models. See their Fig. 4 for the exclusion plot of third generation leptoquark models in mass-coupling plane.
- 205 BHATTACHARYYA 94 limit is from one-loop radiative correction to the leptonic decay width of the Z . $m_H=250$ GeV, $\alpha_s(m_Z)=0.12$, $m_t=180$ GeV, and the electroweak strength of leptoquark coupling are assumed. For leptoquark coupled to $\bar{e}_L t_R$, $\bar{\mu} t$, and $\bar{\tau} t$, see Fig. 2 in BHATTACHARYYA 94B erratum and Fig. 3.
- 206 DAVIDSON 94 gives an extensive list of the bounds on leptoquark-induced four-fermion interactions from π , K , D , B , μ , τ decays and meson mixings, etc. See Table 15 of DAVIDSON 94 for detail.
- 207 KUZNETSOV 94 gives mixing independent bound of the Pati-Salam leptoquark from the cosmological limit on $\pi^0 \rightarrow \bar{\nu}\nu$.
- 208 LEURER 94, LEURER 94B limits are obtained from atomic parity violation and apply to any chiral leptoquark which couples to the first generation with electromagnetic strength. For a nonchiral leptoquark, universality in $\pi_{\ell 2}$ decay provides a much more stringent bound.
- 209 MAHANTA 94 gives bounds of P - and T -violating scalar-leptoquark couplings from atomic and molecular experiments.
- 210 From $(\pi \rightarrow e\nu)/(\pi \rightarrow \mu\nu)$ ratio. SHANKER 82 assumes the leptoquark induced four-fermion coupling $4g^2/M^2 (\bar{\nu}_{eL} u_R) (\bar{d}_L e_R)$ with $g=0.004$ for spin-0 leptoquark and $g^2/M^2 (\bar{\nu}_{eL} \gamma_\mu u_L) (\bar{d}_R \gamma^\mu e_R)$ with $g \approx 0.6$ for spin-1 leptoquark.

MASS LIMITS for Diquarks

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
none 290–420	95	211 ABE	97G CDF	E_6 diquark
none 15–31.7	95	212 ABREU	940 DLPH	SUSY E_6 diquark
211 ABE 97G search for new particle decaying to dijets.				
212 ABREU 940 limit is from $e^+ e^- \rightarrow \bar{c}s c s$. Range extends up to 43 GeV if diquarks are degenerate in mass.				

MASS LIMITS for g_A (axigluon)

Axigluons are massive color-octet gauge bosons in chiral color models and have axial-vector coupling to quarks with the same coupling strength as gluons.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
>365	95	213 DONCHESKI	98 RVUE	$\Gamma(Z \rightarrow \text{hadron})$
none 200–980	95	214 ABE	97G CDF	$p\bar{p} \rightarrow g_A X, X \rightarrow 2 \text{jets}$
none 200–870	95	215 ABE	95N CDF	$p\bar{p} \rightarrow g_A X, g_A \rightarrow q\bar{q}$
none 240–640	95	216 ABE	93G CDF	$p\bar{p} \rightarrow g_A X, g_A \rightarrow 2\text{jets}$
> 50	95	217 CUYPERS	91 RVUE	$\sigma(e^+ e^- \rightarrow \text{hadrons})$
none 120–210	95	218 ABE	90H CDF	$p\bar{p} \rightarrow g_A X, g_A \rightarrow 2\text{jets}$
> 29	95	219 ROBINETT	89 THEO	Partial-wave unitarity
none 150–310	95	220 ALBAJAR	88B UA1	$p\bar{p} \rightarrow g_A X, g_A \rightarrow 2\text{jets}$
> 20		BERGSTROM	88 RVUE	$p\bar{p} \rightarrow \gamma X \text{ via } g_A g$
> 9		221 CUYPERS	88 RVUE	$\gamma \text{ decay}$
> 25		222 DONCHESKI	88B RVUE	$\gamma \text{ decay}$

- 213 DONCHESKI 98 compare α_s derived from low-energy data and that from $\Gamma(Z \rightarrow \text{hadrons})/\Gamma(Z \rightarrow \text{leptons})$.
- 214 ABE 97G search for new particle decaying to dijets.
- 215 ABE 95N assume axigluons decaying to quarks in the Standard Model only.
- 216 ABE 93G assume $\Gamma(g_A) = N\alpha_s m_{g_A}/6$ with $N = 10$.
- 217 CUYPERS 91 compare α_s measured in γ decay and that from R at PEP/PETRA energies.
- 218 ABE 90H assumes $\Gamma(g_A) = N\alpha_s m_{g_A}/6$ with $N = 5$ ($\Gamma(g_A) = 0.09m_{g_A}$). For $N = 10$, the excluded region is reduced to 120–150 GeV.
- 219 ROBINETT 89 result demands partial-wave unitarity of $J = 0$ $t\bar{t} \rightarrow t\bar{t}$ scattering amplitude and derives a limit $m_{g_A} > 0.5 m_t$. Assumes $m_t > 56$ GeV.
- 220 ALBAJAR 88B result is from the nonobservation of a peak in two-jet invariant mass distribution. $\Gamma(g_A) < 0.4 m_{g_A}$ assumed. See also BAGGER 88.
- 221 CUYPERS 88 requires $\Gamma(\gamma \rightarrow g g_A) < \Gamma(\gamma \rightarrow g g g)$. A similar result is obtained by DONCHESKI 88.
- 222 DONCHESKI 88B requires $\Gamma(\gamma \rightarrow g q\bar{q})/\Gamma(\gamma \rightarrow g g g) < 0.25$, where the former decay proceeds via axigluon exchange. A more conservative estimate of < 0.5 leads to $m_{g_A} > 21$ GeV.
-

X^0 (Heavy Boson) Searches in Z Decays

Searches for radiative transition of Z to a lighter spin-0 state X^0 decaying to hadrons, a lepton pair, a photon pair, or invisible particles as shown in the comments. The limits are for the product of branching ratios.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
		223 BARATE 98U	ALEP	$X^0 \rightarrow \ell\bar{\ell}, q\bar{q}, gg, \gamma\gamma, \nu\bar{\nu}$
		224 ACCIARRI 97Q	L3	$X^0 \rightarrow$ invisible particle(s)
		225 ACTON 93E	OPAL	$X^0 \rightarrow \gamma\gamma$
		226 ABREU 92D	DLPH	$X^0 \rightarrow$ hadrons
		227 ADRIANI 92F	L3	$X^0 \rightarrow$ hadrons
		228 ACTON 91	OPAL	$X^0 \rightarrow$ anything
$<1.1 \times 10^{-4}$	95	229 ACTON 91B	OPAL	$X^0 \rightarrow e^+ e^-$
$<9 \times 10^{-5}$	95	229 ACTON 91B	OPAL	$X^0 \rightarrow \mu^+ \mu^-$
$<1.1 \times 10^{-4}$	95	229 ACTON 91B	OPAL	$X^0 \rightarrow \tau^+ \tau^-$
$<2.8 \times 10^{-4}$	95	230 ADEVA 91D	L3	$X^0 \rightarrow e^+ e^-$
$<2.3 \times 10^{-4}$	95	230 ADEVA 91D	L3	$X^0 \rightarrow \mu^+ \mu^-$
$<4.7 \times 10^{-4}$	95	231 ADEVA 91D	L3	$X^0 \rightarrow$ hadrons
$<8 \times 10^{-4}$	95	232 AKRAWY 90J	OPAL	$X^0 \rightarrow$ hadrons
223 BARATE 98U obtain limits on $B(Z \rightarrow \gamma X^0)B(X^0 \rightarrow \ell\bar{\ell}, q\bar{q}, gg, \gamma\gamma, \nu\bar{\nu})$. See their Fig. 17.				
224 See Fig. 4 of ACCIARRI 97Q for the upper limit on $B(Z \rightarrow \gamma X^0; E_\gamma > E_{\min})$ as a function of E_{\min} .				
225 ACTON 93E give $\sigma(e^+ e^- \rightarrow X^0 \gamma) \cdot B(X^0 \rightarrow \gamma\gamma) < 0.4$ pb (95% CL) for $m_{X^0} = 60 \pm 2.5$ GeV. If the process occurs via s -channel γ exchange, the limit translates to $\Gamma(X^0) \cdot B(X^0 \rightarrow \gamma\gamma)^2 < 20$ MeV for $m_{X^0} = 60 \pm 1$ GeV.				
226 ABREU 92D give $\sigma_Z \cdot B(Z \rightarrow \gamma X^0) \cdot B(X^0 \rightarrow \text{hadrons}) < (3-10)$ pb for $m_{X^0} = 10-78$ GeV. A very similar limit is obtained for spin-1 X^0 .				

- 227 ADRIANI 92F search for isolated γ in hadronic Z decays. The limit $\sigma_Z \cdot B(Z \rightarrow \gamma X^0) \cdot B(X^0 \rightarrow \text{hadrons}) < (2-10) \text{ pb}$ (95%CL) is given for $m_{X^0} = 25-85 \text{ GeV}$.
- 228 ACTON 91 searches for $Z \rightarrow Z^* X^0$, $Z^* \rightarrow e^+ e^-$, $\mu^+ \mu^-$, or $\nu \bar{\nu}$. Excludes any new scalar X^0 with $m_{X^0} < 9.5 \text{ GeV}/c$ if it has the same coupling to ZZ^* as the MSM Higgs boson.
- 229 ACTON 91B limits are for $m_{X^0} = 60-85 \text{ GeV}$.
- 230 ADEVA 91D limits are for $m_{X^0} = 30-89 \text{ GeV}$.
- 231 ADEVA 91D limits are for $m_{X^0} = 30-86 \text{ GeV}$.
- 232 AKRAWY 90J give $\Gamma(Z \rightarrow \gamma X^0) \cdot B(X^0 \rightarrow \text{hadrons}) < 1.9 \text{ MeV}$ (95%CL) for $m_{X^0} = 32-80 \text{ GeV}$. We divide by $\Gamma(Z) = 2.5 \text{ GeV}$ to get product of branching ratios. For nonresonant transitions, the limit is $B(Z \rightarrow \gamma q\bar{q}) < 8.2 \text{ MeV}$ assuming three-body phase space distribution.
-

MASS LIMITS for a Heavy Neutral Boson Coupling to $e^+ e^-$

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
none 55-61		233 ODAKA	89 VNS	$\Gamma(X^0 \rightarrow e^+ e^-) \cdot B(X^0 \rightarrow \text{hadrons}) \gtrsim 0.2 \text{ MeV}$
>45	95	234 DERRICK	86 HRS	$\Gamma(X^0 \rightarrow e^+ e^-) = 6 \text{ MeV}$
>46.6	95	235 ADEVA	85 MRKJ	$\Gamma(X^0 \rightarrow e^+ e^-) = 10 \text{ keV}$
>48	95	235 ADEVA 236 BERGER	85 MRKJ 85B PLUT	$\Gamma(X^0 \rightarrow e^+ e^-) = 4 \text{ MeV}$
none 39.8-45.5		237 ADEVA	84 MRKJ	$\Gamma(X^0 \rightarrow e^+ e^-) = 10 \text{ keV}$
>47.8	95	237 ADEVA	84 MRKJ	$\Gamma(X^0 \rightarrow e^+ e^-) = 4 \text{ MeV}$
none 39.8-45.2		237 BEHREND	84C CELL	
>47	95	237 BEHREND	84C CELL	$\Gamma(X^0 \rightarrow e^+ e^-) = 4 \text{ MeV}$
233		ODAKA 89 looked for a narrow or wide scalar resonance in $e^+ e^- \rightarrow \text{hadrons}$ at $E_{\text{cm}} = 55.0-60.8 \text{ GeV}$.		
234		DERRICK 86 found no deviation from the Standard Model Bhabha scattering at $E_{\text{cm}} = 29 \text{ GeV}$ and set limits on the possible scalar boson $e^+ e^-$ coupling. See their figure 4 for excluded region in the $\Gamma(X^0 \rightarrow e^+ e^-)$ - m_{X^0} plane. Electronic chiral invariance requires a parity doublet of X^0 , in which case the limit applies for $\Gamma(X^0 \rightarrow e^+ e^-) = 3 \text{ MeV}$.		
235		ADEVA 85 first limit is from 2γ , $\mu^+ \mu^-$, hadrons assuming X^0 is a scalar. Second limit is from $e^+ e^-$ channel. $E_{\text{cm}} = 40-47 \text{ GeV}$. Supersedes ADEVA 84.		
236		BERGER 85B looked for effect of spin-0 boson exchange in $e^+ e^- \rightarrow e^+ e^-$ and $\mu^+ \mu^-$ at $E_{\text{cm}} = 34.7 \text{ GeV}$. See Fig. 5 for excluded region in the m_{X^0} - $\Gamma(X^0)$ plane.		
237		ADEVA 84 and BEHREND 84C have $E_{\text{cm}} = 39.8-45.5 \text{ GeV}$. MARK-J searched X^0 in $e^+ e^- \rightarrow \text{hadrons}$, 2γ , $\mu^+ \mu^-$, $e^+ e^-$ and CELLO in the same channels plus τ pair. No narrow or broad X^0 is found in the energy range. They also searched for the effect of X^0 with $m_X > E_{\text{cm}}$. The second limits are from Bhabha data and for spin-0 singlet. The same limits apply for $\Gamma(X^0 \rightarrow e^+ e^-) = 2 \text{ MeV}$ if X^0 is a spin-0 doublet. The second limit of BEHREND 84C was read off from their figure 2. The original papers also list limits in other channels.		

Search for X^0 Resonance in e^+e^- Collisions

The limit is for $\Gamma(X^0 \rightarrow e^+e^-) \cdot B(X^0 \rightarrow f)$, where f is the specified final state.

Spin 0 is assumed for X^0 .

VALUE (keV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<10 ³	95	238 ABE	93C VNS	$\Gamma(ee)$
<(0.4–10)	95	239 ABE	93C VNS	$f = \gamma\gamma$
<(0.3–5)	95	240,241 ABE	93D TOPZ	$f = \gamma\gamma$
<(2–12)	95	240,241 ABE	93D TOPZ	$f = \text{hadrons}$
<(4–200)	95	241,242 ABE	93D TOPZ	$f = ee$
<(0.1–6)	95	241,242 ABE	93D TOPZ	$f = \mu\mu$
<(0.5–8)	90	243 STERNER	93 AMY	$f = \gamma\gamma$
238 Limit is for $\Gamma(X^0 \rightarrow e^+e^-)$ $m_{X^0} = 56\text{--}63.5$ GeV for $\Gamma(X^0) = 0.5$ GeV.				
239 Limit is for $m_{X^0} = 56\text{--}61.5$ GeV and is valid for $\Gamma(X^0) \ll 100$ MeV. See their Fig. 5 for limits for $\Gamma = 1, 2$ GeV.				
240 Limit is for $m_{X^0} = 57.2\text{--}60$ GeV.				
241 Limit is valid for $\Gamma(X^0) \ll 100$ MeV. See paper for limits for $\Gamma = 1$ GeV and those for $J = 2$ resonances.				
242 Limit is for $m_{X^0} = 56.6\text{--}60$ GeV.				
243 STERNER 93 limit is for $m_{X^0} = 57\text{--}59.6$ GeV and is valid for $\Gamma(X^0) < 100$ MeV. See their Fig. 2 for limits for $\Gamma = 1, 3$ GeV.				

Search for X^0 Resonance in $e p$ Collisions

VALUE	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •			
244 CHEKANOV 02B ZEUS	ZEUS	$X \rightarrow jj$	
244 CHEKANOV 02B search for photoproduction of X decaying into dijets in $e p$ collisions. See their Fig. 5 for the limit on the photoproduction cross section.			

Search for X^0 Resonance in Two-Photon Process

The limit is for $\Gamma(X^0) \cdot B(X^0 \rightarrow \gamma\gamma)^2$. Spin 0 is assumed for X^0 .

VALUE (MeV)	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<2.6	95	245 ACTON	93E OPAL	$m_{X^0} = 60 \pm 1$ GeV
<2.9	95	BUSKULIC	93F ALEP	$m_{X^0} \sim 60$ GeV

245 ACTON 93E limit for a $J = 2$ resonance is 0.8 MeV.

Search for X^0 Resonance in $e^+e^- \rightarrow X^0\gamma$

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •			
246 ABBIENDI	03D OPAL	$X^0 \rightarrow \gamma\gamma$	
247 ABREU	00Z DLPH	X^0 decaying invisibly	
248 ADAM	96C DLPH	X^0 decaying invisibly	

- 246 ABBIENDI 03D measure the $e^+ e^- \rightarrow \gamma\gamma\gamma$ cross section at $\sqrt{s}=181\text{--}209$ GeV. The upper bound on the production cross section, $\sigma(e^+ e^- \rightarrow X^0\gamma)$ times the branching ratio for $X^0 \rightarrow \gamma\gamma$, is less than 0.03 pb at 95%CL for X^0 masses between 20 and 180 GeV. See their Fig. 9b for the limits in the mass-cross section plane.
- 247 ABREU 00Z is from the single photon cross section at $\sqrt{s}=183, 189$ GeV. The production cross section upper limit is less than 0.3 pb for X^0 mass between 40 and 160 GeV. See their Fig. 4 for the limit in mass-cross section plane.
- 248 ADAM 96C is from the single photon production cross at $\sqrt{s}=130, 136$ GeV. The upper bound is less than 3 pb for X^0 masses between 60 and 130 GeV. See their Fig. 5 for the exact bound on the cross section $\sigma(e^+ e^- \rightarrow \gamma X^0)$.
-

Search for X^0 Resonance in $Z \rightarrow f\bar{f}X^0$

The limit is for $B(Z \rightarrow f\bar{f}X^0) \cdot B(X^0 \rightarrow F)$ where f is a fermion and F is the specified final state. Spin 0 is assumed for X^0 .

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •				
$<3.7 \times 10^{-6}$	95	249 ABREU 250 ABREU 251 ABREU	96T DLPH 96T DLPH 96T DLPH	$f=e,\mu,\tau; F=\gamma\gamma$ $f=\nu; F=\gamma\gamma$ $f=q; F=\gamma\gamma$
$<6.8 \times 10^{-6}$	95	250 ACTON	93E OPAL	$f=e,\mu,\tau; F=\gamma\gamma$
$<5.5 \times 10^{-6}$	95	250 ACTON	93E OPAL	$f=q; F=\gamma\gamma$
$<3.1 \times 10^{-6}$	95	250 ACTON	93E OPAL	$f=\nu; F=\gamma\gamma$
$<6.5 \times 10^{-6}$	95	250 ACTON	93E OPAL	$f=e,\mu; F=\ell\bar{\ell}, q\bar{q}, \nu\bar{\nu}$
$<7.1 \times 10^{-6}$	95	250 BUSKULIC 252 ADRIANI	93F ALEP 92F L3	$f=e,\mu; F=\ell\bar{\ell}, q\bar{q}, \nu\bar{\nu}$ $f=q; F=\gamma\gamma$

249 ABREU 96T obtain limit as a function of m_{X^0} . See their Fig. 6.

250 Limit is for m_{X^0} around 60 GeV.

251 ABREU 96T obtain limit as a function of m_{X^0} . See their Fig. 15.

252 ADRIANI 92F give $\sigma_Z \cdot B(Z \rightarrow q\bar{q}X^0) \cdot B(X^0 \rightarrow \gamma\gamma) < (0.75\text{--}1.5)$ pb (95%CL) for $m_{X^0} = 10\text{--}70$ GeV. The limit is 1 pb at 60 GeV.

Search for X^0 Resonance in $p\bar{p} \rightarrow WX^0$

<u>VALUE (MeV)</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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• • • We do not use the following data for averages, fits, limits, etc. • • •

253 ABE	97W CDF	$X^0 \rightarrow b\bar{b}$	
• • • We do not use the following data for averages, fits, limits, etc. • • •			

253 ABE 97W search for X^0 production associated with W in $p\bar{p}$ collisions at $E_{cm}=1.8$ TeV. The 95%CL upper limit on the production cross section times the branching ratio for $X^0 \rightarrow b\bar{b}$ ranges from 14 to 19 pb for X^0 mass between 70 and 120 GeV. See their Fig. 3 for upper limits of the production cross section as a function of m_{X^0} .

Heavy Particle Production in Quarkonium Decays

Limits are for branching ratios to modes shown.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •				

$<1.5 \times 10^{-5}$	90	254 BAEST	95	CLE2	$\gamma(1S) \rightarrow X^0 \gamma$, $m_{X^0} < 5 \text{ GeV}$
$<3 \times 10^{-5}$ – 6×10^{-3}	90	255 BAEST	95	CLE2	$\gamma(1S) \rightarrow X^0 \bar{X}^0 \gamma$, $m_{X^0} < 3.9 \text{ GeV}$
$<5.6 \times 10^{-5}$	90	256 ANTREASYAN 90C	CBAL		$\gamma(1S) \rightarrow X^0 \gamma$, $m_{X^0} < 7.2 \text{ GeV}$
		257 ALBRECHT	89	ARG	

254 BAEST 95 two-body limit is for pseudoscalar X^0 . The limit becomes $< 10^{-4}$ for $m_{X^0} < 7.7 \text{ GeV}$.

255 BAEST 95 three-body limit is for phase-space photon energy distribution and angular distribution same as for $\gamma \rightarrow g g \gamma$.

256 ANTREASYAN 90C assume that X^0 does not decay in the detector.

257 ALBRECHT 89 give limits for $B(\gamma(1S), \gamma(2S) \rightarrow X^0 \gamma) \cdot B(X^0 \rightarrow \pi^+ \pi^-, K^+ K^-)$, $p\bar{p}$ for $m_{X^0} < 3.5 \text{ GeV}$.

REFERENCES FOR Searches for Heavy Bosons Other Than Higgs Bosons

ABAZOV	06A	PL B636 183	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06L	PL B640 230	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06N	PL B641 423	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABDALLAH	06C	EPJ C45 589	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABULENCIA	06L	PRL 96 211801	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06M	PRL 96 211802	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06T	PR D73 051102R	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABAZOV	05H	PR D71 071104R	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABULENCIA	05A	PRL 95 252001	A. Abulencia <i>et al.</i>	(CDF Collab.)
ACOSTA	05I	PR D71 112001	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	05P	PR D72 051107R	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	05R	PRL 95 131801	D. Acosta <i>et al.</i>	(CDF Collab.)
AKTAS	05B	PL B629 9	A. Aktas <i>et al.</i>	(H1 Collab.)
CHEKANOV	05	PL B610 212	S. Chekanov <i>et al.</i>	(HERA ZEUS Collab.)
CHEKANOV	05A	EPJ C44 463	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
CYBURT	05	ASP 23 313	R.H. Cyburt <i>et al.</i>	
ABAZOV	04A	PRL 92 221801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	04C	PR D69 111101R	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	04G	EPJ C33 173	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03D	EPJ C26 331	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03R	EPJ C31 281	G. Abbiendi <i>et al.</i>	(OPAL)
ACOSTA	03B	PRL 90 081802	D. Acosta <i>et al.</i>	(CDF Collab.)
ADLOFF	03	PL B568 35	C. Adloff <i>et al.</i>	(H1 Collab.)
BARGER	03B	PR D67 075009	V. Barger, P. Langacker, H. Lee	
CHANG	03	PR D68 111101R	M.-C. Chang <i>et al.</i>	(BELLE Collab.)
CHEKANOV	03B	PR D68 052004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
ABAZOV	02	PRL 88 191801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	02B	PL B526 233	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
AFFOLDER	02C	PRL 88 071806	T. Affolder <i>et al.</i>	(CDF Collab.)
CHEKANOV	02	PR D65 092004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
CHEKANOV	02B	PL B531 9	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
MUECK	02	PR D65 085037	A. Mueck, A. Pilaftsis, R. Rueckl	
ABAZOV	01B	PRL 87 061802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	01D	PR D64 092004	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ADLOFF	01C	PL B523 234	C. Adloff <i>et al.</i>	(H1 Collab.)
AFFOLDER	01I	PRL 87 231803	T. Affolder <i>et al.</i>	(CDF Collab.)
BREITWEG	01	PR D63 052002	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
CHEUNG	01B	PL B517 167	K. Cheung	
THOMAS	01	NP A694 559	E. Thomas <i>et al.</i>	
ABBIENDI	00M	EPJ C13 15	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	00C	PRL 84 2088	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	00	PRL 84 5716	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	00S	PL B485 45	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00Z	EPJ C17 53	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	00P	PL B489 81	M. Acciari <i>et al.</i>	(L3 Collab.)
ADLOFF	00	PL B479 358	C. Adloff <i>et al.</i>	(H1 Collab.)
AFFOLDER	00K	PRL 85 2056	T. Affolder <i>et al.</i>	(CDF Collab.)

BARATE	00I	EPJ C12 183	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARGER	00	PL B480 149	V. Barger, K. Cheung	
BREITWEG	00E	EPJ C16 253	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
CHAY	00	PR D61 035002	J. Chay, K.Y. Lee, S. Nam	
CHO	00	MPL A15 311	G. Cho	
CORNET	00	PR D61 037701	F. Cornet, M. Relano, J. Rico	
DELGADO	00	JHEP 0001 030	A. Delgado, A. Pomarol, M. Quiros	
ERLER	00	PRL 84 212	J. Erler, P. Langacker	
GABRIELLI	00	PR D62 055009	E. Gabrielli	
RIZZO	00	PR D61 016007	T.G. Rizzo, J.D. Wells	
ROSNER	00	PR D61 016006	J.L. Rosner	
ZARNECKI	00	EPJ C17 695	A. Zarnecki	
ABBIENDI	99	EPJ C6 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	99J	PRL 83 2896	B. Abbott <i>et al.</i>	(D0 Collab.)
ABREU	99G	PL B446 62	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACKERSTAFF	99D	EPJ C8 3	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ADLOFF	99	EPJ C11 447	C. Adloff <i>et al.</i>	(H1 Collab.)
Also		EPJ C14 553 (erratum)	C. Adloff <i>et al.</i>	(H1 Collab.)
CASALBUONI	99	PL B460 135	R. Casalbuoni <i>et al.</i>	
CZAKON	99	PL B458 355	M. Czakon, J. Gluza, M. Zralek	
ERLER	99	PL B456 68	J. Erler, P. Langacker	
MARCIANO	99	PR D60 093006	W. Marciano	
MASIP	99	PR D60 096005	M. Masip, A. Pomarol	
NATH	99	PR D60 116004	P. Nath, M. Yamaguchi	
STRUMIA	99	PL B466 107	A. Strumia	
ABBOTT	98E	PRL 80 2051	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	98J	PRL 81 38	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	98S	PRL 81 4806	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	98V	PRL 81 5742	F. Abe <i>et al.</i>	(CDF Collab.)
ACCIARRI	98J	PL B433 163	M. Acciari <i>et al.</i>	(L3 Collab.)
ACKERSTAFF	98V	EPJ C2 441	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
BARATE	98U	EPJ C4 571	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARENBOIM	98	EPJ C1 369	G. Barenboim	
CHO	98	EPJ C5 155	G. Cho, K. Hagiwara, S. Matsumoto	
CONRAD	98	RMP 70 1341	J.M. Conrad, M.H. Shaevitz, T. Bolton	
DONCHESKI	98	PR D58 097702	M.A. Doncheski, R.W. Robinett	
GROSS-PILCH..	98	hep-ex/9810015	C. Gross-Pilcher, G. Landsberg, M. Paterno	
ABE	97F	PRL 78 2906	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	97G	PR D55 R5263	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	97S	PRL 79 2192	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	97W	PRL 79 3819	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	97X	PRL 79 4327	F. Abe <i>et al.</i>	(CDF Collab.)
ACCIARRI	97Q	PL B412 201	M. Acciari <i>et al.</i>	(L3 Collab.)
ARIMA	97	PR D55 19	T. Arima <i>et al.</i>	(VENUS Collab.)
BARENBOIM	97	PR D55 4213	G. Barenboim <i>et al.</i>	(VALE, IFIC)
DEANDREA	97	PL B409 277	A. Deandrea	(MARS)
DERRICK	97	ZPHY C73 613	M. Derrick <i>et al.</i>	(ZEUS Collab.)
GROSSMAN	97	PR D55 2768	Y. Grossman, Z. Ligeti, E. Nardi	(REHO, CIT)
JADACH	97	PL B408 281	S. Jadach, B.F.L. Ward, Z. Was	(CERN, INPK+)
STAHL	97	ZPHY C74 73	A. Stahl, H. Voss	(BONN)
ABACHI	96C	PRL 76 3271	S. Abachi <i>et al.</i>	(D0 Collab.)
ABACHI	96D	PL B385 471	S. Abachi <i>et al.</i>	(D0 Collab.)
ABREU	96T	ZPHY C72 179	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ADAM	96C	PL B380 471	W. Adam <i>et al.</i>	(DELPHI Collab.)
AID	96B	PL B369 173	S. Aid <i>et al.</i>	(H1 Collab.)
ALLET	96	PL B383 139	M. Allet <i>et al.</i>	(VILL, LEUV, LOUV, WISC)
ABACHI	95E	PL B358 405	S. Abachi <i>et al.</i>	(D0 Collab.)
ABE	95M	PRL 74 2900	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	95N	PRL 74 3538	F. Abe <i>et al.</i>	(CDF Collab.)
BALEST	95	PR D51 2053	R. Balest <i>et al.</i>	(CLEO Collab.)
KUZNETSOV	95	PRL 75 794	I.A. Kuznetsov <i>et al.</i>	(PNPI, KIAE, HARV+)
KUZNETSOV	95B	PAN 58 2113	A.V. Kuznetsov, N.V. Mikheev	(YARO)
		Translated from YAF 58	2228.	
MIZUKOSHI	95	NP B443 20	J.K. Mizukoshi, O.J.P. Eboli, M.C. Gonzalez-Garcia	
ABREU	940	ZPHY C64 183	P. Abreu <i>et al.</i>	(DELPHI Collab.)
BHATTACH...	94	PL B336 100	G. Bhattacharyya, J. Ellis, K. Sridhar	(CERN)
Also		PL B338 522 (erratum)	G. Bhattacharyya, J. Ellis, K. Sridhar	(CERN)
BHATTACH...	94B	PL B338 522 (erratum)	G. Bhattacharyya, J. Ellis, K. Sridhar	(CERN)
DAVIDSON	94	ZPHY C61 613	S. Davidson, D. Bailey, B.A. Campbell	(CFPA+)
KUZNETSOV	94	PL B329 295	A.V. Kuznetsov, N.V. Mikheev	(YARO)

KUZNETSOV	94B	JETPL 60 315 Translated from ZETFP	I.A. Kuznetsov <i>et al.</i>	(PNPI, KIAE, HARV+)
LEURER	94	PR D50 536	M. Leurer	(REHO)
LEURER	94B	PR D49 333	M. Leurer	(REHO)
Also		PRL 71 1324	M. Leurer	(REHO)
MAHANTA	94	PL B337 128	U. Mahanta	(MEHTA)
SEVERIJNS	94	PRL 73 611 (erratum)	N. Severijns <i>et al.</i>	(LOUV, WISC, LEUV+)
VILAIN	94B	PL B332 465	P. Vilain <i>et al.</i>	(CHARM II Collab.)
ABE	93C	PL B302 119	K. Abe <i>et al.</i>	(VENUS Collab.)
ABE	93D	PL B304 373	T. Abe <i>et al.</i>	(TOPAZ Collab.)
ABE	93G	PRL 71 2542	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	93J	PL B316 620	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACTON	93E	PL B311 391	P.D. Acton <i>et al.</i>	(OPAL Collab.)
ADRIANI	93M	PRPL 236 1	O. Adriani <i>et al.</i>	(L3 Collab.)
ALITTI	93	NP B400 3	J. Alitti <i>et al.</i>	(UA2 Collab.)
BHATTACH...	93	PR D47 R3693	G. Bhattacharyya <i>et al.</i>	(CALC, JADA, ICTP+)
BUSKULIC	93F	PL B308 425	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
DERRICK	93	PL B306 173	M. Derrick <i>et al.</i>	(ZEUS Collab.)
RIZZO	93	PR D48 4470	T.G. Rizzo	(ANL)
SEVERIJNS	93	PRL 70 4047	N. Severijns <i>et al.</i>	(LOUV, WISC, LEUV+)
Also		PRL 73 611 (erratum)	N. Severijns <i>et al.</i>	(LOUV, WISC, LEUV+)
STERNER	93	PL B303 385	K.L. Sterner <i>et al.</i>	(AMY Collab.)
ABREU	92D	ZPHY C53 555	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ADRIANI	92F	PL B292 472	O. Adriani <i>et al.</i>	(L3 Collab.)
DECAMP	92	PRPL 216 253	D. Decamp <i>et al.</i>	(ALEPH Collab.)
IMAZATO	92	PRL 69 877	J. Imazato <i>et al.</i>	(KEK, INUS, TOKY+)
MISHRA	92	PRL 68 3499	S.R. Mishra <i>et al.</i>	(COLU, CHIC, FNAL+)
POLAK	92B	PR D46 3871	J. Polak, M. Zralek	(SILES)
ACTON	91	PL B268 122	D.P. Acton <i>et al.</i>	(OPAL Collab.)
ACTON	91B	PL B273 338	D.P. Acton <i>et al.</i>	(OPAL Collab.)
ADEVA	91D	PL B262 155	B. Adeva <i>et al.</i>	(L3 Collab.)
AQUINO	91	PL B261 280	M. Aquino, A. Fernandez, A. Garcia	(CINV, PUEB)
COLANGELO	91	PL B253 154	P. Colangelo, G. Nardulli	(BARI)
CUPERS	91	PL B259 173	F. Cuypers, A.F. Falk, P.H. Frampton	(DURH, HARV+)
FARAGGI	91	MPL A6 61	A.E. Faraggi, D.V. Nanopoulos	(TAMU)
POLAK	91	NP B363 385	J. Polak, M. Zralek	(SILES)
RIZZO	91	PR D44 202	T.G. Rizzo	(WISC, ISU)
WALKER	91	APJ 376 51	T.P. Walker <i>et al.</i>	(HSCA, OSU, CHIC+)
ABE	90F	PL B246 297	K. Abe <i>et al.</i>	(VENUS Collab.)
ABE	90H	PR D41 1722	F. Abe <i>et al.</i>	(CDF Collab.)
AKRAWY	90J	PL B246 285	M.Z. Akrawy <i>et al.</i>	(OPAL Collab.)
ANTREASYAN	90C	PL B251 204	D. Antreasyan <i>et al.</i>	(Crystal Ball Collab.)
GONZALEZ-G...	90D	PL B240 163	M.C. Gonzalez-Garcia, J.W.F. Valle	(VALE)
GRIFOLS	90	NP B331 244	J.A. Grifols, E. Masso	(BARC)
GRIFOLS	90D	PR D42 3293	J.A. Grifols, E. Masso, T.G. Rizzo	(BARC, CERN+)
KIM	90	PL B240 243	G.N. Kim <i>et al.</i>	(AMY Collab.)
LOPEZ	90	PL B241 392	J.L. Lopez, D.V. Nanopoulos	(TAMU)
ALBAJAR	89	ZPHY C44 15	C. Albajar <i>et al.</i>	(UA1 Collab.)
ALBRECHT	89	ZPHY C42 349	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
BARBIERI	89B	PR D39 1229	R. Barbieri, R.N. Mohapatra	(PISA, UMD)
LANGACKER	89B	PR D40 1569	P. Langacker, S. Uma Sankar	(PENN)
ODAKA	89	JPSJ 58 3037	S. Odaka <i>et al.</i>	(VENUS Collab.)
ROBINETT	89	PR D39 834	R.W. Robinett	(PSU)
ALBAJAR	88B	PL B209 127	C. Albajar <i>et al.</i>	(UA1 Collab.)
BAGGER	88	PR D37 1188	J. Bagger, C. Schmidt, S. King	(HARV, BOST)
BALKE	88	PR D37 587	B. Balke <i>et al.</i>	(LBL, UCB, COLO, NWES+)
BERGSTROM	88	PL B212 386	L. Bergstrom	(STOH)
CUPERS	88	PRL 60 1237	F. Cuypers, P.H. Frampton	(UNCCH)
DONCHESKI	88	PL B206 137	M.A. Doncheski, H. Grotch, R. Robinett	(PSU)
DONCHESKI	88B	PR D38 412	M.A. Doncheski, H. Grotch, R.W. Robinett	(PSU)
ANSARI	87D	PL B195 613	R. Ansari <i>et al.</i>	(UA2 Collab.)
BARTEL	87B	ZPHY C36 15	W. Bartel <i>et al.</i>	(JADE Collab.)
BEHREND	86B	PL B178 452	H.J. Behrend <i>et al.</i>	(CELLO Collab.)
DERRICK	86	PL 166B 463	M. Derrick <i>et al.</i>	(HRS Collab.)
Also		PR D34 3286	M. Derrick <i>et al.</i>	(HRS Collab.)
JODIDIO	86	PR D34 1967	A. Jodidio <i>et al.</i>	(LBL, NWES, TRIU)
Also		PR D37 237 (erratum)	A. Jodidio <i>et al.</i>	(LBL, NWES, TRIU)
MOHAPATRA	86	PR D34 909	R.N. Mohapatra	(UMD)
ADEVA	85	PL 152B 439	B. Adeva <i>et al.</i>	(Mark-J Collab.)
BERGER	85B	ZPHY C27 341	C. Berger <i>et al.</i>	(PLUTO Collab.)
STOKER	85	PRL 54 1887	D.P. Stoker <i>et al.</i>	(LBL, NWES, TRIU)

ADEVA	84	PRL 53 134	B. Adeva <i>et al.</i>	(Mark-J Collab.)
BEHREND	84C	PL 140B 130	H.J. Behrend <i>et al.</i>	(CELLO Collab.)
BERGSMA	83	PL 122B 465	F. Bergsma <i>et al.</i>	(CHARM Collab.)
CARR	83	PRL 51 627	J. Carr <i>et al.</i>	(LBL, NWES, TRIU)
BEALL	82	PRL 48 848	G. Beall, M. Bander, A. Soni	(UCI, UCLA)
SHANKER	82	NP B204 375	O. Shanker	(TRIU)
