

## PRODUCTION AND DECAY OF $b$ -FLAVORED HADRONS

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The  $b$  quark belongs to the third generation of quarks and is the weak-doublet partner of the  $t$  quark. The existence of the third-generation quark doublet was proposed in 1973 by Kobayashi and Maskawa [1] in their model of the quark mixing matrix (“CKM” matrix), and confirmed four years later by the first observation of a  $b\bar{b}$  meson [2]. In the KM model,  $CP$  violation is explained within the Standard Model (SM) by an irreducible phase of the  $3 \times 3$  unitary matrix. The regular pattern of the three lepton and quark families is one of the most intriguing puzzles in particle physics. The existence of families gives rise to many of the free parameters in the SM, including the fermion masses, and the elements of the CKM matrix.

Since the  $b$  quark is the lighter element of the third-generation quark doublet, the decays of  $b$ -flavored hadrons occur via generation-changing processes through this matrix. Because of this, and the fact that the CKM matrix is close to a  $3 \times 3$  unit matrix, many interesting features such as loop and box diagrams, flavor oscillations, as well as large  $CP$  asymmetries, can be observed in the weak decays of  $b$ -flavored hadrons.

The CKM matrix is parameterized by three real parameters and one complex phase. This complex phase can become a source of  $CP$  violation in  $B$  meson decays. A crucial milestone was the first observation of  $CP$  violation in the  $B$  meson system in 2001, by the BaBar [3] and Belle [4] collaborations. They measured a large value for the parameter  $\sin 2\beta$  ( $= \sin 2\phi_1$ ) [5], almost four decades after the discovery of a small  $CP$  asymmetry in neutral kaons. A more detailed discussion of the CKM matrix and  $CP$  violation can be found elsewhere in this *Review* [6,7].

Recent developments in the physics of  $b$ -hadrons include the observation of direct  $CP$  violation, results for rare higher-order weak decays, investigations of heavier  $b$ -hadrons ( $B_s$ ,  $B_c$ , baryons, excited states), measurement of the  $B_s$ -mixing

frequency, increasingly accurate determinations of the CKM matrix parameters.

The structure of this mini-review is organized as follows. After a brief description of theory and terminology, we discuss  $b$ -quark production and current results on spectroscopy and lifetimes of  $b$ -flavored hadrons. We then discuss some basic properties of  $B$ -meson decays, followed by summaries of hadronic, rare, and electroweak penguin decays of  $B$ -mesons. There are separate mini-reviews for  $B\bar{B}$  mixing [8] and the extraction of the CKM matrix elements  $V_{cb}$  and  $V_{ub}$  from  $B$ -meson decays [9] in this *Review*.

**Theory and terminology:** The ground states of  $b$ -flavored hadrons decay via weak interactions. In most hadrons, the  $b$ -quark is accompanied by light-partner quarks ( $d$ ,  $u$ , or  $s$ ), and the decay modes are well described by the decay of the  $b$  quark (spectator model) [10]. The dominant decay mode of a  $b$  quark is  $b \rightarrow cW^{*-}$  (referred to as a “tree” or “spectator” decay), where the virtual  $W$  materializes either into a pair of leptons  $\ell\bar{\nu}$  (“semileptonic decay”), or into a pair of quarks which then hadronizes. The decays in which the spectator quark combines with one of the quarks from  $W^*$  to form one of the final state hadrons are suppressed by a factor  $\sim (1/3)^2$ , because the colors of the two quarks from different sources must match (“color-suppression”).

Many aspects of  $B$  decays can be understood through the Heavy Quark Effective Theory (HQET) [11]. This has been particularly successful for semileptonic decays. For further discussion of HQET, see for instance Ref. 12. For hadronic decays, one typically uses effective Hamiltonian calculations that rely on a perturbative expansion with Wilson coefficients. In addition, some form of the factorization hypothesis is commonly used, where, in analogy with semileptonic decays, two-body hadronic decays of  $B$  mesons are expressed as the product of two independent hadronic currents, one describing the formation of a charm meson (in case of the dominant  $b \rightarrow cW^{*-}$  decays), and the other the hadronization of the remaining  $\bar{u}d$  (or  $\bar{c}s$ ) system from the virtual  $W^-$ . Qualitatively, for a  $B$  decay with a large energy release, the  $\bar{u}d$  pair (produced as a color singlet) travels

fast enough to leave the interaction region without influencing the charm meson. This is known to work well for the dominant spectator decays [13]. There are several common implementations of these ideas for hadronic  $B$  decays, the most common of which are QCD factorization (QCDF) [14], perturbative QCD (pQCD) [15], and soft collinear effective theory (SCET) [16].

The transition  $b \rightarrow u$  is suppressed by  $|V_{ub}/V_{cb}|^2 \sim (0.1)^2$  relative to  $b \rightarrow c$  transitions. The transition  $b \rightarrow s$  is a flavor-changing neutral-current (FCNC) process, and although not allowed in the SM as a tree-process, can occur via more complex loop diagrams (denoted “penguin” decays). The rates for such processes are comparable or larger than CKM-suppressed  $b \rightarrow u$  processes. Penguin processes involving  $b \rightarrow d$  transitions are also possible, and have been observed [17,18]. Other decay processes discussed in this *Review* include  $W$ -exchange (a  $W$  is exchanged between initial-state quarks), penguin annihilation (the gluon from a penguin loop attaches to the spectator quark, similar to an exchange diagram), and pure-annihilation (the initial quarks annihilate to a virtual  $W$ , which then decays).

***Production and spectroscopy:*** The bound states of a  $\bar{b}$  antiquark and a  $u$ ,  $d$ ,  $s$ , or  $c$  quark are referred to as the  $B_u$  ( $B^+$ ),  $B_d$  ( $B^0$ ),  $B_s$ , and  $B_c$  mesons, respectively. The  $B_c$  is the heaviest of the ground-state  $b$ -flavored mesons, and the most difficult to produce: it was observed for the first time in the semileptonic mode by CDF in 1998 [19], but its mass was accurately determined only in 2006, from the fully reconstructed mode  $B_c^+ \rightarrow J/\psi\pi^+$  [20].

The first excited meson is called the  $B^*$  meson, while  $B^{**}$  is the generic name for the four orbitally excited ( $L = 1$ )  $B$ -meson states that correspond to the  $P$ -wave mesons in the charm system,  $D^{**}$ . Excited states of the  $B_s$  meson are similarly named  $B_s^*$  and  $B_s^{**}$ . Of the possible bound  $\bar{b}b$  states, the  $\Upsilon$  series (S-wave) and the  $\chi_b$  (P-wave) are well studied. The pseudoscalar ground state  $\eta_b$  also has been observed by BaBar [21] (and confirmed by CLEO [22]), indirectly through the decay  $\Upsilon(3S) \rightarrow \gamma\eta_b$ . See Ref. 23 for classification and naming of these and other states.

Experimental studies of  $b$  decays have been performed in  $e^+e^-$  collisions at the  $\Upsilon(4S)$  (ARGUS, CLEO, Belle, BaBar) and  $\Upsilon(5S)$  (CLEO, Belle) resonances, as well as at higher energies, at the  $Z$  resonance (SLC, LEP) and in  $p\bar{p}$  collisions (Tevatron). The  $e^+e^- \rightarrow b\bar{b}$  production cross-section at the  $Z$ ,  $\Upsilon(4S)$ , and  $\Upsilon(5S)$  resonances are about  $6.6 \text{ nb}$ ,  $1.1 \text{ nb}$ , and  $0.3 \text{ nb}$  respectively. High-energy hadron collisions produce  $b$ -flavored hadrons of all species with much larger cross-sections:  $\sigma(p\bar{p} \rightarrow bX, |\eta| < 1) \sim 30 \text{ } \mu\text{b}$  at the Tevatron ( $\sqrt{s} = 1.96 \text{ TeV}$ ), and even higher at the energies of the LHC  $pp$  collider (up to a factor of ten at  $\sqrt{s} = 14 \text{ TeV}$ ).

BaBar and Belle have accumulated respectively  $560 \text{ fb}^{-1}$  and  $1020 \text{ fb}^{-1}$  of data, of which  $433 \text{ fb}^{-1}$  and  $710 \text{ fb}^{-1}$  respectively are at the  $\Upsilon(4S)$  resonance; CDF and D0 have currently accumulated about  $10 \text{ fb}^{-1}$  each. At the LHC, CMS and ATLAS have collected  $5 \text{ fb}^{-1}$  of data and LHCb has collected about  $1 \text{ fb}^{-1}$ . These numbers indicate that the majority of  $b$ -quarks have been produced in hadron collisions, but the large backgrounds cause the hadron collider experiments to have lower selection efficiency. Only the few decay modes for which triggering and reconstruction are easiest have been studied so far in hadron collisions. These have included final states with leptons, and exclusive modes with all charged particles in the final state. In contrast, detectors operating at  $e^+e^-$  colliders (“B-Factories”) have a high efficiency for most decays, and have provided large samples of a rich variety of decays of  $B^0$  and  $B^+$  mesons.

In hadron collisions, most production happens as  $b\bar{b}$  pairs, either via  $s$ -channel production or gluon-splitting, with a smaller fraction of single  $b$ -quarks produced by flavor excitation. The total  $b$ -production cross section is an interesting test of our understanding of QCD processes. For many years, experimental measurements have been several times higher than predictions. With improved measurements [24], more accurate input parameters, and more advanced calculations [25], the discrepancy between theory and data is now much reduced, although the presence of inconsistencies among existing measurements makes further studies desirable.

Each quark of a  $b\bar{b}$  pair produced in hadron collisions hadronizes separately and incoherently from the other, but it is still possible, although difficult, to obtain a statistical indication of the charge of a produced  $b/\bar{b}$  quark (“flavor tag” or “charge tag”) from the accompanying particles produced in the hadronization process, or from the decay products of the other quark. The momentum spectrum of produced  $b$ -quarks typically peaks near the  $b$ -quark mass, and extends to much higher momenta, dropping by about a decade for every ten GeV. This implies typical decay lengths of the order of a millimeter; the resolution for the decay vertex must be more precise than this to resolve the fast oscillations of  $B_s$  mesons.

In  $e^+e^-$  colliders, since the  $B$  mesons are very slow in the  $\Upsilon(4S)$  rest frame, asymmetric beam energies are used to boost the decay products to improve the precision of time-dependent measurements that are crucial for the study of  $CP$  violation. At KEKB, the boost is  $\beta\gamma = 0.43$ , and the typical  $B$ -meson decay length is dilated from  $\approx 20 \mu m$  to  $\approx 200 \mu m$ . PEP-II uses a slightly larger boost,  $\beta\gamma = 0.55$ . The two  $B$  mesons produced in  $\Upsilon(4S)$  decay are in a coherent quantum state, which makes it easier than in hadron collisions to infer the charge state of one  $B$  meson from observation of the other; however, the coherence also requires determination of the decay time of both mesons, rather than just one, in order to perform time-dependent  $CP$ -violation measurements.

For the measurement of branching fractions, the initial composition of the data sample must be known. The  $\Upsilon(4S)$  resonance decays predominantly to  $B^0\bar{B}^0$  and  $B^+B^-$ ; the current experimental upper limit for non- $B\bar{B}$  decays of the  $\Upsilon(4S)$  is less than 4% at the 95% confidence level (CL) [26]. The only known modes of this category are decays to lower  $\Upsilon$  states and a pion pair, observed with branching fractions of order  $10^{-4}$  [27]. The ratio  $f_+/f_0$  of the fractions of charged to neutral  $B$  productions from  $\Upsilon(4S)$  decays has been measured by CLEO, BaBar, and Belle in various ways. They typically use pairs of isospin-related decays of  $B^+$  and  $B^0$ , such that it can be assumed that  $\Gamma(B^+ \rightarrow x^+) = \Gamma(B^0 \rightarrow x^0)$ . In this way, the

ratio of the number of events observed in these modes is proportional to  $(f_+\tau_+)/ (f_0\tau_0)$  [28–31]. BaBar has also performed an independent measurement of  $f_0$  with a different method that does not require isospin symmetry or the value of the lifetime ratio, based on the number of events with one or two reconstructed  $B^0 \rightarrow D^{*-}\ell^+\nu$  decays [32]. The combined result, from the current average of  $\tau_+/\tau_0$ , is  $f_+/f_0 = 1.055 \pm 0.025$  [33]. Though the current  $2.2\sigma$  discrepancy with equal production of  $B^+B^-$  and  $B^0\bar{B}^0$  pairs is somewhat larger than previous averages, we still assume  $f_+/f_0 = 1$  in this mini-review except where explicitly stated otherwise. This assumption is also supported by the near equality of the  $B^+$  and  $B^0$  masses: our fit of CLEO, ARGUS, and CDF measurements yields  $m(B^0) = 5279.50 \pm 0.33$  MeV/ $c^2$ ,  $m(B^+) = 5279.13 \pm 0.31$  MeV/ $c^2$ , and  $m(B^0) - m(B^+) = 0.37 \pm 0.24$  MeV/ $c^2$ . The latest measurement from the LHCb agrees well with those and further improves precision [34].

CLEO and Belle have also collected some data at the  $\Upsilon(5S)$  resonance [35,36]. Belle has accumulated more than  $100 \text{ fb}^{-1}$  at this resonance. This resonance does not provide the simple final states of the  $\Upsilon(4S)$ : there are seven possible final states with a pair of non-strange  $B$  mesons and three with a pair of strange  $B$  mesons ( $B_s^*\bar{B}_s^*$ ,  $B_s^*\bar{B}_s$ , and  $B_s\bar{B}_s^*$ ). The fraction of events with a pair of  $B_s$  mesons over the total number of events with a pair of  $b$ -flavored hadrons has been measured to be  $f_s[\Upsilon(5S)] = 0.199 \pm 0.030$ , of which 90% is  $B_s^*\bar{B}_s^*$  events. A few branching fractions of the  $B_s$  have been measured in this way; if the precision of  $f_s$  were improved, they would become the most accurate. Belle has observed a few new  $B_s$  modes that are difficult to reconstruct in hadron colliders and the most precise mass measurement of the  $B_s^*$  meson has been obtained [36,37]. However, the small boost of  $B_s$  mesons produced in this way prevents resolution of their fast oscillations for time-dependent measurements; these are only accessible in hadron collisions or at the  $Z$  peak.

In high-energy collisions, the produced  $b$  or  $\bar{b}$  quarks can hadronize with different probabilities into the full spectrum of  $b$ -hadrons, either in their ground or excited states. Table 1

shows the measured fractions  $f_d$ ,  $f_u$ ,  $f_s$ , and  $f_{\text{baryon}}$  of  $B^0$ ,  $B^+$ ,  $B_s^0$ , and  $b$  baryons, respectively, in an unbiased sample of weakly decaying  $b$  hadrons produced at the  $Z$  resonance or in  $p\bar{p}$  collisions [33]. The results were obtained from a fit where the sum of the fractions were constrained to equal 1.0, neglecting production of  $B_c$  mesons. The observed yields of  $B_c$  mesons at the Tevatron [19] yields  $f_c = 0.2\%$ , in agreement with expectations [38], and well below the current experimental uncertainties in the other fractions.

**Table 1:** Fractions of weakly-decaying  $b$ -hadron species in  $Z \rightarrow b\bar{b}$  decay, in  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8$  TeV and combination of fractions in  $Z \rightarrow b\bar{b}$  decay at Tevatron,  $p\bar{p}$  and  $pp$  collisions at LHC.

$b$ hadron	Fraction at $Z$ [%]	Fraction at $p\bar{p}$ [%]	Combined [%]
$B^+, B^0$	$40.3 \pm 0.9$	$33.9 \pm 3.9$	$40.1 \pm 0.8$
$B_s$	$10.3 \pm 0.9$	$11.1 \pm 1.4$	$10.5 \pm 0.6$
$b$ baryons	$9.0 \pm 1.5$	$21.2 \pm 6.9$	$9.3 \pm 1.6$

The combined values assume identical hadronization in  $p\bar{p}$  collisions and in  $Z$  decay. These could in principle differ, because of the different momentum distributions of the  $b$ -quark in these processes; the sample used in the  $p\bar{p}$  measurements has momenta close to the  $b$  mass, rather than  $m_Z/2$ . A test of the agreement between production fractions may be given by comparison of values of the average time-integrated mixing probability parameter  $\bar{\chi} = f_d\chi_d + f_s\chi_s$  [8]. This is an important input in the determination of the world-averages of production fractions. The current measurements of  $\bar{\chi}$  from LEP and the Tevatron differ by  $1.8\sigma$  [33]. This slight discrepancy increases the uncertainty in the combined fractions in Table 1. It should be noted that the combination is not well defined as both the CDF and LHCb experiments observe a significant dependence of the  $\Lambda_b$  production fraction on transverse momentum. With the availability of large samples of  $b$ -flavored mesons and baryons at  $p\bar{p}$  colliders, the limited knowledge of these fractions has

become an important limiting factor in the determination of their branching fractions.

Excited  $B$ -meson states have been observed by CLEO, LEP, CUSB, D0, and CDF. The current world average of the  $B^*-B$  mass difference is  $45.78 \pm 0.35$  MeV/ $c^2$ . Evidence for  $B^{**}$  ( $L=1$ ) production has been initially obtained at LEP [39], as a broad resonance in the mass of an inclusively reconstructed bottom hadron candidate combined with a charged pion from the primary vertex. Detailed results from exclusive modes have been obtained at the Tevatron, allowing separation of the narrow states  $B_1$  and  $B_2^*$  and also a measurement of the  $B_2^*$  width [40].

Also the narrow  $B_s^{**}$  states, first sighted by OPAL as a single broad enhancement in the  $B^+K$  mass spectrum [41], have now been clearly observed and separately measured at the Tevatron [42]:  $M(B_{s1}) = 5829.4 \pm 0.7$  MeV/ $c^2$  (CDF) and  $M(B_{s2}^*) = 5839.7 \pm 0.7$  MeV/ $c^2$  (CDF),  $M(B_{s2}^*) = 5839.6 \pm 1.1 \pm 0.7$  MeV/ $c^2$  (D0).

Baryon states containing a  $b$  quark are labeled according to the same scheme used for non- $b$  baryons, with the addition of a  $b$  subscript [23]. For many years, the only well-established  $b$  baryon was the  $\Lambda_b^0$  (quark composition  $udb$ ), with only indirect evidence for  $\Xi_b$  ( $dsb$ ) production from LEP [43]. This situation has changed dramatically in the past few years due to the large samples being accumulated at the Tevatron and LHCb. Clear signals of four strongly-decaying baryon states,  $\Sigma_b^+$ ,  $\Sigma_b^{*+}$  ( $uub$ ),  $\Sigma_b^-$ ,  $\Sigma_b^{*-}$  ( $ddb$ ) have been obtained by CDF in  $\Lambda_b^0\pi^\pm$  final states [44]. The strange bottom baryon  $\Xi_b^\pm$  was observed in the exclusive mode  $\Xi_b^\pm \rightarrow J/\psi\Xi^\pm$  by D0 [45], and CDF [46]. More recently CDF has also observed the  $\Xi_b$  in the  $\Xi_c\pi$  final state [47]. The relative production of  $\Xi_b$  and  $\Lambda_b$  baryons has been found to be consistent with the  $B_s$  to  $B_d$  production ratio [45]. Observation of the doubly-strange bottom baryon  $\Omega_b^-$  has been published by both D0 [48] and CDF [49]. However the masses measured by the two experiments show a large discrepancy. The resolution appears to be provided by LHCb; they recently presented a preliminary measurement of the  $\Omega_b^-$  mass consistent with the CDF measurement [50]. Apart from the discrepancy



on  $\Omega_b^-$  mass, the masses of all these new baryons have been measured to a precision of a few  $\text{MeV}/c^2$ , and found to be in agreement with predictions from HQET.

**Lifetimes:** Precise lifetimes are key in extracting the weak parameters that are important for understanding the role of the CKM matrix in  $CP$  violation, such as the determination of  $V_{cb}$  and  $B_s\bar{B}_s$  mixing parameters. In the naive spectator model, the heavy quark can decay only via the external spectator mechanism, and thus, the lifetimes of all mesons and baryons containing  $b$  quarks would be equal. Non-spectator effects, such as the interference between contributing amplitudes, modify this simple picture and give rise to a lifetime hierarchy for  $b$ -flavored hadrons similar to the one in the charm sector. However, since the lifetime differences are expected to scale as  $1/m_Q^2$ , where  $m_Q$  is the mass of the heavy quark, the variations in the  $b$  system are expected to be only 10% or less [51]. We expect:

$$\tau(B^+) \geq \tau(B^0) \approx \tau(B_s) > \tau(\Lambda_b^0) \gg \tau(B_c^+). \quad (1)$$

For the  $B_c^+$ , both quarks decay weakly, so the lifetime is much shorter.

Measurements of the lifetimes of the different  $b$ -flavored hadrons thus provide a means to determine the importance of non-spectator mechanisms in the  $b$  sector. Over the past decade, the precision of silicon vertex detectors and the increasing availability of fully-reconstructed samples has resulted in much-reduced statistical and systematic uncertainties ( $\sim 1\%$ ). The averaging of precision results from different experiments is a complex task that requires careful treatment of correlated systematic uncertainties; the world averages given in Table 2 have been determined by the Heavy Flavor Averaging Group (HFAG) [33].

**Table 2:** Summary of inclusive and exclusive world-average  $b$ -hadron lifetime measurements. For the two  $B_s$  averages, see text below.

Particle	Lifetime [ps]
$B^+$	$1.641 \pm 0.008$
$B^0$	$1.519 \pm 0.007$
$B_s$ (flavor-specific)	$1.463 \pm 0.032$
$B_s$ ( $1/\Gamma_s$ )	$1.495 \pm 0.015$
$B_c^+$	$0.453 \pm 0.041$
$\Lambda_b^0$	$1.425 \pm 0.032$
$\Xi_b^-$	$1.56^{+0.27}_{-0.25}$
$\Omega_b^-$	$1.13^{+0.53}_{-0.40}$
$\Xi_b$ mixture	$1.49^{+0.19}_{-0.18}$
$b$ -baryon mixture	$1.382 \pm 0.029$
$b$ -hadron mixture	$1.568 \pm 0.009$

The short  $B_c^+$  lifetime is in good agreement with predictions [52]. For precision comparisons with theory, lifetime ratios are more sensitive. Experimentally we find:

$$\frac{\tau_{B^+}}{\tau_{B^0}} = 1.079 \pm 0.007, \quad \frac{\tau_{B_s}}{\tau_{B^0}} = 0.984 \pm 0.011,$$

$$\frac{\tau_{\Lambda_b}}{\tau_{B^0}} = 0.938 \pm 0.022,$$

while theory makes the following predictions [51,53]

$$\frac{\tau_{B^+}}{\tau_{B^0}} = 1.06 \pm 0.02, \quad \frac{\tau_{B_s}}{\tau_{B^0}} = 1.00 \pm 0.01, \quad \frac{\tau_{\Lambda_b}}{\tau_{B^0}} = 0.88 \pm 0.05.$$

The ratio of  $B^+$  to  $B^0$  lifetimes has a precision of better than 1%, and is significantly different from 1.0, in agreement with predictions [51]. The ratio of  $B_s$  to  $B^0$  lifetimes is expected to be very close to 1.0; while there used to be mild tension between experiment and theory, the discrepancy is disappearing with newer measurements with large samples of fully reconstructed  $B_s$  decays [54]. The  $\Lambda_b$  lifetime has a history of discrepancies. Predictions were higher than data before the introduction of higher-order effects lowered them. The precision of the measurements has recently been improved by more than a factor of two by two CDF measurements [57,58]. The

measurements are in marginal agreement with each other and previous measurements; the new world average is somewhat larger than the theoretical predictions. The most significant discrepancy comes from the CDF measurement in  $\Lambda_b \rightarrow J/\psi\Lambda$  channel which differs by  $3.3\sigma$  from the average of all other measurements. With more data available at both D0 and CDF and large samples available at LHCb, new results will hopefully resolve this discrepancy in the near future.

Neutral  $B$  mesons are two-component systems similar to neutral kaons, with a light (L) and a heavy (H) mass eigenstate, and independent decay widths  $\Gamma_L$  and  $\Gamma_H$ . The SM predicts a non-zero width difference  $\Delta\Gamma = \Gamma_L - \Gamma_H > 0$  for both  $B_s$  and  $B_d$ . For  $B_d$ ,  $\Delta\Gamma_d/\Gamma_d$  is expected to be  $\sim 0.2\%$ . Analysis of BaBar and DELPHI data on  $CP$ -specific modes of the  $B^0$  yield a combined result:  $\Delta\Gamma_d/\Gamma_d = 0.015 \pm 0.018$  [33]. The issue is much more interesting for the  $B_s$ , since the SM expectation for  $\Delta\Gamma_s/\Gamma_s$  is of order 10%. This potentially non-negligible difference requires care when defining the  $B_s$  lifetime. As indicated in Table 2, two different lifetimes are defined for the  $B_s$  meson: one is defined as  $1/\Gamma_s$ , where  $\Gamma_s$  is the average width of the two mass eigenstates  $(\Gamma_L + \Gamma_H)/2$ ; the other is obtained from “flavor-specific” (*e.g.*, semileptonic) decays and depends both on  $\Gamma_s$  and  $\Delta\Gamma_s$ . Experimentally, the quantity  $\Delta\Gamma_s$  can be accessed by measuring lifetimes in decays into  $CP$  eigenstates, which in the standard model are expected to be close approximations to the mass eigenstates. This has been done with the  $J/\psi\phi$  mode, where the two  $CP$  eigenstates are distinguished by angular distributions, and in  $B_s \rightarrow K^+K^-$  or  $B_s \rightarrow J/\psi f_0(980)$  which are  $CP$ -eigenstates. The current experimental information is dominated by measurements on the  $J/\psi\phi$  mode performed by CDF, D0 and LHCb experiments. By appropriately combining all published measurements of  $J/\psi\phi$  lifetimes and flavor-specific lifetimes, the HFAG group obtains a world-average  $\Delta\Gamma_s/\Gamma_s = 0.092^{+0.051}_{-0.054}$  [33], which is compatible with zero; the latest theoretical predictions yield  $\Delta\Gamma_s/\Gamma_s = 0.133 \pm 0.032$  [59], in agreement with measurements within the large uncertainties on both. From the theoretical point of view, the best quantity to use is  $\Delta\Gamma_s/\Delta M_s$ , which is

much less affected by hadronic uncertainties [59]. Exploiting the very accurate measurement of  $\Delta M_s$  now available [60], this can be turned into a SM prediction with an uncertainty of only 20%:  $\Delta\Gamma_s/\Gamma_s = 0.137 \pm 0.027$ . This is likely to be of importance in future comparisons, as the experimental precision improves with the growth of Tevatron samples. Further improvements are coming from lifetime measurements in the  $CP$ -eigenstates such as  $B_s \rightarrow K^+K^-$  [61] and  $B_s \rightarrow J/\psi f_0(980)$  [62], and alternative (model-dependent) determinations via the  $B_s \rightarrow D_s^{(*)+} D_s^{(*)-}$  branching fraction [63].

The width difference  $\Delta\Gamma_s$  is connected to the  $B_s$  mixing phase  $\phi_s$  by  $\Delta\Gamma_s = \Gamma_{12} \cos \phi_s$ , where  $\Gamma_{12}$  is the off-diagonal term of the decay matrix [6,8,59]. The early measurements by CDF [64] and D0 [65] have produced CL contours in the  $(\phi_s, \Delta\Gamma)$  plane, and both observe a mild deviation, in the same direction, from the expectation of the Standard model of the phase  $\phi_s$  near  $\Delta\Gamma = 0$ . The possibility of a large value of  $\phi_s$  has attracted significant interest, as it would be very clean evidence for the existence of new sources of  $CP$  violation beyond the standard model. However the latest measurements from CDF [66], D0 [67] and LHCb [68], which provide significant improvements over initial measurements, show good agreement with the standard model. The LHCb experiment also used the decay  $B_s \rightarrow J/\psi f_0(980)$  to measure  $\phi_s$  [69]. While this measurement is not as precise as the one from  $B_s \rightarrow J/\psi\phi$ , it does not require analysis of angular distributions, which simplifies the analysis. It should be noted that all above measurements have a two-fold ambiguity in their results. We can resolve this ambiguity using the interference between the decays to  $J/\psi\phi$  and  $J/\psi K^+K^-$ , where  $K^+K^-$  is in relative S-wave state. This has been used by LHCb experiment to determine the sign of the  $\Delta\Gamma_s$  to be positive [70].

***B meson decay properties:*** Semileptonic  $B$  decays  $B \rightarrow X_c \ell \nu$  and  $B \rightarrow X_u \ell \nu$  provide an excellent way to measure the magnitude of the CKM elements  $|V_{cb}|$  and  $|V_{ub}|$  respectively, because the strong interaction effects are much simplified due to the two leptons in the final state. Both exclusive and inclusive decays can be used, and the nature of uncertainties are

quite complementary. For exclusive decay analysis, knowledge of the form factors for the exclusive hadronic system  $X_{c(u)}$  is required. For inclusive analysis, it is usually necessary to restrict the available phase-space of the decay products to suppress backgrounds; subsequently uncertainties are introduced in the extrapolation to the full phase-space. Moreover, restriction to a small corner of the phase-space may result in breakdown of the operator-product expansion scheme, thus making theoretical calculations unreliable. A more detailed discussion of  $B$  semileptonic decays and the extraction of  $|V_{cb}|$  and  $|V_{ub}|$  is given elsewhere in this *Review* [9].

On the other hand, hadronic decays of  $B$  are complicated because of strong interaction effects caused by the surrounding cloud of light quarks and gluons. While this complicates the extraction of CKM matrix elements, it also provides a great opportunity to study perturbative and non-perturbative QCD, hadronization, and Final State Interaction (FSI) effects. Pure-penguin decays were first established by the observation of  $B \rightarrow K^* \gamma$  [71]. Some observed decay modes such as  $B^0 \rightarrow D_s^- K^+$ , may be interpreted as evidence of a  $W$ -exchange process [72]. The evidence for the decay  $B^+ \rightarrow \tau^+ \nu$  from Belle [73] and BaBar [74] is the first sign of a pure annihilation decay. There is growing evidence that penguin annihilation processes may be important in decays with two vector mesons in the final state [75].

***Hadronic decays:*** Most of the hadronic  $B$  decays involve  $b \rightarrow c$  transition at the quark level, resulting in a charmed hadron or charmonium in the final state. Other types of hadronic decays are very rare and will be discussed separately in the next section. The experimental results on hadronic  $B$  decays have steadily improved over the past few years, and the measurements have reached sufficient precision to challenge our understanding of the dynamics of these decays. With the good neutral particle detection and hadron identification capabilities of  $B$ -factory detectors, a substantial fraction of hadronic  $B$  decay events can be fully reconstructed. Because of the kinematic constraint of  $\Upsilon(4S)$ , the energy sum of the final-state particles of a  $B$  meson decay is always equal to one half of the

total energy in the center of mass frame. As a result, the two variables,  $\Delta E$  (energy difference) and  $M_B$  ( $B$  candidate mass with a beam-energy constraint) are very effective for suppressing combinatorial background both from  $\Upsilon(4S)$  and  $e^+e^- \rightarrow q\bar{q}$  continuum events. In particular, the energy-constraint in  $M_B$  improves the signal resolution by almost an order of magnitude.

The kinematically clean environment of  $B$  meson decays provides an excellent opportunity to search for new states. For instance, quark-level  $b \rightarrow c\bar{c}s$  decays have been used to search for new charmonium and charm-strange mesons and study their properties in detail. In 2003, BaBar discovered a new narrow charm-strange state  $D_{sJ}^*(2317)$  [76], and CLEO observed a similar state  $D_{sJ}(2460)$  [77]. The properties of these new states were studied in the  $B$  meson decays,  $B \rightarrow DD_{sJ}^*(2317)$  and  $B \rightarrow DD_{sJ}(2460)$  by Belle [78]. Further studies of  $D_{sJ}^{(*)}$  meson production in  $B$  decays have been made by Belle [79] and BaBar [80]. Now these charm-strange meson states are identified as  $D_{s0}^*(2317)$  and  $D_{s1}(2460)$ , respectively.

More recently, Belle observed a new  $D_{sJ}$  meson produced in  $B^+ \rightarrow \bar{D}^0 D_{sJ} \rightarrow \bar{D}^0 D^0 K^+$  [81]. Combined with a subsequent measurement by BaBar [82], the mass and width of this state are determined to be  $2709_{-6}^{+9}$  MeV/ $c^2$  and  $125 \pm 30$  MeV, respectively. An analysis of the helicity angle distribution determines its spin-parity to be  $1^-$ .

A variety of exotic particles have been discovered in  $B$  decays. Belle found the  $X(3872)$  state [83], which is confirmed by CDF [84] and BaBar [85]. Analyzing their full  $\Upsilon(4S)$  data sample, Belle finds a new upper limit on the width of  $X(3872)$  to be  $\Gamma_{X(3872)} < 1.2$  MeV [86], improving on the existing limit by nearly a factor of 2. Radiative decays of  $X(3872)$  can play a crucial role in understanding the nature of the particle. For example, in the molecular model the decay of  $X(3872)$  to  $\psi'\gamma$  is expected to be highly suppressed in comparison to the decay to  $J/\psi\gamma$  [87]. BaBar has seen the evidence for the decay to  $J/\psi\gamma$  [88]. The ratio  $R \equiv \mathcal{B}(X(3872) \rightarrow \psi'\gamma)/\mathcal{B}(X(3872) \rightarrow J/\psi\gamma)$  is measured to be  $3.4 \pm 1.4$  by BaBar [89], while Belle obtains  $R < 2.1$  at 90% CL [90].

Belle has observed a near-threshold enhancement in the  $J/\psi\omega$  invariant mass for  $B \rightarrow J/\psi\omega K$  decays [91]. BaBar has studied  $B \rightarrow J/\psi\pi^+\pi^-K$ , finding an excess of  $J/\psi\pi^+\pi^-$  events with a mass just above  $4.2 \text{ GeV}/c^2$ ; this is consistent with the  $Y(4260)$  that was observed by BaBar in ISR (Initial State Radiation) events [93]. A Belle study of  $B \rightarrow \psi'K\pi^\pm$  [94] finds a state called  $X(4430)^\pm$  that decays to  $\psi'\pi^\pm$ . Since it is charged, it could not be a charmonium state. This state was searched for by BaBar with similar sensitivity but was not found [95]. In a Dalitz plot analysis of  $\bar{B}^0 \rightarrow \chi_{c1}K^-\pi^+$ , Belle has observed two resonance-like structures in the  $\chi_{c1}\pi^+$  mass distribution [96], labelled as  $X(4050)^\pm$  and  $X(4250)^\pm$  in this *Review*, while no evidence is found by BaBar in a search with similar sensitivity [97].

The hadronic decays  $\bar{B}^0 \rightarrow D^{(*)0}h^0$ , where  $h^0$  stands for light neutral mesons such as  $\pi^0, \eta^{(\prime)}, \rho^0, \omega$ , proceed through color-suppressed diagrams, hence they provide useful tests on the factorization models. Both Belle and BaBar have made comprehensive measurements of such color-suppressed hadronic decays of  $\bar{B}^0$  [98].

Information on  $B_s$  and  $\Lambda_b$  decays is limited, though improving with recent studies of large samples at the Tevatron and LHC experiments. Recent additions are decays of  $B_s \rightarrow J/\psi f_0(980)$  [62,99],  $B_s \rightarrow J/\psi f_2'(1525)$  [100], and  $\Lambda_b \rightarrow \Lambda_c\pi^+\pi^-\pi^-$  [101]. For the later, not only the total rate is measured, but also structure involving decays through excited  $\Lambda_c$  and  $\Sigma_c$  baryons.

There have been hundreds of publications on hadronic  $B$  decays to open-charm and charmonium final states mostly from the  $B$ -factory experiments. These results are nicely summarized in a recent report by HFAG [33].

**Rare  $B$  decays:** All  $B$ -meson decays that do not occur through the  $b \rightarrow c$  transition are usually called rare  $B$  decays. These include both semileptonic and hadronic  $b \rightarrow u$  decays that are suppressed at leading order by the small CKM matrix element  $V_{ub}$ , as well as higher-order  $b \rightarrow s(d)$  processes such as electroweak and gluonic penguin decays.

Charmless  $B$  meson decays into two-body hadronic final states such as  $B \rightarrow \pi\pi$  and  $K\pi$  are experimentally clean, and provide good opportunities to probe new physics and search for indirect and direct  $CP$  violations. Since the final state particles in these decays tend to have larger momenta than average  $B$  decay products, the event environment is cleaner than for  $b \rightarrow c$  decays. Branching fractions are typically around  $10^{-5}$ . Over the past decade, many such modes have been observed by BaBar, Belle, and CLEO. More recently, comparable samples of the modes with all charged final particles have been reconstructed in  $p\bar{p}$  collisions by CDF by triggering on the impact parameter of the charged tracks. This has also allowed observation of charmless decays of the  $B_s$ , in final states such as  $\phi\phi$  [102],  $K^+K^-$  [103], and  $K^-\pi^+$  [104], and of charmless decays of the  $\Lambda_b^0$  baryon [104]. Charmless  $B_s$  modes are related to corresponding  $B^0$  modes by U-spin symmetry, and are determined by similar amplitudes. Combining the observables from  $B_s$  and  $B^0$  modes is a further way of eliminating hadronic uncertainties and extracting relevant CKM information [105].

Because of relatively high-momenta for final state particles, the dominant source of background in  $e^+e^-$  collisions is  $q\bar{q}$  continuum events; sophisticated background suppression techniques exploiting event shape variables are essential for these analyses. In hadron collisions, the dominant background comes from QCD or partially reconstructed heavy flavors, and is similarly suppressed by a combination of kinematic and isolation requirements. The results are in general consistent among the experiments.

BaBar [106] and Belle [107] have observed the decays  $B^+ \rightarrow \bar{K}^0 K^+$  and  $B^0 \rightarrow K^0 \bar{K}^0$ . The world-average branching fractions are  $\mathcal{B}(B^0 \rightarrow K^0 \bar{K}^0) = (0.96_{-0.18}^{+0.20}) \times 10^{-6}$  and  $\mathcal{B}(B^+ \rightarrow \bar{K}^0 K^+) = (1.36 \pm 0.27) \times 10^{-6}$ . These are the first observations of hadronic  $b \rightarrow d$  transitions, with significance  $> 5\sigma$  for all four measurements.  $CP$  asymmetries have even been measured for these modes, though with large errors.

Most rare decay modes including  $B^0 \rightarrow K^+ \pi^-$  have contributions from both  $b \rightarrow u$  tree and  $b \rightarrow sg$  penguin processes.



If the size of the two contributions are comparable, the interference between them may result in direct  $CP$  violation, seen experimentally as a charge asymmetry in the decay rate measurement. BaBar [108], Belle [109], and CDF [103] have measured the direct  $CP$  violating asymmetry in  $B^0 \rightarrow K^+\pi^-$  decays. The BaBar and Belle measurements constitute observation of direct  $CP$  violation with a significance of more than  $5\sigma$ . The world average for this quantity is now rather precise,  $-0.098 \pm 0.013$ . There are sum rules [110] that relate the decay rates and decay-rate asymmetries between the four  $K\pi$  charge states. The experimental measurements of the other three modes are not yet precise enough to test these sum rules.

There is now evidence for direct  $CP$  violation in three other decays:  $B^+ \rightarrow \rho^0 K^+$  [111],  $B^+ \rightarrow \eta K^+$  [112], and  $B^0 \rightarrow \eta K^{*0}$  [113]. The significance is typically 3–4 $\sigma$ , though the significance for the  $B^+ \rightarrow \eta K^+$  decay is now nearly  $5\sigma$  with the recent Belle measurement [112]. In at least the first two cases, a large direct  $CP$  violation might be expected since the penguin amplitude is suppressed so the tree and penguin amplitudes may have comparable magnitudes.

The decay  $B^0 \rightarrow \pi^+\pi^-$  can be used to extract the CKM angle  $\alpha$ . This is complicated by the presence of significant contributions from penguin diagrams. An isospin analysis [114] can be used to untangle the penguin complications. The decay  $B^0 \rightarrow \pi^0\pi^0$ , which is now measured by both BaBar and Belle, is crucial in this analysis. Unfortunately the amount of penguin pollution in the  $B \rightarrow \pi\pi$  system is rather large. In the past few years, measurements in the  $B^0 \rightarrow \rho\rho$  system have produced more precise values of  $\alpha$ , since penguin amplitudes are generally smaller for decays with vector mesons. An important ingredient in the analysis is the  $B^0 \rightarrow \rho^0\rho^0$  branching fraction. The average of measurements from BaBar and Belle BaBar [115] yields a branching fraction of  $(0.73 \pm 0.28) \times 10^{-6}$ . This is only 3% of the  $\rho^+\rho^-$  branching fraction, much smaller than the corresponding ratio in the  $\pi\pi$  system.

The decay  $B \rightarrow a_1\pi$  has been seen by BaBar. An analysis of the time evolution of this decay [116] together with measurements of other related decays has been used to measure

the CKM angle  $\alpha$  [117] in agreement with the more precise measurements from the  $\rho\rho$  system.

Since  $B \rightarrow \rho\rho$  has two vector mesons in the final state, the  $CP$  eigenvalue of the final state depends on the longitudinal polarization fraction  $f_L$  for the decay. Therefore, a measurement of  $f_L$  is needed to extract the CKM angle  $\alpha$ . Both BaBar and Belle have measured  $f_L$  for the decays  $\rho^+\rho^-$  and  $\rho^+\rho^0$  and in both cases the measurements show  $f_L > 0.9$ , making a complete angular analysis unnecessary.

By analyzing the angular distributions of the  $B$  decays to two vector mesons, we can learn a lot about both weak- and strong-interaction dynamics in  $B$  decays. Decays that are penguin-dominated surprisingly have values of  $f_L$  near 0.5. The list of such decays has now grown to include  $B \rightarrow \phi K^*$ ,  $B \rightarrow \rho K^*$ , and  $B \rightarrow \omega K^*$ . The reasons for this "polarization puzzle" are not fully understood. A detailed description of the angular analysis of  $B$  decays to two vector mesons can be found in a separate mini-review [118] in this *Review*.

There has been substantial progress in measurements of many other rare- $B$  decays. The decay  $B \rightarrow \eta' K$  stood out as the largest rare- $B$  decay for many years. The reasons for the large rate are now largely understood [14,119]. However, there are now measurements of several 3-body or quasi-3-body modes with similarly large branching fractions. States seen so far include  $K\pi\pi$  (three charge states) [120],  $KKK$  (four charge states) [121], and  $K^*\pi\pi$  (two charged states) [122]. Many of these analyses now include Dalitz plot treatments with many intermediate resonances. There has also been an observation of the decay  $B^+ \rightarrow K^+K^-\pi^+$  by BaBar [123], noteworthy because an even number of kaons is typically indicative of suppressed  $b \rightarrow d$  transitions as discussed above.

Belle [73] and BaBar [74] have found evidence for  $B^+ \rightarrow \tau^+\nu$ ; the average branching fraction, with a significance of nearly  $5\sigma$  is  $(165 \pm 34) \times 10^{-6}$ . This is somewhat larger than, though consistent with, the value expected in the SM. This is the first observation of a pure annihilation decay. A substantial region of parameter space of charged Higgs mass vs.  $\tan\beta$  is excluded by the measurements of this mode.

**Electroweak penguin decays:** More than a decade has passed since the CLEO experiment first observed an exclusive radiative  $b \rightarrow s\gamma$  transition,  $B \rightarrow K^*(892)\gamma$  [71], thus providing the first evidence for the one-loop FCNC electromagnetic penguin decay. Using much larger data samples, both Belle and BaBar have updated this analysis [124] with an average branching fraction  $\mathcal{B}(B^0 \rightarrow K^{*0}\gamma) = (43.3 \pm 1.5) \times 10^{-6}$ , and have added several new decay modes such as  $B \rightarrow K_1\gamma$ ,  $K_2^*(1430)\gamma$ , *etc.* [125]. With a sample of  $24 \text{ fb}^{-1}$  at  $\Upsilon(5S)$ , Belle observed the radiative penguin decay of  $B_s \rightarrow \phi\gamma$  with a branching fraction  $(57_{-19}^{+22}) \times 10^{-6}$  [126].

Compared to  $b \rightarrow s\gamma$ , the  $b \rightarrow d\gamma$  transitions such as  $B \rightarrow \rho\gamma$ , are suppressed by the small CKM element  $V_{td}$ . Both Belle and BaBar have observed these decays [17,18]. The world average  $\mathcal{B}(B \rightarrow (\rho, \omega)\gamma) = (1.28 \pm 0.21) \times 10^{-6}$ . This can be used to calculate  $|V_{td}/V_{ts}|$  [127]; the measured values are  $0.233_{-0.032}^{+0.033}$  from BaBar [18] and  $0.195_{-0.024}^{+0.025}$  from Belle [17].

The observed radiative penguin branching fractions can constrain a large class of SM extensions [128]. However, due to the uncertainties in the hadronization, only the inclusive  $b \rightarrow s\gamma$  rate can be reliably compared with theoretical calculations. This rate can be measured from the endpoint of the inclusive photon spectrum in  $B$  decay. By combining the measurements of  $B \rightarrow X_s\gamma$  from CLEO, BaBar, and Belle experiments [129,130], HFAG obtains the new average:  $\mathcal{B}(B \rightarrow X_s\gamma) = (3.55 \pm 0.24 \pm 0.09) \times 10^{-4}$  [33] for  $E_\gamma \geq 1.6 \text{ GeV}$ . Consistent results have been reported by ALEPH for inclusive  $b$ -hadrons produced at the  $Z$ . The measured branching fraction can be compared to theoretical calculations. Recent calculations of  $\mathcal{B}(b \rightarrow s\gamma)$  at NNLO level predict the values of  $(3.15 \pm 0.23) \times 10^{-4}$  [131] and  $(2.98 \pm 0.26) \times 10^{-4}$  [132], where the latter is calculated requiring  $E_\gamma \geq 1.6 \text{ GeV}$ .

The  $CP$  asymmetry in  $b \rightarrow s\gamma$  is extensively studied theoretically both in the SM and beyond [133]. According to the SM, the  $CP$  asymmetry in  $b \rightarrow s\gamma$  is smaller than 1%, but some non-SM models allow significantly larger  $CP$  asymmetry ( $\sim 10\%$ ) without altering the inclusive branching fraction. The current world average is  $A_{CP} = -0.012 \pm 0.028$ , again dominated

by BaBar and Belle [134]. In addition to the  $CP$  asymmetry, BaBar also measured the isospin asymmetry  $\Delta_{0-} = 0.06 \pm 0.17$  in  $b \rightarrow s\gamma$  by measuring the companion  $B$  with full reconstruction in the hadronic decay modes [135].

In addition, all three experiments have measured the inclusive photon energy spectrum for  $b \rightarrow s\gamma$ , and by analyzing the shape of the spectrum they obtain the first and second moments for photon energies. Belle has measured these moments covering the widest range in the photon energy ( $1.7 < E_\gamma < 2.8$  GeV) [130]. These results can be used to extract non-perturbative HQET parameters that are needed for precise determination of the CKM matrix element  $V_{ub}$ .

Additional information on FCNC processes can be obtained from  $B \rightarrow X_s \ell^+ \ell^-$  decays, which are mediated by electroweak penguin and  $W$ -box diagrams. Their branching fractions have been measured by Belle [136], BaBar [137], and CDF [138]. Average branching fractions over all charged and neutral modes have been determined from BaBar and Belle data for  $B \rightarrow K \ell^+ \ell^-$ :  $(0.45 \pm 0.04) \times 10^{-6}$  and for  $B \rightarrow K^*(892) \ell^+ \ell^-$ :  $(1.08 \pm 0.11) \times 10^{-6}$ , consistent with the SM expectation.  $B$ -factory experiments also measured the branching fractions for inclusive  $B \rightarrow X_s \ell^+ \ell^-$  decays [139], with an average of  $(3.66_{-0.77}^{+0.76}) \times 10^{-6}$  [140]. Recently corresponding decays of  $B_s$  and  $\Lambda_b$  were observed [138,141]. Branching fraction for the decay  $B_s \rightarrow \phi \mu^+ \mu^-$  is measured to be  $(1.47 \pm 0.24 \pm 0.46) \times 10^{-6}$  and for the decay  $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$  to be  $(1.73 \pm 0.42 \pm 0.55) \times 10^{-6}$ . Excitement was generated by measurements of forward-backward asymmetry in  $B \rightarrow K^*(892) \ell^+ \ell^-$  decays, which exhibited mild tension with the standard model in earlier measurements. The most recent measurements by CDF [142] and LHCb [143] agree with standard model, suggesting that the earlier discrepancy was mainly due to statistical fluctuations.

Finally the decays  $B_{(s)}^0 \rightarrow e^+ e^-$  and  $\mu^+ \mu^-$  are interesting since they only proceed at second order in weak interactions in the SM, but may have large contributions from supersymmetric loops, proportional to  $(\tan\beta)^6$ . Experiments at Tevatron,  $B$ -factories and now also LHC have obtained results that exclude a portion of the region allowed by SUSY models. The most

stringent limits in these modes are obtained by LHCb. The limits in the  $\mu^+\mu^-$  mode are:  $< 1.4 \times 10^{-8}$  and  $< 3.2 \times 10^{-9}$  at 95% confidence level, respectively, for  $B_s$  and  $B^0$  [144]. For the  $B_s$  mode, the result is about factor of five above SM predictions [145]. It should be noted, that the most recent search by CDF observes an excess above expected background [146]. While the branching fraction for decay  $B_s \rightarrow \mu^+\mu^-$  of  $(1.8_{-0.9}^{+1.1}) \times 10^{-8}$  is extracted, CDF concludes that most plausible explanation for the excess is a statistical fluctuation. The limits for the  $e^+e^-$  modes are:  $< 2.8 \times 10^{-7}$  and  $< 8.3 \times 10^{-8}$ , respectively, for  $B_s$  and  $B^0$  [147]. There are also limits for lepton flavor-violating channels  $B_{(s)}^0 \rightarrow e^+\mu^-$ , which are around  $10^{-7}$  [147].

**Summary and Outlook:** The study of  $B$  mesons continues to be one of the most productive fields in particle physics. With the two asymmetric  $B$ -factory experiments Belle and BaBar, we now have a combined data sample of well over  $1 \text{ ab}^{-1}$ .  $CP$  violation has been firmly established in many decays of  $B$  mesons. Evidence for direct  $CP$  violation has been observed. Many rare decays resulting from hadronic  $b \rightarrow u$  transitions and  $b \rightarrow s(d)$  penguin decays have been observed, and the emerging pattern is still full of surprises. Despite the remarkable successes of the  $B$ -factory experiments, many fundamental questions in the flavor sector remain unanswered.

At Fermilab, CDF and D0 each has accumulated about  $10 \text{ fb}^{-1}$ , which is the equivalent of about  $10^{12}$   $b$ -hadrons produced. In spite of the low trigger efficiency of hadronic experiments, a selection of modes have been reconstructed in large quantities, giving a start to a program of studies on  $B_s$  and  $b$ -flavored baryons, in which a first major step has been the determination of the  $B_s$  oscillation frequency.

As Tevatron and  $B$ -factories stop their taking data, the new experiments at the LHC have become very active. The LHC accelerator performed very well in 2011. The general purpose experiments ATLAS and CMS collected about  $5 \text{ fb}^{-1}$  while LHCb collected about  $1 \text{ fb}^{-1}$ . LHCb, which is almost fully dedicated to studies of  $b$ - and  $c$ -hadrons, has a very large data sample. Of particular note is the sensitivity of the LHC

experiments for the decay  $B_s \rightarrow \mu^+ \mu^-$  which is expected to approach the standard model level in 2012.

In addition, two projects for next generation high-luminosity  $B$ -factories at KEK and Frascati are approved. Their aim to increase samples to  $\sim 50 \text{ ab}^{-1}$  will make it possible to explore the indirect evidence of new physics beyond the SM in the heavy-flavor particles ( $b$ ,  $c$ , and  $\tau$ ), in a way that is complementary to the LHC.

These experiments promise a rich spectrum of rare and precise measurements that have the potential to fundamentally affect our understanding of the SM and  $CP$ -violating phenomena.

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5. Currently two different notations ( $\phi_1, \phi_2, \phi_3$ ) and ( $\alpha, \beta, \gamma$ ) are used in the literature for CKM unitarity angles. In this mini-review, we use the latter notation following the other mini-reviews in this *Review*. The two notations are related by  $\phi_1 = \beta$ ,  $\phi_2 = \alpha$  and  $\phi_3 = \gamma$ .
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7. See the “CKM Quark Mixing Matrix,” by A. Cecucci, Z. Ligeti, and Y. Sakai, in this *Review*.
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