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# See the related review(s): Z Boson

### Z MASS

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). The fit is performed using the Z mass and width, the Z hadronic pole cross section, the ratios of hadronic to leptonic partial widths, and the Z pole forward-backward lepton asymmetries. This set is believed to be most free of correlations.

The Z-boson mass listed here corresponds to the mass parameter in a Breit-Wigner distribution with mass dependent width. The value is 34 MeV greater than the real part of the position of the pole (in the energy-squared plane) in the Z-boson propagator. Also the LEP experiments have generally assumed a fixed value of the  $\gamma-Z$  interferences term based on the standard model. Keeping this term as free parameter leads to a somewhat larger error on the fitted Z mass. See ACCIARRI 00Q and ABBIENDI 04G for a detailed investigation of both these issues.

VALUE (GeV)	EVTS	DOCUMENT ID	TECN	COMMENT				
91.1876±0.0021 OUR FIT								
$91.1852 \pm 0.0030$	4.57M	<sup>1</sup> ABBIENDI 01	1A OPAL	$E_{\rm cm}^{ee}=$ 88–94 GeV				
$91.1863 \pm 0.0028$	4.08M	<sup>2</sup> ABREU 00	OF DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV				
$91.1898 \pm 0.0031$	3.96M	<sup>3</sup> ACCIARRI 00	OC L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV				
$91.1885 \pm 0.0031$	4.57M	<sup>4</sup> BARATE 00	OC ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV				
• • • We do not use t	he following	data for averages, fits,	limits, etc.	• • •				
$91.084 \pm 0.107$		<sup>5</sup> ANDREEV 18	8a H1	e <sup>±</sup> p				
$91.1872 \pm 0.0033$		<sup>6</sup> ABBIENDI 04	4g OPAL	$E_{cm}^{ee} = LEP1 +$				
		7		130–209 GeV				
$91.272 \pm 0.032 \pm 0.032$	33	<sup>7</sup> ACHARD 04	4C L3	$E_{Cm}^{ee} = 183209 \; GeV$				
$91.1875 \pm 0.0039$	3.97M	<sup>8</sup> ACCIARRI 00	0Q L3	$E_{cm}^{ee} = LEP1 +$				
		2		130–189 GeV				
$91.151 \pm 0.008$		<sup>9</sup> MIYABAYASHI 9	5 TOPZ	$E_{\rm cm}^{ee}$ = 57.8 GeV				
$91.74 \pm 0.28 \pm 0.93$	3 156	<sup>10</sup> ALITTI 92	2b UA2	E <sup>pp</sup> <sub>cm</sub> = 630 GeV				
90.9 ±0.3 ±0.2	188	<sup>11</sup> ABE 89	9c CDF	$E^{p\overline{p}}_{cm}$ = 1.8 TeV				
$91.14 \pm 0.12$	480	<sup>12</sup> ABRAMS 89	9b MRK2	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89–93 GeV				
93.1 ±1.0 ±3.0	24	<sup>13</sup> ALBAJAR 89	9 UA1	$E_{\rm cm}^{p\overline{p}}$ = 546,630 GeV				

<sup>1</sup> ABBIENDI 01A error includes approximately 2.3 MeV due to statistics and 1.8 MeV due to LEP energy uncertainty.

 $^{2}$  The error includes 1.6 MeV due to LEP energy uncertainty.

<sup>3</sup>The error includes 1.8 MeV due to LEP energy uncertainty.

<sup>4</sup> BARATE 00C error includes approximately 2.4 MeV due to statistics, 0.2 MeV due to experimental systematics, and 1.7 MeV due to LEP energy uncertainty.

- <sup>5</sup> ANDREEV 18A obtain this result in a combined electroweak and QCD analysis using all deep-inelastic  $e^+p$  and  $e^-p$  neutral current and charged current scattering cross sections published by the H1 Collaboration, including data with longitudinally polarized lepton beams.
- <sup>6</sup> ABBIENDI 04G obtain this result using the S-matrix formalism for a combined fit to their cross section and asymmetry data at the Z peak and their data at 130-209 GeV. The authors have corrected the measurement for the 34 MeV shift with respect to the Breit-Wigner fits.
- <sup>7</sup> ACHARD 04C select  $e^+e^- \rightarrow Z\gamma$  events with hard initial-state radiation. Z decays to  $q \overline{q}$  and muon pairs are considered. The fit results obtained in the two samples are found consistent to each other and combined considering the uncertainty due to ISR modelling as fully correlated.
- <sup>8</sup> ACCIARRI 00Q interpret the *s*-dependence of the cross sections and lepton forwardbackward asymmetries in the framework of the S-matrix formalism. They fit to their cross section and asymmetry data at high energies, using the results of S-matrix fits to Z-peak data (ACCIARRI 00C) as constraints. The 130–189 GeV data constrains the  $\gamma/Z$ interference term. The authors have corrected the measurement for the 34.1 MeV shift with respect to the Breit-Wigner fits. The error contains a contribution of  $\pm 2.3$  MeV due to the uncertainty on the  $\gamma Z$  interference.
- <sup>9</sup> MIYABAYASHI 95 combine their low energy total hadronic cross-section measurement with the ACTON 93D data and perform a fit using an S-matrix formalism. As expected, this result is below the mass values obtained with the standard Breit-Wigner parametrization.
- <sup>10</sup> Enters fit through W/Z mass ratio given in the W Particle Listings. The ALITTI 92B systematic error (±0.93) has two contributions: one (±0.92) cancels in  $m_W/m_Z$  and one (±0.12) is noncancelling. These were added in quadrature.

 $^{11}$  First error of ABE 89 is combination of statistical and systematic contributions; second is mass scale uncertainty.

 $^{12}$  ABRAMS 89B uncertainty includes 35 MeV due to the absolute energy measurement.

 $^{13}$ ALBAJAR 89 result is from a total sample of 33  $Z 
ightarrow e^+e^-$  events.

			,				
VALUE	E (GeV)		EVTS	DOCUMENT ID		TECN	COMMENT
2.495	$52 \pm 0.002$	3 OUR F	TIT				
2.494	$8 \pm 0.004$	1	4.57M	<sup>1</sup> ABBIENDI	01A	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
2.487	$6 \pm 0.004$	1	4.08M	<sup>2</sup> ABREU	00F	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
2.502	$4 \pm 0.004$	2	3.96M	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> = 88–94 GeV
2.495	$1 \pm 0.004$	3	4.57M	<sup>4</sup> BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{ee} = 88-94 {\rm GeV}$
• • •	We do i	not use t	he followir	ng data for average	s, fits,	limits, e	etc. • • •
2.494	$3 \pm 0.004$	1		<sup>5</sup> ABBIENDI	<b>0</b> 4G	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = LEP1 + 130–209 GeV
2.502	$25 \pm 0.004$	1	3.97M	<sup>6</sup> ACCIARRI	00Q	L3	$E_{cm}^{ee} = LEP1 +$
2.50	$\pm 0.21$	$\pm 0.06$		<sup>7</sup> ABREU	<b>96</b> R		130–189 GeV E <sup>ee</sup> = 91.2 GeV
3.8	$\pm 0.8$	$\pm 1.0$	188	ABE	89C	CDF	$E_{\rm cm}^{p\overline{p}}$ = 1.8 TeV
2.42	$^{+0.45}_{-0.35}$		480	<sup>8</sup> ABRAMS	<b>89</b> B	MRK2	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89–93 GeV
2.7	$^{+1.2}_{-1.0}$	$\pm 1.3$	24	<sup>9</sup> ALBAJAR	89	UA1	$E_{\rm cm}^{p\overline{p}}$ = 546,630 GeV
2.7	$\pm 2.0$	$\pm 1.0$	25	<sup>10</sup> ANSARI	87	UA2	E <sup>pp</sup> <sub>cm</sub> = 546,630 GeV
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## Z WIDTH

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

- <sup>1</sup>ABBIENDI 01A error includes approximately 3.6 MeV due to statistics, 1 MeV due to event selection systematics, and 1.3 MeV due to LEP energy uncertainty.
- $^{2}$  The error includes 1.2 MeV due to LEP energy uncertainty.
- <sup>3</sup>The error includes 1.3 MeV due to LEP energy uncertainty.
- <sup>4</sup> BARATE 00C error includes approximately 3.8 MeV due to statistics, 0.9 MeV due to experimental systematics, and 1.3 MeV due to LEP energy uncertainty.
- <sup>5</sup> ABBIENDI 04G obtain this result using the S-matrix formalism for a combined fit to their cross section and asymmetry data at the Z peak and their data at 130–209 GeV. The authors have corrected the measurement for the 1 MeV shift with respect to the Breit-Wigner fits.
- <sup>6</sup> ACCIARRI 00Q interpret the s-dependence of the cross sections and lepton forwardbackward asymmetries in the framework of the S-matrix formalism. They fit to their cross section and asymmetry data at high energies, using the results of S-matrix fits to Z-peak data (ACCIARRI 00C) as constraints. The 130–189 GeV data constrains the  $\gamma/Z$ interference term. The authors have corrected the measurement for the 0.9 MeV shift with respect to the Breit-Wigner fits.
- <sup>7</sup>ABREU 96R obtain this value from a study of the interference between initial and final state radiation in the process  $e^+e^- \rightarrow Z \rightarrow \mu^+\mu^-$ .
- $^{8}$  ABRAMS 89B uncertainty includes 50 MeV due to the miniSAM background subtraction  $_{\circ}$  error.
- <sup>9</sup> ALBAJAR 89 result is from a total sample of 33  $Z \rightarrow e^+e^-$  events.
- <sup>10</sup> Quoted values of ANSARI 87 are from direct fit. Ratio of Z and W production gives either  $\Gamma(Z) < (1.09 \pm 0.07) \times \Gamma(W)$ , CL = 90% or  $\Gamma(Z) = (0.82 + 0.19 \pm 0.06) \times \Gamma(W)$ . Assuming Standard-Model value  $\Gamma(W) = 2.65$  GeV then gives  $\Gamma(Z) < 2.89 \pm 0.19$  or  $= 2.17 + 0.50 \pm 0.16$ .

	Mode	Fraction $(\Gamma_i/\Gamma)$	Scale factor/ Confidence level
$\Gamma_1$	$e^+e^-$	[ <i>a</i> ] ( 3.3632±0.0042) %	
Γ2	$\mu^+\mu^-$	[ <i>a</i> ] ( 3.3662±0.0066) %	
Γ <sub>3</sub>	$\tau^+ \tau^-$	[a] ( 3.3696±0.0083) %	
Γ4	$\ell^+ \ell^-$	[a,b] ( 3.3658±0.0023) %	
$\Gamma_5$	$\mu^+\mu^-\mu^+\mu^-$		
Г <sub>6</sub>	$\ell^+ \ell^- \ell^+ \ell^-$	$[c]$ (4.58 $\pm$ 0.26 ) $ imes$ 1	0-6
Γ <sub>7</sub>	invisible	$[a]$ (20.000 $\pm 0.055$ )%	
Г <sub>8</sub>	hadrons	[a] (69.911 $\pm 0.056$ )%	
Г9	$(u\overline{u}+c\overline{c})/2$	(11.6 $\pm$ 0.6 ) %	
$\Gamma_{10}$	$(d\overline{d} + s\overline{s} + b\overline{b})/3$	$(15.6 \pm 0.4)$ %	
$\Gamma_{11}$	<u>c</u>	(12.03 $\pm$ 0.21 ) %	
$\Gamma_{12}$	b b	(15.12 $\pm 0.05$ ) %	
Γ <sub>13</sub>	bbbb	$(3.6 \pm 1.3)  imes 1$	0 <sup>-4</sup>
$\Gamma_{14}$	ggg	< 1.1 %	CL=95%
$\Gamma_{15}$	$\pi^{0}\gamma$	< 2.01 × 1	$0^{-5}$ CL=95%
$\Gamma_{16}$	$\eta\gamma$	< 5.1 × 1	$0^{-5}$ CL=95%
$\Gamma_{17}$	$ ho^{0}\gamma$	< 2.5 × 1	$0^{-5}$ CL=95%
Γ <sub>18</sub>	$\omega \gamma$	< 6.5 × 1	$0^{-4}$ CL=95%
Γ <sub>19</sub>	$\eta^{\prime}(958)\gamma$	< 4.2 × 1	$0^{-5}$ CL=95%

Z DECAY MODES

Γ <sub>20</sub> Γ <sub>21</sub> Γ <sub>22</sub>	$ \phi \gamma \\ \gamma \gamma \\ \pi^0 \pi^0 $	< 9 < 1.46 < 1.52	$\begin{array}{c} \times 10^{-7} \\ \times 10^{-5} \\ \times 10^{-5} \end{array}$	CL=95% CL=95% CL=95%
Γ <sub>23</sub>	$\gamma \gamma \gamma$	< 2.2	imes 10 <sup>-6</sup>	CL=95%
Г <sub>24</sub>	$\pi^{\pm}W^{\mp}$	[d] < 7	$\times 10^{-5}$	CL=95%
Γ <sub>25</sub>	$ ho^{\pm} W^{\mp}$	[d] < 8.3	imes 10 <sup>-5</sup>	CL=95%
Γ <sub>26</sub>	$J/\psi(1S)$ X	$(\begin{array}{cc} 3.51 & +0.23 \\ -0.25 \end{array})$	) × 10 <sup>-3</sup>	S=1.1
Γ <sub>27</sub>	$J/\psi(1S)\gamma$	< 2.3	imes 10 <sup>-6</sup>	CL=95%
Г <sub>28</sub>	$\psi(2S)X$	( $1.60$ $\pm 0.29$	$)  imes 10^{-3}$	
Γ <sub>29</sub>	$\psi(2S)\gamma$	< 4.5	imes 10 <sup>-6</sup>	CL=95%
Г <sub>30</sub>	$J/\psi(1S)\ell^+\ell^-$			
Г <sub>31</sub>	$\chi_{c1}(1P)X$	$(2.9 \pm 0.7)$		
Г <sub>32</sub>	$\chi_{c2}(1P)X$	< 3.2	$\times 10^{-3}$	CL=90%
Г <sub>33</sub>	$arphi(1S)  imes + arphi(2S)  imes \ + arphi(3S)  imes$	$(1.0 \pm 0.5)$	) × 10 <sup>-4</sup>	
Г <sub>34</sub>	$\Upsilon(1S)X$	< 3.4	imes 10 <sup>-6</sup>	CL=95%
Γ <sub>35</sub>	$\Upsilon(1S)\chi$ $\Upsilon(1S)\gamma$	< 2.8	$\times 10 \times 10^{-6}$	CL=95%
Γ <sub>35</sub>	$\Upsilon(2S) X$	< 6.5	$\times 10 \times 10^{-6}$	CL=95%
Γ <sub>37</sub>	$\Upsilon(2S)\chi$ $\Upsilon(2S)\gamma$	< 1.7	$\times 10 \times 10^{-6}$	CL=95%
Γ <sub>38</sub>	$\Upsilon(3S) X$	< 5.4	$\times 10 \times 10^{-6}$	CL=95%
Γ <sub>38</sub>	$\Upsilon(3S)\gamma$	< 4.8	$\times 10^{-6}$	CL=95%
Γ <sub>40</sub>	$(D^0/\overline{D}^0)$ X	(20.7 ±2.0	)%	CL=5570
Γ <sub>41</sub>	$D^{\pm}X$	$(12.2 \pm 1.7)$	) %	
$\Gamma_{42}$	$D^{*}(2010)^{\pm}X$	[d] $(11.4 \pm 1.3)$	) %	
Γ <sub>43</sub>	$D_{s1}(2536)^{\pm}X$	$(3.6 \pm 0.8)$	$) \times 10^{-3}$	
Γ <sub>44</sub>		$(5.8 \pm 2.2)$	$) \times 10^{-3}$	
	$D^{*'}(2629)^{\pm}X$	searched for	) / 20	
Γ <sub>46</sub>	BX			
Γ <sub>47</sub>	<i>B</i> *X			
Γ <sub>48</sub>	B <sup>+</sup> X	$[e]$ ( 6.08 $\pm 0.13$	) %	
Γ <sub>49</sub>	$B_s^0 X$	[e] (1.59 ±0.13		
Г <sub>50</sub>	$B_c^+ X$ $\Lambda_c^+ X$	searched for	,	
Γ <sub>51</sub>	Λ <sup>Ť</sup> X	$(1.54 \pm 0.33)$	) %	
Γ <sub>52</sub>	$\Xi_c^0 X$	seen	,	
	$= \frac{1}{c}$	seen		
	<i>b</i> -baryon X	[e] ( $1.38 \pm 0.22$	) %	
	anomalous $\gamma$ + hadrons	[f] < 3.2	$\times 10^{-3}$	CL=95%
	$e^+e^-\gamma$	[f] < 5.2 [f] < 5.2	$\times$ 10 $\times$ 10 <sup>-4</sup>	CL=95%
$\Gamma_{57}$	$\mu^+\mu^-\gamma$	[f] < 5.6	$\times 10^{-4}$	CL=95%
	$\tau^{\mu}$ $\tau^{\mu}$ $\tau^{-}$ $\gamma$	[f] < 7.3	$\times 10^{-4}$	CL=95%
	$\ell^+\ell^-\gamma\gamma$	[g] < 6.8	$\times 10^{-6}$	CL=95%
	$q \overline{q} \gamma \gamma$	[g] < 5.5	imes 10 <sup>-6</sup>	CL=95%
	$\nu \overline{\nu} \gamma \gamma$	[g] < 3.1	imes 10 <sup>-6</sup>	CL=95%
<b>J 1</b>				

Г <sub>62</sub>	$e^{\pm}\mu^{\mp}$	LF	[d] < 7.5	imes 10 <sup>-7</sup>	CL=95%
Г <sub>63</sub>	$e^{\pm} au^{\mp}$	LF	[d] < 9.8	imes 10 <sup>-6</sup>	CL=95%
Г <sub>64</sub>	$\mu^{\pm} \tau^{\mp}$	LF	[d] < 1.2	imes 10 <sup>-5</sup>	CL=95%
Г <sub>65</sub>	pe	L,B	< 1.8	imes 10 <sup>-6</sup>	CL=95%
Г <sub>66</sub>	$p\mu$	L,B	< 1.8	imes 10 <sup>-6</sup>	CL=95%

- [a] This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06 (Physics Reports (Physics Letters C) 427 257 (2006)).
- [b]  $\ell$  indicates each type of lepton (e,  $\mu$ , and  $\tau$ ), not sum over them.
- [c] Here  $\ell$  indicates e or  $\mu$ .
- [d] The value is for the sum of the charge states or particle/antiparticle states indicated.
- [e] This value is updated using the product of (i) the  $Z \rightarrow bb$  fraction from this listing and (ii) the *b*-hadron fraction in an unbiased sample of weakly decaying *b*-hadrons produced in *Z*-decays provided by the Heavy Flavor Averaging Group (HFLAV, http://www.slac.stanford.edu/xorg/hflav/osc/PDG\_2009/#FRACZ).
- [f] See the Particle Listings below for the  $\gamma$  energy range used in this measurement.
- [g] For  $m_{\gamma \gamma} = (60 \pm 5)$  GeV.

### **Z PARTIAL WIDTHS**

Γ(	(e+	e <sup></sup>	)

Γ1

For the LEP experiments, this parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
83.91±0.12 OUR FIT					
$83.66 \pm 0.20$	137.0K	ABBIENDI	01A	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$83.54 \pm 0.27$	117.8k	ABREU	00F	DLPH	$E_{\rm cm}^{ee}$ = 88–94 GeV
$84.16 \pm 0.22$	124.4k	ACCIARRI	00C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$83.88 \!\pm\! 0.19$		BARATE	00C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$82.89 \!\pm\! 1.20 \!\pm\! 0.89$		<sup>1</sup> ABE	95J	SLD	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.31 GeV

<sup>1</sup> ABE 95J obtain this measurement from Bhabha events in a restricted fiducial region to improve systematics. They use the values 91.187 and 2.489 GeV for the Z mass and total decay width to extract this partial width.

# $\Gamma(\mu^+\mu^-)$

#### Γ2

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
83.99 $\pm$ 0.18 OUR FIT					
$84.03 \pm 0.30$	182.8K	ABBIENDI	01A	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$84.48 \pm 0.40$	157.6k	ABREU	00F	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
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83.95±0.44	113.4k	ACCIARRI	00c L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$84.02 \pm 0.28$		BARATE	00C ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV

# $\Gamma(\tau^+\tau^-)$

Γ3

Γ<sub>4</sub>

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
84.08±0.22 OUR FIT					
$83.94 \!\pm\! 0.41$	151.5K	ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$83.71 \pm 0.58$	104.0k	ABREU	00F	DLPH	<i>E<sup>ee</sup></i> = 88–94 GeV
$84.23 \pm 0.58$	103.0k	ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> = 88–94 GeV
$84.38 \pm 0.31$		BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV

 $\Gamma(\ell^+\ell^-)_{\ell}$  indicates each type of lepton (*e*,  $\mu$ , and  $\tau$ ), not sum over them.

In our fit  $\Gamma(\ell^+ \ell^-)$  is defined as the partial Z width for the decay into a pair of massless charged leptons. This parameter is not directly used in the 5-parameter fit assuming lepton universality but is derived using the fit results. See the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
83.984±0.086 OUR FI	т				
$83.82 \pm 0.15$	471.3K	ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$83.85 \pm 0.17$	379.4k	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.14 \pm 0.17$	340.8k	ACCIARRI	<b>00</b> C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$84.02 \pm 0.15$	500k	BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

## **Γ**(invisible)

Γ7

We use only direct measurements of the invisible partial width using the single photon channel to obtain the average value quoted below. OUR FIT value is obtained as a difference between the total and the observed partial widths assuming lepton universality.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
499.0 $\pm$ 1.5 OUR FIT					
503 $\pm 16$ OUR AVER	RAGE Erro	r includes scale f	actor	of 1.2.	
498 $\pm 12$ $\pm 12$	1791	ACCIARRI	<b>98</b> G	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV
$539$ $\pm 26$ $\pm 17$	410	AKERS	<b>95</b> C	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$450 \pm 34 \pm 34$	258	BUSKULIC	93L	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$540$ $\pm80$ $\pm40$	52	ADEVA	92	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV
$\bullet \bullet \bullet$ We do not use th	e following	data for average	s, fits,	limits, e	etc. • • •
498.1± 2.6		<sup>1</sup> ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> = 88–94 GeV
498.1± 3.2		<sup>1</sup> ABREU	00F	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
499.1± 2.9		<sup>1</sup> ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> = 88–94 GeV
499.1± 2.5		<sup>1</sup> BARATE	<b>00</b> C	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>This is an indirect determination of  $\Gamma(invisible)$  from a fit to the visible Z decay modes.

### Γ(hadrons)

## Г8

This parameter is not directly used in the 5-parameter fit assuming lepton universality, but is derived using the fit results. See the note "The Z boson" and ref. LEP-SLC 06.

VALUE (MeV)	EVTS	DOCUMENT ID		TECN	COMMENT
1744.4 $\pm$ 2.0 OUR FIT					
$1745.4 \pm 3.5$	4.10M	ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$1738.1 \pm 4.0$	3.70M	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1751.1 \pm 3.8$	3.54M	ACCIARRI	<b>00</b> C	L3	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
$1744.0 \pm 3.4$	4.07M	BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{ee}=$ 88–94 GeV
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# **Z BRANCHING RATIOS**

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

# $\Gamma(\mu^+\mu^-)/\Gamma(e^+e^-)$

 $\Gamma_2/\Gamma_1$ 

This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06. VALUE DOCUMENT ID

### 1.0009±0.0028 OUR FIT

$\Gamma( au^+ au^-)/\Gamma(e^+e^-)$	DOCUMENT ID	TECN	Соммент	′Γ <sub>1</sub>
1.0020±0.0032 OUR AVERAGE	DOCOMENT ID	<u></u>	COMMENT	
$1.02 \pm 0.06$	<sup>1</sup> AAIJ	18AR LHCB	$E^{pp}_{cm} = 8 \text{ TeV}$	
$1.0019 \pm 0.0032$	<sup>2</sup> LEP-SLC	06	$E_{\rm cm}^{ee} = 88-94  {\rm GeV}$	

<sup>1</sup>AAIJ 18AR obtain the result from the ratio of the measured  $pp \rightarrow Z + X$  cross sections in the corresponding Z decay channels.

 $^{2}$  This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

$\Gamma( au^+ au^-)/\Gamma(\mu^+\mu^-)$			$\Gamma_3/\Gamma_2$
VALUE	DOCUMENT ID	TECN	COMMENT
$1.0010\pm0.0026$ OUR AVERAGE			
$1.01 \pm 0.05$	<sup>1</sup> AAIJ	18AR LHCB	$E_{\rm cm}^{pp} = 8 { m TeV}$
$1.0010 \pm 0.0026$	<sup>2</sup> LEP-SLC	06	$E_{\rm cm}^{ee} = 88$ –94 GeV

<sup>1</sup>AAIJ 18AR obtain the result from the ratio of the measured  $pp \rightarrow Z + X$  cross sections in the corresponding Z decay channels.

 $^{2}$  This parameter is not directly used in the overall fit but is derived using the fit results; see the note "The Z boson" and ref. LEP-SLC 06.

# $\Gamma(\ell^+\ell^-\ell^+\ell^-)/\Gamma_{total}$

### $\Gamma_6/\Gamma$

Here  $\ell$  indicates either *e* or  $\mu$ . The branching fractions in this node are given within the phase-space defined by the requirements that (i) the 4-lepton invariant mass is between 80 GeV and 100 GeV, and (ii) any opposite-sign same-flavor lepton pair has a di-lepton invariant mass larger than 4 GeV.

OUR EVALUATION below is presented in RAINBOLT 19, assessing and including correlated systematic uncertainties between the measurements.

	EVTS	DOCUMENT ID		TECN	COMMENT	_
4.58±0.26 OUR EVALU	ATION					_
$4.83^{+0.23}_{-0.22}{}^{+0.35}_{-0.32}$	509	<sup>1</sup> SIRUNYAN	18bt	CMS	$E^{pp}_{ m cm}=13~ m TeV$	
$\begin{array}{rrrr} 4.9 & +0.8 & +0.4 \\ & -0.7 & -0.2 \end{array}$	39	<sup>2</sup> KHACHATRY	. <b>16</b> CC	CMS	$E^{pp}_{ m cm}=13~ m TeV$	
$4.31\!\pm\!0.34\!\pm\!0.17$	172	AAD	14N	ATLS	$E^{pp}_{cm} =$ 7, 8 TeV	
$4.6 \begin{array}{c} +1.0 \\ -0.9 \end{array} \pm 0.2$	28	<sup>3</sup> CHATRCHYAN	12bn	CMS	$E_{\rm cm}^{pp} = 7 { m TeV}$	

<sup>1</sup>SIRUNYAN 18BT report the  $Z \rightarrow 4\ell$  branching fraction =  $(4.83 + 0.23 + 0.32 \pm 0.08 \pm 0.12) \times 10^{-6}$  where the uncertainties are statistical partematic due to the provide the second statistical partematic due to the second statistical pa

 $0.12) \times 10^{-6}$ , where the uncertainties are statistical, systematic, due to theory, and luminosity. The last three have been added in quadrature to obtain the total systematic error.

 $^{\rm error.}_{\rm 2}$  KHACHATRYAN 16CC reports  $(4.9 \substack{+0.8 \\ -0.7 \ -0.2 \ -0.1 \ -0.1}) \times 10^{-6}$  value, where the uncertainties are statistical, systematic, theory, and due to luminosity. We have combined uncertainties in quadrature.

<sup>3</sup>CHATRCHYAN 12BN reports  $(4.2^{+0.9}_{-0.8} \pm 0.2) \times 10^{-6}$  value. Their result (both central value and uncertainties) is scaled up by 10% to account for the different phase-space definition used here (see RAINBOLT 19).

$\Gamma(hadrons)/\Gamma(e^+e^-)$					$\Gamma_8/\Gamma_1$
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
20.804 $\pm$ 0.050 OUR FIT					
$20.902 \pm 0.084$	137.0K	<sup>1</sup> ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.88~\pm~0.12$	117.8k	ABREU	00F	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.816 \pm 0.089$	124.4k	ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.677 \pm \ 0.075$		<sup>2</sup> BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
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• • • We do not use the following data for averages, fits, limits, etc. • • •

27.0	$^{+11.7}_{-8.8}$	12	<sup>3</sup> ABRAMS	<b>89</b> D	MRK2	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89–93 GeV
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<sup>1</sup> ABBIENDI 01A error includes approximately 0.067 due to statistics, 0.040 due to event selection systematics, 0.027 due to the theoretical uncertainty in *t*-channel prediction, and 0.014 due to LEP energy uncertainty.

<sup>2</sup> BARATE 00C error includes approximately 0.062 due to statistics, 0.033 due to experimental systematics, and 0.026 due to the theoretical uncertainty in *t*-channel prediction.
 <sup>3</sup> ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

### $\Gamma(\text{hadrons})/\Gamma(\mu^+\mu^-)$

HTTP://PDG.LBL.GOV

### $\Gamma_8/\Gamma_2$

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OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
20.785 $\pm$ 0.033 OUR FIT					
$20.811 \!\pm\! 0.058$	182.8K	<sup>1</sup> ABBIENDI	01A	OPAL	$E_{cm}^{ee}$ = 88–94 GeV
$20.65 \pm 0.08$	157.6k	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$20.861 \!\pm\! 0.097$	113.4k	ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.799 \!\pm\! 0.056$		<sup>2</sup> BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV
$\bullet \bullet \bullet$ We do not use the f	ollowing da	ata for averages, fi	ts, lim	its, etc.	• • •
$18.9 \begin{array}{c} +7.1 \\ -5.3 \end{array}$	13	<sup>3</sup> ABRAMS	<b>89</b> D	MRK2	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 89–93 GeV

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- <sup>1</sup> ABBIENDI 01A error includes approximately 0.050 due to statistics and 0.027 due to event selection systematics.
- $^2\,{\rm BARATE}$  00C error includes approximately 0.053 due to statistics and 0.021 due to experimental systematics.
- <sup>3</sup>ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

# $\Gamma(\text{hadrons})/\Gamma(\tau^+\tau^-)$

## Г8/Г3

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06).

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
20.764 $\pm$ 0.045 OUR FIT					
$20.832 \!\pm\! 0.091$	151.5K	<sup>1</sup> ABBIENDI	01A	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.84 \pm 0.13$	104.0k	ABREU	00F	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.792 \!\pm\! 0.133$	103.0k	ACCIARRI	00C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$20.707 \!\pm\! 0.062$		<sup>2</sup> BARATE	00C	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
				•	

• • • We do not use the following data for averages, fits, limits, etc. • • •

15.2 
$$^{+4.8}_{-3.9}$$
 21 <sup>3</sup> ABRAMS 89D MRK2  $E_{cm}^{ee}$  = 89–93 GeV

 $^1\,{\sf ABBIENDI}$  01A error includes approximately 0.055 due to statistics and 0.071 due to event selection systematics.

- $^2$  BARATE 00C error includes approximately 0.054 due to statistics and 0.033 due to experimental systematics.
- $^3$  ABRAMS 89D have included both statistical and systematic uncertainties in their quoted errors.

# $\Gamma(hadrons)/\Gamma(\ell^+\ell^-)$

## $\Gamma_8/\Gamma_4$

 $\ell$  indicates each type of lepton (e,  $\mu$ , and  $\tau$ ), not sum over them.

Our fit result is obtained requiring lepton universality.

VALUE	EVTS	DOCUMENT ID		TECN	COMMENT		
20.767±0.025 OUR	FIT						
$20.823 \!\pm\! 0.044$	471.3K	<sup>1</sup> ABBIENDI	01A	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV		
$20.730 \!\pm\! 0.060$	379.4k	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$20.810 \!\pm\! 0.060$	340.8k	ACCIARRI	00C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$20.725 \!\pm\! 0.039$	500k	<sup>2</sup> BARATE	<b>00</b> C	ALEP	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV		
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$							

18.9  $^{+3.6}_{-3.2}$  46 ABRAMS 89B MRK2  $E_{cm}^{ee}$  = 89–93 GeV

<sup>1</sup> ABBIENDI 01A error includes approximately 0.034 due to statistics and 0.027 due to event selection systematics.

<sup>2</sup> BARATE 00C error includes approximately 0.033 due to statistics, 0.020 due to experimental systematics, and 0.005 due to the theoretical uncertainty in t-channel prediction.

# $\Gamma((u\overline{u}+c\overline{c})/2)/\Gamma(hadrons)$

## Γ<sub>9</sub>/Γ<sub>8</sub>

This quantity is the branching ratio of  $Z \rightarrow$  "up-type" quarks to  $Z \rightarrow$  hadrons. Except ACKERSTAFF 97T the values of  $Z \rightarrow$  "up-type" and  $Z \rightarrow$  "down-type" branchings are extracted from measurements of  $\Gamma$ (hadrons), and  $\Gamma(Z \rightarrow \gamma + \text{jets})$  where  $\gamma$  is a high-energy (>5 or 7 GeV) isolated photon. As the experiments use different procedures and slightly different values of  $M_Z$ ,  $\Gamma$ (hadrons) and  $\alpha_s$  in their extraction procedures, our average has to be taken with caution.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.166 \pm 0.009$ OUR AVERAGE				
$0.172 \substack{+\ 0.011 \\ -\ 0.010}$	<sup>1</sup> ABBIENDI	04E	OPAL	$E_{ m cm}^{ee}=91.2~{ m GeV}$
$0.160\!\pm\!0.019\!\pm\!0.019$	<sup>2</sup> ACKERSTAFF	97⊤	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.137 \substack{+\ 0.038 \\ -\ 0.054}$	<sup>3</sup> ABREU	95X	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.137 \pm 0.033$	<sup>4</sup> ADRIANI	93	L3	$E_{\rm cm}^{ee}$ = 91.2 GeV

<sup>1</sup> ABBIENDI 04E select photons with energy > 7 GeV and use  $\Gamma$ (hadrons) = 1744.4 ± 2.0 MeV and  $\alpha_s = 0.1172 \pm 0.002$  to obtain  $\Gamma_u = 300 + \frac{19}{-18}$  MeV.

<sup>2</sup> ACKERSTAFF 97T measure  $\Gamma_{u\overline{u}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}}) = 0.258 \pm 0.031 \pm 0.032$ . To obtain this branching ratio authors use  $R_c + R_b = 0.380 \pm 0.010$ . This measurement is fully negatively correlated with the measurement of  $\Gamma_{d\overline{d}}, s\overline{s}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})$  given in the next data block.

<sup>3</sup>ABREU 95X use  $M_Z = 91.187 \pm 0.009$  GeV,  $\Gamma$ (hadrons) = 1725 ± 12 MeV and  $\alpha_s = 0.123 \pm 0.005$ . To obtain this branching ratio we divide their value of  $C_{2/3} = 0.91 + 0.25$ by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.66 \pm 0.05$ .

<sup>4</sup> ADRIANI 93 use  $M_Z = 91.181 \pm 0.022$  GeV,  $\Gamma$ (hadrons) = 1742 ± 19 MeV and  $\alpha_s = 0.125 \pm 0.009$ . To obtain this branching ratio we divide their value of  $C_{2/3} = 0.92 \pm 0.22$  by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.720 \pm 0.076$ .

# $\Gamma((d\overline{d}+s\overline{s}+b\overline{b})/3)/\Gamma(hadrons)$

## $\Gamma_{10}/\Gamma_8$

This quantity is the branching ratio of  $Z \rightarrow$  "down-type" quarks to  $Z \rightarrow$  hadrons. Except ACKERSTAFF 97T the values of  $Z \rightarrow$  "up-type" and  $Z \rightarrow$  "down-type" branchings are extracted from measurements of  $\Gamma$ (hadrons), and  $\Gamma(Z \rightarrow \gamma + \text{jets})$  where  $\gamma$  is a high-energy (>5 or 7 GeV) isolated photon. As the experiments use different procedures and slightly different values of  $M_Z$ ,  $\Gamma$ (hadrons) and  $\alpha_s$  in their extraction procedures, our average has to be taken with caution.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.223 \pm 0.006$ OUR AVERAGE				
$0.218 \pm 0.007$				$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$0.230\!\pm\!0.010\!\pm\!0.010$	<sup>2</sup> ACKERSTAFF	<b>97</b> T	OPAL	$E_{\rm cm}^{ee}=$ 88–94 GeV
$0.243 \substack{+\ 0.036 \\ -\ 0.026}$	<sup>3</sup> ABREU	95X	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.243 \pm 0.022$	<sup>4</sup> ADRIANI	93	L3	$E_{\rm cm}^{ee}=$ 91.2 GeV

<sup>1</sup> ABBIENDI 04E select photons with energy > 7 GeV and use  $\Gamma$ (hadrons) = 1744.4 ± 2.0 MeV and  $\alpha_s = 0.1172 \pm 0.002$  to obtain  $\Gamma_d = 381 \pm 12$  MeV.

<sup>2</sup> ACKERSTAFF 97<sup>T</sup> measure  $\Gamma_{d\overline{d},s\overline{s}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}}) = 0.371 \pm 0.016 \pm 0.016$ . To obtain this branching ratio authors use  $R_c+R_b = 0.380 \pm 0.010$ . This measurement is fully negatively correlated with the measurement of  $\Gamma_{u\overline{u}}/(\Gamma_{d\overline{d}}+\Gamma_{u\overline{u}}+\Gamma_{s\overline{s}})$  presented in the previous data block.

- <sup>3</sup>ABREU 95X use  $M_Z = 91.187 \pm 0.009$  GeV,  $\Gamma$ (hadrons) = 1725 ± 12 MeV and  $\alpha_s = 0.123 \pm 0.005$ . To obtain this branching ratio we divide their value of  $C_{1/3} = 1.62 \substack{+0.24 \\ -0.17}$  by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.66 \pm 0.05$ .
- <sup>4</sup> ADRIANI 93 use  $M_Z = 91.181 \pm 0.022$  GeV, Γ(hadrons) = 1742 ± 19 MeV and  $\alpha_s = 0.125 \pm 0.009$ . To obtain this branching ratio we divide their value of  $C_{1/3} = 1.63 \pm 0.15$  by their value of  $(3C_{1/3} + 2C_{2/3}) = 6.720 \pm 0.076$ .

# $R_{c} = \Gamma(c\overline{c})/\Gamma(\text{hadrons})$

### $\Gamma_{11}/\Gamma_8$

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_c = 0.1723$  for  $m_t = 174.3$  GeV and  $M_H = 150$  GeV.

VALUE	DOCUMENT ID		TECN	COMMENT
$0.1721 \pm 0.0030$ OUR FIT				
$0.1744 \!\pm\! 0.0031 \!\pm\! 0.0021$	<sup>1</sup> ABE	05F	SLD	E <sup>ee</sup> <sub>cm</sub> =91.28 GeV
$0.1665 \!\pm\! 0.0051 \!\pm\! 0.0081$	<sup>2</sup> ABREU		DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.1698 \!\pm\! 0.0069$	<sup>3</sup> BARATE	<b>00</b> B	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.180\ \pm 0.011\ \pm 0.013$	<sup>4</sup> ACKERSTAFF	98E	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
$0.167\ \pm 0.011\ \pm 0.012$	<sup>5</sup> ALEXANDER	<b>96</b> R	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
• • • We do not use the fo	llowing data for a	verage	es, fits, l	imits, etc. • • •
$0.1623 \!\pm\! 0.0085 \!\pm\! 0.0209$	<sup>6</sup> ABREU	<b>95</b> D	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

- <sup>1</sup> ABE 05F use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $c\overline{c}$  events using a double tag method. The single *c*-tag is obtained with a neural network trained to perform flavor discrimination using as input several signatures (corrected secondary vertex mass, vertex decay length, multiplicity and total momentum of the hemisphere). A multitag approach is used, defining 4 regions of the output value of the neural network and  $R_c$  is extracted from a simultaneous fit to the count rates of the 4 different tags. The quoted systematic error includes an uncertainty of  $\pm 0.0006$  due to the uncertainty on  $R_b$ .
- <sup>2</sup> ABREU 00 obtain this result properly combining the measurement from the  $D^{*+}$  production rate ( $R_c = 0.1610 \pm 0.0104 \pm 0.0077 \pm 0.0043$  (BR)) with that from the overall charm counting ( $R_c = 0.1692 \pm 0.0047 \pm 0.0063 \pm 0.0074$  (BR)) in  $c \overline{c}$  events. The systematic error includes an uncertainty of  $\pm 0.0054$  due to the uncertainty on the charmed hadron branching fractions.
- <sup>3</sup>BARATE 00B use exclusive decay modes to independently determine the quantities  $R_c \times f(c \rightarrow X)$ ,  $X=D^0$ ,  $D^+$ ,  $D_s^+$ , and  $\Lambda_c$ . Estimating  $R_c \times f(c \rightarrow \Xi_c / \Omega_c) = 0.0034$ , they simply sum over all the charm decays to obtain  $R_c = 0.1738 \pm 0.0047 \pm 0.0088 \pm 0.0075$ (BR). This is combined with all previous ALEPH measurements (BARATE 98T and BUSKULIC 94G,  $R_c = 0.1681 \pm 0.0054 \pm 0.0062$ ) to obtain the quoted value.
- <sup>4</sup> ACKERSTAFF 98E use an inclusive/exclusive double tag. In one jet  $D^{*\pm}$  mesons are exclusively reconstructed in several decay channels and in the opposite jet a slow pion (opposite charge inclusive  $D^{*\pm}$ ) tag is used. The *b* content of this sample is measured by the simultaneous detection of a lepton in one jet and an inclusively reconstructed  $D^{*\pm}$  meson in the opposite jet. The systematic error includes an uncertainty of  $\pm 0.006$ due to the external branching ratios.
- <sup>5</sup> ALEXANDER 96R obtain this value via direct charm counting, summing the partial contributions from  $D^0$ ,  $D^+$ ,  $D_s^+$ , and  $\Lambda_c^+$ , and assuming that strange-charmed baryons account for the 15% of the  $\Lambda_c^+$  production. An uncertainty of  $\pm 0.005$  due to the uncertainties in the charm hadron branching ratios is included in the overall systematics. <sup>6</sup> ABREU 95D perform a maximum likelihood fit to the combined p and  $p_T$  distributions
- of single and dilepton samples. The second error includes an uncertainty of  $\pm 0.0124$  due to models and branching ratios.

# $R_b = \Gamma(b\overline{b})/\Gamma(\text{hadrons})$

# $\Gamma_{12}/\Gamma_8$

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_b$ =0.21581 for  $m_t$ =174.3 GeV and  $M_H$ =150 GeV.

VALUE	DOCUMENT ID		TECN	<u>COMMENT</u>
0.21629±0.00066 OUR FIT				
$0.21594 \pm 0.00094 \pm 0.00075$	<sup>1</sup> ABE	05F	SLD	<i>E<sup>ee</sup></i> =91.28 GeV
$0.2174 \ \pm 0.0015 \ \pm 0.0028$	<sup>2</sup> ACCIARRI	00	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 89–93 GeV
$0.2178\ \pm 0.0011\ \pm 0.0013$	<sup>3</sup> ABBIENDI	<b>99</b> B	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.21634 \pm 0.00067 \pm 0.00060$	<sup>4</sup> ABREU	<b>99</b> B	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.2159\ \pm 0.0009\ \pm 0.0011$	<sup>5</sup> BARATE	97F	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$\bullet$ $\bullet$ We do not use the following	ng data for averag	es, fit	s, limits,	etc. • • •
$0.2145 \ \pm 0.0089 \ \pm 0.0067$	<sup>6</sup> ABREU	<b>95</b> D	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.219 \pm 0.006 \pm 0.005$	<sup>7</sup> BUSKULIC	<b>94</b> G	ALEP	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.251 \ \pm 0.049 \ \pm 0.030$	<sup>8</sup> JACOBSEN	91	MRK2	$E_{\rm cm}^{ee} = 91  { m GeV}$

<sup>1</sup> ABE 05F use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $b\bar{b}$  events using a double tag method. The single b-tag is obtained with a neural network trained to perform flavor discrimination using as input several signatures (corrected secondary vertex mass, vertex decay length, multiplicity and total momentum of the hemisphere; the key tag is obtained requiring the secondary vertex corrected mass to be above the D-meson mass). ABE 05F obtain  $R_b = 0.21604 \pm 0.00098 \pm 0.00074$  where the systematic error includes an uncertainty of  $\pm 0.00012$  due to the uncertainty on  $R_c$ . The value reported here is obtained properly combining with ABE 98D. The quoted systematic error includes an uncertainty of  $\pm 0.00012$  due to the uncertainty on  $R_c$ .

<sup>2</sup> ACCIARRI 00 obtain this result using a double-tagging technique, with a high  $p_T$  lepton tag and an impact parameter tag in opposite hemispheres.

- <sup>3</sup>ABBIENDI 99B tag  $Z \rightarrow b\overline{b}$  decays using leptons and/or separated decay vertices. The *b*-tagging efficiency is measured directly from the data using a double-tagging technique.
- <sup>4</sup> ABREU 99B obtain this result combining in a multivariate analysis several tagging methods (impact parameter and secondary vertex reconstruction, complemented by event shape variables). For  $R_c$  different from its Standard Model value of 0.172,  $R_b$  varies as  $-0.024 \times (R_c - 0.172)$ .
- <sup>5</sup> BARATE 97F combine the lifetime-mass hemisphere tag (BARATE 97E) with event shape information and lepton tag to identify  $Z \rightarrow b\overline{b}$  candidates. They further use *c*- and *u d s*-selection tags to identify the background. For  $R_c$  different from its Standard Model value of 0.172,  $R_b$  varies as  $-0.019 \times (R_c - 0.172)$ .
- <sup>6</sup>ABREU 95D perform a maximum likelihood fit to the combined p and  $p_T$  distributions of single and dilepton samples. The second error includes an uncertainty of  $\pm 0.0023$  due to models and branching ratios.
- <sup>7</sup> BUSKULIC 94G perform a simultaneous fit to the p and  $p_T$  spectra of both single and dilepton events.

<sup>8</sup> JACOBSEN 91 tagged  $b\overline{b}$  events by requiring coincidence of  $\geq$  3 tracks with significant impact parameters using vertex detector. Systematic error includes lifetime and decay uncertainties (±0.014).

$\Gamma(b\overline{b}b\overline{b})/\Gamma(hadrons)$					Γ <sub>13</sub> /Γ <sub>8</sub>
VALUE (units $10^{-4}$ )	DOCUMENT ID		TECN	COMMENT	
5.2 $\pm$ 1.9 OUR AVERAGE					
$3.6 \pm 1.7 \pm 2.7$	<sup>1</sup> ABBIENDI	<b>01</b> G	OPAL	$E_{cm}^{ee} = 88-94$	GeV
$6.0 \pm 1.9 \pm 1.4$	<sup>2</sup> ABREU	<b>99</b> U	DLPH	$E_{cm}^{ee} = 88-94$	GeV

<sup>1</sup>ABBIENDI 01G use a sample of four-jet events from hadronic Z decays. To enhance the  $b\overline{b}b\overline{b}$  signal, at least three of the four jets are required to have a significantly detached secondary vertex.

<sup>2</sup>ABREU 990 force hadronic Z decays into 3 jets to use all the available phase space and require a b tag for every jet. This decay mode includes primary and secondary 4b production, e.g, from gluon splitting to  $b\overline{b}$ .

$\Gamma(ggg)/\Gamma(hadrons)$					Г <sub>14</sub> /Г <sub>8</sub>
VALUE	CL%	DOCUMENT ID		TECN	COMMENT
$< 1.6 \times 10^{-2}$	95	<sup>1</sup> ABREU	96s	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

 $^{1}$  This branching ratio is slightly dependent on the jet-finder algorithm. The value we quote is obtained using the JADE algorithm, while using the DURHAM algorithm ABREU 96S obtain an upper limit of  $1.5 \times 10^{-2}$ .

# $\Gamma(\pi^0 \gamma) / \Gamma_{\text{total}}$

 $\Gamma_{15}/\Gamma$ 

					=• /
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
<2.01 × 10 <sup>-5</sup>	95	AALTONEN	14E	CDF	${\cal E}^{{m p}{\overline{m p}}}_{ m cm}=1.96~{ m TeV}$
$< 5.2 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>95</b> G	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$<$ 5.5 $ imes$ 10 $^{-5}$	95	ABREU	<b>94</b> B	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$< 2.1 \times 10^{-4}$	95	DECAMP	92	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$< 1.4  imes 10^{-4}$	95	AKRAWY	91F	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup> This limit is for both decay modes  $Z \rightarrow \pi^0 \gamma / \gamma \gamma$  which are indistinguishable in ACCIA-RRI 95G.

$\Gamma(\eta\gamma)/\Gamma_{total}$					Г <sub>16</sub> /Г
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
$< 7.6  imes 10^{-5}$	95	ACCIARRI	<b>95</b> G	L3	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
$< 8.0  imes 10^{-5}$	95	ABREU	<b>94</b> B	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
<5.1 × 10 <sup>-5</sup>	95	DECAMP	92	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$< 2.0 \times 10^{-4}$	95	AKRAWY	91F	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

$\Gamma( ho^0\gamma)/\Gamma_{ ext{total}}$						Г <sub>17</sub> /Г
VALUE	<u>CL%</u>	<u>EVTS</u>	DOCUMENT ID	TECN	COMMENT	
<2.5 × 10 <sup>-5</sup>	95	12.5k	<sup>1</sup> AABOUD	18AU ATLS	$E^{pp}_{ m cm}=13~{ m TeV}$	

<sup>1</sup>AABOUD 18AU search for the  $Z \rightarrow \rho \gamma$  decay mode where the  $\rho$  is identified through its decay  $\rho \rightarrow \pi^+ \pi^-$ . In the data corresponding to 32.3 fb<sup>-1</sup>, 12,583 events are selected for 635 < m( $\pi^+ \pi^-$ ) < 915 MeV.

$\Gamma(\omega\gamma)/\Gamma_{ ext{total}}$					Г <sub>18</sub> /Г
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
<6.5 × 10 <sup>-4</sup>	95	ABREU	<b>94</b> B	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV
$\Gamma(\eta'(958)\gamma)/\Gamma_{total}$					Г <sub>19</sub> /Г
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
$<4.2 \times 10^{-5}$	95	DECAMP	92	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

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$\Gammaig(\phi\gammaig)/\Gamma_{ ext{total}}$						Г <sub>20</sub> /Г		
VALUE	<u>CL%</u>	EVTS	DOCUMENT ID		TECN	COMMENT		
<9 × 10 <sup>-7</sup>	95	3.3k	<sup>1</sup> AABOUD	18AU	ATLS	$E^{pp}_{ m cm}=$ 13 TeV		
$\bullet \bullet \bullet$ We do not	use the	e follow	ing data for average	s, fits,	limits, e	etc. • • •		
$< \! 8.3  imes 10^{-6}$	95	1.0k	<sup>2</sup> AABOUD	16K	ATLS	$E^{pp}_{cm} = 13 \; { m TeV}$		
<sup>1</sup> AABOUD 18AU search for the $Z \rightarrow \phi \gamma$ decay mode where the $\phi$ is identified through its decay $\phi \rightarrow K^+ K^-$ . In the data corresponding to 32.3 fb <sup>-1</sup> , 3,364 events are selected for 1012 < m( $K^+ K^-$ ) < 1028 MeV. <sup>2</sup> AABOUD 16K search for the $Z \rightarrow \phi \gamma$ decay mode where the $\phi$ is identified through its decay into $K^+ K^-$ . In the data corresponding to a total luminosity of 2.7 fb <sup>-1</sup> , 1065 events are selected and their $K^+ K^- \gamma$ invariant mass spectrum is analyzed.								
$\Gamma(\gamma \gamma)/\Gamma_{total}$	would	violate	the Landau-Yang th	orem		Γ <sub>21</sub> /Γ		
VALUE		CL%	-			COMMENT		
<1.46 × 10 <sup>-5</sup>		95	AALTONEN					
$<$ 5.2 $ imes$ 10 $^{-5}$		95	<sup>1</sup> ACCIARRI	<b>95</b> G	L3	E <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$<$ 5.5 $ imes$ 10 $^{-5}$		95	ABREU	<b>94</b> B	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV		
$< 1.4  imes 10^{-4}$		95	AKRAWY	91F	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
<sup>1</sup> This limit is fo RRI 95G.	r both	decay r	modes $Z \to \pi^0 \gamma / \gamma$	$\gamma$ whic	ch are in	distinguishable in ACCIA-		
						F /F		

$\Gamma(\pi^{0}\pi^{0})/\Gamma_{total}$					Г <sub>22</sub> /Г
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
<1.52 × 10 <sup>-5</sup>	95	AALTONEN	14E	CDF	$E_{Cm}^{p\overline{p}}=1.96\;TeV$
$\Gamma(\gamma\gamma\gamma)/\Gamma_{ m total}$					Г <sub>23</sub> /Г
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
<2.2 × 10 <sup>-6</sup>	95	AAD	16L	ATLS	$E^{pp}_{cm} = 8 \text{ TeV}$
$\bullet \bullet \bullet$ We do not use the	e following o	data for averages	s, fits,	limits, e	etc. • • •
$< 1.0 \times 10^{-5}$	95	<sup>1</sup> ACCIARRI	<b>95</b> C	L3	<i>E<sup>ee</sup></i> = 88–94 GeV
$< 1.7 \times 10^{-5}$	95	<sup>1</sup> ABREU	<b>94</b> B	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$< 6.6  imes 10^{-5}$	95	AKRAWY	91F	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>Limit derived in the context of composite Z model.

#### $\Gamma(\pi^{\pm}W^{\mp})/\Gamma_{\text{total}}$ $\Gamma_{24}/\Gamma$ The value is for the sum of the charge states indicated. VALUE DOCUMENT ID TECN COMMENT <u>CL%</u> $< 7 \times 10^{-5}$ 92 ALEP $E_{cm}^{ee} = 88-94 \text{ GeV}$ 95 DECAMP $\Gamma(\rho^{\pm} W^{\mp})/\Gamma_{\text{total}}$ The value is for the sum of the charge states indicated. $\Gamma_{25}/\Gamma$ VALUE DOCUMENT ID TECN COMMENT <u>CL%</u> $< 8.3 \times 10^{-5}$ 92 ALEP *E*<sup>ee</sup><sub>cm</sub>= 88–94 GeV 95 DECAMP

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$\Gamma(J/\psi(1S)X)/\Gamma_{total}$					Г <sub>26</sub> /Г		
VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID		TECN	COMMENT		
$3.51^{+0.23}_{-0.25}$ OUR AVERA	GE Erro	r includes scale fa	ctor o	of 1.1.			
$3.21 {\pm} 0.21 {+} 0.19 \\ -0.28$	553	<sup>1</sup> ACCIARRI	99F	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$3.9 \ \pm 0.2 \ \pm 0.3$	511	<sup>2</sup> ALEXANDER	<b>96</b> B	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV		
$3.73\!\pm\!0.39\!\pm\!0.36$	153	<sup>3</sup> ABREU	<b>94</b> P	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV		
<sup>1</sup> ACCIARRI 99F combine $\mu^+ \mu^-$ and $e^+ e^- J/\psi(1S)$ decay channels. The branching ratio for prompt $J/\psi(1S)$ production is measured to be $(2.1 \pm 0.6 \pm 0.4 \substack{+0.4 \\ -0.2}$ (theor.))×10 <sup>-4</sup> <sup>2</sup> ALEXANDER 96B identify $J/\psi(1S)$ from the decays into lepton pairs. $(4.8 \pm 2.4)\%$ of this branching ratio is due to prompt $J/\psi(1S)$ production (ALEXANDER 96N). <sup>3</sup> Combining $\mu^+ \mu^-$ and $e^+ e^-$ channels and taking into account the common systematii errors. $(7.7 \substack{+6.3 \\ -5.4})\%$ of this branching ratio is due to prompt $J/\psi(1S)$ production. $\Gamma(J/\psi(1S)\gamma)/\Gamma_{total}$							
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT		
<2.3 × 10 <sup>-0</sup>	95	<sup>1</sup> AABOUD	18BL	ATLS	$E_{\rm cm}^{pp} = 13 { m TeV}$		
• • • We do not use the	e following	data for averages	s, fits,	limits, e	etc. ● ● ●		
$< 2.6  imes 10^{-6}$	95	<sup>2</sup> AAD	151	ATLS	$E^{pp}_{cm} = 8 \text{ TeV}$		
isolated photon of <i>p</i>	T > 35(2)	25) GeV and a mι	uon w	ith <i>pT</i>	Two triggers were used: $> 18(24)$ GeV. The $J/\psi$ azimuthal angle between		

is detected via its dimuon decay and it is required that the azimuthal angle between the photon and the  $J/\psi$  in the plane transverse to the beam direction is  $> \pi/2$ . The number of observed/expected background events is  $92/89 \pm 6$  in the dimuon mass range 2.9–3.3 GeV leading to the quoted 95% C.L. limit.

<sup>2</sup>AAD 151 use events with the highest  $p_T$  muon in the pair required to have  $p_T > 20$ GeV, the dimuon mass required to be within 0.2 GeV of the  $J/\psi(1S)$  mass and it's transverse momentum required to be > 36 GeV. The photon is also required to have it's  $p_T > 36$  GeV.

# $\Gamma(a/a(2S)X)/\Gamma$

$\Gamma(\psi(2S)X)/\Gamma_{total}$					Г <sub>28</sub> /Г
VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID		TECN	COMMENT
$1.60\pm0.29$ OUR AVE	RAGE				
$1.6 \ \pm 0.5 \ \pm 0.3$	39	<sup>1</sup> ACCIARRI	97J	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1.6 \ \pm 0.3 \ \pm 0.2$	46.9	<sup>2</sup> ALEXANDER	<b>96</b> B	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$1.60\!\pm\!0.73\!\pm\!0.33$	5.4	<sup>3</sup> ABREU	<b>94</b> P	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>ACCIARRI 97J measure this branching ratio via the decay channel  $\psi(2S) 
ightarrow \ell^+ \ell^-$  ( $\ell$  $= \mu$ , e).

 $^2\,{\sf ALEXANDER}$  96B measure this branching ratio via the decay channel  $\psi(2S)$   $\rightarrow$  $J/\psi \pi^+ \pi^-$ , with  $J/\psi \rightarrow \ell^+ \ell^-$ .

<sup>3</sup>ABREU 94P measure this branching ratio via decay channel  $\psi(2S) \rightarrow J/\psi \pi^+ \pi^-$ , with  $J/\psi \rightarrow \mu^+ \mu^-$ .

Citation: M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018) and 2019 update

$\Gammaig(\psi(2S)\gammaig)/\Gamma_{total}$					Г <sub>29</sub> /Г
VALUE	CL%	DOCUMENT ID	TECN	COMMENT	
<4.5 × 10 <sup>-6</sup>	95	<sup>1</sup> AABOUD	18BL ATLS	$E^{pp}_{cm} = 13 \text{ TeV}$	

<sup>1</sup> AABOUD 18BL study  $Z \rightarrow \psi(2S)\gamma$  in 13 TeV pp interactions. Two triggers were used: isolated photon of  $p_T > 35(25)$  GeV and a muon with  $p_T > 18(24)$  GeV. The  $\psi(2S)$ is detected via its dimuon decay and it is required that the azimuthal angle between the photon and the  $\psi(2S)$  in the plane transverse to the beam direction is  $> \pi/2$ . The number of observed/expected background events is  $43/42 \pm 5$  in the dimuon mass range 3.5-3.9 GeV leading to the quoted 95% C.L. limit.

$\Gammaig(J/\psi(1S)\ell^+\ell^-ig)/\Gammaig(\mu^+\mu^-\mu^+\mu^-ig)$					
VALUE	DOCUMENT ID	TECN	COMMENT		
0.67±0.18±0.05	<sup>1</sup> SIRUNYAN	18DZ CMS	<i>pp</i> at 13 TeV		

<sup>1</sup> SIRUNYAN 18DZ observe the decay  $Z \to \Psi \ell^+ \ell^-$  in pp collisions at  $\sqrt{s} = 13$  TeV, where  $\Psi$  includes  $J/\psi$  as well as  $\psi(2S) \to J/\psi X$ , and  $\ell^+ \ell^-$  represents an electron or muon pair while the  $J/\psi$  is detected via its  $\mu^+ \mu^-$  decay channel. To reduce systematic errors they determine the ratio of the branching fraction of this decay to that of  $Z \to \mu^+ \mu^- \mu^+ \mu^-$  within phase-space cuts imposed on lepton transverse momentum and pseudo rapidity, dilepton invariant mass, and  $J/\psi$  transverse momentum. The number of selected  $\Psi \mu^+ \mu^- (\Psi e^+ e^-)$  candidate events is 29 (18). Analyzing the  $\mu^+ \mu^-$  and  $\mu^+ \mu^- \ell^+ \ell^-$  invariant mass distributions, a yield of  $13.0 \pm 3.9$  ( $11.2 \pm 3.4$ ) events for the  $\Psi \mu^+ \mu^- (\Psi e^+ e^-)$  mode is obtained. The ratio of the branching fractions is determined as  $0.67 \pm 0.18 \pm 0.05$  within the selected phase-space cuts. Assuming extrapolation to full phase space cancels in the ratio, and using their measured value of  $B(Z \to \mu^+ \mu^- \mu^+ \mu^-) = (1.20 \pm 0.08) \times 10^{-6}$ , they estimate  $B(Z \to J/\psi \ell^+ \ell^-) = 8 \times 10^{-7}$ .

$\Gamma(\chi_{c1})$	(1 <i>P</i> )	)X),	/Γ <sub>total</sub>
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 $\Gamma_{31}/\Gamma$ 

(/////////	Juli				<b>U</b> =/
VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID		TECN	COMMENT
$2.9\pm0.7$ OUR AVER	AGE				
$2.7\!\pm\!0.6\!\pm\!0.5$	33	<sup>1</sup> ACCIARRI	97J	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$5.0\!\pm\!2.1\!+\!1.5_{-0.9}$	6.4	<sup>2</sup> ABREU	<b>94</b> P	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup> ACCIARRI 97J measure this branching ratio via the decay channel  $\chi_{c1} \rightarrow J/\psi + \gamma$ , with  $J/\psi \rightarrow \ell^+ \ell^-$  ( $\ell = \mu$ , e). The  $M(\ell^+ \ell^- \gamma) - M(\ell^+ \ell^-)$  mass difference spectrum is fitted with two gaussian shapes for  $\chi_{c1}$  and  $\chi_{c2}$ .

<sup>2</sup> This branching ratio is measured via the decay channel  $\chi_{c1} \rightarrow J/\psi + \gamma$ , with  $J/\psi \rightarrow \mu^+ \mu^-$ .

$\Gamma(\chi_{c2}(1P)X)/\Gamma_{total}$					Г <sub>32</sub> /Г
VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<3.2 × 10 <sup>-3</sup>	90	<sup>1</sup> ACCIARRI	97J	L3	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>ACCIARRI 97J derive this limit via the decay channel  $\chi_{c2} \rightarrow J/\psi + \gamma$ , with  $J/\psi \rightarrow \ell^+ \ell^-$  ( $\ell = \mu$ , e). The  $M(\ell^+ \ell^- \gamma) - M(\ell^+ \ell^-)$  mass difference spectrum is fitted with two gaussian shapes for  $\chi_{c1}$  and  $\chi_{c2}$ .

$\Gamma(\Upsilon(1S) \times +\Upsilon(2S))$	) X + 7(3	3 <i>S</i> )X)/Γ <sub>total</sub>	Г <sub>33</sub> /Г	$\Gamma = (\Gamma_{34} + \Gamma_{36} + \Gamma_{38}) / \Gamma$
VALUE (units $10^{-4}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
$1.0 \pm 0.4 \pm 0.22$	6.4	<sup>1</sup> ALEXANDER 96	F OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup> ALEXANDER 96F identify the  $\Upsilon$  (which refers to any of the three lowest bound states) through its decay into  $e^+e^-$  and  $\mu^+\mu^-$ . The systematic error includes an uncertainty of  $\pm 0.2$  due to the production mechanism.

# $\Gamma(\Upsilon(1S)X)/\Gamma_{total}$

Γ <sub>34</sub>	/Γ
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VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<3.4 × 10 <sup>-6</sup>	95	<sup>1</sup> AAD	151	ATLS	$E^{pp}_{cm}=8$ TeV
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$					
$< 4.4  imes 10^{-5}$	95	<sup>2</sup> ACCIARRI	99F	L3	<i>E<sup>ee</sup></i> = 88–94 GeV

<sup>1</sup> AAD 15I use events with the highest  $p_T$  muon in the pair required to have  $p_T > 20$  GeV, the dimuon mass required to be in the range 8–12 GeV and it's transverse momentum required to be > 36 GeV. The photon is also required to have it's  $p_T > 36$  GeV.

<sup>2</sup>ACCIARRI 99F search for  $\Upsilon(1S)$  through its decay into  $\ell^+ \ell^-$  ( $\ell = e$  or  $\mu$ ).

# $\Gamma(\Upsilon(1S)\gamma)/\Gamma_{\text{total}}$

VALUE	<u>CL%</u>	DOCUMENT ID	TECN	COMMENT	
<2.8 × 10 <sup>-6</sup>	95	<sup>1</sup> AABOUD	18BL ATLS	$E^{pp}_{cm} = 13 \text{ TeV}$	

<sup>1</sup> AABOUD 18BL study  $Z \rightarrow \Upsilon(1S)\gamma$  in 13 TeV pp interactions. Two triggers were used: isolated photon of  $p_T > 35(25)$  GeV and a muon with  $p_T > 18(24)$  GeV. The  $\Upsilon(1S)$ is detected via its dimuon decay and it is required that the azimuthal angle between the photon and the  $\Upsilon(1S)$  in the plane transverse to the beam direction is  $> \pi/2$ . The number of observed/expected background events is  $115/126 \pm 8$  in the dimuon mass range 9.0–10.0 GeV leading to the quoted 95% C.L. limit.

# $\Gamma(\Upsilon(2S)X)/\Gamma_{total}$

 $-(m(\alpha c))/r$ 

 $\Gamma_{36}/\Gamma$ 

/ -

 $\Gamma_{35}/\Gamma$ 

						,
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
$< 6.5 \times 10^{-6}$	95	<sup>1</sup> AAD	151	ATLS	$E^{pp}_{cm}=$ 8 TeV	
• • • We do not use the	e following	data for averages	s, fits,	limits, e	etc. • • •	
$< \! 13.9  imes 10^{-5}$	95	<sup>2</sup> ACCIARRI	<b>97</b> R	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 Ge	V

<sup>1</sup> AAD 15I use events with the highest  $p_T$  muon in the pair required to have  $p_T > 20$  GeV, the dimuon mass required to be in the range 8–12 GeV and it's transverse momentum required to be > 36 GeV. The photon is also required to have it's  $p_T > 36$  GeV.

<sup>2</sup>ACCIARRI 97R search for  $\Upsilon(2S)$  through its decay into  $\ell^+ \ell^-$  ( $\ell = e$  or  $\mu$ ).

$(1(25)\gamma)/(total)$					1 37/1
VALUE	CL%	DOCUMENT ID	TECN	COMMENT	
<1.7 × 10 <sup>-6</sup>	95	<sup>1</sup> AABOUD	18bl ATLS	${\cal E}^{m{pp}}_{\sf cm}=13\;{\sf TeV}$	

<sup>1</sup> AABOUD 18BL study  $Z \rightarrow \Upsilon(2S)\gamma$  in 13 TeV pp interactions. Two triggers were used: isolated photon of  $p_T > 35(25)$  GeV and a muon with  $p_T > 18(24)$  GeV. The  $\Upsilon(2S)$ is detected via its dimuon decay and it is required that the azimuthal angle between the photon and the  $\Upsilon(2S)$  in the plane transverse to the beam direction is  $> \pi/2$ . The number of observed/expected background events is  $106/121 \pm 8$  in the dimuon mass range 9.5–10.5 GeV leading to the quoted 95% C.L. limit.

$\Gamma(\Upsilon(3S)X)/\Gamma_{total}$						Г <sub>38</sub> /Г
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
<5.4 × 10 <sup>-6</sup>	95	<sup>1</sup> AAD	151	ATLS	$E^{pp}_{cm} = 8 \text{ TeV}$	
• • • We do not use the	e following	data for averages	s, fits,	limits,	etc. • • •	
$< 9.4  imes 10^{-5}$	95	<sup>2</sup> ACCIARRI	<b>97</b> R	L3	<i>E<sup>ee</sup></i> = 88–94 G	eV

<sup>1</sup> AAD 15I use events with the highest  $p_T$  muon in the pair required to have  $p_T > 20$  GeV, the dimuon mass required to be in the range 8–12 GeV and it's transverse momentum required to be > 36 GeV. The photon is also required to have it's  $p_T > 36$  GeV.

<sup>2</sup>ACCIARRI 97R search for  $\Upsilon(3S)$  through its decay into  $\ell^+\ell^-$  ( $\ell = e$  or  $\mu$ ).

Г(1	r(35)	γ)/ſ	total
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Г<sub>39</sub>/Г

 $\Gamma_{40}/\Gamma_8$ 

 $\Gamma_{42}/\Gamma_8$ 

VALUE	CL%	DOCUMENT ID	TECN	COMMENT	
<4.8 × 10 <sup>-6</sup>	95	<sup>1</sup> AABOUD	18BL ATLS	$E^{pp}_{cm} = 13 \text{ TeV}$	

<sup>1</sup> AABOUD 18BL study  $Z \rightarrow \Upsilon(3S)\gamma$  in 13 TeV pp interactions. Two triggers were used: isolated photon of  $p_T > 35(25)$  GeV and a muon with  $p_T > 18(24)$  GeV. The  $\Upsilon(3S)$ is detected via its dimuon decay and it is required that the azimuthal angle between the photon and the  $\Upsilon(3S)$  in the plane transverse to the beam direction is  $> \pi/2$ . The number of observed/expected background events is  $112/113 \pm 8$  in the dimuon mass range 10.0–11.0 GeV leading to the quoted 95% C.L. limit.

$\Gamma((D^{o}/D^{o})X)/\Gamma(ha)$	adrons)				
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
0.296±0.019±0.021	369	<sup>1</sup> ABREU	93I	DLPH	$E_{\rm cm}^{ee} = 88-94 {\rm Ge}$

<sup>1</sup> The  $(D^0/\overline{D}^0)$  states in ABREU 931 are detected by the  $K\pi$  decay mode. This is a corrected result (see the erratum of ABREU 931).

$\Gamma(D^{\pm}X)/\Gamma(hadrons)$	5)				Γ <sub>41</sub> /Γ <sub>8</sub>
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$0.174 {\pm} 0.016 {\pm} 0.018$	539	<sup>1</sup> ABREU	931	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup> The  $D^{\pm}$  states in ABREU 93I are detected by the  $K \pi \pi$  decay mode. This is a corrected result (see the erratum of ABREU 93I).

# $\Gamma(D^*(2010)^{\pm}X)/\Gamma(hadrons)$

. . . .

I he value is for the sum of the charge states indicated.						
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT	
0.163±0.019 OUR AVE	RAGE	Error includes scale	factor	of 1.3.		
$0.155 \!\pm\! 0.010 \!\pm\! 0.013$	358	<sup>1</sup> ABREU	931	DLPH	$E_{\rm cm}^{ee}$ = 88–94 GeV	
$0.21 \pm 0.04$	362	<sup>2</sup> DECAMP	91J	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV	

 ${}^{1}D^{*}(2010)^{\pm}$  in ABREU 93I are reconstructed from  $D^{0}\pi^{\pm}$ , with  $D^{0} \rightarrow K^{-}\pi^{+}$ . The new CLEO II measurement of  $B(D^{*\pm} \rightarrow D^{0}\pi^{\pm}) = (68.1 \pm 1.6)$  % is used. This is a corrected result (see the erratum of ABREU 93I).

<sup>2</sup> DECAMP 91J report  $B(D^*(2010)^+ \rightarrow D^0 \pi^+) B(D^0 \rightarrow K^- \pi^+) \Gamma(D^*(2010)^{\pm}X)$ /  $\Gamma(hadrons) = (5.11 \pm 0.34) \times 10^{-3}$ . They obtained the above number assuming  $B(D^0 \rightarrow K^- \pi^+) = (3.62 \pm 0.34 \pm 0.44)\%$  and  $B(D^*(2010)^+ \rightarrow D^0 \pi^+) = (55 \pm 4)\%$ . We have rescaled their original result of  $0.26 \pm 0.05$  taking into account the new CLEO II branching ratio  $B(D^*(2010)^+ \rightarrow D^0 \pi^+) = (68.1 \pm 1.6)\%$ .

#### $\Gamma(D_{s1}(2536)^{\pm}X)/\Gamma(hadrons)$ $\Gamma_{43}/\Gamma_8$ $D_{s1}(2536)^{\pm}$ is an expected orbitally-excited state of the $D_s$ meson. DOCUMENT ID TECN COMMENT VALUE (%) EVTS 02B ALEP $E_{cm}^{ee} = 88-94 \text{ GeV}$ <sup>1</sup> HEISTER 92 $0.52 \pm 0.09 \pm 0.06$ <sup>1</sup>HEISTER 02B reconstruct this meson in the decay modes $D_{s1}(2536)^{\pm} \rightarrow D^{*\pm} \kappa^0$ and $D_{s1}(2536)^{\pm} \rightarrow D^{*0} K^{\pm}$ . The quoted branching ratio assumes that the decay width of the $D_{s1}(2536)$ is saturated by the two measured decay modes.

# $\Gamma(D_{sJ}(2573)^{\pm}X)/\Gamma(hadrons)$

$D_{sJ}^{}(2573)^{\pm}$ is	an expected	orbitally-excited a	state o	of the D	<sub>s</sub> meson.
VALUE (%)	EVTS	DOCUMENT ID		TECN	COMMENT
$0.83 {\pm} 0.29 {+} 0.07 {-} 0.13$	64	<sup>1</sup> HEISTER	<b>02</b> B	ALEP	<i>E<sup>ee</sup></i> = 88–94 GeV

<sup>1</sup>HEISTER 02B reconstruct this meson in the decay mode  $D^*_{s2}(2573)^{\pm} \rightarrow D^0 K^{\pm}$ . The quoted branching ratio assumes that the detected decay mode represents 45% of the full decay width.

# $\Gamma(D^{*'}(2629)^{\pm}X)/\Gamma(hadrons)$

$D^{*'}(2629)^{\pm}$ is a predicted radial excitation of the $D^{*}(2010)^{\pm}$ meson	n.
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VALUE	<u>DOCUMENT ID</u>	TECN	<u>COMMENT</u>
searched for	1 ABBIENDI 01M	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
1 ABBIENDI 01N searched	for the decay mode $D^*$	*/(2620)±	$ \rightarrow D^{*\pm}\pi^{\pm}\pi^{-}$ with

ABBIENDI 01N searched for the decay mode  $D^{*'}(2629)^{\pm} \rightarrow D^{*\pm}\pi^{+}\pi^{-}$  with  $D^{*+} \rightarrow D^0 \pi^+$ , and  $D^0 \rightarrow K^- \pi^+$ . They quote a 95% CL limit for  $Z \rightarrow D^{*\prime}(2629)^{\pm} \times B(D^{*\prime}(2629)^+ \rightarrow D^{*+} \pi^+ \pi^-) < 3.1 \times 10^{-3}$ .

# $\Gamma(B^*X)/[\Gamma(BX) + \Gamma(B^*X)]$

 $\Gamma_{47}/(\Gamma_{46}+\Gamma_{47})$ As the experiments assume different values of the b-baryon contribution, our average should be taken with caution.

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
$0.75 \pm 0.04$ OUR AVE	RAGE				
$0.760 \!\pm\! 0.036 \!\pm\! 0.083$		<sup>1</sup> ACKERSTAFF	<b>9</b> 7M	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.771\!\pm\!0.026\!\pm\!0.070$		<sup>2</sup> BUSKULIC	<b>96</b> D	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.72 \ \pm 0.03 \ \pm 0.06$		<sup>3</sup> ABREU	<b>95</b> R	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$0.76\ \pm 0.08\ \pm 0.06$	1378	<sup>4</sup> ACCIARRI	<b>95</b> B	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV

 $^1$ ACKERSTAFF 97M use an inclusive B reconstruction method and assume a (13.2  $\pm$ 4.1)% b-baryon contribution. The value refers to a b-flavored meson mixture of  $B_u$ ,  $B_d$ , and  $B_{s}$ .

 $^2$ BUSKULIC 96D use an inclusive reconstruction of B hadrons and assume a (12.2  $\pm$ 4.3)% b-baryon contribution. The value refers to a b-flavored mixture of  $B_{\mu\nu}$ ,  $\dot{B}_{d}$ , and В<sub>s</sub>.

 $^3$  ABREU 95R use an inclusive *B*-reconstruction method and assume a  $(10\pm4)\%$  *b*-baryon contribution. The value refers to a *b*-flavored meson mixture of  $B_{\mu}$ ,  $B_{d}$ , and  $B_{s}$ .

 $^4$  ACCIARRI 95B assume a 9.4% *b*-baryon contribution. The value refers to a *b*-flavored mixture of  $B_{\mu}$ ,  $B_{d}$ , and  $B_{s}$ .

HTTP://PDG.LBL.GOV

 $\Gamma_{45}/\Gamma_8$ 

 $\Gamma_{44}/\Gamma_8$ 

## $\Gamma(B^+X)/\Gamma(hadrons)$

### $\Gamma_{48}/\Gamma_8$

"OUR EVALUATION" is obtained using our current values for  $f(\overline{b} \rightarrow B^+)$  and  $R_b = \Gamma(b\overline{b})/\Gamma(hadrons)$ . We calculate  $\Gamma(B^+ X)/\Gamma(hadrons) = R_b \times f(\overline{b} \rightarrow B^+)$ . The decay fraction  $f(\overline{b} \rightarrow B^+)$  was provided by the Heavy Flavor Averaging Group (HFLAV, http://www.slac.stanford.edu/xorg/hflav/osc/PDG\_2009/#FRACZ). DOCUMENT ID TECN COMMENT

#### $0.0869 \pm 0.0019$ OUR EVALUATION

 $0.0887 \pm 0.0030$ 

VALUE

<sup>1</sup> ABDALLAH 03K DLPH  $E_{cm}^{ee} = 88-94$  GeV

<sup>1</sup> ABDALLAH 03K measure the production fraction of  $B^+$  mesons in hadronic Z decays  $f(B^+) = (40.99 \pm 0.82 \pm 1.11)\%$ . The value quoted here is obtained multiplying this production fraction by our value of  $R_b = \Gamma(\overline{b}b)/\Gamma(hadrons)$ .

# $\Gamma(B_s^0 X) / \Gamma(hadrons)$

#### Γ<sub>49</sub>/Γ<sub>8</sub>

"OUR EVALUATION" is obtained using our current values for  $f(\overline{b} \rightarrow B_s^0)$  and  $R_b = \Gamma(b\overline{b})/\Gamma(hadrons)$ . We calculate  $\Gamma(B_s^0)/\Gamma(hadrons) = R_b \times f(\overline{b} \rightarrow B_s^0)$ . The decay fraction  $f(\overline{b} \rightarrow B_s^0)$  was provided by the Heavy Flavor Averaging Group (HFLAV, http://www.slac.stanford.edu/xorg/hflav/osc/PDG\_2009/#FRACZ).

VALUE	DOCOMENTID	TLCN	COMMENT				
0.0227±0.0019 OUR EVALUATION							
seen	<sup>1</sup> ABREU	92м DLPH	I <i>E<sup>ee</sup></i> = 88–94 GeV				
seen	<sup>2</sup> ACTON	92N OPAL	<i>E<sup>ee</sup></i> = 88–94 GeV				
seen	<sup>3</sup> BUSKULIC	92E ALEF	2 <i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV				
$^1$ ABREU 92M reported value is	$\Gamma(B_s^0 X) * B(B_s^0 \rightarrow$	$D_{s} \mu \nu_{\mu} X$ )	$*B(D_{s} \rightarrow \phi \pi)/\Gamma(hadrons)$				

$$= (18 \pm 8) \times 10^{-1}$$

5

<sup>2</sup> ACTON 92N find evidence for  $B_s^0$  production using  $D_s$ - $\ell$  correlations, with  $D_s^+ \to \phi \pi^+$ and  $K^*(892)K^+$ . Assuming  $R_b$  from the Standard Model and averaging over the e and  $\mu$  channels, authors measure the product branching fraction to be  $f(\overline{b} \to B_s^0) \times B(B_s^0 \to D_s^- \ell^+ \nu_\ell X) \times B(D_s^- \to \phi \pi^-) = (3.9 \pm 1.1 \pm 0.8) \times 10^{-4}$ .

<sup>3</sup>BUSKULIC 92E find evidence for  $B_s^0$  production using  $D_s$ - $\ell$  correlations, with  $D_s^+ \rightarrow \phi \pi^+$  and  $K^*(892)K^+$ . Using  $B(D_s^+ \rightarrow \phi \pi^+) = (2.7 \pm 0.7)\%$  and summing up the e and  $\mu$  channels, the weighted average product branching fraction is measured to be  $B(\overline{b} \rightarrow B_s^0) \times B(B_s^0 \rightarrow D_s^- \ell^+ \nu_\ell X) = 0.040 \pm 0.011 \substack{+0.010 \\ -0.012}$ .

# $\Gamma(B_c^+ X)/\Gamma(hadrons)$

 $\Gamma_{50}/\Gamma_8$ 

VALUE	DOCUMENT ID	TECN	COMMENT
searched for	<sup>1</sup> ACKERSTAFF 98	BO OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
searched for	<sup>2</sup> ABREU 97	7e DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
searched for	<sup>3</sup> BARATE 97	7H ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV

<sup>1</sup> ACKERSTAFF 980 searched for the decay modes  $B_c \rightarrow J/\psi \pi^+$ ,  $J/\psi a_1^+$ , and  $J/\psi \ell^+ \nu_{\ell}$ , with  $J/\psi \rightarrow \ell^+ \ell^-$ ,  $\ell = e, \mu$ . The number of candidates (background) for the three decay modes is 2 (0.63 ± 0.2), 0 (1.10 ± 0.22), and 1 (0.82 ± 0.19) respectively. Interpreting the  $2B_c \rightarrow J/\psi \pi^+$  candidates as signal, they report  $\Gamma(B_c^+ X) \times B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) = (3.8^{+5.0}_{-2.4} \pm 0.5) \times 10^{-5}$ . Interpreted as background, the 90% CL bounds are  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(hadrons) < 1.06 \times 10^{-4}$ ,  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+)/\Gamma(B_c^+ X)$ 

 $J/\psi a_1^+)/\Gamma(hadrons) < 5.29 \times 10^{-4}, \ \Gamma(B_c^+X)*B(B_c \rightarrow J/\psi \ell^+ \nu_\ell)/\Gamma(hadrons) < 6.96 \times 10^{-5}.$ 

<sup>2</sup> ABREU 97E searched for the decay modes  $B_c \rightarrow J/\psi \pi^+$ ,  $J/\psi \ell^+ \nu_\ell$ , and  $J/\psi (3\pi)^+$ , with  $J/\psi \rightarrow \ell^+ \ell^-$ ,  $\ell = e,\mu$ . The number of candidates (background) for the three decay modes is 1 (1.7), 0 (0.3), and 1 (2.3) respectively. They report the following 90% CL limits:  $\Gamma(B_c^+X)*B(B_c \rightarrow J/\psi \pi^+)/\Gamma(\text{hadrons}) < (1.05-0.84) \times 10^{-4}$ ,  $\Gamma(B_c^+X)*B(B_c \rightarrow J/\psi \pi^+)/\Gamma(\text{hadrons}) < (1.05-0.84) \times 10^{-4}$ ,  $\Gamma(B_c^+X)*B(B_c \rightarrow J/\psi \pi^+)/\Gamma(\text{hadrons}) < (1.05-0.84) \times 10^{-4}$ ,  $\Gamma(B_c^+X)*B(B_c \rightarrow J/\psi \ell \nu_\ell)/\Gamma(\text{hadrons}) < (5.8-5.0) \times 10^{-5}$ ,  $\Gamma(B_c^+X)*B(B_c \rightarrow J/\psi (3\pi)^+)/\Gamma(\text{hadrons}) < (1.75 \times 10^{-4})$ , where the ranges are due to the predicted  $B_c$  lifetime (0.4-1.4) ps. <sup>3</sup> BARATE 97H searched for the decay modes  $B_c \rightarrow J/\psi \pi^+$  and  $J/\psi \ell^+ \nu_\ell$  with

 $J/\psi \rightarrow \ell^+ \ell^-$ ,  $\ell = e,\mu$ . The number of candidates (background) for the two decay modes is 0 (0.44) and 2 (0.81) respectively. They report the following 90% CL limits:  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+) / \Gamma(hadrons) < 3.6 \times 10^{-5}$  and  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \pi^+) / \Gamma(hadrons) < 3.6 \times 10^{-5}$  and  $\Gamma(B_c^+ X) * B(B_c \rightarrow J/\psi \ell^+ \nu_\ell) / \Gamma(hadrons) < 5.2 \times 10^{-5}$ .

 $\Gamma(\Lambda_c^+ X) / \Gamma(hadrons)$ 

 $\Gamma_{51}/\Gamma_8$ 

VALUE	DOCUMENT ID		TECN	COMMENT
$0.022\pm0.005$ OUR AVERAGE				
$0.024 \pm 0.005 \pm 0.006$	<sup>1</sup> ALEXANDER	<b>96</b> R	OPAL	$E_{\rm cm}^{ee} = 88-94  { m GeV}$
$0.021\!\pm\!0.003\!\pm\!0.005$	<sup>2</sup> BUSKULIC	96Y	ALEP	$E_{\rm cm}^{ee} =$ 88–94 GeV

<sup>1</sup>ALEXANDER 96R measure  $R_b \times f(b \rightarrow \Lambda_c^+ X) \times B(\Lambda_c^+ \rightarrow pK^-\pi^+) = (0.122 \pm 0.023 \pm 0.010)\%$  in hadronic Z decays; the value quoted here is obtained using our best value  $B(\Lambda_c^+ \rightarrow pK^-\pi^+) = (5.0 \pm 1.3)\%$ . The first error is the total experiment's error and the second error is the systematic error due to the branching fraction uncertainty. <sup>2</sup>BUSKULIC 96Y obtain the production fraction of  $\Lambda_c^+$  baryons in hadronic Z decays  $f(b \rightarrow \Lambda_c^+ X) = 0.110 \pm 0.014 \pm 0.006$  using  $B(\Lambda_c^+ \rightarrow pK^-\pi^+) = (4.4 \pm 0.6)\%$ ; we have rescaled using our best value  $B(\Lambda_c^+ \rightarrow pK^-\pi^+) = (5.0 \pm 1.3)\%$  obtaining  $f(b \rightarrow \Lambda_c^+ X) = 0.097 \pm 0.013 \pm 0.025$  where the first error is their total experiment's error and the second error is the systematic error due to the branching fraction uncertainty. The value quoted here is obtained multiplying this production fraction by our value of  $R_b = \Gamma(b \overline{b})/\Gamma(hadrons)$ .

#### $\Gamma(\Xi_c^0 X) / \Gamma(hadrons)$ $\Gamma_{52}/\Gamma_8$ DOCUMENT ID TECN COMMENT We do not use the following data for averages, fits, limits, etc. • • <sup>1</sup> ABDALLAH 05C DLPH $E_{cm}^{ee} = 88-94 \text{ GeV}$ seen <sup>1</sup>ABDALLAH 05C searched for the charmed strange baryon $\Xi_c^0$ in the decay channel $\Xi_c^0 \rightarrow \Xi^- \pi^+ (\Xi^- \rightarrow \Lambda \pi^-)$ . The production rate is measured to be $f_{\Xi_c^0} \times B(\Xi_c^0 \rightarrow \Xi_c^0)$ $arepsilon^-\pi^+)=$ (4.7 $\pm$ 1.4 $\pm$ 1.1) imes 10 $^{-4}$ per hadronic Z decay. $\Gamma(\Xi_b X)/\Gamma(hadrons)$ $\Gamma_{53}/\Gamma_8$ Here $\Xi_b$ is used as a notation for the strange *b*-baryon states $\Xi_b^-$ and $\Xi_b^0$ . DOCUMENT ID \_\_\_\_\_ TECN \_\_\_\_\_ COMMENT VALUE • We do not use the following data for averages, fits, limits, etc. • • • <sup>1</sup> ABDALLAH 05C DLPH $E_{cm}^{ee} = 88-94 \text{ GeV}$ seen 96T ALEP $E_{cm}^{ee} = 88-94 \text{ GeV}$ 95V DLPH $E_{cm}^{ee} = 88-94 \text{ GeV}$ <sup>2</sup> BUSKULIC seen <sup>3</sup> ABREU seen

- $^1$ ABDALLAH 05C searched for the beauty strange baryon  $arepsilon_b$  in the inclusive semileptonic decay channel  $\Xi_b \rightarrow \Xi^- \ell^- \overline{\nu}_\ell X$ . Evidence for the  $\Xi_b$  production is seen from the observation of  $\Xi^{\mp}$  production accompanied by a lepton of the same sign. From the excess of "right-sign" pairs  $\Xi^{\mp} \ell^{\mp}$  compared to "wrong-sign" pairs  $\Xi^{\mp} \ell^{\pm}$  the production rate is measured to be B( $b \rightarrow \Xi_b$ ) × B( $\Xi_b \rightarrow \Xi^- \ell^- X$ ) = (3.0 ± 1.0 ± 0.3) × 10<sup>-4</sup> per lepton species, averaged over electrons and muons.
- $^2$  BUSKULIC 96T investigate  $\varXi$ -lepton correlations and find a significant excess of "rightsign" pairs  $\Xi^{\mp} \ell^{\mp}$  compared to "wrong-sign" pairs  $\Xi^{\mp} \ell^{\pm}$ . This excess is interpreted as evidence for  $\Xi_b$  semileptonic decay. The measured product branching ratio is B( $b \rightarrow$  $\Xi_b$ ) × B( $\Xi_b \rightarrow X_c X \ell^- \overline{\nu}_\ell$ ) × B( $X_c \rightarrow \Xi^- X'$ ) = (5.4 ± 1.1 ± 0.8) × 10<sup>-4</sup> per lepton species, averaged over electrons and muons, with  $X_c$  a charmed baryon.
- <sup>3</sup>ABREU 95V observe an excess of "right-sign" pairs  $\Xi^{\pm}\ell^{\pm}$  compared to "wrong-sign" pairs  $\Xi^{\mp}\ell^{\pm}$  in jets: this excess is interpreted as evidence for the beauty strange baryon  $\Xi_b$  production, with  $\Xi_b \to \Xi^- \ell^- \overline{\nu}_\ell X$ . They find that the probability for this signal to come from non *b*-baryon decays is less than  $5 \times 10^{-4}$  and that  $\Lambda_b$  decays can account for less than 10% of these events. The  $\Xi_b$  production rate is then measured to be B( $b \rightarrow$  $(\Xi_h) \times B(\Xi_h \rightarrow \Xi^- \ell^- X) = (5.9 \pm 2.1 \pm 1.0) \times 10^{-4}$  per lepton species, averaged over electrons and muons.

### $\Gamma(b-baryon X)/\Gamma(hadrons)$

 $\Gamma_{54}/\Gamma_8$ 

 $\Gamma_{55}/\Gamma$ 

"OUR EVALUATION" is obtained using our current values for f(  $b \rightarrow b$ -baryon) and  $R_b = \Gamma(b \overline{b})/\Gamma(hadrons)$ . We calculate  $\Gamma(b-baryon X)/\Gamma(hadrons) = R_b \times f(b \rightarrow baryon X)$ *b*-baryon). The decay fraction  $f(b \rightarrow b$ -baryon) was provided by the Heavy Flavor Averaging Group (HFLAV, http://www.slac.stanford.edu/xorg/hflav/osc/PDG\_2009). VALUE DOCUMENT ID TECN COMMENT

#### 0.0197±0.0032 OUR EVALUATION $0.0221 \pm 0.0015 \pm 0.0058$

<sup>1</sup> BARATE 98v ALEP *E*<sup>ee</sup><sub>cm</sub>= 88–94 GeV  $^1$  BARATE 98V use the overall number of identified protons in b-hadron decays to measure  $f(b \rightarrow b\text{-baryon}) = 0.102 \pm 0.007 \pm 0.027$ . They assume  $BR(b\text{-baryon} \rightarrow pX) =$ 

 $(58 \pm 6)\%$  and BR $(B_5^0 \rightarrow pX) = (8.0 \pm 4.0)\%$ . The value quoted here is obtained multiplying this production fraction by our value of  $R_b = \Gamma(b\overline{b})/\Gamma(hadrons)$ .

### $\Gamma(\text{anomalous } \gamma + \text{hadrons}) / \Gamma_{\text{total}}$

Limits on additional sources of prompt photons beyond expectations for final-state bremsstrahlung.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<3.2 × 10 <sup>-3</sup>	95	<sup>1</sup> AKRAWY	90J	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>AKRAWY 90J report  $\Gamma(\gamma X) < 8.2$  MeV at 95%CL. They assume a three-body  $\gamma q \overline{q}$ distribution and use  $E(\gamma) > 10$  GeV.

$\Gamma(e^+ e^- \gamma) / \Gamma_{ m total}$					Г	<sub>56</sub> /Г
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
<5.2 × 10 <sup>-4</sup>	95	<sup>1</sup> ACTON	<b>91</b> B	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV	

<sup>1</sup>ACTON 91B looked for isolated photons with E>2% of beam energy (> 0.9 GeV).

$\Gamma(\mu^+\mu^-\gamma)/\Gamma_{ ext{total}}$						Г <sub>57</sub> /Г
VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
<5.6 × 10 <sup>-4</sup>	95	<sup>1</sup> ACTON	<b>91</b> B	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV	,

<sup>1</sup> ACTON 91B looked for isolated photons with E>2% of beam energy (> 0.9 GeV).

Citation: M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018) and 2019 update

$\Gamma(\tau^+\tau^-\gamma)/\Gamma_{total}$					Г <sub>58</sub> /Г	
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT	
<7.3 × 10 <sup>-4</sup>	95	<sup>1</sup> ACTON	<b>91</b> B	OPAL	$E_{\rm Cm}^{ee} = 91.2  {\rm GeV}$	
<sup>1</sup> ACTON 91B looke	ed for isolat	ed photons with E>	2% c	of beam	energy ( $> 0.9$ GeV).	
$\Gamma(\ell^+\ell^-\gamma\gamma)/\Gamma_{total}$ The value is the		$\ell={\it e},\mu, au.$			Г <sub>59</sub> /Г	
VALUE	<u>CL%</u>	DOCUMENT ID				
<6.8 × 10 <sup>—6</sup>	95	<sup>1</sup> ACTON	93E	OPAL	$E_{\rm cm}^{ee}=$ 88–94 GeV	
$^{1}$ For $m_{\gamma\gamma}=$ 60 $\pm$	5 GeV.					
$\Gamma(q\overline{q}\gamma\gamma)/\Gamma_{ m total}$					Г <sub>60</sub> /Г	
VALUE	<u>CL%</u>	DOCUMENT ID			COMMENT	
<5.5 × 10 <sup>-6</sup>	95	<sup>1</sup> ACTON	93E	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV	
$^{1}$ For $m_{\gamma\gamma}=$ 60 $\pm$	5 GeV.					
$\Gamma ig(  u \overline{ u} \gamma \gamma ig) / \Gamma_{total}$					Г <sub>61</sub> /Г	
VALUE	<u>CL%</u>	<u>DOCUMENT ID</u> <sup>1</sup> ACTON		TECN	COMMENT	
<3.1 × 10 <sup>-6</sup>	95	<sup>1</sup> ACTON	93E	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV	
$^1$ For $m_{\gamma\gamma}=$ 60 $\pm$	5 GeV.					
	-	ber conservation. T	he va	alue is f	<b>Γ<sub>62</sub>/Γ</b> or the sum of the charge	
states indicated VALUE	CL%	DOCUMENT ID		TECN	COMMENT	
<7.5 × 10 <sup>-7</sup>	95	AAD			$E_{\rm cm}^{pp} = 8  {\rm TeV}$	
$< 2.5 \times 10^{-6}$	95	ABREU			$E_{\rm cm}^{ee}$ = 88–94 GeV	
$< 1.7 \times 10^{-6}$	95	AKERS	95W	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV	
$< 0.6 \times 10^{-5}$	95	ADRIANI	931	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV	
${<}2.6 imes10^{-5}$	95	DECAMP	92	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV	
$ \Gamma(e^{\pm}\mu^{\mp})/\Gamma(e^{+}e^{-}) \qquad \qquad \Gamma_{62}/\Gamma_{1} $ Test of lepton family number conservation. The value is for the sum of the charge states indicated.						
VALUE	<u>CL%</u>	DOCUMENT ID	TE	<u>CN</u> CC	DMMENT	
<0.07	90	ALBAJAR 89	UA	$A1 E_{0}^{l}$	o <del>p</del> cm= 546,630 GeV	

 $\Gamma(e^{\pm}\tau^{\mp})/\Gamma_{\text{total}}$ Test of lepton family number conservation. The value is for the sum of the charge

VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT	
${<}5.8 imes10^{-5}$	95	AABOUD	18CN	ATLS	$E^{pp}_{ m cm}=13~{ m TeV}$	
$< 2.2 \times 10^{-5}$	95	ABREU	<b>97</b> C	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 88–94 GeV	
<9.8 × 10 <sup>-6</sup>	95	AKERS	95W	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV	
$< 1.3 \times 10^{-5}$	95	ADRIANI	931	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV	
$< 1.2 \times 10^{-4}$	95	DECAMP	92	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV	

# $\Gamma(\mu^{\pm}\tau^{\mp})/\Gamma_{\text{total}}$

# Г<sub>64</sub>/Г

Test of lepton family number conservation. The value is for the sum of the charge states indicated.

oraroo marcaroar					
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
${<}1.3\times10^{-5}$	95	AABOUD	18CN	ATLS	$E^{pp}_{cm}=$ 8, 13 TeV
<1.2 × 10 <sup>-5</sup>	95	ABREU	<b>97</b> C	DLPH	$E_{\rm cm}^{ee}$ = 88–94 GeV
$< 1.7 \times 10^{-5}$	95	AKERS	95W	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$< 1.9 \times 10^{-5}$	95	ADRIANI	931	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV
$< 1.0 \times 10^{-4}$	95	DECAMP	92	ALEP	$E_{\rm cm}^{ee}=$ 88–94 GeV

# $\Gamma(pe)/\Gamma_{total}$

Γ<sub>65</sub>/Γ

Test of baryon number and lepton number conservations. Charge conjugate states are implied.

VALUE	CL%	DOCUMENT ID		TECN	COMMENT
<1.8 × 10 <sup>-6</sup>	95	<sup>1</sup> ABBIENDI	991	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>ABBIENDI 991 give the 95%CL limit on the partial width  $\Gamma(Z^0 \rightarrow pe) < 4.6$  KeV and we have transformed it into a branching ratio.

# $\Gamma(\rho\mu)/\Gamma_{\text{total}}$

# Г<sub>66</sub>/Г

Test of baryon number and lepton number conservations. Charge conjugate states are implied.

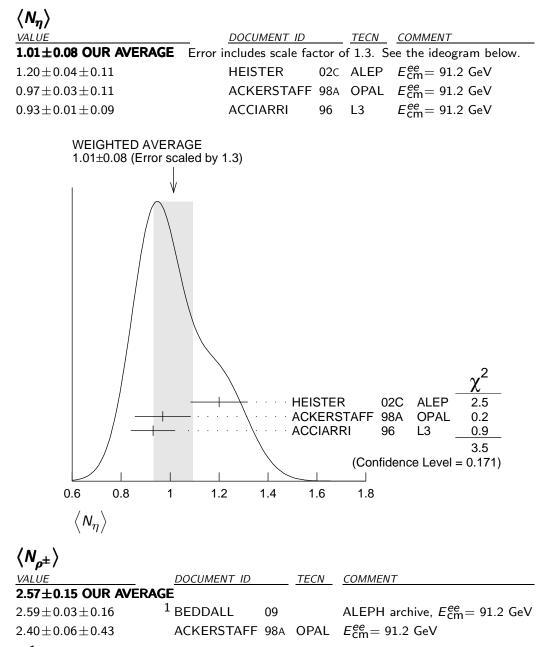
VALUE	<u>CL%</u>	DOCUMENT ID		TECN	COMMENT
<1.8 × 10 <sup>-6</sup>	95	<sup>1</sup> ABBIENDI	991	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
					-

<sup>1</sup>ABBIENDI 991 give the 95%CL limit on the partial width  $\Gamma(Z^0 \rightarrow p\mu) < 4.4$  KeV and we have transformed it into a branching ratio.

### AVERAGE PARTICLE MULTIPLICITIES IN HADRONIC Z DECAY

Summed over particle and antiparticle, when appropriate.

$\langle N_{\gamma} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$20.97 \pm 0.02 \pm 1.15$	ACKERSTAFF	98A	OPAL	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$\langle N_{\pi^{\pm}} \rangle$			TECN	COMMENT
<u>VALUE</u> 17.03 ±0.16 OUR AVERAGE	DOCUMENT ID		<u>TECN</u>	COMMENT
$17.007 \pm 0.209$	ABE	04C	SLD	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
$17.26\ \pm 0.10\ \pm 0.88$	ABREU	98L	DLPH	$E_{\rm cm}^{ee}=$ 91.2 GeV
$17.04 \pm 0.31$	BARATE	98v	ALEP	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$17.05 \pm 0.43$	AKERS	<b>94</b> P	OPAL	$E_{\rm Cm}^{ee}=$ 91.2 GeV
$\langle N_{\pi^0} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
9.76±0.26 OUR AVERAGE				
$9.55 \pm 0.06 \pm 0.75$	ACKERSTAFF	98A	OPAL	$E_{\rm cm}^{ee}$ = 91.2 GeV
$9.63\!\pm\!0.13\!\pm\!0.63$	BARATE	97J	ALEP	$E_{\rm cm}^{ee}$ = 91.2 GeV
$9.90\!\pm\!0.02\!\pm\!0.33$	ACCIARRI	96	L3	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$9.2\ \pm 0.2\ \pm 1.0$	ADAM	96	DLPH	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
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<sup>1</sup> BEDDALL 09 analyse 3.2 million hadronic Z decays as archived by ALEPH collaboration and report a value of  $2.59 \pm 0.03 \pm 0.15 \pm 0.04$ . The first error is statistical, the second systematic, and the third arises from extrapolation to full phase space. We combine the systematic errors in quadrature.

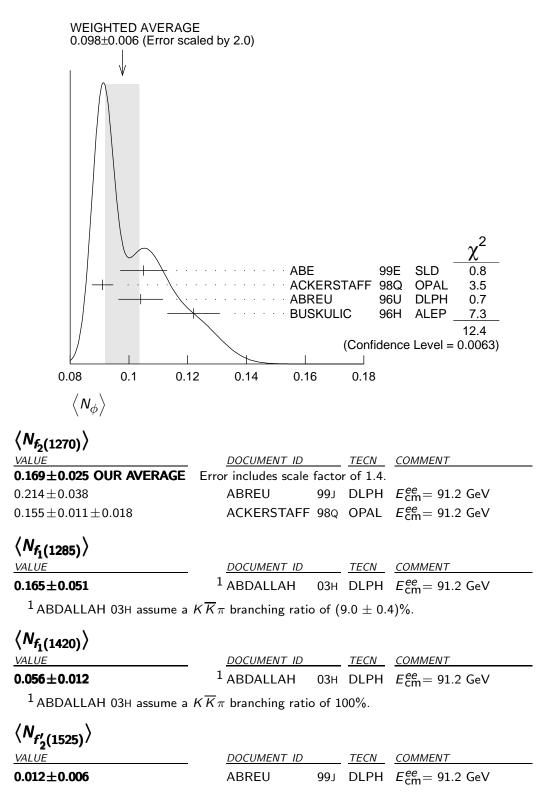
# $\left< \textit{N}_{\rho^0} \right>$

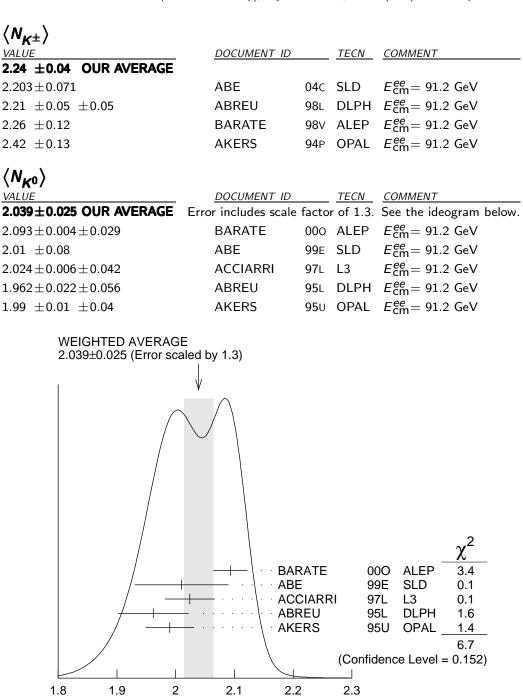
VALUE	DOCUMENT ID	TECN	COMMENT
$1.24\pm0.10$ OUR AVERAGE	Error includes scale fa	ctor of 1.1.	
$1.19 \pm 0.10$	ABREU	99J DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
$1.45\!\pm\!0.06\!\pm\!0.20$	BUSKULIC	96H ALEP	$E_{\rm cm}^{ee}$ = 91.2 GeV

Citation: M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018) and 2019 update

$\langle N_{\omega} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$1.02\pm0.06$ OUR AVERAGE				- 20
$1.00 \pm 0.03 \pm 0.06$	HEISTER			$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$1.04 \pm 0.04 \pm 0.14$	ACKERSTAFF	98A	OPAL	<i>E</i> <sup><i>ee</i></sup> <sub>cm</sub> = 91.2 GeV
$1.17\!\pm\!0.09\!\pm\!0.15$	ACCIARRI	<b>97</b> D	L3	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$\langle N_{\eta'} \rangle$				
VALUE	DOCUMENT ID			COMMENT
0.17 $\pm$ 0.05 OUR AVERAGE	Error includes scale			
$0.14 \pm 0.01 \pm 0.02$				$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$0.25 \pm 0.04$	<sup>1</sup> ACCIARRI			$E_{\rm Cm}^{ee}=$ 91.2 GeV
$\bullet \bullet \bullet$ We do not use the follow	ing data for averages	, fits,	limits, e	etc. • • •
$0.068\!\pm\!0.018\!\pm\!0.016$	<sup>2</sup> BUSKULIC	<b>92</b> D	ALEP	$E_{\rm cm}^{ee}=$ 91.2 GeV
<sup>1</sup> ACCIARRI 97D obtain this $\gamma' \rightarrow \rho^0 \gamma$ . <sup>2</sup> BUSKULIC 92D obtain this		the tv	vo decay	channels $\eta'  o \pi^+ \pi^- \eta$
- BUSKULIC 92D obtain this				
$\langle N_{f_0(980)} \rangle$				
⟨N <sub>f0</sub> (980)⟩ <sub>VALUE</sub>	DOCUMENT ID		TECN	COMMENT
$\langle N_{f_0(980)} \rangle$	DOCUMENT ID			
⟨N <sub>f0</sub> (980)⟩ <sub>VALUE</sub>				<u>COMMENT</u> E <sup>ce</sup> <sub>cm</sub> = 91.2 GeV
⟨ <i>N<sub>f0(980)</sub></i> ⟩ <u>VALUE</u> 0.147±0.011 OUR AVERAGE	<u>DOCUMENT ID</u> ABREU	99J	DLPH	
<pre>     ⟨N<sub>f0(980)</sub>⟩     <sub>VALUE</sub>     0.147±0.011 OUR AVERAGE     0.164±0.021 </pre>	<u>DOCUMENT ID</u> ABREU	99J	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
<pre></pre>	<u>DOCUMENT ID</u> ABREU	99J 98Q	DLPH OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
<pre></pre>	<u>DOCUMENT ID</u> ABREU ACKERSTAFF <u>DOCUMENT ID</u>	99J 98Q	DLPH OPAL <u>TECN</u>	E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV E <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$ \frac{\langle N_{f_0(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.141 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_0(980)} \pm \rangle}{\overset{VALUE}{0.27 \pm 0.04 \pm 0.10}} \\ \frac{\langle N_{\phi} \rangle}{\langle N_{\phi} \rangle} $	<u>DOCUMENT ID</u> ABREU ACKERSTAFF <u>DOCUMENT ID</u> ACKERSTAFF	99J 98Q 98A	DLPH OPAL <u>TECN</u> OPAL	$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{\rm cm}^{ee}} = 91.2 \text{ GeV}$
$\frac{\langle N_{f_{0}(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.141 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_{0}(980)} \pm \rangle}{\overset{VALUE}{0.27 \pm 0.04 \pm 0.10}} \\ \frac{\langle N_{\phi} \rangle}{\overset{VALUE}{0.27 \pm 0.04 \pm 0.$	DOCUMENT ID ABREU ACKERSTAFF <u>DOCUMENT ID</u> ACKERSTAFF	99J 98Q 98A	DLPH OPAL <u>TECN</u> OPAL <u>TECN</u>	$E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $\underline{COMMENT}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $\underline{COMMENT}$
$\frac{\langle N_{f_0(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.164 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_0(980)} \pm \rangle}{VALUE} \\ 0.27 \pm 0.04 \pm 0.10 \\ \frac{\langle N_{\phi} \rangle}{VALUE} \\ 0.098 \pm 0.006 \text{ OUR AVERAGE} \\ \end{array}$	DOCUMENT ID ABREU ACKERSTAFF DOCUMENT ID ACKERSTAFF DOCUMENT ID Error includes scale	99J 98Q 98A 98A	DLPH OPAL <u>TECN</u> OPAL <u>TECN</u> r of 2.0.	$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{\rm cm}^{ee}} = 91.2 \text{ GeV}$ $\frac{COMMENT}{See \text{ the ideogram below.}}$
$\frac{\langle N_{f_0(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.164 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_0(980)} \pm \rangle}{V_{a_0(980)} \pm \rangle} \\ \frac{\langle N_{a_0(980)} \pm 0.004 \pm 0.100}{V_{ALUE}} \\ \frac{\langle N_{\phi} \rangle}{V_{ALUE}} \\ 0.098 \pm 0.006 \text{ OUR AVERAGE} \\ 0.105 \pm 0.008 \\ \end{array}$	DOCUMENT ID ABREU ACKERSTAFF DOCUMENT ID ACKERSTAFF DOCUMENT ID Error includes scale ABE	99J 98Q 98A 98A facto 99E	DLPH OPAL <u>TECN</u> OPAL <u>TECN</u> r of 2.0. SLD	$E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ $\frac{COMMENT}{See \text{ the ideogram below.}}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$
$\frac{\langle N_{f_0(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.141 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_0(980)} \pm \rangle}{\overset{VALUE}{0.27 \pm 0.04 \pm 0.10}} \\ \frac{\langle N_{\phi} \rangle}{\overset{VALUE}{0.098 \pm 0.006 \text{ OUR AVERAGE}}} \\ 0.105 \pm 0.008 \\ 0.091 \pm 0.002 \pm 0.003 \\ \end{array}$	DOCUMENT ID ABREU ACKERSTAFF DOCUMENT ID ACKERSTAFF DOCUMENT ID Error includes scale ABE ACKERSTAFF	99J 98Q 98A 98A facto 99E 98Q	DLPH OPAL <u>TECN</u> OPAL TECN of 2.0. SLD OPAL	$E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$
$\frac{\langle N_{f_0(980)} \rangle}{\overset{VALUE}{0.147 \pm 0.011 \text{ OUR AVERAGE}}} \\ 0.164 \pm 0.021 \\ 0.164 \pm 0.007 \pm 0.011 \\ \frac{\langle N_{a_0(980)} \pm \rangle}{V_{a_0(980)} \pm \rangle} \\ \frac{\langle N_{a_0(980)} \pm 0.004 \pm 0.100}{V_{ALUE}} \\ \frac{\langle N_{\phi} \rangle}{V_{ALUE}} \\ 0.098 \pm 0.006 \text{ OUR AVERAGE} \\ 0.105 \pm 0.008 \\ \end{array}$	DOCUMENT ID ABREU ACKERSTAFF DOCUMENT ID ACKERSTAFF DOCUMENT ID Error includes scale ABE	99J 98Q 98A 98A facto 99E 98Q 96U	DLPH OPAL <u>TECN</u> OPAL TECN r of 2.0. SLD OPAL DLPH	$E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$ $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ $\frac{COMMENT}{See \text{ the ideogram below.}}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$

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$$\langle N_{K^0} \rangle$$
  
 $\langle N_{K^*(892)\pm} \rangle$ 

<u>VALUE</u> 0.72 ±0.05 OUR AVERAGE	DOCUMENT ID	TECN	COMMENT
$\begin{array}{c} 0.712 \pm 0.031 \pm 0.059 \\ 0.72 \pm 0.02 \pm 0.08 \end{array}$	ABREU ACTON		$E_{cm}^{ee}$ = 91.2 GeV $E_{cm}^{ee}$ = 91.2 GeV

Citation: M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018) and 2019 update

Citation. IN. Tanabasin et al. (Farticle L	vata Group), r nys. r	EV. D	56, 05000	
$\langle N_{\kappa^{*}(892)^{0}} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.739 \pm 0.022$ OUR AVERAGE				
$0.707 \pm 0.041$	ABE			$E_{\rm cm}^{ee}=$ 91.2 GeV
$0.74 \pm 0.02 \pm 0.02$	ACKERSTAFF	<b>97</b> S	OPAL	$E_{\rm cm}^{ee}$ = 91.2 GeV
$0.77 \pm 0.02 \pm 0.07$	ABREU	<b>96</b> U	DLPH	$E_{\rm cm}^{ee}$ = 91.2 GeV
$0.83 \pm 0.01 \pm 0.09$	BUSKULIC	96H	ALEP	$E_{\rm cm}^{ee}$ = 91.2 GeV
$0.97\ \pm 0.18\ \pm 0.31$	ABREU	93	DLPH	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$\langle N_{\kappa_2^*(1430)} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.073±0.023	ABREU	99J	DLPH	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$\bullet \bullet \bullet$ We do not use the following of	data for averages	s, fits,	limits, e	etc. • • •
$0.19\ \pm 0.04\ \pm 0.06$	<sup>1</sup> AKERS	95X	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
$^1$ AKERS 95X obtain this value fo	r <i>x</i> < 0.3.			
$\langle N_{D^{\pm}} \rangle$				
<u>VALUE</u> 0.187±0.020 OUR AVERAGE Erro	DOCUMENT ID			
				•
$0.170 \pm 0.009 \pm 0.014$				$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$
$0.251 \pm 0.026 \pm 0.025$ $0.199 \pm 0.019 \pm 0.024$	BUSKULIC <sup>1</sup> ABREU			$E_{cm}^{ee} = 91.2 \text{ GeV}$ $E_{cm}^{ee} = 91.2 \text{ GeV}$
	ADREU	951	DLFH	$E_{cm} = 91.2 \text{ GeV}$
<sup>1</sup> See ABREU 95 (erratum).				
WEIGHTED AVERAGE				
0.187±0.020 (Error scaled	by 1.5)			
$\wedge$				
				2
				_χ
		XAN		SR OPAL 1.1
		SKULI REU	C 94 93	4J ALEP 3.1 3I DLPH 0.2
	AD	10	9.	4.3
	$\sim$		(Confid	lence Level = 0.114)
0.1 0.15 0.2 0.2		25	 0.4	
	25 0.3 0	.35	0.4	
$\left< \textit{N}_{D^{\pm}} \right>$				

Citation: M. Tanabashi et al. (Particle Data Group), Phys. Rev. D 98, 030001 (2018) and 2019 update

<b>(Ν<sub>D</sub>0)</b> VALUE	DOCUMENT ID		TECN	COMMENT
$0.462 \pm 0.026$ OUR AVERAGE				
$0.465 \!\pm\! 0.017 \!\pm\! 0.027$	ALEXANDER	<b>96</b> R	OPAL	$E_{\rm cm}^{ee}=$ 91.2 GeV
$0.518 \!\pm\! 0.052 \!\pm\! 0.035$	BUSKULIC	94J		$E_{\rm Cm}^{ee} =$ 91.2 GeV
$0.403 \!\pm\! 0.038 \!\pm\! 0.044$	<sup>1</sup> ABREU	931	DLPH	$E_{\rm cm}^{ee}=$ 91.2 GeV
$^{1}$ See ABREU 95 (erratum).				
$\langle N_{D^{\pm}} \rangle$				
S VALUE	DOCUMENT ID		TECN	COMMENT
$0.131 \pm 0.010 \pm 0.018$	ALEXANDER	<b>96</b> R	OPAL	$E_{\rm cm}^{ee}=$ 91.2 GeV
$\langle N_{D^*(2010)^{\pm}} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
<b>0.183 ±0.008 OUR AVERAG</b> 0.1854±0.0041±0.0091		00-		<i>E<sup>ee</sup></i> = 91.2 GeV
$0.1854 \pm 0.0041 \pm 0.0091$ $0.187 \pm 0.015 \pm 0.013$	BUSKULIC			$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $E_{\rm cm}^{ee} = 91.2 \text{ GeV}$
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	<sup>2</sup> ABREU			$E_{\rm cm}^{ee} = 91.2 \text{ GeV}$ $E_{\rm cm}^{ee} = 91.2 \text{ GeV}$
$^{1}$ ACKERSTAFF 98E system				
branching ratios $B(D^{*+} \rightarrow 0.0012)$	,			
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)^+} \rangle$	,			
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ VALUE (units 10 <sup>-3</sup> )	DOCUMENT ID			COMMENT
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)^+} \rangle$	DOCUMENT ID			
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ VALUE (units 10 <sup>-3</sup> )	<u>DOCUMENT ID</u> wing data for average	s, fits,	limits, e	
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> ••• We do not use the follow	DOCUMENT ID wing data for averages $^1$ ACKERSTAFF this value for $x > 0.6$	s, fits, 97w	limits, e	etc. • • • • $E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> • • We do not use the follow $2.9^{+0.7}_{-0.6} \pm 0.2$ <sup>1</sup> ACKERSTAFF 97W obtain	<u>DOCUMENT ID</u> wing data for averages <sup>1</sup> ACKERSTAFF this value for x> 0.6 D* K final states.	s, fits, 97w and v	limits, o OPAL vith the	etc. • • • $E_{Cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)}+\rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> • • We do not use the follow 2.9 <sup>+0.7</sup> ±0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*}\rangle$ <u>VALUE</u>	$\frac{DOCUMENT \ ID}{1}$ wing data for averages <sup>1</sup> ACKERSTAFF this value for $x > 0.6$ $D^* K$ final states. <u>DOCUMENT ID</u>	s, fits, 97w and v	Iimits, o OPAL vith the <u>TECN</u>	etc. • • • $E_{Cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)}+\rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> ••• We do not use the follow $2.9^{+0.7}_{-0.6}\pm0.2$ <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*}\rangle$ <u>VALUE</u> 0.28±0.01±0.03	$\frac{DOCUMENT \ ID}{1}$ wing data for averages $1 \text{ ACKERSTAFF}$ this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1}$ ABREU	s, fits, 97w and v 95R	limits, o OPAL vith the <u>TECN</u> DLPH	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> • • We do not use the follow $2.9^{+0.7}_{-0.6} \pm 0.2$ <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*} \rangle$	$\frac{DOCUMENT \ ID}{1}$ wing data for averages $1 \text{ ACKERSTAFF}$ this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1}$ ABREU	s, fits, 97w and v 95R	limits, o OPAL vith the <u>TECN</u> DLPH	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)}+\rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> ••• We do not use the follow 2.9 <sup>+0.7</sup> $\pm$ 0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*}\rangle$ <u>VALUE</u> 0.28±0.01±0.03 <sup>1</sup> ABREU 95R quote this val $\langle N_{J/\psi(1S)}\rangle$	DOCUMENT ID wing data for averages 1 ACKERSTAFF this value for $x > 0.6$ $D^* K$ final states. $\frac{DOCUMENT ID}{1}$ ABREU ue for a flavor-average	s, fits, 97W and v 95R ed exc	imits, o OPAL vith the <u>TECN</u> DLPH	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee} = 91.2 \text{ GeV}}$ te.
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)+} \rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> • • We do not use the follow $2.9^{+0.7}_{-0.6} \pm 0.2$ <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*} \rangle$ <u>VALUE</u> 0.28 \pm 0.01 \pm 0.03 <sup>1</sup> ABREU 95R quote this val $\langle N_{J/\psi(1S)} \rangle$ <u>VALUE</u>	$\frac{DOCUMENT \ ID}{M}$ wing data for averages $\frac{1}{ACKERSTAFF}$ this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1 \ ABREU}$ ue for a flavor-average $\frac{DOCUMENT \ ID}{DOCUMENT \ ID}$	s, fits, 97W and v 95R ed exc	imits, o OPAL vith the <u>TECN</u> DLPH tited stat	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ te. $\underline{COMMENT}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)}+ \rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> ••• We do not use the follow 2.9 <sup>+0.7</sup> <sub>-0.6</sub> ±0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*} \rangle$ <u>VALUE</u> 0.28±0.01±0.03 <sup>1</sup> ABREU 95R quote this val $\langle N_{J/\psi(1S)} \rangle$ <u>VALUE</u> 0.0056±0.0003±0.0004	$\frac{DOCUMENT \ ID}{Ming}$ wing data for averages 1 ACKERSTAFF this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1 \ ABREU}$ ue for a flavor-average $\frac{DOCUMENT \ ID}{1 \ ALEXANDER}$	s, fits, 97W and v 95R ed exc 96B	imits, o OPAL vith the <u>TECN</u> DLPH cited stat	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee} = 91.2 \text{ GeV}}$ te. $\frac{COMMENT}{E_{cm}^{ee} = 91.2 \text{ GeV}}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). (N <sub>D<sub>s1</sub>(2536)+) VALUE (units 10<sup>-3</sup>) • • We do not use the follow 2.9<sup>+0.7</sup>±0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> (N<sub>B*</sub>) <u>VALUE</u> 0.28±0.01±0.03 <sup>1</sup> ABREU 95R quote this val (N<sub>J</sub>/ψ(1S)) <u>VALUE</u> 0.0056±0.0003±0.0004 <sup>1</sup> ALEXANDER 96B identify</sub>	$\frac{DOCUMENT \ ID}{Ming}$ wing data for averages 1 ACKERSTAFF this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1 \ ABREU}$ ue for a flavor-average $\frac{DOCUMENT \ ID}{1 \ ALEXANDER}$	s, fits, 97W and v 95R ed exc 96B	imits, o OPAL vith the <u>TECN</u> DLPH cited stat	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee} = 91.2 \text{ GeV}}$ te. $\frac{COMMENT}{E_{cm}^{ee} = 91.2 \text{ GeV}}$
0.0012. <sup>2</sup> See ABREU 95 (erratum). $\langle N_{D_{s1}(2536)}+\rangle$ <u>VALUE (units 10<sup>-3</sup>)</u> ••• We do not use the follow 2.9 <sup>+0.7</sup> <sub>-0.6</sub> ±0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> $\langle N_{B^*}\rangle$ <u>VALUE</u> 0.28±0.01±0.03 <sup>1</sup> ABREU 95R quote this val $\langle N_{J/\psi(1S)}\rangle$ <u>VALUE</u> 0.0056±0.0003±0.0004 <sup>1</sup> ALEXANDER 96B identify $\langle N_{\psi(2S)}\rangle$	$\frac{DOCUMENT \ ID}{Ming data for averages}$ wing data for averages $1_{ACKERSTAFF}$ this value for x> 0.6 D* K final states. $\frac{DOCUMENT \ ID}{1_{ABREU}}$ ue for a flavor-average $\frac{DOCUMENT \ ID}{1_{ALEXANDER}}$ $J/\psi(1S) from the definition of t$	s, fits, 97W and v 95R ed exc 96B ecays i	imits, o OPAL vith the <u>TECN</u> DLPH cited stat	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ te. $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ on pairs.
0.0012. <sup>2</sup> See ABREU 95 (erratum). (N <sub>D<sub>s1</sub>(2536)+) VALUE (units 10<sup>-3</sup>) • • We do not use the follow 2.9<sup>+0.7</sup>±0.2 <sup>1</sup> ACKERSTAFF 97W obtain width is saturated by the <i>L</i> (N<sub>B*</sub>) <u>VALUE</u> 0.28±0.01±0.03 <sup>1</sup> ABREU 95R quote this val (N<sub>J</sub>/ψ(1S)) <u>VALUE</u> 0.0056±0.0003±0.0004 <sup>1</sup> ALEXANDER 96B identify</sub>	DOCUMENT ID         wing data for averages         1 ACKERSTAFF         this value for $x > 0.6$ D* K final states.         DOCUMENT ID         1 ABREU         ue for a flavor-average         DOCUMENT ID         1 ALEXANDER $J/\psi(1S)$ from the design	s, fits, 97W and v 95R ed exc 96B ecays i	limits, o OPAL vith the <u>TECN</u> DLPH tited stat <u>TECN</u> OPAL into lept	etc. • • • $E_{cm}^{ee} = 91.2 \text{ GeV}$ assumption that its dec $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ te. $\frac{COMMENT}{E_{cm}^{ee}} = 91.2 \text{ GeV}$ on pairs.

$\langle N_{\rho} \rangle$				
VALUE	DOCUMENT ID	7	TECN	COMMENT
1.046±0.026 OUR AVERAGE				
$1.054 \pm 0.035$	ABE			$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$1.08 \pm 0.04 \pm 0.03$	ABREU			$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$1.00 \pm 0.07$	BARATE			$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$0.92 \pm 0.11$	AKERS	94P C	OPAL	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$\langle N_{\Delta(1232)^{++}} \rangle$				
VALUE 0.087±0.033 OUR AVERAGE	<u>DOCUMENT ID</u> Error includes scale			COMMENT
$0.079 \pm 0.009 \pm 0.011$	ABREU			$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$0.22 \pm 0.04 \pm 0.04$	ALEXANDER	95D C	JPAL	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$\langle N_A \rangle$				
VALUE	DOCUMENT ID			
0.388±0.009 OUR AVERAGE	Error includes scale	factor o	of 1.7.	See the ideogram below.
$0.404 \pm 0.002 \pm 0.007$	BARATE	000 A	<b>ALEP</b>	$E_{\rm cm}^{ee}$ = 91.2 GeV
$0.395 \pm 0.022$	ABE	99e S	SLD	$E_{\rm cm}^{ee}$ = 91.2 GeV
$0.364 \pm 0.004 \pm 0.017$	ACCIARRI	97∟ L	_3	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$0.374 \pm 0.002 \pm 0.010$	ALEXANDER	97D C	DPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.2 GeV
$0.357 \pm 0.003 \pm 0.017$	ABREU	93L D	DLPH	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
WEIGHTED AVERAGE 0.388±0.009 (Error scale		CIARRI EXANDE REU	93	E SLD 0.1 L L3 1.9 D OPAL 1.9
			`	
0.3 0.35	0.4 0.45		0.5	
$\langle N_A \rangle$				

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⟨ <i>N<sub>A(1520)</sub>⟩</i> <sup>VALUE</sup> 0.0224±0.0027 OUR AVERAGE	DOCUMENT ID		<u>TECN</u>	COMMENT
$0.029 \ \pm 0.005 \ \pm 0.005$	ABREU	<b>00</b> P	DLPH	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$0.0213\!\pm\!0.0021\!\pm\!0.0019$	ALEXANDER	<b>97</b> D	OPAL	$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$\langle N_{\Sigma^+} \rangle$ VALUE	DOCUMENT ID		TECN	COMMENT
0.107±0.010 OUR AVERAGE	DOCOMENT		TLCN	COMMENT
$0.114\!\pm\!0.011\!\pm\!0.009$	ACCIARRI	L00	L3	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$0.099\!\pm\!0.008\!\pm\!0.013$	ALEXANDER	97E	OPAL	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$\langle N_{\Sigma^{-}} \rangle$ $\frac{VALUE}{0.082 \pm 0.007} \text{ OUR AVERACE}$	DOCUMENT ID		TECN	COMMENT
<u>VALUE</u> 0.082±0.007 OUR AVERAGE		00P		
<u>VALUE</u> 0.082±0.007 OUR AVERAGE 0.081±0.002±0.010	ABREU		DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
$\frac{V_{ALUE}}{0.082 \pm 0.007 \text{ OUR AVERAGE}}$ $0.081 \pm 0.002 \pm 0.010$ $0.083 \pm 0.006 \pm 0.009$ $\left< N_{\Sigma^{+}+\Sigma^{-}} \right>$ $\frac{V_{ALUE}}{0.181 \pm 0.018 \text{ OUR AVERAGE}}$	ABREU ALEXANDER <u>DOCUMENT ID</u>	97e	DLPH OPAL <u>TECN</u>	E <sup>ee</sup> <sub>Cm</sub> = 91.2 GeV E <sup>ee</sup> <sub>Cm</sub> = 91.2 GeV <u>COMMENT</u>
$\frac{V_{ALUE}}{0.082 \pm 0.007 \text{ OUR AVERAGE}}$ $0.081 \pm 0.002 \pm 0.010$ $0.083 \pm 0.006 \pm 0.009$ $\left< N_{\Sigma^{+}+\Sigma^{-}} \right>$ $\frac{V_{ALUE}}{0.181 \pm 0.018 \text{ OUR AVERAGE}}$	ABREU ALEXANDER <u>DOCUMENT ID</u>	97e	DLPH OPAL <u>TECN</u>	$E_{cm}^{ee}$ = 91.2 GeV $E_{cm}^{ee}$ = 91.2 GeV

 $^1\,\text{We}$  have combined the values of  $\langle \textit{N}_{\Sigma^+}\rangle$  and  $\langle \textit{N}_{\Sigma^-}\rangle$  from ALEXANDER 97E adding the statistical and systematic errors of the two final states separately in quadrature. If isospin symmetry is assumed this value becomes 0.174  $\pm$  0.010  $\pm$  0.015.

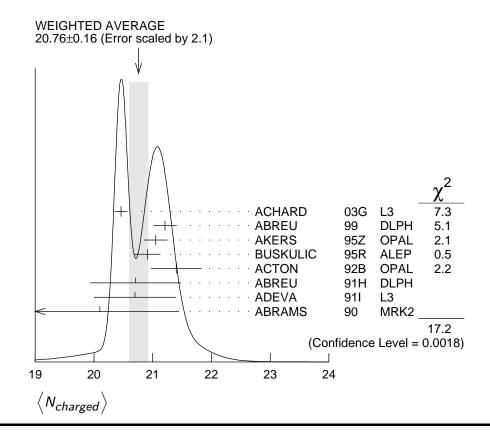
$\langle N_{50} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.076 \pm 0.010$ OUR AVERAGE				
$0.095 \!\pm\! 0.015 \!\pm\! 0.013$	ACCIARRI	00J	L3	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$0.071\!\pm\!0.012\!\pm\!0.013$	ALEXANDER	97E	OPAL	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$0.070 \pm 0.010 \pm 0.010$	ADAM	<b>96</b> B	DLPH	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$\langle N_{(\Sigma^++\Sigma^-+\Sigma^0)/3} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.084 \pm 0.005 \pm 0.008$	ALEXANDER	97E	OPAL	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$\langle N_{\Sigma(1385)^+}  angle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.0239 \pm 0.0009 \pm 0.0012$	ALEXANDER	<b>97</b> D	OPAL	$E_{\rm cm}^{ee}$ = 91.2 GeV
$\langle N_{\Sigma(1385)^{-}} \rangle$	DOCUMENT ID		TECN	COMMENT
VALUE	DOCUMENT ID		<u>TECN</u>	<u>COMMENT</u>
$0.0240 \pm 0.0010 \pm 0.0014$	ALEXANDER	<b>97</b> D	OPAL	$E_{\rm Cm}^{ee} = 91.2  { m GeV}$

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$\langle N_{\Sigma(1385)^++\Sigma(1385)^-} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.046 ±0.004 OUR AVERAGE	Error includes sca			•
$0.0479 \!\pm\! 0.0013 \!\pm\! 0.0026$	ALEXANDER	<b>97</b> D	OPAL	$E_{\rm cm}^{ee} =$ 91.2 GeV
$0.0382\!\pm\!0.0028\!\pm\!0.0045$	ABREU	<b>95</b> 0	DLPH	$E_{\rm Cm}^{ee} = 91.2  {\rm GeV}$
⟨ <b>N</b> <sub>=</sub> -⟩ <sub>VALUE</sub>	DOCUMENT ID		TECN	COMMENT
0.0258±0.0009 OUR AVERAGE	DOCOMENT ID		TLCN	COMMENT
$0.0247 \pm 0.0009 \pm 0.0025$	ABDALLAH	06E	DLPH	$E_{\rm cm}^{ee} = 91.2 \; { m GeV}$
$0.0259 \pm 0.0004 \pm 0.0009$				$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$\langle N_{\equiv(1530)^0} \rangle$				
VALUE	DOCUMENT ID			
	Error includes sca			
$0.0045 \pm 0.0005 \pm 0.0006$	ABDALLAH			$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$0.0068 \pm 0.0005 \pm 0.0004$	ALEXANDER	<b>97</b> D	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV
$\langle N_{\Omega^{-}} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
0.00164±0.00028 OUR AVERAGE				- 00
$0.0018 \pm 0.0003 \pm 0.0002$				$E_{\rm cm}^{ee} = 91.2 {\rm GeV}$
$0.0014 \pm 0.0002 \pm 0.0004$	ADAM	<b>96</b> B	DLPH	$E_{\rm cm}^{ee}$ = 91.2 GeV
$\langle N_{\Lambda^+} \rangle$				
VALUE	DOCUMENT ID		TECN	COMMENT
$0.078 \pm 0.012 \pm 0.012$	ALEXANDER	<b>96</b> R	OPAL	$E_{\rm cm}^{ee}$ = 91.2 GeV
$\langle N_{\overline{D}} \rangle$				
VALUE (units $10^{-6}$ )	DOCUMENT ID		TECN	COMMENT
• • We do not use the following of				
-	<sup>1</sup> SCHAEL			$E_{ m cm}^{ee}=91.2~{ m GeV}$
<sup>1</sup> SCHAEL 06A obtain this anti- anti-deuteron momentum range	deuteron produc from 0.62 to 1.0	tion r 03 Ge <sup>v</sup>	ate per V/c.	hadronic $Z$ decay in the
$\langle N_{charged} \rangle$				

\' <b>"</b> charged /			
VALUE	DOCUMENT ID	TECN	COMMENT
20.76±0.16 OUR AVERAGE	Error includes scale facto	or of 2.1.	See the ideogram below.
$20.46 \!\pm\! 0.01 \!\pm\! 0.11$	ACHARD 03	G <b>L3</b>	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$21.21\!\pm\!0.01\!\pm\!0.20$	ABREU 99	DLPH	$E_{\rm cm}^{ee}$ = 91.2 GeV
$21.05 \pm 0.20$	AKERS 95	z OPAL	$E_{\rm cm}^{ee}$ = 91.2 GeV
$20.91 \!\pm\! 0.03 \!\pm\! 0.22$	BUSKULIC 95	r ALEP	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$21.40 \pm 0.43$	ACTON 92	b OPAL	$E_{\rm cm}^{ee} = 91.2  {\rm GeV}$
$20.71\!\pm\!0.04\!\pm\!0.77$	ABREU 91	h DLPH	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$20.7 \pm 0.7$	ADEVA 91	I L3	$E_{\rm cm}^{ee} = 91.2  { m GeV}$
$20.1 \ \pm 1.0 \ \pm 0.9$	ABRAMS 90	MRK2	$E_{\rm cm}^{ee}=$ 91.1 GeV

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### **Z HADRONIC POLE CROSS SECTION**

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). This quantity is defined as

$$\sigma_h^0 = \frac{12\pi}{M_7^2} \frac{\Gamma(e^+ e^-) \Gamma(\text{hadrons})}{\Gamma_7^2}$$

It is one of the parameters used in the Z lineshape fit.

VALUE	(nb)	EVTS	DOCUMENT ID		TECN	COMMENT		
41.541±0.037 OUR FIT								
41.501	$1\!\pm\!0.055$	4.10M	<sup>1</sup> ABBIENDI	01A	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
41.578	$8 \pm 0.069$	3.70M	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
41.535	$5\pm0.055$	3.54M	ACCIARRI	<b>00</b> C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
41.559	$9 \pm 0.058$	4.07M	<sup>2</sup> BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
• • •	We do not use	the following	g data for average	s, fits,	limits, e	etc. • • •		
42	$\pm 4$	450	ABRAMS	<b>89</b> B	MRK2	$E_{\rm cm}^{ee}$ = 89.2–93.0 GeV		
<sup>1</sup> A	$^1$ ABBIENDI 01A error includes approximately 0.031 due to statistics, 0.033 due to event							

selection systematics, 0.029 due to uncertainty in luminosity measurement, and 0.011 due to LEP energy uncertainty.

 $^2$  BARATE 00C error includes approximately 0.030 due to statistics, 0.026 due to experimental systematics, and 0.025 due to uncertainty in luminosity measurement.

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### **Z VECTOR COUPLINGS**

These quantities are the effective vector couplings of the Z to charged leptons and quarks. Their magnitude is derived from a measurement of the Z lineshape and the forward-backward lepton asymmetries as a function of energy around the Z mass. The relative sign among the vector to axial-vector couplings is obtained from a measurement of the Z asymmetry parameters,  $A_e$ ,  $A_{\mu}$ , and  $A_{\tau}$ . By convention the sign of  $g_A^e$  is fixed to be negative (and opposite to that of  $g^{\nu_e}$  obtained using  $\nu_e$  scattering measurements). For the light quarks, the sign of the couplings is assigned consistently with this assumption. The LEP/SLD-based fit values quoted below correspond to global nine- or five-parameter fits to lineshape, lepton forward-backward asymmetry, and  $A_e$ ,  $A_{\mu}$ , and  $A_{\tau}$  measurements. See the note "The Z boson" and ref. LEP-SLC 06 for details. Where  $p\overline{p}$  and ep data is quoted, OUR FIT value corresponds to a weighted average of this with the LEP/SLD fit result.

# g<sub>V</sub>

VALUE	<u>EVTS</u>	DOCUMENT ID	TECN	I COMMENT
$-0.03817 \pm 0.00047$ OUR FI	т			
$-0.058$ $\pm 0.016$ $\pm 0.007$	5026	<sup>1</sup> ACOSTA	05м CDF	$E_{\rm cm}^{p\overline{p}}$ = 1.96 TeV
$-0.0346 \pm 0.0023$	137.0K	<sup>2</sup> ABBIENDI	010 OPA	L $E_{\rm cm}^{ee} = 88-94$ GeV
$-0.0412 \pm 0.0027$	124.4k	<sup>3</sup> ACCIARRI	00C L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0400 \pm 0.0037$		BARATE	00C ALE	$P  E_{CM}^{ee} = 88-94  GeV$
$-0.0414\ \pm 0.0020$		<sup>4</sup> ABE	95J SLD	$E_{\rm cm}^{ee}$ = 91.31 GeV

<sup>1</sup>ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to  $e^+e^-$ , assuming the quark couplings are as predicted by the standard model. Higher order radiative corrections have not been taken into account.

<sup>2</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^3\,{\rm ACCIARRI}$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<sup>4</sup> ABE 95J obtain this result combining polarized Bhabha results with the  $A_{LR}$  measurement of ABE 94C. The Bhabha results alone give  $-0.0507 \pm 0.0096 \pm 0.0020$ .

$g_V^{\mu}$							
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT		
$-0.0367 \pm 0.0023$ OU	R FIT						
$-0.0388\substack{+0.0060\\-0.0064}$	182.8K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV		
$-0.0386 \!\pm\! 0.0073$	113.4k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV		
$-0.0362 \pm 0.0061$		BARATE	<b>00</b> C	ALEP	<i>E<sup>ee</sup></i> = 88–94 GeV		
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$							
$-0.0413 \pm 0.0060$	66143	<sup>3</sup> ABBIENDI	01K	OPAL	<i>E<sup>ee</sup></i> = 89–93 GeV		

<sup>1</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^2\,{\rm ACCIARRI}$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

 $^3$ ABBIENDI 01K obtain this from an angular analysis of the muon pair asymmetry which takes into account effects of initial state radiation on an event by event basis and of initial-final state interference.

$g_V^{\tau}$					
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$-0.0366 \pm 0.0010$ OUR	FIT				
$-0.0365 \!\pm\! 0.0023$	151.5K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.0384 \pm 0.0026$	103.0k	<sup>2</sup> ACCIARRI	00C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.0361\!\pm\!0.0068$		BARATE	<b>00</b> C	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

<sup>2</sup>ACCIARRI 00C use their measurement of the au polarization in addition to forwardbackward lepton asymmetries.

δV					
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$-0.03783 \pm 0.00041$ C	OUR FIT				
$-0.0358 \pm 0.0014$	471.3K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.0397 \pm 0.0020$	379.4k	<sup>2</sup> ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0397 \pm 0.0017$	340.8k	<sup>3</sup> ACCIARRI	<b>00</b> C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.0383\ \pm 0.0018$	500k	BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

<sup>1</sup>ABBIENDI 010 use their measurement of the au polarization in addition to the lineshape and forward-backward lepton asymmetries.

<sup>2</sup>Using forward-backward lepton asymmetries.

<sup>3</sup>ACCIARRI 00C use their measurement of the  $\tau$  polarization in addition to forwardbackward lepton asymmetries.

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VALUE	EVTS	DOCUMENT ID		TECN	COMMENT	
0.266±0.034 OUR AVERAGE						
$0.270 \!\pm\! 0.037$		<sup>1</sup> ANDREEV	18A	H1	$e^{\pm}p$	
$0.201\!\pm\!0.112$	156k	<sup>2</sup> ABAZOV	<b>11</b> D	D0	$E_{ m cm}^{p\overline{p}}=1.97~ m TeV$	
$\begin{array}{c} 0.24 & +0.28 \\ -0.11 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{\rm Cm}^{ee}=$ 88–94 GeV	
$0.399^{+0.152}_{-0.188}{\pm}0.066$	5026	<sup>4</sup> ACOSTA	<b>0</b> 5M	CDF	$E_{cm}^{p\overline{p}}$ = 1.96 TeV	
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$						
$0.14 \begin{array}{c} +0.09 \\ -0.09 \end{array}$		<sup>5</sup> ABRAMOWIC	Z16A	ZEUS		
$0.144 \substack{+ 0.066 \\ - 0.058}$		<sup>6</sup> ABT	16			
$0.27 \pm 0.13$	1500	<sup>7</sup> AKTAS	06	H1	$e^{\pm} p  ightarrow ~ \overline{ u}_e( u_e) X, \ \sqrt{s} pprox 300 \; { m GeV}$	

<sup>1</sup>ANDREEV 18A obtain this result in a combined electroweak and QCD analysis using all deep-inelastic  $e^+p$  and  $e^-p$  neutral current and charged current scattering cross sections published by the H1 Collaboration, including data with longitudinally polarized lepton beams.

<sup>2</sup>ABAZOV 11D study  $p\overline{p} \rightarrow Z/\gamma^* e^+ e^-$  events using 5 fb<sup>-1</sup> data at  $\sqrt{s} = 1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T > 25$  GeV, at least one electron in the central region and the di-electron mass in the range 50-1000 GeV. From the forward-backward asymmetry, determined as a function of

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the di-electron mass, they derive the axial and vector couplings of the *u*- and *d*- quarks and the value of  $\sin^2 \theta_{eff}^{\ell} = 0.2309 \pm 0.0008(\text{stat}) \pm 0.0006(\text{syst})$ .

- <sup>3</sup>LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.
- <sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.
- <sup>5</sup> ABRAMOWICZ 16A determine the  $Z^0$  couplings to *u* and *d*-quarks using the ZEUS polarised data from Run II together with the unpolarised data from both ZEUS and H1 Collaborations for Run I and unpolarised H1 data from Run II.
- <sup>6</sup>ABT 16 determine the  $Z^0$  couplings to u- and d-quarks using the same techniques and data as ABRAMOWICZ 16A but additionally use the published H1 polarised data.
- $^7$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  30,000 GeV<sup>2</sup>) and charged current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  15,000 GeV<sup>2</sup>) differential cross sections. In the determination of the *u*-quark couplings the electron and *d*-quark couplings are fixed to their standard model values.

<b>BV</b> VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT	_
$-0.38 \begin{array}{c} +0.04 \\ -0.05 \end{array}$ our an	<b>VERAGE</b>					
$-0.488 \pm 0.092$		<sup>1</sup> ANDREEV	-		$e^{\pm}p$	
$-0.351\!\pm\!0.251$	156k	<sup>2</sup> ABAZOV	<b>11</b> D	D0	$E^{p\overline{p}}_{ m cm}=$ 1.97 TeV	
$-0.33 \begin{array}{c} +0.05 \\ -0.07 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{\rm cm}^{ee}=$ 88–94 GeV	
$-0.226^{+0.635}_{-0.290}{\pm}0.090$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{\rm cm}^{p\overline{p}}$ = 1.96 TeV	
$\bullet \bullet \bullet$ We do not use th	e following	g data for averages	s, fits,	limits, e	etc. • • •	
$-0.41 \ \begin{array}{c} +0.25 \\ -0.20 \end{array}$		<sup>5</sup> ABRAMOWIC	<b>Z16</b> A	ZEUS		
$-0.503 \substack{+0.171 \\ -0.103}$		<sup>6</sup> ABT	16			
$-0.33 \pm 0.33$	1500	<sup>7</sup> AKTAS	06	H1	$e^{\pm} p  ightarrow ~ \overline{ u}_e( u_e) X$ ,	

- <sup>1</sup> ANDREEV 18A obtain this result in a combined electroweak and QCD analysis using all deep-inelastic  $e^+ p$  and  $e^- p$  neutral current and charged current scattering cross sections published by the H1 Collaboration, including data with longitudinally polarized lepton beams.
- <sup>2</sup>ABAZOV 11D study  $p\overline{p} \rightarrow Z/\gamma^* e^+ e^-$  events using 5 fb<sup>-1</sup> data at  $\sqrt{s} = 1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T > 25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the *u* and *d* quarks and the value of  $\sin^2\theta_{eff}^{\ell} = 0.2309 \pm 0.0008(\text{stat}) \pm 0.0006(\text{syst})$ .

<sup>3</sup>LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.

<sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.

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 $\sqrt{s} \approx 300 \text{ GeV}$ 

- <sup>5</sup> ABRAMOWICZ 16A determine the  $Z^0$  couplings to *u* and *d*-quarks using the ZEUS polarised data from Run II together with the unpolarised data from both ZEUS and H1 Collaborations for Run I and unpolarised H1 data from Run II.
- <sup>6</sup>ABT 16 determine the  $Z^0$  couplings to u- and d-quarks using the same techniques and \_data as ABRAMOWICZ 16A but additionally use the published H1 polarised data.
- <sup>7</sup>AKTAS 06 fit the neutral current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  30,000 GeV<sup>2</sup>) and charged current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  15,000 GeV<sup>2</sup>) differential cross sections. In the determination of the *d*-quark couplings the electron and *u*-quark couplings are fixed to their standard model values.

## Z AXIAL-VECTOR COUPLINGS

These quantities are the effective axial-vector couplings of the Z to charged leptons and quarks. Their magnitude is derived from a measurement of the Z lineshape and the forward-backward lepton asymmetries as a function of energy around the Z mass. The relative sign among the vector to axial-vector couplings is obtained from a measurement of the Z asymmetry parameters,  $A_e$ ,  $A_{\mu}$ , and  $A_{\tau}$ . By convention the sign of  $g_A^e$  is fixed to be negative (and opposite to that of  $g^{\nu_e}$  obtained using  $\nu_e$  scattering measurements). For the light quarks, the sign of the couplings is assigned consistently with this assumption. The LEP/SLD-based fit values quoted below correspond to global nine- or five-parameter fits to lineshape, lepton forward-backward asymmetry, and  $A_e$ ,  $A_{\mu}$ , and  $A_{\tau}$  measurements. See the note "The Z boson" and ref. LEP-SLC 06 for details. Where  $p\overline{p}$  and ep data is quoted, OUR FIT value corresponds to a weighted average of this with the LEP/SLD fit result.

# g<sub>A</sub>

VALUE	EVTS	DOCUMENT ID	TE	ECN	COMMENT
$-0.50111 \pm 0.00035$ OUR Fi	Т				
$-0.528$ $\pm 0.123$ $\pm 0.059$	5026	<sup>1</sup> ACOSTA	05м CI	DF	$E_{ m cm}^{p\overline{p}}$ = 1.96 TeV
$-0.50062\!\pm\!0.00062$	137.0K	<sup>2</sup> ABBIENDI	010 O	PAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$-0.5015 \pm 0.0007$	124.4k	<sup>3</sup> ACCIARRI	00C L3	3	$E_{\rm cm}^{ee}$ = 88–94 GeV
$-0.50166 \pm 0.00057$		BARATE	00C AI	LEP	$E_{\rm cm}^{ee}$ = 88–94 GeV
$-0.4977 \pm 0.0045$		<sup>4</sup> ABE	95J SL	D	$E_{\rm cm}^{ee} = 91.31 {\rm GeV}$

<sup>1</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q\overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to  $e^+e^-$ , assuming the quark couplings are as predicted by the standard model. Higher order radiative corrections have not been taken into account.

<sup>2</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^3\,\rm ACCIARRI$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<sup>4</sup> ABE 95J obtain this result combining polarized Bhabha results with the  $A_{LR}$  measurement of ABE 94C. The Bhabha results alone give  $-0.4968 \pm 0.0039 \pm 0.0027$ .

$g^{\mu}_{A}$					
VALUE	EVTS	DOCUMENT ID		TECN	<u>COMMENT</u>
$-0.50120\pm0.00054$ O	UR FIT				
$-0.50117 \!\pm\! 0.00099$	182.8K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.5009 \pm 0.0014$	113.4k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.50046 \!\pm\! 0.00093$		BARATE	<b>00</b> C	ALEP	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
• • • We do not use t	he followin	g data for average	s, fits,	limits, e	etc. ● ● ●

 $-0.520 \pm 0.015$  66143 <sup>3</sup> ABBIENDI 01K OPAL  $E_{cm}^{ee} = 89-93$  GeV

<sup>1</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^2\,{\rm ACCIARRI}$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<sup>3</sup>ABBIENDI 01K obtain this from an angular analysis of the muon pair asymmetry which takes into account effects of initial state radiation on an event by event basis and of initial-final state interference.

## $g_A^{\tau}$

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VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$-0.50204 \pm 0.00064$ O	UR FIT				
$-0.50165 \pm 0.00124$	151.5K	<sup>1</sup> ABBIENDI	010	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$-0.5023\ \pm 0.0017$	103.0k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$-0.50216\!\pm\!0.00100$		BARATE	<b>00</b> C	ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV

<sup>1</sup>ABBIENDI 010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^2\,{\rm ACCIARRI}$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<i>B</i> <sub>A</sub>					
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
-0.50123±0.00026 OUR FIT					
$-0.50089 \!\pm\! 0.00045$	471.3K	<sup>1</sup> ABBIENDI	010	OPAL	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.5007 \pm 0.0005$	379.4k	ABREU	00F	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.50153 \!\pm\! 0.00053$	340.8k	<sup>2</sup> ACCIARRI	<b>00</b> C	L3	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV
$-0.50150 \pm 0.00046$	500k	BARATE	<b>00</b> C	ALEP	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV

 $^{1}\,{\rm ABBIENDI}$  010 use their measurement of the  $\tau$  polarization in addition to the lineshape and forward-backward lepton asymmetries.

 $^2\,{\rm ACCIARRI}$  00C use their measurement of the  $\tau$  polarization in addition to forward-backward lepton asymmetries.

<i>B</i> <sup><i>U</i></sup> <u>VALUE</u> 0.519 <sup>+0.028</sup> OUR AVE	<u>EVTS</u>	DOCUMENT ID		<u>TECN</u>	COMMENT	_
$0.548 \pm 0.036$		<sup>1</sup> ANDREEV	18A		$e^{\pm} p$	
$0.501\!\pm\!0.110$	156k	<sup>2</sup> ABAZOV	<b>11</b> D	D0	$E_{ m cm}^{p\overline{p}}=$ 1.97 TeV	
$0.47 \begin{array}{c} +0.05 \\ -0.33 \end{array}$		<sup>3</sup> LEP-SLC	06		$E_{\rm cm}^{ee}=$ 88–94 GeV	
$0.441^{+0.207}_{-0.173}{\pm}0.067$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{\sf cm}^{p\overline{p}}$ = 1.96 TeV	

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.50 \begin{array}{c} +0.12 \\ -0.05 \end{array}$		<sup>5</sup> ABRAMOWIC	Z16A	ZEUS	
$0.532 \substack{+ 0.107 \\ - 0.063}$		<sup>6</sup> ABT	16		
$0.57\ \pm 0.08$	1500	<sup>7</sup> AKTAS	06	H1	$e^{\pm}p \rightarrow \overline{\nu}_{e}(\nu_{e})X,$

<sup>1</sup> ANDREEV 18A obtain this result in a combined electroweak and QCD analysis using all deep-inelastic  $e^+ p$  and  $e^- p$  neutral current and charged current scattering cross sections published by the H1 Collaboration, including data with longitudinally polarized lepton beams.

- <sup>2</sup>ABAZOV 11D study  $p\overline{p} \rightarrow Z/\gamma^* e^+ e^-$  events using 5 fb<sup>-1</sup> data at  $\sqrt{s} = 1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T > 25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the *u* and *d* quarks and the value of  $\sin^2 \theta_{eff}^{\ell} = 0.2309 \pm 0.0008(\text{stat}) \pm 0.0006(\text{syst})$ .
- <sup>3</sup>LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.
- <sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.
- <sup>5</sup> ABRAMOWICZ 16A determine the  $Z^0$  couplings to *u* and *d*-quarks using the ZEUS polarised data from Run II together with the unpolarised data from both ZEUS and H1 Collaborations for Run I and unpolarised H1 data from Run II.
- <sup>6</sup>ABT 16 determine the  $Z^0$  couplings to u- and d-quarks using the same techniques and data as ABRAMOWICZ 16A but additionally use the published H1 polarised data.
- $^7$  AKTAS 06 fit the neutral current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  30,000 GeV<sup>2</sup>) and charged current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  15,000 GeV<sup>2</sup>) differential cross sections. In the determination of the *u*-quark couplings the electron and *d*-quark couplings are fixed to their standard model values.

<b>gd</b> VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$-0.527 + 0.040_{-0.028}$ OUR AV		DOCOMENT ID		<u>TLCN</u>	
$-0.619 \pm 0.108$		<sup>1</sup> ANDREEV	18A		$e^{\pm}p$
$-0.497 \pm 0.165$	156k	<sup>2</sup> ABAZOV	<b>11</b> D	D0	$E_{ m cm}^{p\overline{p}}=1.97~ m TeV$
$-0.52 \   {+0.05 \atop -0.03}$		<sup>3</sup> LEP-SLC	06		$E_{\rm cm}^{ee}=$ 88–94 GeV
$-0.016^{+0.346}_{-0.536}{\pm}0.091$	5026	<sup>4</sup> ACOSTA	05м	CDF	$E_{cm}^{p\overline{p}}$ = 1.96 TeV
$\bullet \bullet \bullet$ We do not use th	e following	g data for averages	s, fits,	limits, e	etc. • • •
$-0.56 \ \begin{array}{c} +0.41 \\ -0.15 \end{array}$		<sup>5</sup> ABRAMOWIC	Z16A	ZEUS	
$-0.409\substack{+0.373 \\ -0.213}$		<sup>6</sup> <sub>ABT</sub>	16		
$-0.80 \pm 0.24$	1500	<sup>7</sup> AKTAS	06	H1	$e^{\pm} p  ightarrow rac{\overline{ u}_{e}( u_{e})X,}{\sqrt{s}pprox 300 \; { m GeV}}$

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- <sup>1</sup>ANDREEV 18A obtain this result in a combined electroweak and QCD analysis using all deep-inelastic  $e^+p$  and  $e^-p$  neutral current and charged current scattering cross sections published by the H1 Collaboration, including data with longitudinally polarized lepton beams.
- <sup>2</sup>ABAZOV 11D study  $p\overline{p} \rightarrow Z/\gamma^* e^+ e^-$  events using 5 fb<sup>-1</sup> data at  $\sqrt{s} = 1.96$  TeV. The candidate events are selected by requiring two isolated electromagnetic showers with  $E_T > 25$  GeV, at least one electron in the central region and the di-electron mass in the range 50–1000 GeV. From the forward-backward asymmetry, determined as a function of the di-electron mass, they derive the axial and vector couplings of the *u* and *d* quarks and the value of  $\sin^2\theta_{eff}^{\ell} = 0.2309 \pm 0.0008(\text{stat}) \pm 0.0006(\text{syst})$ .
- <sup>3</sup>LEP-SLC 06 is a combination of the results from LEP and SLC experiments using light quark tagging. s- and d-quark couplings are assumed to be identical.
- <sup>4</sup> ACOSTA 05M determine the forward-backward asymmetry of  $e^+e^-$  pairs produced via  $q \overline{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$  in 15 M( $e^+e^-$ ) effective mass bins ranging from 40 GeV to 600 GeV. These results are used to obtain the vector and axial-vector couplings of the Z to the light quarks, assuming the electron couplings are as predicted by the Standard Model. Higher order radiative corrections have not been taken into account.
- <sup>5</sup> ABRAMOWICZ 16A determine the  $Z^0$  couplings to *u* and *d*-quarks using the ZEUS polarised data from Run II together with the unpolarised data from both ZEUS and H1 Collaborations for Run I and unpolarised H1 data from Run II.
- <sup>6</sup>ABT 16 determine the  $Z^0$  couplings to u- and d-quarks using the same techniques and data as ABRAMOWICZ 16A but additionally use the published H1 polarised data.
- <sup>7</sup>AKTAS 06 fit the neutral current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  30,000 GeV<sup>2</sup>) and charged current (1.5  $\leq$  Q<sup>2</sup>  $\leq$  15,000 GeV<sup>2</sup>) differential cross sections. In the determination of the *d*-quark couplings the electron and *u*-quark couplings are fixed to their standard model values.

### Z COUPLINGS TO NEUTRAL LEPTONS

Averaging over neutrino species, the invisible Z decay width determines the effective neutrino coupling  $g^{\nu_{\ell}}$ . For  $g^{\nu_{e}}$  and  $g^{\nu_{\mu}}$ ,  $\nu_{e}e$  and  $\nu_{\mu}e$  scattering results are combined with  $g^{e}_{A}$  and  $g^{e}_{V}$  measurements at the Z mass to obtain  $g^{\nu_{e}}$  and  $g^{\nu_{\mu}}$  following NOVIKOV 93C.

## gνℓ

VALUE	DOCUMENT ID		COMMENT		
0.50076±0.00076	<sup>1</sup> LEP-SLC	06	$E_{\rm cm}^{ee}=$ 88–94 GeV		

<sup>1</sup> From invisible Z-decay width.

## $g^{\nu_e}$

VALUE	DOCUMENT ID	TECN	<u>СОММ</u>	ENT
$0.528 \pm 0.085$	<sup>1</sup> VILAIN 94	CHM2	From	$ u_{\mu} e$ and $ u_{e} e$ scattering
$^1$ VILAIN 94 derive this $1.05 \substack{+0.15 \\ -0.18}.$	value from their va	lue of $g^{l}$	${}^{\prime\mu}$ and	their ratio $g^{ u_e}/g^{ u_\mu} =$
$g^{ u_{\mu}}$				
VALUE	DOCUMENT I	D	TECN	COMMENT
0.502±0.017	<sup>1</sup> VILAIN	94	CHM2	From $ u_{\mu} e$ scattering

 $^1$  VILAIN 94 derive this value from their measurement of the couplings  $g_{\it A}^{e\,\nu_\mu}=-0.503\pm$ 0.017 and  $g_V^{e 
u \mu} = -$  0.035  $\pm$  0.017 obtained from  $u_\mu e$  scattering. We have re-evaluated

this value using the current PDG values for  $g^e_A$  and  $g^e_V$ .

## Z ASYMMETRY PARAMETERS

For each fermion-antifermion pair coupling to the Z these quantities are defined as

$$A_{f} = rac{2g_{V}^{f}g_{A}^{f}}{(g_{V}^{f})^{2} + (g_{A}^{f})^{2}}$$

where  $g_V^f$  and  $g_A^f$  are the effective vector and axial-vector couplings. For their relation to the various lepton asymmetries see the note "The Z boson" and ref. LEP-SLC 06.

## Ae

Using polarized beams, this quantity can also be measured as  $(\sigma_I - \sigma_R)/(\sigma_I + \sigma_R)$ , where  $\sigma_L$  and  $\sigma_R$  are the  $e^+e^-$  production cross sections for Z bosons produced with left-handed and right-handed electrons respectively.

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT
$0.1515 \pm 0.0019$ OUR AVER/	AGE				
$0.1454 \pm 0.0108 \pm 0.0036$	144810	<sup>1</sup> ABBIENDI	010	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV
$0.1516 \!\pm\! 0.0021$	559000	<sup>2</sup> ABE	<b>01</b> B	SLD	$E_{\rm cm}^{ee} = 91.24  { m GeV}$
$0.1504 \pm 0.0068 \pm 0.0008$		<sup>3</sup> HEISTER	01	ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV
$0.1382\!\pm\!0.0116\!\pm\!0.0005$	105000	<sup>4</sup> ABREU	00e	DLPH	$E_{\rm cm}^{ee}$ = 88–94 GeV
$0.1678 \!\pm\! 0.0127 \!\pm\! 0.0030$	137092	<sup>5</sup> ACCIARRI	98H	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV
$0.162\ \pm 0.041\ \pm 0.014$	89838	<sup>6</sup> ABE	97	SLD	$E_{\rm Cm}^{ee} = 91.27  { m GeV}$
$0.202 \ \pm 0.038 \ \pm 0.008$		<sup>7</sup> ABE	95J	SLD	$E_{\rm Cm}^{ee}=$ 91.31 GeV

<sup>1</sup>ABBIENDI 010 fit for  $A_e$  and  $A_{\tau}$  from measurements of the  $\tau$  polarization at varying  $\tau$  production angles. The correlation between  $A_e$  and  $A_{\tau}$  is less than 0.03.

 $^2$  ABE 01B use the left-right production and left-right forward-backward decay asymmetries in leptonic Z decays to obtain a value of 0.1544  $\pm$  0.0060. This is combined with leftright production asymmetry measurement using hadronic Z decays (ABE 00B) to obtain the quoted value.

- $^3$ HEISTER 01 obtain this result fitting the au polarization as a function of the polar production angle of the  $\tau$ .
- $^4$ ABREU 00E obtain this result fitting the au polarization as a function of the polar au production angle. This measurement is a combination of different analyses (exclusive  $\tau$  decay modes, inclusive hadronic 1-prong reconstruction, and a neural network analysis).
- $^5$  Derived from the measurement of forward-backward au polarization asymmetry.
- <sup>6</sup>ABE 97 obtain this result from a measurement of the observed left-right charge asymmetry,  $A_{O}^{\rm obs}$  = 0.225  $\pm$  0.056  $\pm$  0.019, in hadronic Z decays. If they combine this value of  $A_Q^{obs}$  with their earlier measurement of  $A_{LR}^{obs}$  they determine  $A_e$  to be 0.1574  $\pm$  0.0197  $\pm$  0.0067 independent of the beam polarization.

<sup>7</sup>ABE 95J obtain this result from polarized Bhabha scattering.

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# $A_{\mu}$

This quantity is directly extracted from a measurement of the left-right forwardbackward asymmetry in  $\mu^+\mu^-$  production at SLC using a polarized electron beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter A<sub>o</sub>.

VALUE	EVTS	DOCUMENT IL	7	TECN	COMMENT			
0.142±0.015	16844	<sup>1</sup> ABE	<b>01</b> B	SLD	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.24 GeV			
● ● ● We do not use the following data for averages, fits, limits, etc. ● ●								
		•						

15BT ATLS  $E^{pp}_{cm} =$  7 TeV  $^{2}$  AAD 1.7M  $0.153 \pm 0.012$ 

 $^1 {\sf ABE}$  01B obtain this direct measurement using the left-right production and left-right forward-backward polar angle asymmetries in  $\mu^+\mu^-$  decays of the Z boson obtained with a polarized electron beam.

<sup>2</sup>AAD 15BT study  $pp \rightarrow Z \rightarrow \ell^+ \ell^-$  events where  $\ell$  is an electron or a muon in the dilepton mass region 70–1000 GeV. The background in the Z peak region is estimated to be < 1% for the muon channel. The muon asymmetry parameter is derived from the measured forward-backward asymmetry assuming the value of the quark asymmetry parameter from the SM. For this reason it is not used in the average.

## $A_{\tau}$

The LEP Collaborations derive this quantity from the measurement of the  $\tau$  polarization in  $Z \rightarrow \tau^+ \tau^-$ . The SLD Collaboration directly extracts this quantity from its measured left-right forward-backward asymmetry in  $Z \rightarrow \tau^+ \tau^-$  produced using a polarized  $e^-$  beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter  $A_e$ .

VALUE	<u>EVTS</u>	DOCUMENT ID		TECN	COMMENT		
0.143 ±0.004 OUR AVERAGE							
$0.1456\!\pm\!0.0076\!\pm\!0.0057$	144810	<sup>1</sup> ABBIENDI	010	OPAL	$E_{\rm cm}^{ee}$ = 88–94 GeV		
$0.136\ \pm 0.015$	16083	<sup>2</sup> ABE	<b>01</b> B	SLD	$E_{\rm cm}^{ee}$ = 91.24 GeV		
$0.1451\!\pm\!0.0052\!\pm\!0.0029$		<sup>3</sup> HEISTER	01	ALEP	$E_{\rm cm}^{ee}$ = 88–94 GeV		
$0.1359 \!\pm\! 0.0079 \!\pm\! 0.0055$	105000	<sup>4</sup> ABREU	00E	DLPH	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 88–94 GeV		
$0.1476 \!\pm\! 0.0088 \!\pm\! 0.0062$	137092	ACCIARRI	98H	L3	$E_{\rm cm}^{ee}$ = 88–94 GeV		

<sup>1</sup>ABBIENDI 010 fit for  $A_e$  and  $A_{\tau}$  from measurements of the  $\tau$  polarization at varying  $\tau$  production angles. The correlation between  $A_e$  and  $A_{\tau}$  is less than 0.03.

- $^2$ ABE 01B obtain this direct measurement using the left-right production and left-right forward-backward polar angle asymmetries in  $\tau^+\tau^-$  decays of the Z boson obtained with a polarized electron beam.
- $^3\,{\sf HEISTER}$  01 obtain this result fitting the  $\tau$  polarization as a function of the polar production angle of the  $\tau$ .
- $^4$ ABREU 00E obtain this result fitting the au polarization as a function of the polar au production angle. This measurement is a combination of different analyses (exclusive au decay modes, inclusive hadronic 1-prong reconstruction, and a neural network analysis).

#### A۶

The SLD Collaboration directly extracts this quantity by a simultaneous fit to four measured s-quark polar angle distributions corresponding to two states of  $e^-$  polarization (positive and negative) and to the  $K^+K^-$  and  $K^{\pm}K^0_S$  strange particle tagging modes in the hadronic final states.

VALUE	EVTS	DOCUMENT ID TECN		DOCUMENT ID TECN COMMENT	
0.895±0.066±0.062	2870	<sup>1</sup> ABE	<b>00</b> D	SLD	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.2 GeV

<sup>1</sup>ABE 00D tag  $Z \rightarrow s\bar{s}$  events by an absence of B or D hadrons and the presence in each hemisphere of a high momentum  $K^{\pm}$  or  $K^{0}_{S}$ .

#### Ac

This quantity is directly extracted from a measurement of the left-right forwardbackward asymmetry in  $c\overline{c}$  production at SLC using polarized electron beam. This double asymmetry eliminates the dependence on the *Z*-*e*-*e* coupling parameter  $A_e$ . OUR FIT is obtained by a simultaneous fit to several *c*- and *b*-quark measurements as explained in the note "The *Z* boson" and ref. LEP-SLC 06.

VALUE	<u>DOCUMENT ID</u>		TECN	COMMENT
0.670 ±0.027 OUR FIT				
$0.6712 \!\pm\! 0.0224 \!\pm\! 0.0157$	<sup>1</sup> ABE	05	SLD	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.24 GeV
$\bullet \bullet \bullet$ We do not use the followin	g data for average	s, fits,	limits, o	etc. • • •
$0.583 \ \pm 0.055 \ \pm 0.055$	<sup>2</sup> ABE	<b>0</b> 2G	SLD	<i>E</i> <sup>ee</sup> <sub>cm</sub> = 91.24 GeV
$0.688 \pm 0.041$	<sup>3</sup> ABE	<b>01</b> C	SLD	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.25 GeV

<sup>1</sup> ABE 05 use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $c\overline{c}$  events tagging on the invariant mass of reconstructed secondary decay vertices. The charge of the underlying c-quark is obtained with an algorithm that takes into account the net charge of the vertex as well as the charge of tracks emanating from the vertex and identified as kaons. This yields (9970 events)  $A_c = 0.6747 \pm 0.0290 \pm 0.0233$ . Taking into account all correlations with earlier results reported in ABE 02G and ABE 01C, they obtain the quoted overall SLD result.

<sup>2</sup> ABE 02G tag *b* and *c* quarks through their semileptonic decays into electrons and muons. A maximum likelihood fit is performed to extract simultaneously  $A_b$  and  $A_c$ .

<sup>3</sup> ABE 01C tag  $Z \rightarrow c\overline{c}$  events using two techniques: exclusive reconstruction of  $D^{*+}$ ,  $D^+$ and  $D^0$  mesons and the soft pion tag for  $D^{*+} \rightarrow D^0 \pi^+$ . The large background from D mesons produced in  $b\overline{b}$  events is separated efficiently from the signal using precision vertex information. When combining the  $A_c$  values from these two samples, care is taken to avoid double counting of events common to the two samples, and common systematic errors are properly taken into account.

#### Ab

This quantity is directly extracted from a measurement of the left-right forwardbackward asymmetry in  $b\overline{b}$  production at SLC using polarized electron beam. This double asymmetry eliminates the dependence on the Z-e-e coupling parameter  $A_e$ . OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
0.923 ±0.020 OUR FIT	-				
$0.9170 \!\pm\! 0.0147 \!\pm\! 0.0145$		<sup>1</sup> ABE	05	SLD	$E_{\rm Cm}^{ee}$ = 91.24 GeV
$\bullet \bullet \bullet$ We do not use the	following	data for averages	, fits, li	mits, etc	C. ● ● ●
$0.907\ \pm 0.020\ \pm 0.024$	48028	<sup>2</sup> ABE	03F	SLD	<i>E<sup>ee</sup></i> <sub>cm</sub> = 91.24 GeV
$0.919\ \pm 0.030\ \pm 0.024$		<sup>3</sup> ABE	<b>0</b> 2G	SLD	$E_{\rm cm}^{ee}$ = 91.24 GeV
$0.855\ \pm 0.088\ \pm 0.102$	7473	<sup>4</sup> ABE	99L	SLD	E <sup>ee</sup> <sub>cm</sub> = 91.27 GeV

<sup>1</sup> ABE 05 use hadronic Z decays collected during 1996–98 to obtain an enriched sample of  $b \overline{b}$  events tagging on the invariant mass of reconstructed secondary decay vertices. The charge of the underlying b-quark is obtained with an algorithm that takes into account the net charge of the vertex as well as the charge of tracks emanating from the vertex and identified as kaons. This yields (25917 events)  $A_b = 0.9173 \pm 0.0184 \pm 0.0173$ . Taking into account all correlations with earlier results reported in ABE 03F, ABE 02G and ABE 99L, they obtain the quoted overall SLD result.

 $^2$  ABE 03F obtain an enriched sample of  $b\overline{b}$  events tagging on the invariant mass of a 3-dimensional topologically reconstructed secondary decay. The charge of the underlying b quark is obtained using a self-calibrating track-charge method. For the 1996–1998 data sample they measure  $A_b=0.906\pm0.022\pm0.023$ . The value quoted here is obtained combining the above with the result of ABE 98I (1993–1995 data sample).

- <sup>3</sup>ABE 02G tag *b* and *c* quarks through their semileptonic decays into electrons and muons. A maximum likelihood fit is performed to extract simultaneously  $A_b$  and  $A_c$ .
- <sup>4</sup> ABE 99L obtain an enriched sample of  $b\overline{b}$  events tagging with an inclusive vertex mass cut. For distinguishing *b* and  $\overline{b}$  quarks they use the charge of identified  $K^{\pm}$ .

# TRANSVERSE SPIN CORRELATIONS IN $Z \rightarrow \tau^+ \tau^-$

The correlations between the transverse spin components of  $\tau^+ \tau^-$  produced in Z decays may be expressed in terms of the vector and axial-vector couplings:

$$C_{TT} = \frac{|g_A^{\tau}|^2 - |g_V^{\tau}|^2}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2}$$
$$C_{TN} = -2 \frac{|g_A^{\tau}| |g_V^{\tau}|}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2} \sin(\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}})$$

 $C_{TT}$  refers to the transverse-transverse (within the collision plane) spin correlation and  $C_{TN}$  refers to the transverse-normal (to the collision plane) spin correlation.

The longitudinal  $\tau$  polarization  $P_{\tau}$  (=  $-A_{\tau}$ ) is given by:

$$P_{\tau} = -2 \frac{|g_A^{\tau}| |g_V^{\tau}|}{|g_A^{\tau}|^2 + |g_V^{\tau}|^2} \cos(\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}})$$

Here  $\Phi$  is the phase and the phase difference  $\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}}$  can be obtained using both the measurements of  $C_{TN}$  and  $P_{\tau}$ .

VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$1.01\pm0.12$ OUR AVERA	GE				
$0.87 \!\pm\! 0.20 \!+\! 0.10 \!-\! 0.12$	9.1k	ABREU	<b>97</b> G	DLPH	$E_{\rm cm}^{ee}$ = 91.2 GeV
$1.06\!\pm\!0.13\!\pm\!0.05$	120k	BARATE	<b>97</b> D	ALEP	$E_{\rm cm}^{ee}$ = 91.2 GeV
C <sub>TN</sub>					
VALUE	EVTS	DOCUMENT ID		TECN	COMMENT
$0.08 {\pm} 0.13 {\pm} 0.04$	120k <sup>1</sup>	BARATE	<b>97</b> D	ALEP	$E_{\rm cm}^{ee}$ = 91.2 GeV
_					

<sup>1</sup>BARATE 97D combine their value of  $C_{TN}$  with the world average  $P_{\tau} = -0.140 \pm 0.007$  to obtain  $\tan(\Phi_{g_V^{\tau}} - \Phi_{g_A^{\tau}}) = -0.57 \pm 0.97$ .

# FORWARD-BACKWARD $e^+e^- \rightarrow f\overline{f}$ CHARGE ASYMMETRIES

These asymmetries are experimentally determined by tagging the respective lepton or quark flavor in  $e^+\,e^-$  interactions. Details of heavy flavor (c- or b-quark) tagging at LEP are described in the note on "The Z boson" and ref. LEP-SLC 06. The Standard Model predictions for LEP data have been (re)computed using the ZFITTER package (version 6.36) with input parameters  $M_Z=91.187~{\rm GeV},~M_{\rm top}=174.3~{\rm GeV},~M_{\rm Higgs}=150~{\rm GeV},~\alpha_s=0.119,~\alpha^{(5)}~(M_Z)=1/128.877$  and the Fermi constant  $G_F=1.16637\times 10^{-5}~{\rm GeV}^{-2}$  (see the note on "The Z boson" for references).

For non-LEP data the Standard Model predictions are as given by the authors of the respective publications.



# $A_{FB}^{(0,e)}$ CHARGE ASYMMETRY IN $e^+e^- \rightarrow e^+e^- -$

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_e^2$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
1.45±0.25 OUR FIT					
$0.89 \pm 0.44$	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$1.71 \pm 0.49$	1.57	91.2	ABREU	00F	DLPH
$1.06 \pm 0.58$	1.57	91.2	ACCIARRI	<b>00</b> C	L3
$1.88 \pm 0.34$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP

<sup>1</sup>ABBIENDI 01A error includes approximately 0.38 due to statistics, 0.16 due to event selection systematics, and 0.18 due to the theoretical uncertainty in *t*-channel prediction. <sup>2</sup> BARATE 00C error includes approximately 0.31 due to statistics, 0.06 due to experimental

systematics, and 0.13 due to the theoretical uncertainty in *t*-channel prediction.

# — $A^{(0,\mu)}_{FB}$ CHARGE ASYMMETRY IN $e^+e^- ightarrow \mu^+\mu^-$ -

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_eA_\mu$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID	TECN
1.69 $\pm$ 0.13 OUR FIT			1	
$1.59\pm$ 0.23	1.57	91.2	<sup>1</sup> ABBIENDI 01A	OPAL
$1.65\pm$ 0.25	1.57	91.2	ABREU 00F	DLPH
$1.88\pm$ 0.33	1.57	91.2	ACCIARRI 00C	L3
$1.71\pm~0.24$	1.57	91.2	<sup>2</sup> BARATE 00C	ALEP
$\bullet$ • • We do not use the follow	wing data fo	r averages,	fits, limits, etc. • • •	
9 ±30	-1.3	20	<sup>3</sup> ABREU 95м	DLPH
$7 \pm 26$	-8.3	40	<sup>3</sup> ABREU 95м	DLPH
$-11 \pm 33$	-24.1	57	<sup>3</sup> ABREU 95м	DLPH
$-62 \pm 17$	-44.6	69	<sup>3</sup> ABREU 95м	DLPH
$-56 \pm 10$	-63.5	79	<sup>3</sup> ABREU 95м	DLPH
$-13$ $\pm$ 5	-34.4	87.5	<sup>3</sup> ABREU 95м	DLPH
$-29.0 \ \begin{array}{c} + \ 5.0 \\ - \ 4.8 \end{array} \pm 0.5$	-32.1	56.9	<sup>4</sup> ABE 901	VNS
$-$ 9.9 $\pm$ 1.5 $\pm$ 0.5	-9.2	35	HEGNER 90	JADE
$0.05 \pm 0.22$	0.026	91.14	<sup>5</sup> ABRAMS 89D	MRK2
$-43.4 \pm 17.0$	-24.9	52.0	<sup>6</sup> BACALA 89	AMY
$-11.0 \pm 16.5$	-29.4	55.0	<sup>6</sup> BACALA 89	AMY
$-30.0 \pm 12.4$	-31.2	56.0	<sup>6</sup> BACALA 89	AMY
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$-46.2 \pm 14.9$	-33.0	57.0	<sup>6</sup> BACALA	89	AMY
$-29 \pm 13$	-25.9	53.3	ADACHI	88C	TOPZ
$+$ 5.3 $\pm$ 5.0 $\pm$ 0.5	-1.2	14.0	ADEVA	88	MRKJ
$-10.4~\pm~1.3~\pm0.5$	-8.6	34.8	ADEVA	88	MRKJ
$-12.3~\pm~5.3~\pm0.5$	-10.7	38.3	ADEVA	88	MRKJ
$-15.6~\pm~3.0~\pm0.5$	-14.9	43.8	ADEVA	88	MRKJ
$-$ 1.0 $\pm$ 6.0	-1.2	13.9	BRAUNSCH	<b>88</b> D	TASS
$-$ 9.1 $\pm$ 2.3 $\pm 0.5$	-8.6	34.5	BRAUNSCH	<b>88</b> D	TASS
$-10.6 \ + \ 2.2 \ \pm 0.5$	-8.9	35.0	BRAUNSCH	<b>88</b> D	TASS
$-17.6 \ \begin{array}{c} + & 4.4 \\ - & 4.3 \end{array} \pm 0.5$	-15.2	43.6	BRAUNSCH	<b>88</b> D	TASS
$-$ 4.8 $\pm$ 6.5 $\pm 1.0$	-11.5	39	BEHREND	87C	CELL
$-18.8~\pm~4.5~\pm1.0$	-15.5	44	BEHREND	87C	CELL
$+$ 2.7 $\pm$ 4.9	-1.2	13.9	BARTEL	86C	JADE
$-11.1~\pm~1.8~\pm1.0$	-8.6	34.4	BARTEL	86C	JADE
$-17.3$ $\pm$ 4.8 $\pm1.0$	-13.7	41.5	BARTEL	86C	JADE
$-22.8~\pm~5.1~\pm1.0$	-16.6	44.8	BARTEL	86C	JADE
$-$ 6.3 $\pm$ 0.8 $\pm$ 0.2	-6.3	29	ASH	85	MAC
$-$ 4.9 $\pm$ 1.5 $\pm$ 0.5	-5.9	29	DERRICK	85	HRS
$-$ 7.1 $\pm$ 1.7	-5.7	29	LEVI	83	MRK2
$-16.1~\pm~3.2$	-9.2	34.2	BRANDELIK	82C	TASS
1					

<sup>1</sup>ABBIENDI 01A error is almost entirely on account of statistics.

<sup>2</sup> BARATE 00C error is almost entirely on account of statistics.

<sup>3</sup> ABREU 95M perform this measurement using radiative muon-pair events associated with high-energy isolated photons.

<sup>4</sup>ABE 901 measurements in the range 50  $\leq \sqrt{s} \leq$  60.8 GeV.

<sup>5</sup>ABRAMS 89D asymmetry includes both  $9 \mu^+ \mu^-$  and 15  $\tau^+ \tau^-$  events.

<sup>6</sup>BACALA 89 systematic error is about 5%.

# - $A^{(0, au)}_{FB}$ CHARGE ASYMMETRY IN $e^+e^- \rightarrow \tau^+\tau^-$ -------

OUR FIT is obtained using the fit procedure and correlations as determined by the LEP Electroweak Working Group (see the note "The Z boson" and ref. LEP-SLC 06). For the Z peak, we report the pole asymmetry defined by  $(3/4)A_eA_{\tau}$  as determined by the nine-parameter fit to cross-section and lepton forward-backward asymmetry data.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
1.88 $\pm$ 0.17 OUR FIT					
$1.45\pm$ 0.30	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$2.41\pm$ 0.37	1.57	91.2	ABREU	00F	DLPH
$2.60\pm$ 0.47	1.57	91.2	ACCIARRI	<b>00</b> C	L3
$1.70\pm~0.28$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP
$\bullet \bullet \bullet$ We do not use the follow	wing data fo	or averages,	, fits, limits, etc. $ullet$	• •	
$-32.8 \ + \ 6.4 \ \pm 1.5$	-32.1	56.9	<sup>3</sup> ABE	901	VNS
$-$ 8.1 $\pm$ 2.0 $\pm$ 0.6	-9.2	35	HEGNER	90	JADE
$-18.4 \pm 19.2$	-24.9	52.0	<sup>4</sup> BACALA	89	AMY
$-17.7 \pm 26.1$	-29.4	55.0	<sup>4</sup> BACALA	89	AMY
$-45.9 \pm 16.6$	-31.2	56.0	<sup>4</sup> BACALA	89	AMY
$-49.5 \pm 18.0$	-33.0	57.0	<sup>4</sup> BACALA	89	AMY
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$-20 \pm 14$	-25.9	53.3	ADACHI	88C	TOPZ
$-10.6~\pm~3.1~\pm1.5$	-8.5	34.7	ADEVA	88	MRKJ
$-$ 8.5 $\pm$ 6.6 $\pm 1.5$	-15.4	43.8	ADEVA	88	MRKJ
$-$ 6.0 $\pm$ 2.5 $\pm1.0$	8.8	34.6	BARTEL	85F	JADE
$-11.8$ $\pm$ 4.6 $\pm1.0$	14.8	43.0	BARTEL	85F	JADE
$-$ 5.5 $\pm$ 1.2 $\pm$ 0.5	-0.063	29.0	FERNANDEZ	85	MAC
$-$ 4.2 $\pm$ 2.0	0.057	29	LEVI	83	MRK2
$-10.3~\pm~5.2$	-9.2	34.2	BEHREND	82	CELL
$-$ 0.4 $\pm$ 6.6	-9.1	34.2	BRANDELIK	82C	TASS

<sup>1</sup>ABBIENDI 01A error includes approximately 0.26 due to statistics and 0.14 due to event selection systematics.

- <sup>2</sup>BARATE 00C error includes approximately 0.26 due to statistics and 0.11 due to experimental systematics.
- <sup>3</sup>ABE 901 measurements in the range 50  $\leq \sqrt{s} \leq$  60.8 GeV. <sup>4</sup>BACALA 89 systematic error is about 5%.

# ------ $A_{FB}^{(0,\ell)}$ CHARGE ASYMMETRY IN $e^+e^- \rightarrow \ell^+\ell^-$ -------

For the Z peak, we report the pole asymmetry defined by  $(3/4)A_\ell^2$  as determined by the five-parameter fit to cross-section and lepton forwardbackward asymmetry data assuming lepton universality. For details see the note "The Z boson" and ref. LEP-SLC 06.

ASYMMETRY (%)	STD. MODEL	√ <i>s</i> (GeV)	DOCUMENT ID		TECN
$1.71\pm0.10$ OUR FIT					
$1.45 \pm 0.17$	1.57	91.2	<sup>1</sup> ABBIENDI	01A	OPAL
$1.87 \pm 0.19$	1.57	91.2	ABREU	00F	DLPH
$1.92 \pm 0.24$	1.57	91.2	ACCIARRI	00C	L3
$1.73 \pm 0.16$	1.57	91.2	<sup>2</sup> BARATE	<b>00</b> C	ALEP

<sup>1</sup>ABBIENDI 01A error includes approximately 0.15 due to statistics, 0.06 due to event selection systematics, and 0.03 due to the theoretical uncertainty in *t*-channel prediction. <sup>2</sup> BARATE 00C error includes approximately 0.15 due to statistics, 0.04 due to experimental systematics, and 0.02 due to the theoretical uncertainty in *t*-channel prediction.

 $----- A_{FB}^{(0,u)}$  CHARGE ASYMMETRY IN  $e^+e^- \rightarrow u\overline{u}$ 

4.0±6.7±2.8	7.2	91.2	1 ACKERSTAFF 97T	OPAL
ASYMMETRY (%)	STD. MODEL	√ <i>s</i> (GeV)	DOCUMENT ID	TECN

<sup>1</sup>ACKERSTAFF 97<sup>T</sup> measure the forward-backward asymmetry of various fast hadrons made of light quarks. Then using SU(2) isospin symmetry and flavor independence for down and strange quarks authors solve for the different quark types.

 $- A_{FB}^{(0,s)} \text{ CHARGE ASYMMETRY IN } e^+ e^- \rightarrow s\overline{s} - \cdots$ 

The *s*-quark asymmetry is derived from measurements of the forward-backward asymmetry of fast hadrons containing an *s* quark.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID	TECN
9.8 $\pm$ 1.1 OUR AVERAGE				
$10.08\!\pm\!1.13\!\pm\!0.40$	10.1	91.2	<sup>1</sup> ABREU 00B	DLPH
$6.8\ \pm 3.5\ \pm 1.1$	10.1	91.2	<sup>2</sup> ACKERSTAFF 97T	OPAL

<sup>1</sup> ABREU 00B tag the presence of an *s* quark requiring a high-momentum-identified charged kaon. The *s*-quark pole asymmetry is extracted from the charged-kaon asymmetry taking the expected *d*- and *u*-quark asymmetries from the Standard Model and using the measured values for the *c*- and *b*-quark asymmetries.

<sup>2</sup> ACKERSTAFF 97T measure the forward-backward asymmetry of various fast hadrons made of light quarks. Then using SU(2) isospin symmetry and flavor independence for down and strange quarks authors solve for the different quark types. The value reported here corresponds then to the forward-backward asymmetry for "down-type" quarks.

# $----- A_{FB}^{(0,c)} \text{ CHARGE ASYMMETRY IN } e^+ e^- \rightarrow c \overline{c} ---$

OUR FIT, which is obtained by a simultaneous fit to several c- and bquark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06, refers to the **Z pole** asymmetry. The experimental values, on the other hand, correspond to the measurements carried out at the respective energies.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
7.07± 0.35 OUR FIT			1		
$6.31 \pm \ 0.93 \pm 0.65$	6.35	91.26	<sup>1</sup> ABDALLAH	04F	DLPH
$5.68 \pm 0.54 \pm 0.39$	6.3	91.25	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
$6.45 \pm 0.57 \pm 0.37$	6.10	91.21	<sup>3</sup> HEISTER	02н	ALEP
$6.59 \pm 0.94 \pm 0.35$	6.2	91.235	<sup>4</sup> ABREU	99Y	DLPH
$6.3~\pm~0.9~\pm0.3$	6.1	91.22	<sup>5</sup> BARATE	980	ALEP
$6.3~\pm~1.2~\pm0.6$	6.1	91.22	<sup>6</sup> ALEXANDER	<b>97</b> C	OPAL
$8.3~\pm~3.8~\pm2.7$	6.2	91.24	<sup>7</sup> ADRIANI	<b>9</b> 2D	L3
$\bullet$ $\bullet$ We do not use the follow	wing data for	<sup>,</sup> averages, f	its, limits, etc. •	• •	
$3.1~\pm~3.5~\pm0.5$	- 3.5	89.43	<sup>1</sup> ABDALLAH	04F	DLPH
11.0 $\pm$ 2.8 $\pm$ 0.7	12.3	92.99	<sup>1</sup> ABDALLAH	04F	DLPH
$-$ 6.8 $\pm$ 2.5 $\pm$ 0.9	-3.0	89.51	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
14.6 $\pm$ 2.0 $\pm$ 0.8	12.2	92.95	<sup>2</sup> ABBIENDI	<b>03</b> P	OPAL
$-12.4 \ \pm 15.9 \ \pm 2.0$	-9.6	88.38	<sup>3</sup> HEISTER	02H	ALEP
$-$ 2.3 $\pm$ 2.6 $\pm$ 0.2	-3.8	89.38	<sup>3</sup> HEISTER	02H	ALEP
$-$ 0.3 $\pm$ 8.3 $\pm$ 0.6	0.9	90.21	<sup>3</sup> HEISTER	02H	ALEP
10.6 $\pm$ 7.7 $\pm$ 0.7	9.6	92.05	<sup>3</sup> HEISTER	02H	ALEP
11.9 $\pm$ 2.1 $\pm$ 0.6	12.2	92.94	<sup>3</sup> HEISTER	02H	ALEP
$12.1\ \pm 11.0\ \pm 1.0$	14.2	93.90	<sup>3</sup> HEISTER	02H	ALEP
$-$ 4.96 $\pm$ 3.68 $\pm$ 0.53	- 3.5	89.434	<sup>4</sup> ABREU	99Y	DLPH
$11.80\pm~3.18\pm0.62$	12.3	92.990	<sup>4</sup> ABREU	99Y	DLPH
$-$ 1.0 $\pm$ 4.3 $\pm$ 1.0	-3.9	89.37	<sup>5</sup> BARATE	980	ALEP

$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	12.3 - 3.4 12.4	89.45 93.00	<sup>6</sup> ALEXANDER <sup>6</sup> ALEXANDER	97C 97C	OPAL
$egin{array}{rl} -12.9 \ \pm \ 7.8 \ \pm 5.5 \ 7.7 \ \pm 13.4 \ \pm 5.0 \end{array}$	-13.6 -22.1	35 43	BEHREND BEHREND		CELL CELL
$\begin{array}{rrrrr} -12.8 \ \pm \ 4.4 \ \pm 4.1 \\ -10.9 \ \pm 12.9 \ \pm 4.6 \\ -14.9 \ \pm \ 6.7 \end{array}$	-13.6 -23.2 -13.3	35 44 35	ELSEN ELSEN OULD-SAADA	90 90 89	JADE JADE JADE

<sup>1</sup>ABDALLAH 04F tag *b*- and *c*-quarks using semileptonic decays combined with charge flow information from the hemisphere opposite to the lepton. Enriched samples of  $c\overline{c}$  and  $b\overline{b}$  events are obtained using lifetime information.

- <sup>2</sup>ABBIENDI 03P tag heavy flavors using events with one or two identified leptons. This allows the simultaneous fitting of the *b* and *c* quark forward-backward asymmetries as well as the average  $B^0-\overline{B}^0$  mixing.
- <sup>3</sup> HEISTER 02H measure simultaneously b and c quark forward-backward asymmetries using their semileptonic decays to tag the quark charge. The flavor separation is obtained with a discriminating multivariate analysis.
- <sup>4</sup>ABREU 99Y tag  $Z \rightarrow b\overline{b}$  and  $Z \rightarrow c\overline{c}$  events by an exclusive reconstruction of several D meson decay modes ( $D^{*+}$ ,  $D^0$ , and  $D^+$  with their charge-conjugate states).
- <sup>5</sup> BARATE 980 tag  $Z \rightarrow c\overline{c}$  events requiring the presence of high-momentum reconstructed  $D^{*+}$ ,  $D^+$ , or  $D^0$  mesons.
- <sup>6</sup>ALEXANDER 97C identify the *b* and *c* events using a  $D/D^*$  tag.
- <sup>7</sup>ADRIANI 92D use both electron and muon semileptonic decays.

 $----- A_{FB}^{(0,b)}$  CHARGE ASYMMETRY IN  $e^+e^- \rightarrow b\overline{b}$  ----

OUR FIT, which is obtained by a simultaneous fit to several c- and bquark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06, refers to the **Z pole** asymmetry. The experimental values, on the other hand, correspond to the measurements carried out at the respective energies.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
9.92 $\pm$ 0.16 OUR FIT					
$9.58 \pm \ 0.32 \pm \ 0.14$	9.68	91.231	<sup>1</sup> ABDALLAH	05	DLPH
$10.04 \pm \ 0.56 \pm \ 0.25$	9.69	91.26	<sup>2</sup> ABDALLAH	04F	DLPH
$9.72 \pm \ 0.42 \pm \ 0.15$	9.67	91.25	<sup>3</sup> ABBIENDI	<b>03</b> P	OPAL
$9.77 \pm \ 0.36 \pm \ 0.18$	9.69	91.26	<sup>4</sup> ABBIENDI	021	OPAL
$9.52\pm~0.41\pm~0.17$	9.59	91.21	<sup>5</sup> HEISTER	02H	ALEP
$10.00 \pm \ 0.27 \pm \ 0.11$	9.63	91.232	<sup>6</sup> HEISTER	<b>01</b> D	ALEP
$7.62 \pm \ 1.94 \pm \ 0.85$	9.64	91.235	<sup>7</sup> ABREU	99Y	DLPH
$9.60\pm~0.66\pm~0.33$	9.69	91.26	<sup>8</sup> ACCIARRI	<b>99</b> D	L3
$9.31 \pm \ 1.01 \pm \ 0.55$	9.65	91.24	<sup>9</sup> ACCIARRI	<b>98</b> U	L3
9.4 $\pm$ 2.7 $\pm$ 2.2	9.61	91.22	<sup>10</sup> ALEXANDER	<b>97</b> C	OPAL
$\bullet \bullet \bullet$ We do not use the follow	wing data fo	r averages,	fits, limits, etc. $\bullet$	• •	
$6.37 \pm \ 1.43 \pm \ 0.17$	5.8	89.449	<sup>1</sup> ABDALLAH	05	DLPH
$10.41 \pm \ 1.15 \pm \ 0.24$	12.1	92.990	<sup>1</sup> ABDALLAH	05	DLPH
$6.7 ~\pm~ 2.2 ~\pm~ 0.2$	5.7	89.43	<sup>2</sup> ABDALLAH	04F	DLPH
11.2 $\pm$ 1.8 $\pm$ 0.2	12.1	92.99	<sup>2</sup> ABDALLAH	04F	DLPH
$4.7~\pm~1.8~\pm~0.1$	5.9	89.51	<sup>3</sup> ABBIENDI	<b>03</b> P	OPAL
10.3 $\pm$ 1.5 $\pm$ 0.2	12.0	92.95	<sup>3</sup> ABBIENDI	<b>03</b> P	OPAL
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$5.82 \pm \ 1.53 \pm \ 0.12$	5.9	89.50	<sup>4</sup> ABBIENDI	021	OPAL
$12.21 \pm \ 1.23 \pm \ 0.25$	12.0	92.91	<sup>4</sup> ABBIENDI	021	OPAL
$-13.1 \ \pm 13.5 \ \pm \ 1.0$	3.2	88.38	<sup>5</sup> HEISTER	02н	ALEP
5.5 $\pm$ 1.9 $\pm$ 0.1	5.6	89.38	<sup>5</sup> HEISTER	02н	ALEP
$-$ 0.4 $\pm$ 6.7 $\pm$ 0.8	7.5	90.21	<sup>5</sup> HEISTER	02н	ALEP
11.1 $\pm$ 6.4 $\pm$ 0.5	11.0	92.05	<sup>5</sup> HEISTER	02н	ALEP
10.4 $\pm$ 1.5 $\pm$ 0.3	12.0	92.94	<sup>5</sup> HEISTER	02н	ALEP
13.8 $\pm$ 9.3 $\pm$ 1.1	12.9	93.90	<sup>5</sup> HEISTER	02н	ALEP
$4.36 \pm \ 1.19 \pm \ 0.11$	5.8	89.472	<sup>6</sup> HEISTER	<b>01</b> D	ALEP
$11.72\pm~0.97\pm~0.11$	12.0	92.950	<sup>6</sup> HEISTER	<b>01</b> D	ALEP
$5.67 \pm \ 7.56 \pm \ 1.17$	5.7	89.434	<sup>7</sup> ABREU	<b>99</b> Y	DLPH
$8.82\pm~6.33\pm~1.22$	12.1	92.990	<sup>7</sup> ABREU	99Y	DLPH
$6.11\pm~2.93\pm~0.43$	5.9	89.50	<sup>8</sup> ACCIARRI	<b>99</b> D	L3
$13.71\pm~2.40\pm~0.44$	12.2	93.10	<sup>8</sup> ACCIARRI	<b>99</b> D	L3
$4.95 \pm 5.23 \pm 0.40$	5.8	89.45	<sup>9</sup> ACCIARRI	<b>98</b> U	L3
$11.37 \pm \ 3.99 \pm \ 0.65$	12.1	92.99	<sup>9</sup> ACCIARRI	<b>98</b> U	L3
$-$ 8.6 $\pm 10.8$ $\pm$ 2.9	5.8	89.45	<sup>10</sup> ALEXANDER	<b>97</b> C	OPAL
$-$ 2.1 $\pm$ 9.0 $\pm$ 2.6	12.1	93.00	<sup>10</sup> ALEXANDER	<b>97</b> C	OPAL
$-71$ $\pm 34$ $+$ $\frac{7}{8}$	- 58	58.3	SHIMONAKA	91	TOPZ
$-22.2 \pm 7.7 \pm 3.5$	-26.0	35	BEHREND	<b>90</b> D	CELL
$-49.1 \ \pm 16.0 \ \pm \ 5.0$	- 39.7	43	BEHREND	<b>90</b> D	CELL
$-28 \pm 11$	-23	35	BRAUNSCH	90	TASS
$-16.6 \pm 7.7 \pm 4.8$	-24.3	35	ELSEN	90	JADE
$-33.6 \pm 22.2 \pm 5.2$	- 39.9	44	ELSEN	90	JADE
$3.4 \pm 7.0 \pm 3.5$	-16.0	29.0	BAND	89	MAC
$-72 \pm 28 \pm 13$	- 56	55.2	SAGAWA	89	AMY

<sup>1</sup> ABDALLAH 05 obtain an enriched samples of  $b\overline{b}$  events using lifetime information. The quark (or antiquark) charge is determined with a neural network using the secondary vertex charge, the jet charge and particle identification.

<sup>2</sup>ABDALLAH 04F tag *b*- and *c*-quarks using semileptonic decays combined with charge flow information from the hemisphere opposite to the lepton. Enriched samples of  $c\overline{c}$  and  $b\overline{b}$  events are obtained using lifetime information.

- <sup>3</sup>ABBIENDI 03P tag heavy flavors using events with one or two identified leptons. This allows the simultaneous fitting of the *b* and *c* quark forward-backward asymmetries as well as the average  $B^0-\overline{B}^0$  mixing.
- <sup>4</sup>ABBIENDI 02I tag  $Z^0 \rightarrow b\overline{b}$  decays using a combination of secondary vertex and lepton tags. The sign of the *b*-quark charge is determined using an inclusive tag based on jet, vertex, and kaon charges.

<sup>5</sup> HEISTER 02H measure simultaneously b and c quark forward-backward asymmetries using their semileptonic decays to tag the quark charge. The flavor separation is obtained with a discriminating multivariate analysis.

<sup>6</sup> HEISTER 01D tag  $Z \rightarrow b\overline{b}$  events using the impact parameters of charged tracks complemented with information from displaced vertices, event shape variables, and lepton identification. The *b*-quark direction and charge is determined using the hemisphere charge method along with information from fast kaon tagging and charge estimators of primary and secondary vertices. The change in the quoted value due to variation of  $A_{FR}^c$ 

and  $R_b$  is given as +0.103 ( $A_{FB}^c$  - 0.0651) -0.440 ( $R_b$  - 0.21585).

<sup>7</sup> ABREU 99Y tag  $Z \rightarrow b\overline{b}$  and  $\overline{Z} \rightarrow c\overline{c}$  events by an exclusive reconstruction of several D meson decay modes ( $D^{*+}$ ,  $D^{0}$ , and  $D^{+}$  with their charge-conjugate states).

<sup>8</sup> ACCIARRI 99D tag  $Z \rightarrow b\overline{b}$  events using high p and p<sub>T</sub> leptons. The analysis determines simultaneously a mixing parameter  $\chi_b = 0.1192 \pm 0.0068 \pm 0.0051$  which is used to correct the observed asymmetry.

<sup>9</sup> ACCIARRI 980 tag  $Z \rightarrow b\overline{b}$  events using lifetime and measure the jet charge using the hemisphere charge.

<sup>10</sup>ALEXANDER 97C identify the *b* and *c* events using a  $D/D^*$  tag.

# CHARGE ASYMMETRY IN $e^+e^- \rightarrow q\overline{q}$

Summed over five lighter flavors.

Experimental and Standard Model values are somewhat event-selection dependent. Standard Model expectations contain some assumptions on  $B^0$ - $\overline{B}^0$  mixing and on other electroweak parameters.

ASYMMETRY (%)	STD. MODEL	$\sqrt{s}$ (GeV)	DOCUMENT ID		TECN
$\bullet$ $\bullet$ $\bullet$ We do not use the following	owing data fo	r averages, f	its, limits, etc. • •	•	
$-$ 0.76 $\pm$ 0.12 $\pm$ 0.15		91.2	<sup>1</sup> ABREU	92ı	DLPH
$4.0 \pm 0.4 \pm 0.63$	4.0	91.3	<sup>2</sup> ACTON	92L	OPAL
$9.1 \ \pm 1.4 \ \pm 1.6$	9.0	57.9	ADACHI	91	TOPZ
$- 0.84 \!\pm\! 0.15 \!\pm\! 0.04$		91	DECAMP	<b>91</b> B	ALEP
$8.3\ \pm 2.9\ \pm 1.9$	8.7	56.6	STUART	90	AMY
$11.4\ \pm 2.2\ \pm 2.1$	8.7	57.6	ABE	89L	VNS
$6.0 \pm 1.3$	5.0	34.8	GREENSHAW	89	JADE
8.2 ±2.9	8.5	43.6	GREENSHAW	89	JADE
-					

 $^{1}$ ABREU 921 has 0.14 systematic error due to uncertainty of quark fragmentation.

<sup>2</sup> ACTON 92L use the weight function method on 259k selected  $Z \rightarrow$  hadrons events. The systematic error includes a contribution of 0.2 due to  $B^{0}-\overline{B}^{0}$  mixing effect, 0.4 due to Monte Carlo (MC) fragmentation uncertainties and 0.3 due to MC statistics. ACTON 92L derive a value of  $\sin^{2}\theta_{W}^{\text{eff}}$  to be 0.2321 ± 0.0017 ± 0.0028.

CHARGE ASYMMETRY IN $p\overline{p} \rightarrow Z \rightarrow e^+ e^-$						
ASYMMETRY (%)	STD. MODEL	$\frac{\sqrt{s}}{(\text{GeV})}$	DOCUMENT ID		TECN	
ullet $ullet$ $ullet$ We do not use the following data for averages, fits, limits, etc. $ullet$ $ullet$						
$5.2{\pm}5.9{\pm}0.4$		91	ABE	91E	CDF	

# ANOMALOUS $ZZ\gamma$ , $Z\gamma\gamma$ , AND ZZV COUPLINGS See the related review(s):

Anomalous  $ZZ\gamma$ ,  $Z\gamma\gamma$ , and ZZV Couplings

# h<sub>i</sub>V

Combining the LEP-2 results taking into account the correlations, the following 95% CL limits are derived [SCHAEL 13A]:

$-0.12 < h_1^Z < +0.11$ ,	$-0.07 < h_2^Z < +0.07$ ,
$-0.19 < h_{\overline{3}}^{\overline{Z}} < +0.06$ ,	$-0.04 < h_{4}^{\overline{Z}} < +0.13$ ,
$-0.05 <  h_1^{\gamma} < +0.05$ ,	$-0.04 < h_2^\gamma \; < +0.02$ ,
$-0.05 < h_3^{\gamma} < +0.00$ ,	$+0.01 < h_4^{\gamma} < +0.05.$

Some of the recent results from the Tevatron and LHC experiments individually surpass the combined LEP-2 results in precision (see below).

VALUE	DOCUMENT ID		TECN	COMMENT
$\bullet$ $\bullet$ We do not use the following	g data for averages	, fits,	limits, e	tc. ● ● ●
	<sup>1</sup> AAD	16Q	ATLS	$E^{pp}_{cm} = 8 \text{ TeV}$
	<sup>2</sup> KHACHATRY	.16AE	CMS	$E_{cm}^{pp}=$ 8 TeV
	<sup>3</sup> KHACHATRY	. <b>15</b> AC	CMS	$E^{pp}_{cm} = 8 \text{ TeV}$
	<sup>4</sup> CHATRCHYAN	<b>14</b> AB	CMS	$E_{cm}^{pp}=$ 7 TeV
	<sup>5</sup> AAD	13AN	ATLS	$E_{ m cm}^{pp}=$ 7 TeV
	<sup>6</sup> CHATRCHYAN	<b>13</b> BI	CMS	$E_{\rm cm}^{pp} =$ 7 TeV
	<sup>7</sup> ABAZOV	12s	D0	$E_{\sf cm}^{p\overline{p}}=1.96\;{\sf TeV}$
	<sup>8</sup> AALTONEN	11s	CDF	$E_{\sf cm}^{p\overline{p}}=1.96\;{\sf TeV}$
	<sup>9</sup> CHATRCHYAN	11M	CMS	$E_{\rm cm}^{pp} =$ 7 TeV
	<sup>10</sup> ABAZOV	09L	D0	$E_{\sf cm}^{p\overline{p}}=1.96\;{\sf TeV}$
	$^{11}$ ABAZOV	<b>07</b> M	D0	$E_{\sf cm}^{p\overline{p}}=1.96\;{\sf TeV}$
	<sup>12</sup> ABDALLAH	<b>07</b> C	DLPH	$E_{\rm cm}^{ee}=$ 183–208 GeV
	<sup>13</sup> ACHARD	04H	L3	$E_{\rm cm}^{ee} = 183$ –208 GeV
	<sup>14</sup> ABBIENDI,G	<b>00</b> C	OPAL	$E_{\rm cm}^{ee} = 189 \; { m GeV}$
	<sup>15</sup> ABBOTT	98M	D0	$E_{\sf cm}^{p\overline{p}}=1.8\;{\sf TeV}$
	<sup>16</sup> ABREU	98K	DLPH	$E_{Cm}^{ee}=$ 161, 172 GeV

- <sup>1</sup> AAD 16Q study  $Z\gamma$  production in pp collisions. In events with no additional jets, 10268 (12738) Z decays to electron (muon) pairs are selected, with an expected background of 1291 ± 340 (1537 ± 408) events, as well as 1039 Z decays to neutrino pairs with an expected background of 450 ± 96 events. Analyzing the photon transverse momentum distribution above 250 GeV (400 GeV) for lepton (neutrino) events, yields the 95% C.L. limits:  $-7.8 \times 10^{-4} < h_3^Z < 8.6 \times 10^{-4}$ ,  $-3.0 \times 10^{-6} < h_4^Z < 2.9 \times 10^{-6}$ ,  $-9.5 \times 10^{-4} < h_3^{\gamma} < 9.9 \times 10^{-4}$ ,  $-3.2 \times 10^{-6} < h_4^{\gamma} < 3.2 \times 10^{-6}$ .
- <sup>2</sup> KHACHATRYAN 16AE determine the  $Z\gamma \rightarrow \nu \overline{\nu} \gamma$  cross section by selecting events with a photon of  $E_T$  > 145 GeV and  $E_T$  > 140 GeV. 630 candidate events are observed with an expected SM background of 269 ± 26. The  $E_T$  spectrum of the photon is used to set 95% C.L. limits as follows:  $-1.5 \times 10^{-3} < h_3^Z < 1.6 \times 10^{-3}, -3.9 \times 10^{-6} < h_4^Z < 4.5 \times 10^{-6}, -1.1 \times 10^{-3} < h_3^{\gamma} < 0.9 \times 10^{-3}, -3.8 \times 10^{-6} < h_4^{\gamma} < 4.3 \times 10^{-6}.$ <sup>3</sup> KHACHATRYAN 15AC study  $Z\gamma$  events in 8 TeV pp interactions, where the Z decays into 2 same-flavor, opposite sign leptons (e or  $\mu$ ) and a photon with  $p_T$  > 15 GeV. The  $n_T$  of a lepton is required to be > 20 GeV/c, their effective mass > 50 GeV and
- The  $p_T$  of a lepton is required to be > 20 GeV/c, their effective mass > 50 GeV, and the photon should have a separation  $\Delta R > 0.7$  with each lepton. The observed  $p_T$  distribution of the photons is used to extract the 95% C.L. limits:  $-3.8 \times 10^{-3} < h_3^Z < 3.7 \times 10^{-3}, -3.1 \times 10^{-5} < h_4^Z < 3.0 \times 10^{-5}, -4.6 \times 10^{-3} < h_3^\gamma < 4.6 \times 10^{-3}, -3.6 \times 10^{-5} < h_4^\gamma < 3.5 \times 10^{-5}$ .
- <sup>4</sup> CHATRCHYAN 14AB measure  $Z\gamma$  production cross section for  $p_T^{\gamma} > 15$  GeV and  $R(\ell\gamma) > 0.7$ , which is the separation between the  $\gamma$  and the final state charged lepton (e or  $\mu$ ) in the azimuthal angle-pseudorapidity ( $\phi \eta$ ) plane. The di-lepton mass is required to be > 50 GeV. After background subtraction the number of  $ee\gamma$  and  $\mu\mu\gamma$  events is determined to be  $3160 \pm 120$  and  $5030 \pm 233$  respectively, compatible with expectations

from the SM. This leads to a 95% CL limits of  $-1 \times 10^{-2} < h_3^{\gamma} < 1 \times 10^{-2}$ ,  $-9 \times 10^{-5} < h_4^{\gamma} < 9 \times 10^{-5}$ ,  $-9 \times 10^{-3} < h_3^Z < 9 \times 10^{-3}$ ,  $-8 \times 10^{-5} < h_4^Z < 8 \times 10^{-5}$ , assuming  $h_1^V$  and  $h_2^V$  have SM values,  $V = \gamma$  or Z.

- <sup>5</sup> AAD 13AN study  $Z\gamma$  production in pp collisions. In events with no additional jet, 1417 (2031) Z decays to electron (muon) pairs are selected, with an expected background of 156 ± 54 (244 ± 64) events, as well as 662 Z decays to neutrino pairs with an expected background of 302 ± 42 events. Analysing the photon  $p_T$  spectrum above 100 GeV yields the 95% C.L. limts:  $-0.013 < h_3^Z < 0.014$ ,  $-8.7 \times 10^{-5} < h_4^Z < 8.7 \times 10^{-5}$ ,  $-0.015 < h_3^{\gamma} < 0.016$ ,  $-9.4 \times 10^{-5} < h_4^{\gamma} < 9.2 \times 10^{-5}$ . Supersedes AAD 12BX.
- <sup>6</sup> CHATRCHYAN 13BI determine the  $Z\gamma \rightarrow \nu \overline{\nu} \gamma$  cross section by selecting events with a photon of  $E_T > 145$  GeV and a  $\not\!\!E_T > 130$  GeV. 73 candidate events are observed with an expected SM background of  $30.2 \pm 6.5$ . The  $E_T$  spectrum of the photon is used to set 95% C.L. limits as follows:  $|h_3^Z| < 2.7 \times 10^{-3}$ ,  $|h_4^Z| < 1.3 \times 10^{-5}$ ,  $|h_3^\gamma| < 2.9 \times 10^{-3}$ ,
- $|h_{A}^{\gamma}| < 1.5 \times 10^{-5}.$
- <sup>7</sup> ABAZOV 12S study  $Z\gamma$  production in  $p\overline{p}$  collisions at  $\sqrt{s} = 1.96$  TeV using 6.2 fb<sup>-1</sup> of data where the Z decays to electron (muon) pairs and the photon has at least 10 GeV of transverse momentum. In data, 304 (308) di-electron (di-muon) events are observed with an expected background of 255  $\pm$  16 (285  $\pm$  24) events. Based on the photon  $p_T$  spectrum, and including also earlier data and the  $Z \rightarrow \nu\overline{\nu}$  decay mode (from ABAZOV 09L), the following 95% C.L. limits are reported:  $|h_{03}^Z| < 0.026$ ,  $|h_{04}^Z| < 0.0013$ ,  $|h_{03}^\gamma| < 0.027$ ,  $|h_{04}^\gamma| < 0.0014$  for a form factor scale of  $\Lambda = 1.5$  TeV.
- <sup>8</sup> AALTONEN 11S study  $Z\gamma$  events in  $p\overline{p}$  interactions at  $\sqrt{s} = 1.96$  TeV with integrated luminosity 5.1 fb<sup>-1</sup> for  $Z \rightarrow e^+e^-/\mu^+\mu^-$  and 4.9 fb<sup>-1</sup> for  $Z \rightarrow \nu\overline{\nu}$ . For the charged lepton case, the two leptons must be of the same flavor with the transverse momentum/energy of one > 20 GeV and the other > 10 GeV. The isolated photon must have  $E_T > 50$  GeV. They observe 91 events with 87.2  $\pm$  7.8 events expected from standard model processes. For the  $\nu\overline{\nu}$  case they require solitary photons with  $E_T > 25$  GeV and observe 85 events with standard model expectation of 85.9  $\pm$  5.6 events. Taking the form factor  $\Lambda = 1.5$  TeV they derive 95% C.L. limits as  $|h_3^{\gamma}, Z| < 0.022$  and  $|h_4^{\gamma}, Z| < 0.0009$ .
- <sup>9</sup> CHATRCHYAN 11M study  $Z\gamma$  production in pp collisions at  $\sqrt{s} = 7$  TeV using 36 pb<sup>-1</sup> pp data, where the Z decays to  $e^+e^-$  or  $\mu^+\mu^-$ . The total cross sections are measured for photon transverse energy  $E_T^{\gamma} > 10$  GeV and spatial separation from charged leptons in the plane of pseudo rapidity and azimuthal angle  $\Delta R(\ell,\gamma) > 0.7$  with the dilepton invariant mass requirement of  $M_{\ell\ell} > 50$  GeV. The number of  $e^+e^-\gamma$  and  $\mu^+\mu^-\gamma$  candidates is 81 and 90 with estimated backgrounds of  $20.5 \pm 2.5$  and  $27.3 \pm 3.2$  events respectively. The 95% CL limits for  $ZZ\gamma$  couplings are  $-0.05 < h_3^{\gamma} < 0.06$  and  $-0.0005 < h_4^{\gamma} < 0.0005$ , and for  $Z\gamma\gamma$  couplings are  $-0.07 < h_3^{\gamma} < 0.07$  and  $-0.0005 < h_4^{\gamma} < 0.0006$ .
- <sup>10</sup> ABAZOV 09L study  $Z\gamma$ ,  $Z \rightarrow \nu \overline{\nu}$  production in  $p\overline{p}$  collisions at 1.96 TeV C.M. energy. They select 51 events with a photon of transverse energy  $E_T$  larger than 90 GeV, with an expected background of 17 events. Based on the photon  $E_T$  spectrum and including also Z decays to charged leptons (from ABAZOV 07M), the following 95% CL limits are reported:  $|h_{30}^{\gamma}| < 0.033$ ,  $|h_{40}^{\gamma}| < 0.0017$ ,  $|h_{30}^{Z}| < 0.033$ ,  $|h_{40}^{\gamma}| < 0.0017$ .
- <sup>11</sup> ABAZOV 07M use 968  $p\overline{p} \rightarrow e^+e^-/\mu^+\mu^-\gamma X$  candidates, at 1.96 TeV center of mass energy, to tag  $p\overline{p} \rightarrow Z\gamma$  events by requiring  $E_T(\gamma) > 7$  GeV, lepton-gamma separation  $\Delta R_{\ell\gamma} > 0.7$ , and di-lepton invariant mass > 30 GeV. The cross section is in agreement with the SM prediction. Using these  $Z\gamma$  events they obtain 95% C.L. limits on each

 $h_i^V$ , keeping all others fixed at their SM values. They report:  $-0.083 < h_{30}^Z < 0.082$ ,  $-0.0053 < h_{40}^Z < 0.0054$ ,  $-0.085 < h_{30}^\gamma < 0.084$ ,  $-0.0053 < h_{40}^\gamma < 0.0054$ , for the form factor scale  $\Lambda = 1.2$  TeV.

- <sup>12</sup> Using data collected at  $\sqrt{s} = 183-208$ , ABDALLAH 07C select 1,877  $e^+e^- \rightarrow Z\gamma$ events with  $Z \rightarrow q\overline{q}$  or  $\nu\overline{\nu}$ , 171  $e^+e^- \rightarrow ZZ$  events with  $Z \rightarrow q\overline{q}$  or lepton pair (except an explicit  $\tau$  pair), and 74  $e^+e^- \rightarrow Z\gamma^*$  events with a  $q\overline{q}\mu^+\mu^-$  or  $q\overline{q}e^+e^$ signature, to derive 95% CL limits on  $h_i^V$ . Each limit is derived with other parameters set to zero. They report:  $-0.23 < h_1^Z < 0.23$ ,  $-0.30 < h_3^Z < 0.16$ ,  $-0.14 < h_1^{\gamma} < 0.14$ ,  $-0.049 < h_3^{\gamma} < 0.044$ .
- <sup>13</sup> ACHARD 04H select 3515  $e^+e^- \rightarrow Z\gamma$  events with  $Z \rightarrow q \overline{q}$  or  $\nu \overline{\nu}$  at  $\sqrt{s} = 189-209$  GeV to derive 95% CL limits on  $h_i^V$ . For deriving each limit the other parameters are fixed at zero. They report:  $-0.153 < h_1^Z < 0.141$ ,  $-0.087 < h_2^Z < 0.079$ ,  $-0.220 < h_3^Z < 0.112$ ,  $-0.068 < h_4^Z < 0.148$ ,  $-0.057 < h_1^\gamma < 0.057$ ,  $-0.050 < h_2^\gamma < 0.023$ ,  $-0.059 < h_3^\gamma < 0.004$ ,  $-0.004 < h_4^\gamma < 0.042$ .
- <sup>14</sup> ABBIENDI,G 00C study  $e^+e^- \rightarrow Z\gamma$  events (with  $Z \rightarrow q\overline{q}$  and  $Z \rightarrow \nu\overline{\nu}$ ) at 189 GeV to obtain the central values (and 95% CL limits) of these couplings:  $h_1^Z = 0.000 \pm 0.100 \ (-0.190, 0.190), \ h_2^Z = 0.000 \pm 0.068 \ (-0.128, 0.128), \ h_3^Z = -0.074^{+0.102}_{-0.103} \ (-0.269, 0.119), \ h_4^Z = 0.046 \pm 0.068 \ (-0.084, 0.175), \ h_1^{\gamma} = 0.000 \pm 0.061 \ (-0.115, 0.115), \ h_2^{\gamma} = 0.000 \pm 0.041 \ (-0.077, 0.077), \ h_3^{\gamma} = -0.080^{+0.039}_{-0.041} \ (-0.164, -0.006), \ h_4^{\gamma} = 0.064^{+0.033}_{-0.030} \ (+0.007, +0.134).$  The results are derived assuming that only one coupling at a time is different from zero.
- <sup>15</sup> ABBOTT 98M study  $p\overline{p} \to Z\gamma + X$ , with  $Z \to e^+e^-$ ,  $\mu^+\mu^-$ ,  $\overline{\nu}\nu$  at 1.8 TeV, to obtain 95% CL limits at  $\Lambda = 750 \text{ GeV}$ :  $|h_{30}^Z| < 0.36$ ,  $|h_{40}^Z| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ), and  $|h_{30}^{\gamma}| < 0.37$ ,  $|h_{40}^{\gamma}| < 0.05$  (keeping  $h_i^Z = 0$ ). Limits on the *CP*-violating couplings are  $|h_{10}^Z| < 0.36$ ,  $|h_{20}^Z| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ), and  $|h_{10}^{\gamma}| < 0.37$ ,  $|h_{20}^{\gamma}| < 0.05$  (keeping  $h_i^{\gamma} = 0$ ).

<sup>16</sup> ABREU 98K determine a 95% CL upper limit on  $\sigma(e^+e^- \rightarrow \gamma + \text{invisible particles}) < 2.5 \text{ pb using 161 and 172 GeV data}$ . This is used to set 95% CL limits on  $|h_{30}^{\gamma}| < 0.8$  and  $|h_{30}^{Z}| < 1.3$ , derived at a scale  $\Lambda = 1$  TeV and with n = 3 in the form factor representation.

## f

Combining the LEP-2 results taking into account the correlations, the following 95% CL limits are derived [SCHAEL 13A]:

$$\begin{aligned} -0.28 < f_4^Z < +0.32, & -0.34 < f_5^Z < +0.35, \\ -0.17 < f_4^\gamma < +0.19, & -0.35 < f_5^\gamma < +0.32. \end{aligned}$$

Some of the recent results from the Tevatron and LHC experiments individually surpass the combined LEP-2 results in precision (see below).

VALUE	DOCUMENT ID	TECN COMMENT	
$\bullet \bullet \bullet$ We do not use the following	g data for averages	es, fits, limits, etc. • • •	
	<sup>1</sup> AABOUD	CIII	
	<sup>2</sup> SIRUNYAN	18BT CMS $E_{cm}^{pp} = 13 \text{ TeV}$	
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<sup>3</sup> KHACHATRY	. <b>15</b> B	CMS	$E^{pp}_{cm} = 8 \; \text{TeV}$
<sup>4</sup> KHACHATRY	. <b>15</b> BC	CMS	$E^{pp}_{cm}=$ 7, 8 TeV
<sup>5</sup> AAD	13z	ATLS	$E_{\rm cm}^{pp} =$ 7 TeV
<sup>6</sup> CHATRCHYAN	<b>13</b> B	CMS	$E_{\rm cm}^{pp} =$ 7 TeV
<sup>7</sup> SCHAEL	09		$E_{\rm cm}^{ee} = 192209~{ m GeV}$
<sup>8</sup> ABAZOV	08K	D0	$E_{ m cm}^{p\overline{p}}=1.96~ m TeV$
<sup>9</sup> ABDALLAH	<b>07</b> C	DLPH	$E_{\rm cm}^{ee} = 183-208  {\rm GeV}$
<sup>10</sup> ABBIENDI	04C	OPAL	
<sup>11</sup> ACHARD	<b>03</b> D	L3	

- <sup>1</sup>AABOUD 18Q study  $pp \rightarrow ZZ$  events at  $\sqrt{s} = 13$  TeV with  $Z \rightarrow e^+e^-$  or  $Z \rightarrow \mu^+\mu^-$ . The number of events observed in the 4e, 2e 2 $\mu$ , and 4 $\mu$  channels is 249, 465, and 303 respectively. Analysing the  $p_T$  spectrum of the leading Z boson, the following the following 95% C.L. limits are derived in units of  $10^{-4}$ :  $-1.8 < f_4^{\gamma} < 1.8$ ,  $-1.5 < f_4^Z < 1.5$ ,  $-1.8 < f_5^{\gamma} < 1.8$ ,  $-1.5 < f_5^Z < 1.5$ .
- <sup>2</sup>SIRUNYAN 18BT study ppZZ events at  $\sqrt{s} = 13$  TeV with  $Z \rightarrow e^+e^-$  or  $Z \rightarrow \mu^+\mu^-$ . The number of events observed in the 4e,  $2e2\mu$ , and  $4\mu$  channels is 220, 543 and 335 respectively. Analysing the 4-lepton invariant mass spectrum, the following 95% C.L. limits are derived in units of  $10^{-3}$ :  $-1.2 < f_4^{\gamma} < 1.3$ ,  $-1.2 < f_4^Z < 1.0$ ,  $-1.2 < f_5^{\gamma} < 1.3$ ,  $-1.0 < f_5^Z < 1.3$ .
- <sup>3</sup> KHACHATRYAN 15B study ZZ production in 8 TeV pp collisions. In the decay modes  $ZZ \rightarrow 4e, 4\mu, 2e2\mu, 54, 75, 148$  events are observed, with an expected background of  $2.2 \pm 0.9, 1.2 \pm 0.6$ , and  $2.4 \pm 1.0$  events, respectively. Analysing the 4-lepton invariant mass spectrum in the range from 110 GeV to 1200 GeV, the following 95% C.L. limits are obtained:  $|f_A^Z| < 0.004, |f_5^Z| < 0.004, |f_A^\gamma| < 0.005, |f_5^\gamma| < 0.005.$
- <sup>4</sup> KHACHATRYAN 15BC use the cross section measurement of the final state  $pp \rightarrow ZZ \rightarrow 2\ell 2\nu$ , ( $\ell$  being an electron or a muon) at 7 and 8 TeV to put limits on these triple gauge couplings. Effective mass of the charged lepton pair is required to be in the range 83.5–98.5 GeV and the dilepton  $p_T > 45$  GeV. The reduced missing  $E_T$  is required to be > 65 GeV, which takes into account the fake missing  $E_T$  due to detector effects. The numbers of  $e^+e^-$  and  $\mu^+\mu^-$  events selected are 35 and 40 at 7 TeV and 176 and 271 at 8 TeV respectively. The production cross sections so obtained are in agreement with SM predictions. The following 95% C.L. limits are set:  $-0.0028 < f_4^Z < 0.0032$ ,  $-0.0037 < f_4^{\gamma} < 0.0033$ ,  $-0.0029 < f_5^Z < 0.0031$ ,  $-0.0033 < f_5^{\gamma} < 0.0037$ . Combining with previous results (KHACHATRYAN 15B and CHATRCHYAN 13B) which include 7 TeV and 8 TeV data on the final states  $pp \rightarrow ZZ \rightarrow 2\ell 2\ell'$  where  $\ell$  and  $\ell'$  are an electron or a muon, the best limits are  $-0.0022 < f_4^Z < 0.0026$ ,  $-0.0023 < f_5^Z < 0.0023$ ,  $-0.0026 < f_5^{\gamma} < 0.0027$ .
- <sup>5</sup> AAD 13Z study ZZ production in pp collisions at  $\sqrt{s} = 7$  TeV. In the ZZ  $\rightarrow \ell^+ \ell^- \ell'^+ \ell'^-$  final state they observe a total of 66 events with an expected background of  $0.9 \pm 1.3$ . In the ZZ  $\rightarrow \ell^+ \ell^- \nu \nu$  final state they observe a total of 87 events with an expected background of  $46.9 \pm 5.2$ . The limits on anomalous TGCs are determined using the observed and expected numbers of these ZZ events binned in  $p_T^Z$ . The 95% C.L. are as follows: for form factor scale  $\Lambda = \infty$ ,  $-0.015 < f_4^{\gamma} < 0.015$ ,  $-0.013 < f_5^Z < 0.013$ ; for form factor scale  $\Lambda = \infty$

- 3 TeV,  $-0.022 < f_4^{\gamma} < 0.023$ ,  $-0.019 < f_4^Z < 0.019$ ,  $-0.023 < f_5^{\gamma} < 0.023$ ,  $-0.020 < f_5^Z < 0.019$ .
- <sup>6</sup> CHATRCHYAN 13B study ZZ production in pp collisions and select 54 ZZ candidates in the Z decay channel with electrons or muons with an expected background of  $1.4 \pm 0.5$ events. The resulting 95% C.L. ranges are:  $-0.013 < f_4^{\gamma} < 0.015, -0.011 < f_4^{Z} < 0.012, -0.014 < f_5^{\gamma} < 0.014, -0.012 < f_5^{Z} < 0.012.$
- <sup>7</sup> Using data collected in the center of mass energy range 192–209 GeV, SCHAEL 09 select 318  $e^+e^- \rightarrow ZZ$  events with 319.4 expected from the standard model. Using this data they derive the following 95% CL limits:  $-0.321 < f_4^{\gamma} < 0.318$ ,  $-0.534 < f_4^Z < 0.534$ ,  $-0.724 < f_5^{\gamma} < 0.733$ ,  $-1.194 < f_5^Z < 1.190$ .
- <sup>8</sup> ABAZOV 08K search for ZZ and  $Z\gamma^*$  events with  $1 \text{ fb}^{-1} p\overline{p}$  data at  $\sqrt{s} = 1.96 \text{ TeV}$  in (ee)(ee),  $(\mu\mu)(\mu\mu)$ ,  $(ee)(\mu\mu)$  final states requiring the lepton pair masses to be > 30 GeV. They observe 1 event, which is consistent with an expected signal of  $1.71 \pm 0.15$  events and a background of  $0.13 \pm 0.03$  events. From this they derive the following limits, for a form factor ( $\Lambda$ ) value of 1.2 TeV:  $-0.28 < f_{40}^Z < 0.28$ ,  $-0.31 < f_{50}^Z < 0.28$

 $0.29, \, -0.26 < f_{40}^{\gamma} < 0.26, \, -0.30 < f_{50}^{\gamma} < 0.28.$ 

<sup>9</sup> Using data collected at  $\sqrt{s} = 183-208$  GeV, ABDALLAH 07C select 171  $e^+e^- \rightarrow ZZ$ events with  $Z \rightarrow q\bar{q}$  or lepton pair (except an explicit  $\tau$  pair), and 74  $e^+e^- \rightarrow Z\gamma^*$ events with a  $q\bar{q}\mu^+\mu^-$  or  $q\bar{q}e^+e^-$  signature, to derive 95% CL limits on  $f_i^V$ . Each limit is derived with other parameters set to zero. They report:  $-0.40 < f_A^Z < 0.42$ ,

$$-0.38 < f_5^Z < 0.62, -0.23 < f_4^\gamma < 0.25, -0.52 < f_5^\gamma < 0.48.$$

<sup>10</sup> ABBIENDI 04C study ZZ production in  $e^+e^-$  collisions in the C.M. energy range 190–209 GeV. They select 340 events with an expected background of 180 events. Including the ABBIENDI 00N data at 183 and 189 GeV (118 events with an expected background of 65 events) they report the following 95% CL limits:  $-0.45 < f_4^Z < 0.58$ ,

$$-0.94 < f_5^Z < 0.25, -0.32 < f_4^{\gamma} < 0.33, \text{ and } -0.71 < f_5^{\gamma} < 0.59.$$

<sup>11</sup> ACHARD 03D study Z-boson pair production in  $e^+e^-$  collisions in the C.M. energy range 200–209 GeV. They select 549 events with an expected background of 432 events. Including the ACCIARRI 99G and ACCIARRI 990 data (183 and 189 GeV respectively, 286 events with an expected background of 241 events) and the 192–202 GeV ACCIARRI 011 results (656 events, expected background of 512 events), they report the following 95% CL limits:  $-0.48 \le f_4^Z \le 0.46$ ,  $-0.36 \le f_5^Z \le 1.03$ ,  $-0.28 \le f_4^\gamma \le 0.28$ , and  $-0.40 \le f_5^\gamma \le 0.47$ .

## ANOMALOUS W/Z QUARTIC COUPLINGS

## See the related review(s):

Anomalous W/Z Quartic Couplings (QGCs)

# $a_0/\Lambda^2$ , $a_c/\Lambda^2$

Combining published and unpublished preliminary LEP results the following 95% CL intervals for the QGCs associated with the  $ZZ\gamma\gamma$  vertex are derived (CERN-PH-EP/2005-051 or hep-ex/0511027):

$$-0.008 < a_0^Z / \Lambda^2 < +0.021$$
  
$$-0.029 < a_C^Z / \Lambda^2 < +0.039$$

Anomalous Z quartic couplings have also been measured by the Tevatron and LHC experiments. As discussed in the review on "Anomalous W/Z quartic couplings," the coupling parameters in the Anomalous QGC Lagrangian may relate to processes involving only the W or only to the Z or to both. Thus, results on all other AQGCs are reported together in the W listings.

VALUE	DOCUMENT ID	TECN	
• • • We do not use the following d	lata for averages, fits,	limits, etc.	

<sup>1</sup> ABBIENDI	04L	OPAL
<sup>2</sup> HEISTER	04A	ALEP
<sup>3</sup> ACHARD	026	13

- <sup>1</sup>ABBIENDI 04L select 20  $e^+e^- \rightarrow \nu \overline{\nu} \gamma \gamma$  acoplanar events in the energy range 180–209 GeV and 176  $e^+e^- \rightarrow q \bar{q} \gamma \gamma$  events in the energy range 130–209 GeV. These samples are used to constrain possible anomalous  $W^+W^-\gamma\gamma$  and  $ZZ\gamma\gamma$  quartic couplings. Further combining with the  $W^+W^-\gamma$  sample of ABBIENDI 04B the following oneparameter 95% CL limits are obtained:  $-0.007 < a_0^Z/\Lambda^2 < 0.023 \text{ GeV}^{-2}$ ,  $-0.029 < 0.023 \text{ GeV}^{-2}$  $a_c^Z/\Lambda^2 < 0.029 \; {
  m GeV}^{-2}$ ,  $-0.020 < a_0^W/\Lambda^2 < 0.020 \; {
  m GeV}^{-2}$ ,  $-0.052 < a_c^W/\Lambda^2 < 0.020 \; {
  m GeV}^{-2}$ ,  $-0.052 < a_c^W/\Lambda^2 < 0.020 \; {
  m GeV}^{-2}$  $0.037 \text{ GeV}^{-2}$ .
- <sup>2</sup> In the CM energy range 183 to 209 GeV HEISTER 04A select 30  $e^+e^- \rightarrow \nu \overline{\nu} \gamma \gamma$  events with two acoplanar, high energy and high transverse momentum photons. The photonphoton acoplanarity is required to be  $>5^{\circ},~E_{\gamma}/\sqrt{s}~>0.025$  (the more energetic photon having energy > 0.2  $\sqrt{s}),~{\rm p}_{T_{\gamma}}/{\rm E_{beam}}~>$  0.05 and  $\left|\cos\,\theta_{\gamma}\right|~<$  0.94. A likelihood fit to the photon energy and recoil missing mass yields the following one-parameter 95% CL limits:  $-0.012 < a_0^Z/\Lambda^2 < 0.019 \text{ GeV}^{-2}$ ,  $-0.041 < a_c^Z/\Lambda^2 < 0.044 \text{ GeV}^{-2}$ ,  $-0.060 < a_0^W/\Lambda^2 < 0.055 \text{ GeV}^{-2}$ ,  $-0.099 < a_c^W/\Lambda^2 < 0.093 \text{ GeV}^{-2}$ .
- <sup>3</sup>ACHARD 02G study  $e^+e^- \rightarrow Z\gamma\gamma \rightarrow q\overline{q}\gamma\gamma$  events using data at center-of-mass energies from 200 to 209 GeV. The photons are required to be isolated, each with energy >5 GeV and  $|\cos\theta| < 0.97$ , and the di-jet invariant mass to be compatible with that of the Z boson (74–111 GeV). Cuts on Z velocity ( $\beta < 0.73$ ) and on the energy of the most energetic photon reduce the backgrounds due to non-resonant production of the  $q \bar{q} \gamma \gamma$  state and due to ISR respectively, yielding a total of 40 candidate events of which 8.6 are expected to be due to background. The energy spectra of the least energetic photon are fitted for all ten center-of-mass energy values from 130 GeV to 209 GeV (as obtained adding to the present analysis 130–202 GeV data of ACCIARRI 01E, for a total of 137 events with an expected background of 34.1 events) to obtain the fitted values  $a_0/\Lambda^2 = 0.00 \stackrel{+0.02}{_{-0.01}} \text{ GeV}^{-2}$  and  $a_c/\Lambda^2 = 0.03 \stackrel{+0.01}{_{-0.02}} \text{ GeV}^{-2}$ , where the other parameter is kept fixed to its Standard Model value (0). A simultaneous fit to both parameters yields the 95% CL limits  $-0.02 \text{ GeV}^{-2} < a_0/\Lambda^2 < 0.03 \text{ GeV}^{-2}$  and -0.07 $GeV^{-2} < a_c/\Lambda^2 < 0.05 GeV^{-2}$ .

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RAINBOLT	19	PR D99 013004	J.L. Rainbolt, M. Schmitt		(NWES)
AABOUD	18AU	JHEP 1807 127	M. Aaboud <i>et al.</i>	(ATLAS	Collab.)
AABOUD	18BL	PL B786 134	M. Aaboud <i>et al.</i>	ATLAS	Collab.)
AABOUD	18CN	PR D98 092010	M. Aaboud <i>et al.</i>	(ATLAS	Collab.)
AABOUD		PR D97 032005	M. Aaboud <i>et al.</i>	(ATLAS	Collab.)
AAIJ	18AR	JHEP 1809 159	R. Aaij <i>et al.</i>	(LHCb	Collab.)
ANDREEV	18A	EPJ C78 777	V. Andreev et al.	(H1	Collab.)
SIRUNYAN	18BT	EPJ C78 165	A.M. Sirunyan <i>et al.</i>	(CMS	Collab.)
SIRUNYAN	18DZ	PRL 121 141801	A.M. Sirunyan <i>et al.</i>	(CMS	Collab.)
AABOUD	16K	PRL 117 111802	M. Aaboud <i>et al.</i>	(ATLAS	Collab.)
AAD	16L	EPJ C76 210	G. Aad <i>et al.</i>	(ATLAS	Collab.)
AAD	16Q	PR D93 112002	G. Aad <i>et al.</i>	(ATLAS	Collab.)

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ABRAMOWICZ 16A PR D93 092002 PR D94 052007 ABT 16 KHACHATRY... 16AE PL B760 448 KHACHATRY... 16CC PL B763 280 15BT JHEP 1509 049 AAD AAD PRL 114 121801 15I KHACHATRY... 15AC JHEP 1504 164 KHACHATRY... 15B PL B740 250 KHACHATRY... 15BC EPJ C75 511 AAD 14AU PR D90 072010 AAD 14N PRL 112 231806 AALTONEN 14E PRL 112 111803 CHATRCHYAN 14AB PR D89 092005 13AN PR D87 112003 AAD PR D91 119901 (e Also JHEP 1303 128 AAD 13Z CHATRCHYAN 13B JHEP 1301 063 CHATRCHYAN 13BI JHEP 1310 164 SCHAEL PRPL 532 119 13A AAD 12BX PL B717 49 ABAZOV 12S PR D85 052001 CHATRCHYAN 12BN JHEP 1212 034 AALTONEN PRL 107 051802 11S PR D84 012007 ABAZOV 11D CHATRCHYAN 11M PL B701 535 PRL 102 201802 ABAZOV 09L BEDDALL 09 PL B670 300 SCHAEL 09 JHEP 0904 124 PRL 100 131801 ABAZOV 08K ABAZOV PL B653 378 07M ABDALLAH EPJ C51 525 07C ABDALLAH 06E PL B639 179 AKTAS 06 PL B632 35 LEP-SLC PRPL 427 257 06 SCHAEL 06A PL B639 192 ABDALLAH 05 EPJ C40 1 ABDALLAH EPJ C44 299 05C PRL 94 091801 ABE 05 ABE 05F PR D71 112004 PR D71 052002 ACOSTA 05M ABBIENDI 04B PL B580 17 ABBIENDI 04C EPJ C32 303 04E PL B586 167 ABBIENDI ABBIENDI 04G EPJ C33 173 ABBIENDI PR D70 032005 04L ABDALLAH 04F EPJ C34 109 PR D69 072003 ABE 04C ACHARD PL B585 42 04C ACHARD 04H PL B597 119 HEISTER 04A PL B602 31 03P PL B577 18 ABBIENDI PL B569 129 ABDALLAH 03H ABDALLAH 03K PL B576 29 ABE 03F PRL 90 141804 ACHARD PL B572 133 03D ACHARD 03G PL B577 109 ABBIENDI 02I PL B546 29 ABE 02G PRL 88 151801 ACHARD PL B540 43 02G HEISTER 02B PL B526 34 PL B528 19 HEISTER 02C HEISTER 02H EPJ C24 177 EPJ C19 587 ABBIENDI 01A EPJ C18 447 ABBIENDI 01G ABBIENDI 01K PL B516 1 EPJ C20 445 ABBIENDI 01N ABBIENDI 010 EPJ C21 1 PRL 86 1162 ABE 01B ABE 01C PR D63 032005 ACCIARRI 01E PL B505 47 ACCIARRI 011 PL B497 23

	U Abramowicz at al	(7EUS Collab.)
	H. Abramowicz <i>et al.</i>	(ZEUS Collab.)
	I. Abt <i>et al.</i>	(MPIM, OXF, HAMB, DESY)
	V Khachatrian at al	(CMS Collab.)
	V. Khachatryan <i>et al.</i>	
	V. Khachatryan <i>et al.</i>	(CMS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	V. Khachatryan <i>et al.</i>	(CMS Collab.)
	V. Khachatryan <i>et al.</i>	(CMS Collab.)
	V. Khachatryan et al.	(CMS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	T. Aaltonen <i>et al.</i>	(CDF Collab.)
	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
errat.)	G. Aad <i>et al.</i>	(ATLAS Collab.)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
	S. Schael <i>et al.</i>	(ALEPH Collab., DELPHI, L3+)
	G. Aad <i>et al.</i>	(ATLAS Collab.)
	V.M. Abazov <i>et al.</i>	(D0 Collab.)
	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
	T. Aaltonen <i>et al.</i>	(CDF Collab.)
	V.M. Abazov <i>et al.</i>	(D0 Collab.)
	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
	V.M. Abazov et al.	) (D0 Collab.)
	A. Beddall, A. Beddall,	A. Bingul (UGAZ)
	S. Schael <i>et al.</i>	(ALEPH Collab.)
	V.M. Abazov <i>et al.</i>	(D0 Collab.)
	V.M. Abazov <i>et al.</i>	(D0 Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
		· · · · · · · · · · · · · · · · · · ·
	A. Aktas <i>et al.</i>	(H1 Collab.)
	ALEPH, DELPHI, L3, O	PAL, SLD and working groups
	S. Schael <i>et al.</i>	(ALEPH Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	D. Acosta <i>et al.</i>	(CDF Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	J. Abdallah <i>et al.</i>	
		(DELPHI Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	P. Achard <i>et al.</i>	(L3 Collab.)
	P. Achard <i>et al.</i>	(L3 Collab.)
	A. Heister <i>et al.</i>	(ALEPH Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
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	K. Abe <i>et al.</i>	(SLD Collab.)
	P. Achard <i>et al.</i>	(L3 Collab.)
	P. Achard <i>et al.</i>	(L3 Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	P. Achard <i>et al.</i>	(L3 Collab.)
	A. Heister <i>et al.</i>	(ALEPH Collab.)
	A. Heister <i>et al.</i>	(ALEPH Collab.)
	A. Heister <i>et al.</i>	
		(ALEPH Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	
		(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	G. Abbiendi et al.	(OPAL Collab.)
	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	K. Abe <i>et al.</i>	(SLD Collab.)
	M. Acciarri <i>et al.</i>	(L3 Collab.)
	M. Acciarri <i>et al.</i>	(L3 Collab.)

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HEISTER	01	EPJ C20 401	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	01D	EPJ C22 201	A. Heister <i>et al.</i>	(ALEPH Collab.)
ABBIENDI	00N	PL B476 256	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI,G	00C	EPJ C17 553	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABE	00B	PRL 84 5945	K. Abe <i>et al.</i>	(SLD Collab.)
ABE	00D	PRL 85 5059	K. Abe <i>et al.</i>	(SLD Collab.)
ABREU	00	EPJ C12 225	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00B	EPJ C14 613	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00B	EPJ C14 585	P. Abreu <i>et al.</i>	(DELPHI Collab.)
	00E		P. Abreu <i>et al.</i>	
ABREU		EPJ C16 371	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00P	PL B475 429		(DELPHI Collab.)
ACCIARRI	00	EPJ C13 47	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00C	EPJ C16 1	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00J	PL B479 79	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00Q	PL B489 93	M. Acciarri <i>et al.</i>	(L3 Collab.)
BARATE	00B	EPJ C16 597	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	00C	EPJ C14 1	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	000	EPJ C16 613	R. Barate <i>et al.</i>	(ALEPH Collab.)
ABBIENDI	99B	EPJ C8 217	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	991	PL B447 157	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABE	99E	PR D59 052001	K. Abe <i>et al.</i>	(SLD Collab.)
ABE	99L	PRL 83 1902	K. Abe <i>et al.</i>	(SLD Collab.)
ABREU	99	EPJ C6 19	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99B	EPJ C10 415	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99 J	PL B449 364	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99U	PL B462 425	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99Y	EPJ C10 219	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	99D	PL B448 152	M. Acciarri et al.	(L3 Collab.)
ACCIARRI	99F	PL B453 94	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99G	PL B450 281	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	990	PL B465 363	M. Acciarri <i>et al.</i>	(L3 Collab.)
ABBOTT	98M	PR D57 3817	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	98D	PRL 80 660	K. Abe <i>et al.</i>	(SLD Collab.)
ABE	98I	PRL 81 942	K. Abe et al.	(SLD Collab.)
ABREU	98K	PL B423 194	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	98L	EPJ C5 585	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	98C 98G	PL B431 199	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	98G 98H	PL B431 199 PL B429 387	M. Acciarri <i>et al.</i> M. Acciarri <i>et al.</i>	
	98U	PL B439 225	M. Acciarri <i>et al.</i> M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI				(L3 Collab.)
ACKERSTAFF	98A	EPJ C5 411	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	98E	EPJ C1 439	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	980	PL B420 157	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	98Q	EPJ C4 19	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
BARATE	980 00T	PL B434 415	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98T	EPJ C4 557	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98V	EPJ C5 205	R. Barate <i>et al.</i>	(ALEPH Collab.)
ABE	97	PRL 78 17	K. Abe <i>et al.</i>	(SLD Collab.)
ABREU	97C	ZPHY C73 243	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	97E	PL B398 207	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	97G	PL B404 194	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	97D	PL B393 465	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	97J	PL B407 351	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	97L	PL B407 389	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	97R	PL B413 167	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACKERSTAFF		ZPHY C74 413	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	97S	PL B412 210	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF		ZPHY C76 387	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	97W	ZPHY C76 425	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ALEXANDER	97C	ZPHY C73 379	G. Alexander <i>et al.</i>	(OPAL Collab.)
ALEXANDER	97D	ZPHY C73 569	G. Alexander <i>et al.</i>	(OPAL Collab.)
ALEXANDER	97E	ZPHY C73 587	G. Alexander <i>et al.</i>	(OPAL Collab.)
BARATE	97D	PL B405 191	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	97E	PL B401 150	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	97F	PL B401 163	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	97H	PL B402 213	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	97J	ZPHY C74 451	R. Barate <i>et al.</i>	(ALEPH Collab.)
ABREU	96R	ZPHY C72 31	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	96S	PL B389 405	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	96U	ZPHY C73 61	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	96	PL B371 126	M. Acciarri <i>et al.</i>	(L3 Collab.)
ADAM	96	ZPHY C69 561	W. Adam <i>et al.</i>	(DELPHI Collab.)

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	06 P	7DUV C70 271	))/ Adams at al	
ADAM ALEXANDER	96B 96B	ZPHY C70 371 ZPHY C70 197	W. Adam <i>et al.</i> G. Alexander <i>et al.</i>	(DELPHI Collab.)
ALEXANDER	906 96F	PL B370 185	G. Alexander <i>et al.</i> G. Alexander <i>et al.</i>	(OPAL Collab.) (OPAL Collab.)
ALEXANDER	96N	PL B384 343	G. Alexander <i>et al.</i>	(OPAL Collab.)
ALEXANDER	96R	ZPHY C72 1	G. Alexander <i>et al.</i>	(OPAL Collab.)
BUSKULIC	96D	ZPHY C69 393	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
BUSKULIC	96H	ZPHY C69 379	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
BUSKULIC	96T	PL B384 449	D. Buskulic et al.	(ALEPH Collab.)
BUSKULIC	96Y	PL B388 648	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
ABE	95 J	PRL 74 2880	K. Abe <i>et al.</i>	(SLD Collab.)
ABREU	95 05 D		(erratum)P. Abreu <i>et al.</i> P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU ABREU	95D 95L	ZPHY C66 323 ZPHY C65 587	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i>	(DELPHI Collab.) (DELPHI Collab.)
ABREU	95L 95M	ZPHY C65 603	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	950	ZPHY C67 543	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	95R	ZPHY C68 353	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	95V	ZPHY C68 541	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	95W	PL B361 207	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	95X	ZPHY C69 1	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	95B	PL B345 589	M. Acciarri <i>et al.</i>	(L3 Collab.)
acciarri Acciarri	95C 95G	PL B345 609 PL B353 136	M. Acciarri <i>et al.</i> M. Acciarri <i>et al.</i>	(L3 Collab.) (L3 Collab.)
AKERS	95G 95C	ZPHY C65 47	R. Akers <i>et al.</i>	(OPAL Collab.)
AKERS	95U	ZPHY C67 389	R. Akers <i>et al.</i>	(OPAL Collab.)
AKERS	95W	ZPHY C67 555	R. Akers <i>et al.</i>	(OPAL Collab.)
AKERS	95X	ZPHY C68 1	R. Akers <i>et al.</i>	(OPAL Collab.)
AKERS	95Z	ZPHY C68 203	R. Akers <i>et al.</i>	(OPAL Collab.)
ALEXANDER	95D	PL B358 162	G. Alexander et al.	(OPAL Collab.)
BUSKULIC	95R	ZPHY C69 15	D. Buskulic <i>et al.</i>	(ALEPH Collab.)
MIYABAYASHI		PL B347 171	K. Miyabayashi <i>et al.</i>	(TOPAZ Collab.)
ABE ABREU	94C 94B	PRL 73 25 PL B327 386	K. Abe <i>et al.</i> P. Abreu <i>et al.</i>	(SLD Collab.) (DELPHI Collab.)
ABREU	94D 94P	PL B341 109	P. Abreu <i>et al.</i>	(DELPHI Collab.)
AKERS	94P	ZPHY C63 181	R. Akers <i>et al.</i>	(OPAL Collab.)
BUSKULIC	94G	ZPHY C62 179	D. Buskulic <i>et al.</i>	(ÀLEPH Collab.)
BUSKULIC	94J	ZPHY C62 1	D. Buskulic et al.	(ALEPH Collab.)
VILAIN	94	PL B320 203	P. Vilain <i>et al.</i>	(CHARM II Collab.)
ABREU	93	PL B298 236	P. Abreu <i>et al.</i>	(CHARM II Collab.) (DELPHI Collab.)
ABREU ABREU	-	PL B298 236 ZPHY C59 533	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.)
ABREU ABREU Also	93 93I	PL B298 236 ZPHY C59 533 ZPHY C65 709	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> (erratum)P. Abreu <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.)
ABREU ABREU Also ABREU	93 93I 93L	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> (erratum)P. Abreu <i>et al.</i> P. Abreu <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.)
ABREU ABREU Also	93 93I	PL B298 236 ZPHY C59 533 ZPHY C65 709	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> (erratum)P. Abreu <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.)
ABREU ABREU Also ABREU ACTON	93 931 93L 93	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> (erratum)P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> P.D. Acton <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.)
ABREU ABREU Also ABREU ACTON ACTON	93 93I 93L 93 93D 93E 93	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219	P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> (erratum)P. Abreu <i>et al.</i> P. Abreu <i>et al.</i> P.D. Acton <i>et al.</i> P.D. Acton <i>et al.</i> P.D. Acton <i>et al.</i> O. Adriani <i>et al.</i>	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.)
ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI ADRIANI	93 93I 93L 93 93D 93E 93 93I	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (DPAL Collab.) (L3 Collab.) (L3 Collab.)
ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI ADRIANI BUSKULIC	93 93I 93L 93 93D 93E 93 93I 93I 93L	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. D. Buskulic et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.)
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ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI ADRIANI BUSKULIC NOVIKOV ABREU	93 93I 93L 93 93D 93E 93 93I 93L 93C 92I	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520 PL B298 453 PL B277 371	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) M.I. Vysotsky (ITEP) (DELPHI Collab.)
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ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI ADRIANI BUSKULIC NOVIKOV ABREU	93 93I 93L 93 93D 93E 93 93I 93L 93C 92I	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520 PL B298 453 PL B277 371	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DPAL Collab.) (OPAL Collab.) (OPAL Collab.) (DPAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.)
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ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI BUSKULIC NOVIKOV ABREU ABREU ACTON ACTON ACTON ADEVA ADRIANI	93 93I 93L 93D 93E 93 93I 93L 93L 93C 92I 92M 92B 92B 92L 92N 92 92D	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520 PL B298 453 PL B277 371 PL B289 199 ZPHY C53 539 PL B294 436 PL B295 357 PL B275 209 PL B292 454	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. D. Acton et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al. D.P. Acton et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. B. Adeva et al. O. Adriani et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (DPAL Collab.) (CPAL Collab.) (CPAL Collab.) (L3 Collab.) (L3 Collab.)
ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI BUSKULIC NOVIKOV ABREU ABREU ACTON ACTON ACTON ACTON ADEVA ADRIANI ALITTI	93 93I 93L 93 93D 93E 93 93I 93L 93L 93C 92I 92M 92B 92L 92N 92 92D 92B	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520 PL B298 453 PL B277 371 PL B289 199 ZPHY C53 539 PL B294 436 PL B295 357 PL B275 209 PL B292 454 PL B276 354	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. B. Adeva et al. O. Adriani et al. J. Alitti et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DPAL Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (DAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.)
ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI BUSKULIC NOVIKOV ABREU ABREU ABREU ACTON ACTON ACTON ACTON ACTON ADEVA ADRIANI ALITTI BUSKULIC	93 931 932 93 935 938 938 931 932 932 921 922 922 922 922 922 922 922 922 92	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B313 520 PL B298 453 PL B277 371 PL B289 199 ZPHY C53 539 PL B294 436 PL B295 357 PL B275 209 PL B292 454 PL B276 354 PL B292 210	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. B. Adeva et al. O. Adriani et al. J. Alitti et al. D. Buskulic et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (DAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (DPAL Collab.) (CPAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.)
ABREU ABREU Also ABREU ACTON ACTON ACTON ADRIANI BUSKULIC NOVIKOV ABREU ABREU ACTON ACTON ACTON ACTON ACTON ACTON ALITTI BUSKULIC BUSKULIC	93 931 932 93 935 933 931 932 932 921 922 922 922 922 922 922 922 922 92	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B316 427 PL B316 427 PL B298 453 PL B277 371 PL B289 199 ZPHY C53 539 PL B294 436 PL B295 357 PL B275 209 PL B292 454 PL B276 354 PL B292 210 PL B294 145	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. B. Adeva et al. O. Adriani et al. J. Alitti et al. D. Buskulic et al. D. Buskulic et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L4 Collab.) (L4 Collab.)
ABREU ABREU ALso ABREU ACTON ACTON ACTON ADRIANI BUSKULIC NOVIKOV ABREU ABREU ACTON ACTON ACTON ACTON ACTON ACTON ADEVA ADRIANI ALITTI BUSKULIC BUSKULIC DECAMP	93 93I 93L 93D 93E 93 93I 93L 93L 92M 92B 92L 92M 92B 92L 92N 92 92D 92D 92E 92 92 92 92 92	PL B298 236 ZPHY C59 533 ZPHY C65 709 PL B318 249 PL B305 407 ZPHY C58 219 PL B311 391 PL B301 136 PL B316 427 PL B316 427 PL B318 520 PL B298 453 PL B277 371 PL B289 199 ZPHY C53 539 PL B294 436 PL B295 357 PL B295 357 PL B276 354 PL B292 210 PL B294 145 PRPL 216 253	P. Abreu et al. P. Abreu et al. (erratum)P. Abreu et al. P. Abreu et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. O. Adriani et al. O. Adriani et al. O. Adriani et al. D. Buskulic et al. V.A. Novikov, L.B. Okun, P. Abreu et al. D.P. Acton et al. P.D. Acton et al. P.D. Acton et al. P.D. Acton et al. B. Adeva et al. O. Adriani et al. J. Alitti et al. D. Buskulic et al.	(CHÀRM II Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (DELPHI Collab.) (DELPHI Collab.) (DELPHI Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (OPAL Collab.) (DPAL Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (L3 Collab.) (ALEPH Collab.) (ALEPH Collab.)
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