# SEARCHES not in other sections

#### Magnetic Monopole Searches

Isolated supermassive monopole candidate events have not been confirmed. The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$<~1.4\times10^{-16}~{\rm cm^{-2}sr^{-1}s^{-1}}~~{\rm for}~1.1\times10^{-4}<\beta<1$$

#### Supersymmetric Particle Searches

All supersymmetric mass bounds here are model dependent.

The limits assume:

1)  $\widetilde{\chi}_1^0$  is the lightest supersymmetric particle; 2) *R*-parity is conserved, unless stated otherwise;

See the Particle Listings for a Note giving details of supersymmetry.

$$\begin{array}{l} \widetilde{\chi}_i^0 - \text{neutralinos (mixtures of } \widetilde{\gamma}, \ \widetilde{Z}^0, \ \text{and } \widetilde{H}_i^0) \\ \text{Mass } m_{\widetilde{\chi}_1^0} > 0 \ \text{GeV, CL} = 95\% \\ \text{[general MSSM, non-universal gaugino masses]} \\ \text{Mass } m_{\widetilde{\chi}_1^0} > 46 \ \text{GeV, CL} = 95\% \\ \text{[all } \tan\beta, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \text{Mass } m_{\widetilde{\chi}_2^0} > 670 \ \text{GeV, CL} = 95\% \\ \text{[} 3/4\ell + \cancel{E}_T, \ \text{Tn2n3B, } m_{\widetilde{\chi}_1^0} < 200 \text{GeV}] \\ \text{Mass } m_{\widetilde{\chi}_3^0} > 670 \ \text{GeV, CL} = 95\% \\ \text{[} 3/4\ell + \cancel{E}_T, \ \text{Tn2n3B, } m_{\widetilde{\chi}_1^0} < 200 \text{GeV}] \\ \text{Mass } m_{\widetilde{\chi}_3^0} > 116 \ \text{GeV, CL} = 95\% \\ \text{[} 1<\tan\beta < 40, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \\ \widetilde{\chi}_i^{\pm} - \text{charginos (mixtures of } \widetilde{W}^{\pm} \ \text{and } \widetilde{H}_i^{\pm}) \\ \text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 94 \ \text{GeV, CL} = 95\% \\ \text{[} \tan\beta < 40, \ m_{\widetilde{\chi}_1^{\pm}} - m_{\widetilde{\chi}_1^0} > 3 \ \text{GeV, all } m_0] \\ \text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 500 \ \text{GeV, CL} = 95\% \\ \text{[} 2\ell^{\pm} + \cancel{E}_T, \ \text{Tchi1chi1B, } m_{\widetilde{\chi}_1^0} = 0 \ \text{GeV}] \\ \end{array}$$

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\tilde{\chi}^{\pm} — long-lived chargino
       Mass m_{\widetilde{\chi}^{\pm}} > 620 GeV, CL = 95\% [stable \widetilde{\chi}^{\pm}]
\widetilde{\nu} — sneutrino
       Mass m > 41 GeV, CL = 95\% [model independent]
       Mass m > 94 GeV, CL = 95\%
                [CMSSM, 1 \le \tan\beta \le 40, m_{\widetilde{e}_R} - m_{\widetilde{\chi}_1^0} > 10 GeV]
       Mass m > 2300 GeV, CL = 95\% [R-Parity Violating]
               [RPV, \widetilde{\nu}_{\tau} \rightarrow e \mu, \lambda'_{311} = 0.11]
\tilde{e} — scalar electron (selectron)
       Mass m(\tilde{e}_L) > 107 GeV, CL = 95\% [all m_{\tilde{e}_l} - m_{\tilde{\chi}_0^0}]
      Mass m > 410 GeV, CL = 95% [R-Parity Violating] [RPV, \geq 4\ell^{\pm}, \tilde{\ell} \rightarrow I \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow \ell^{\pm}\ell^{\mp}\nu]
\widetilde{\mu} — scalar muon (smuon)
       Mass m > 94 GeV, CL = 95\%
                [CMSSM, 1 \leq 	an\!eta \leq 40, m_{\widetilde{\mu}_R} - m_{\widetilde{\chi}_1^0} > 10 \ {
m GeV}]
       Mass m > 410 GeV, CL = 95\% [R-Parity Violating]
               [RPV, \geq 4\ell^{\pm}, \widetilde{\ell} \rightarrow I\widetilde{\chi}_{1}^{0}, \widetilde{\chi}_{1}^{0} \rightarrow \ell^{\pm}\ell^{\mp}\nu]
\widetilde{	au} — scalar tau (stau)
       Mass m > 81.9 GeV. CL = 95\%
              [m_{\widetilde	au_R}-m_{\widetilde\chi^0_1}\>>15 GeV, all 	heta_	au, B(\widetilde	au	o \;	au\widetilde\chi^0_1)=100%]
       Mass m > 286 GeV, CL = 95\% [long-lived \tilde{\tau}]
\tilde{q} – squarks of the first two quark generations
       Mass m > 1450 \text{ GeV}, CL = 95\%
              [CMSSM, tan\beta = 30, A_0 = -2max(m_0, m_{1/2}), \mu > 0]
       Mass m > 1550 \text{ GeV}, CL = 95\%
              [mass degenerate squarks]
       Mass m > 1050 \text{ GeV}, CL = 95\%
              [single light squark bounds]
       Mass m > 1.600 \times 10^3 GeV, CL = 95\% [R-Parity Violating]
              [RPV, \widetilde{q} \rightarrow q \widetilde{\chi}_1^0, \widetilde{\chi}_1^0 \rightarrow \ell \ell \nu, \lambda_{121}, \lambda_{122} \neq 0, m_{\widetilde{g}} = 2400 \text{GeV}]
\tilde{q} — long-lived squark
       Mass m > 1000, CL = 95%
              [\tilde{t}, \text{ charge-suppressed interaction model}]
       Mass m > 845, CL = 95\% [b, stable, Regge model]
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$$\begin{split} \widetilde{b} & -- \text{ scalar bottom (sbottom)} \\ & \text{Mass } m > \ 1230 \text{ GeV}, \text{ CL} = 95\% \\ & \text{ [jets+} E_T, \text{ Tsbot1, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 307 \text{ GeV}, \text{ CL} = 95\% \quad \text{[R-Parity Violating]} \\ & \text{[RPV, } \widetilde{b} \rightarrow \ t d \text{ or } ts, \lambda_{332}'' \text{ or } \lambda_{331}'' \text{ coupling]} \\ \widetilde{t} & -- \text{ scalar top (stop)} \\ & \text{Mass } m > \ 1120 \text{ GeV}, \text{ CL} = 95\% \\ & \text{[}1\ell + \text{jets+} E_T, \text{ Tstop1, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 610 \text{ GeV}, \text{ CL} = 95\% \quad \text{[R-Parity Violating]} \\ & \text{[RPV, 4 jets, Tstop1RPV, } \lambda_{323}'' \text{ coupling]} \\ \widetilde{g} & -- \text{gluino} \\ & \text{Mass } m > \ 1.860 \times 10^3 \text{ GeV}, \text{ CL} = 95\% \\ & \text{[} \geq 1 \text{ jets+} E_T, \text{ Tglu1A, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 2.260 \times 10^3 \text{ GeV}, \text{ CL} = 95\% \quad \text{[R-Parity Violating]} \\ & \text{[RPV, } \geq 4\ell, \lambda_{12k} \neq 0, m_{\widetilde{\chi}_1^0} > 1000 \text{ GeV]} \end{split}$$

#### Technicolor

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full *Review* (the data listings).

# Quark and Lepton Compositeness, Searches for

## Scale Limits $\Lambda$ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \overline{\psi}_L \gamma_\mu \psi_L \overline{\psi}_L \gamma^\mu \psi_L$$

(with  $g^2/4\pi$  set equal to 1), then we define  $\Lambda \equiv \Lambda_{LL}^{\pm}$ . For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

$$\Lambda_{LL}^+(eeee)$$
 > 8.3 TeV, CL = 95%  $\Lambda_{LL}^-(eeee)$  > 10.3 TeV, CL = 95%

$$\Lambda_{LL}^{+}(ee\mu\mu)$$
 > 8.5 TeV, CL = 95%

 $\Lambda_{LL}^{-}(ee\mu\mu)$  > 9.5 TeV, CL = 95%

 $\Lambda_{LL}^{-}(ee\tau\tau)$  > 7.9 TeV, CL = 95%

 $\Lambda_{LL}^{-}(ee\tau\tau)$  > 7.2 TeV, CL = 95%

 $\Lambda_{LL}^{-}(ee\tau\tau)$  > 7.2 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\ell\ell\ell\ell)$  > 9.1 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\ell\ell\ell\ell)$  > 10.3 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eeqq)$  > 24 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eeqq)$  > 37 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eeuu)$  > 23.3 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eeuu)$  > 12.5 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eedd)$  > 11.1 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eedd)$  > 26.4 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eecc)$  > 9.4 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eecc)$  > 5.6 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eebb)$  > 10.2 TeV, CL = 95%

 $\Lambda_{LL}^{-}(eebb)$  > 10.2 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\mu qq)$  > 20 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\mu qq)$  > 30 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\mu qq)$  > 30 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu qqq)$  > 13.1 none 17.4–29.5 TeV, CL = 95%

 $\Lambda_{LL}^{-}(qqqq)$  > 21.8 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\nu qq)$  > 5.0 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\nu qq)$  > 5.0 TeV, CL = 95%

 $\Lambda_{LL}^{-}(\mu\nu qq)$  > 5.4 TeV, CL = 95%

#### **Excited Leptons**

The limits from  $\ell^{*+}\ell^{*-}$  do not depend on  $\lambda$  (where  $\lambda$  is the  $\ell\ell^{*}$  transition coupling). The  $\lambda$ -dependent limits assume chiral coupling.

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e^{*\pm} — excited electron Mass m>103.2 GeV, CL = 95% (from e^*e^*) Mass m>3.000\times 10^3 GeV, CL = 95% (from e^*e^*) Mass m>356 GeV, CL = 95% (if \lambda_{\gamma}=1)
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u^{*\pm} — excited muon
    Mass m > 103.2 \text{ GeV}, CL = 95\% (from \mu^* \mu^*)
    Mass m > 3.000 \times 10^3 GeV. CL = 95\% (from \mu \mu^*)
\tau^{*\pm} — excited tau
    Mass m > 103.2 \text{ GeV}, CL = 95\% (from \tau^* \tau^*)
    Mass m > 2.500 \times 10^3 GeV. CL = 95\% (from \tau \tau^*)
\nu^* — excited neutrino
    Mass m > 1.600 \times 10^3 GeV, CL = 95\% (from \nu^* \nu^*)
    Mass m > 213 GeV, CL = 95\% (from \nu^* X)
a^* — excited quark
    Mass m > 338 \text{ GeV}, CL = 95\% (from q^* q^*)
    Mass m > 6.000 \times 10^3 GeV, CL = 95\% (from q^* X)
Color Sextet and Octet Particles
Color Sextet Quarks (q_6)
    Mass m > 84 GeV, CL = 95\% (Stable q_6)
Color Octet Charged Leptons (\ell_8)
    Mass m > 86 GeV, CL = 95\% (Stable \ell_8)
Color Octet Neutrinos (\nu_8)
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#### Extra Dimensions

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

Mass m > 110 GeV, CL = 90%  $(\nu_R \rightarrow \nu_R)$ 

### Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

$$R < 30~\mu\text{m}$$
, CL = 95% (direct tests of Newton's law)  $R < 4.8~\mu\text{m}$ , CL = 95% ( $pp \rightarrow jG$ )  $R < 0.16-916~\text{nm}$  (astrophysics; limits depend on technique and assumptions)

#### Constraints on the fundamental gravity scale

$$M_{TT}>9.02$$
 TeV, CL  $=95\%$  ( $p\,p\to$  dijet, angular distribution)  $M_{c}>4.16$  TeV, CL  $=95\%$  ( $p\,p\to\ell\,\overline{\ell}$ )

#### Constraints on the Kaluza-Klein graviton in warped extra dimensions

$$\mathit{M}_G$$
  $>$  4.25 TeV, CL  $=$  95%  $(\mathit{pp} \rightarrow \gamma \gamma)$ 

#### Constraints on the Kaluza-Klein gluon in warped extra dimensions

$$M_{g_{KK}} > 3.8$$
 TeV, CL  $= 95\% \quad (g_{KK} 
ightarrow t \, \overline{t})$