

## 108. $W'$ -Boson Searches

Revised November, 2019 by B.A. Dobrescu (Fermilab) and S. Willocq (U. Massachusetts).

The  $W'$  boson is a massive hypothetical particle of spin 1 and electric charge  $\pm 1$ , which is a color singlet and is predicted in various extensions of the Standard Model (SM).

### 108.1 $W'$ couplings to quarks and leptons

The Lagrangian terms describing the couplings of a  $W'^+$  boson to fermions are given by

$$\frac{W'^+_\mu}{\sqrt{2}} \left[ \bar{u}_i (C_{qij}^R P_R + C_{qij}^L P_L) \gamma^\mu d_j + \bar{\nu}_i (C_{\ell ij}^R P_R + C_{\ell ij}^L P_L) \gamma^\mu e_j \right]. \quad (108.1)$$

Here,  $u, d, \nu$ , and  $e$  are the SM fermions in the mass eigenstate basis,  $i, j = 1, 2, 3$  label the fermion generation, and  $P_{R,L} = (1 \pm \gamma_5)/2$ . The coefficients  $C_{qij}^L, C_{qij}^R, C_{\ell ij}^L$ , and  $C_{\ell ij}^R$  are complex dimensionless parameters. If  $C_{\ell ij}^R \neq 0$ , then the  $i$ th generation includes a right-handed neutrino. Using this notation, the SM  $W$  couplings are  $C_q^L = g V_{\text{CKM}}$ ,  $C_\ell^L = g \approx 0.63$  and  $C_q^R = C_\ell^R = 0$ .

Unitarity considerations imply that the  $W'$  boson is associated with a spontaneously-broken gauge symmetry. This is true even when it is a composite particle (*e.g.*  $\rho^\pm$ -like bound states [1]) if its mass is much smaller than the compositeness scale, or a Kaluza-Klein mode in theories where the  $W$  boson propagates in extra dimensions [2]. The simplest extension of the electroweak gauge group that includes a  $W'$  boson is  $SU(2)_1 \times SU(2)_2 \times U(1)$ , but larger groups are encountered in some theories. A generic property of these gauge theories is that they also include a  $Z'$  boson [3]; the  $W'$ -to- $Z'$  mass ratio is often a free parameter.

A tree-level mass mixing may be induced between the electrically-charged gauge bosons. Upon diagonalization of their mass matrix, the  $W$ -to- $Z$  mass ratio and the couplings of the observed  $W$  boson are shifted from the SM values. Their measurements imply that the mixing angle,  $\theta_+$ , between the gauge eigenstates must be smaller than about  $10^{-2}$ . In certain theories the mixing is negligible (*e.g.*, due to a new parity [4]), even when the  $W'$  mass is near the electroweak scale. Note that  $SU(2)$  gauge invariance suppresses the kinetic mixing between the  $W$  and  $W'$  bosons (in contrast to the case of a  $Z'$  boson [3]).

The  $W'$  coupling to  $WZ$  is fixed by Lorentz and gauge invariances, and to leading order in  $\theta_+$  is given by [5]

$$\frac{g \theta_+ i}{\cos \theta_W} \left[ W'^+_\mu (W^-_\nu Z^{\nu\mu} + Z_\nu W^{-\mu\nu}) + Z^\nu W^{-\mu} W'^+_{\nu\mu} \right] + \text{H.c.}, \quad (108.2)$$

where  $W^{\mu\nu} \equiv \partial^\mu W^\nu - \partial^\nu W^\mu$ , etc. The  $\theta_W$  dependence shown here corrects the one given in Ref. [6], which has been referred to as the Extended Gauge Model by the experimental collaborations. The  $W'$  coupling to  $Wh^0$ , where  $h^0$  is the SM Higgs boson, is

$$-\xi_h g_{W'} M_W W'^+_\mu W^{\mu-} h^0 + \text{H.c.}, \quad (108.3)$$

where  $g_{W'}$  is the gauge coupling of the  $W'$  boson, and the coefficient  $\xi_h$  satisfies  $\xi_h \leq 1$  in simple Higgs sectors [5].

In models based on the ‘‘left-right symmetric’’ gauge group [7],  $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$ , the SM fermions that couple to the  $W$  boson transform as doublets under  $SU(2)_L$  while the other fermions transform as doublets under  $SU(2)_R$ . Consequently, the  $W'$  boson couples primarily to right-handed fermions; its coupling to left-handed fermions arises due to the  $\theta_+$  mixing, so that  $C_q^L$  is proportional to the CKM matrix and its elements are much smaller than the diagonal elements of  $C_q^R$ . Generically,  $C_q^R$  does not need to be proportional to  $V_{\text{CKM}}$ .

There are many other models based on the  $SU(2)_1 \times SU(2)_2 \times U(1)$  gauge symmetry. In the “alternate left-right” model [8], all the couplings shown in Eq. (108.1) vanish, but there are some new fermions such that the  $W'$  boson couples to pairs involving a SM fermion and a new fermion. In the “unified SM” [9], the left-handed quarks are doublets under one  $SU(2)$ , and the left-handed leptons are doublets under a different  $SU(2)$ , leading to a mostly leptophobic  $W'$  boson:  $C_{\ell ij}^L \ll C_{q ij}^L$  and  $C_{\ell ij}^R = C_{q ij}^R = 0$ . Fermions of different generations may also transform as doublets under different  $SU(2)$  gauge groups [10]. In particular, the couplings to third generation quarks may be enhanced [11].

It is also possible that the  $W'$  couplings to SM fermions are highly suppressed. For example, if the quarks and leptons are singlets under one  $SU(2)$  [12], then the couplings are proportional to the tiny mixing angle  $\theta_+$ . Similar suppressions may arise if some vectorlike fermions mix with the SM fermions [13].

Gauge groups that embed the electroweak symmetry, such as  $SU(3)_W \times U(1)$  or  $SU(4)_W \times U(1)$ , also include one or more  $W'$  bosons [14].

## 108.2 Collider searches

At LEP-II,  $W'$  bosons could have been produced in pairs via their photon and  $Z$  couplings. The production cross section is large enough to rule out  $M_{W'} < \sqrt{s}/2 \approx 105$  GeV for most patterns of decay modes.

At hadron colliders,  $W'$  bosons can be detected through resonant pair production of fermions ( $f$  and  $f'$ ) or electroweak bosons with a net electric charge equal to  $\pm 1$ . When  $W'$  has a width much smaller than its mass ( $M_{W'}/\Gamma_{W'} \lesssim 7\%$ ), the contribution of the  $s$ -channel  $W'$  exchange to the total rate for  $pp \rightarrow f\bar{f}'X$ , where  $X$  is any final state, may be approximated by the branching fraction  $B(W' \rightarrow f\bar{f}')$  times the production cross section

$$\sigma(pp \rightarrow W'X) \simeq \frac{\pi}{6s} \sum_{i,j} \left[ (C_{qij}^L)^2 + (C_{qij}^R)^2 \right] w_{ij} \left( M_{W'}^2/s, M_{W'} \right). \quad (108.4)$$

The functions  $w_{ij}$  include the information about proton structure, and are given to leading order in  $\alpha_s$  by

$$w_{ij}(z, \mu) = \int_z^1 \frac{dx}{x} \left[ u_i(x, \mu) \bar{d}_j \left( \frac{z}{x}, \mu \right) + \bar{u}_i(x, \mu) d_j \left( \frac{z}{x}, \mu \right) \right], \quad (108.5)$$

where  $u_i(x, \mu)$  and  $d_i(x, \mu)$  are the parton distributions inside the proton at the factorization scale  $\mu$  and parton momentum fraction  $x$  for the up- and down-type quarks of the  $i$ th generation, respectively. QCD corrections to  $W'$  production are sizable (they also include quark-gluon initial states), but preserve the above factorization of couplings at next-to-leading order [15].

The most commonly studied  $W'$  signal consists of a high-momentum electron or muon and large missing transverse momentum. The signal transverse mass distribution forms a Jacobian peak with its endpoint at  $M_{W'}$  (see Fig. 1 (top) of Ref. [16]). Given that the branching fractions for  $W' \rightarrow e\nu$  and  $W' \rightarrow \mu\nu$  could be very different, the results in these channels should be presented separately. Searches in these channels often implicitly assume that the left-handed couplings vanish (no interference between  $W$  and  $W'$ ), and that the right-handed neutrino is light compared to the  $W'$  boson and escapes the detector. An example of parameter values that satisfy these assumptions is  $C_q^R = gV_{\text{CKM}}$ ,  $C_\ell^R = g$ ,  $C_q^L = C_\ell^L = 0$ , which define a model that preserves lepton universality and predicts the same total cross section as the Sequential SM used in many  $W'$  searches. However, if a  $W'$  boson were discovered and the final state fermions have left-handed helicity, then the effects of  $W - W'$  interference could be observed [17], providing information about the  $W'$  couplings. The effects of the  $W'$  width on interference are discussed in [18].

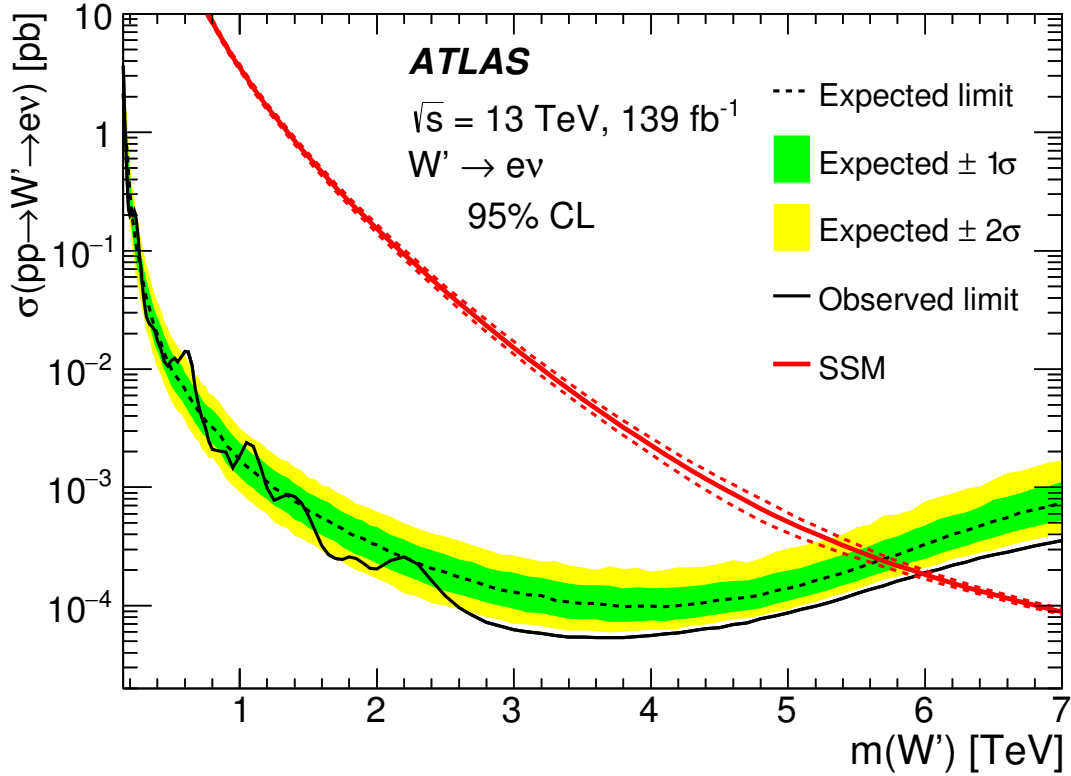


Figure 108.1: Upper limit on  $\sigma(pp \rightarrow W'X) B(W' \rightarrow e\nu)$  from ATLAS [16]. The red line shows the theoretical prediction in the Sequential SM.

In the  $e\nu$  channel, the ATLAS and CMS collaborations set limits on the  $W'$  production cross section times branching fraction (and thus indirectly on the  $W'$  couplings). These limits are set for  $M_{W'}$  in the 0.15 – 7 TeV range and are based on 36–139  $\text{fb}^{-1}$  at  $\sqrt{s} = 13$  TeV [16, 19], as shown in Fig. 108.1 for the most stringent limits. ATLAS sets the strongest mass limit  $M_{W'} > 6.0$  TeV in the Sequential SM (all limits in this mini-review are at the 95% CL). The coupling limits are much weaker for  $M_{W'} < 150$  GeV, a range last explored with the Tevatron at  $\sqrt{s} = 1.8$  TeV [20].

In the  $\mu\nu$  channel, ATLAS and CMS set rate limits for  $M_{W'}$  in the 0.15 – 7 TeV range from the same analyses as mentioned above, with the strongest mass lower limit of 5.1 TeV in the Sequential SM set by ATLAS [16] using 139  $\text{fb}^{-1}$  of  $\sqrt{s} = 13$  TeV data. When combined with the  $e\nu$  channel assuming lepton universality, the upper limit on the  $\sqrt{s} = 13$  TeV cross section times branching fraction to  $\ell\nu$  varies between 0.05 and 2.1 fb for  $M_{W'}$  values in the range between 1 and 6 TeV [16]. Only weak limits on  $W' \rightarrow \mu\nu$  exist for  $M_{W'} < 150$  GeV [21]. Note that masses of the order of the electroweak scale are interesting from a theory point of view, while lepton universality does not necessarily apply to a  $W'$  boson.

Dedicated searches for  $W' \rightarrow \tau\nu$  have been performed by CMS at 8 TeV [22] and both ATLAS and CMS with 36  $\text{fb}^{-1}$  at 13 TeV [23, 24]. Limits are set on  $\sigma \cdot B$  for  $M_{W'}$  between 0.4 and 4 TeV for the former and between 0.4 and 5.6 TeV for the latter. A mass lower limit of 4.0 TeV is set in the Sequential SM and the upper limit on the cross section times branching fraction to  $\tau\nu$  at 13 TeV varies between 1.7 and 12 fb for  $M_{W'}$  values in the range between 1 and 5 TeV [24].

The  $W'$  decay into a charged lepton and a right-handed neutrino,  $\nu_R$ , may also be followed by the  $\nu_R$  decay through a virtual  $W'$  boson into a charged lepton and two quark jets. The CMS [25] and ATLAS [26] searches in the  $eejj$  and  $\mu\mu jj$  channels have set limits on the cross section times

branching fraction as a function of the  $\nu_R$  mass or of  $M_{W'}$ . No requirement is placed on the charge of the lepton pair. A related  $W'$  search in the  $\tau\tau jj$  channel with hadronic  $\tau$  decays was also performed by CMS [27].

The  $t\bar{b}$  channel is particularly important because a  $W'$  boson that couples only to right-handed fermions cannot decay to leptons when the right-handed neutrinos are heavier than  $M_{W'}$ . Additional motivations are provided by a  $W'$  boson with enhanced couplings to the third generation [11], and by a leptophobic  $W'$  boson. The usual signature consists of a leptonically-decaying  $W$  boson and two  $b$ -jets. Recent studies have also incorporated the fully hadronic decay channel for  $M_{W'} \gg m_t$  with the use of jet substructure techniques to tag highly boosted top-jets. For a detailed discussion of this channel, see Ref. [28].

Searches for dijet resonances may be used to set limits on  $W' \rightarrow q\bar{q}'$ . ATLAS [29] and CMS [30] provide similar coverage in the  $\sim 1.5 - 8.0$  TeV mass range with 139 and 137  $\text{fb}^{-1}$  of data, respectively, collected at  $\sqrt{s} = 13$  TeV. Interpretation in terms of  $W'$  decays with 139  $\text{fb}^{-1}$  of 13 TeV data yields a  $W'$  mass lower limit of 4.0 TeV in the Sequential SM [29]. For masses in the range  $\sim 0.5 - 1.5$  TeV, analyses based on jets reconstructed online provide the best sensitivity because they circumvent trigger bandwidth limitations [31,32]. For  $W'$  masses below  $\sim 0.5$  TeV, the best limits are set in novel analyses exploiting boosted technologies and initial state radiation [33–36]. Cross-section limits for  $W'$  masses below  $\sim 1.5$  TeV can be derived from the dijet limits on  $Z'$  bosons summarized in Ref. [3].

In some theories [4] the  $W'$  couplings to SM fermions are suppressed by discrete symmetries.  $W'$  production then occurs in pairs, through a photon or  $Z$  boson. The decay modes are model-dependent and often involve other new particles. The ensuing collider signals arise from cascade decays and typically include missing transverse momentum.

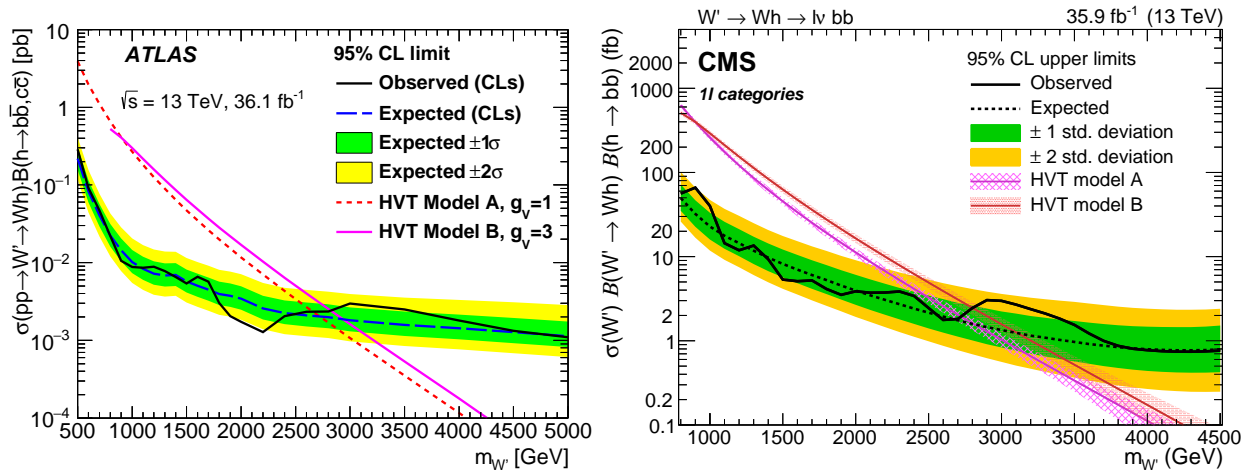


Figure 108.2: Upper limits on  $W'$  production cross section times branching fraction into a  $W$  and a SM Higgs boson decaying into heavy-flavor quarks, from ATLAS [37] (left) and CMS [38] (right).

Searches for  $WZ$  resonances at the LHC have focused on the process  $pp \rightarrow W' \rightarrow WZ$  with the production mainly from  $u\bar{d} \rightarrow W'$  assuming SM-like couplings to quarks. ATLAS and CMS have set the upper limits on the  $W'WZ$  coupling for  $M_{W'}$  in the 0.2 – 5.0 TeV range with a combination of fully leptonic, semi-leptonic and fully hadronic channels with  $\sim 36 \text{fb}^{-1}$  at 13 TeV [39,40] (see also Ref. [28]). The strongest lower limits on the  $W'$  mass are set by ATLAS [41] and CMS [42] at 13 TeV with 139  $\text{fb}^{-1}$  and 77  $\text{fb}^{-1}$ , respectively, in the  $WZ \rightarrow (jj)(jj)$  final state, where the parentheses represent a resonance. The lower limit on  $M_{W'}$  is 3.4 TeV in the context of the Heavy

Vector Triplet (HVT) weakly-coupled scenario A [43]. A fermiophobic  $W'$  boson that couples to  $WZ$  may be produced at hadron colliders in association with a  $Z$  boson, or via  $WZ$  fusion. This would give rise to  $(WZ)Z$  and  $(WZ)jj$  final states [44].

$W'$  bosons have also been searched for in final states with a  $W$  boson and a SM Higgs boson in the channels  $W \rightarrow \ell\nu$  or  $W \rightarrow q\bar{q}'$  and  $h^0 \rightarrow b\bar{b}$  by ATLAS [37, 45] and CMS [38, 46] with  $36 \text{ fb}^{-1}$  at  $\sqrt{s} = 13 \text{ TeV}$ . Cross-section limits are set for  $W'$  masses in the range between 0.5 and 5.0 TeV. The ATLAS and CMS 13 TeV analyses both set the most stringent lower limit on the mass:  $M_{W'} > 2.7 \text{ TeV}$  for the HVT weakly-coupled scenario A, as shown in Fig. 108.2.

### 108.3 Low-energy constraints

The properties of  $W'$  bosons are also constrained by measurements of processes at energies much below  $M_{W'}$ . The bounds on  $W - W'$  mixing [47] are mostly due to the change in  $W$  properties compared to the SM. Limits on deviations in the  $ZWW$  couplings provide a leading constraint for fermiophobic  $W'$  bosons [13].

Constraints arising from low-energy effects of  $W'$  exchange are strongly model-dependent. If the  $W'$  couplings to quarks are not suppressed, then box diagrams involving a  $W$  and a  $W'$  boson contribute to neutral meson-mixing. In the case of  $W'$  couplings to right-handed quarks as in the left-right symmetric model, the limit from  $K_L - K_S$  mixing is severe:  $M_{W'} > 2.9 \text{ TeV}$  for  $C_q^R = gV_{\text{CKM}}$  [48]. However, if no correlation between the  $W'$  and  $W$  couplings is assumed, then the limit on  $M_{W'}$  may be significantly relaxed [49].

$W'$  exchange also contributes at tree level to various low-energy processes. In particular, it would impact the measurement of the Fermi constant  $G_F$  in muon decay, which in turn would change the predictions of many other electroweak processes. A recent test of parity violation in polarized muon decay [50] has set limits of about 600 GeV on  $M_{W'}$ , assuming  $W'$  couplings to right-handed leptons as in left-right symmetric models and a light  $\nu_R$ . There are also  $W'$  contributions to the neutron electric dipole moment,  $\beta$  decays, and other processes [47].

If right-handed neutrinos have Majorana masses, then there are tree-level contributions to neutrinoless double-beta decay, and a limit on  $M_{W'}$  versus the  $\nu_R$  mass may be derived [51]. For  $\nu_R$  masses below a few GeV, the  $W'$  boson contributes to leptonic and semileptonic  $B$  meson decays, so that limits may be placed on various combinations of  $W'$  parameters [49]. For  $\nu_R$  masses below  $\sim 30 \text{ MeV}$ , the most stringent constraints on  $M_{W'}$  are due to the limits on  $\nu_R$  emission from supernovae.

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