SEARCHES not in other sections

Magnetic Monopole Searches

The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$<~1.4\times 10^{-16}~{\rm cm}^{-2}{\rm sr}^{-1}{\rm s}^{-1}~~{\rm for}~1.1\times 10^{-4}<\beta<1$$

Supersymmetric Particle Searches

All supersymmetric mass bounds here are model dependent.

The limits assume:

1) $\widetilde{\chi}_1^0$ is the lightest supersymmetric particle; 2) *R*-parity is conserved, unless stated otherwise;

See the Particle Listings for a Note giving details of supersymmetry.

$$\begin{array}{l} \widetilde{\chi}_i^0 - \text{neutralinos (mixtures of } \widetilde{\gamma}, \ \widetilde{Z}^0, \ \text{and } \widetilde{H}_i^0) \\ \text{Mass } m_{\widetilde{\chi}_1^0} > 0 \ \text{GeV, CL} = 95\% \\ \text{[general MSSM, non-universal gaugino masses]} \\ \text{Mass } m_{\widetilde{\chi}_1^0} > 46 \ \text{GeV, CL} = 95\% \\ \text{[all } \tan\beta, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \text{Mass } m_{\widetilde{\chi}_2^0} > 62.4 \ \text{GeV, CL} = 95\% \\ \text{[1<$} \tan\beta < 40, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \text{Mass } m_{\widetilde{\chi}_3^0} > 99.9 \ \text{GeV, CL} = 95\% \\ \text{[1<$} \tan\beta < 40, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \text{Mass } m_{\widetilde{\chi}_4^0} > 116 \ \text{GeV, CL} = 95\% \\ \text{[1<$} \tan\beta < 40, \ \text{all } m_0, \ \text{all } m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \widetilde{\chi}_i^{\pm} - \text{charginos (mixtures of } \widetilde{W}^{\pm} \ \text{and } \widetilde{H}_i^{\pm}) \\ \text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 94 \ \text{GeV, CL} = 95\% \\ \text{[$\tan\beta$ < 40, $m_{\widetilde{\chi}_1^{\pm}} - m_{\widetilde{\chi}_1^0} > 3 \ \text{GeV, all } m_0]} \\ \text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 810 \ \text{GeV, CL} = 95\% \\ \text{[$\ell^{\pm}\ell^{\mp}$, Tchi1chi1C, $m_{\widetilde{\chi}_1^0} = 0 \ \text{GeV}]} \\ \end{array}$$

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$$\begin{array}{ll} \widetilde{\chi}^{\pm} & - \text{long-lived chargino} \\ \text{Mass } m_{\widetilde{\chi}^{\pm}} > 620 \text{ GeV}, \text{ CL} = 95\% \quad [\text{stable } \widetilde{\chi}^{\pm}] \\ \widetilde{\nu} & - \text{sneutrino} \\ \text{Mass } m > 41 \text{ GeV}, \text{ CL} = 95\% \quad [\text{model independent}] \\ \text{Mass } m > 94 \text{ GeV}, \text{ CL} = 95\% \quad [\text{Rodel independent}] \\ \text{Mass } m > 94 \text{ GeV}, \text{ CL} = 95\% \quad [\text{R-Parity Violating}] \\ \widetilde{\nu}_{\tau} & - e\mu, \lambda_{312} = \lambda_{321} = 0.07, \lambda_{311}' = 0.11] \\ \widetilde{\epsilon} & - \text{ scalar electron (selectron)} \\ \text{Mass } m(\widetilde{\epsilon}_{L}) > 107 \text{ GeV}, \text{ CL} = 95\% \quad [\text{R-Parity Violating}] \\ \mathbb{E} & + \ell^{\pm}, \ell$$

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Mass m > 1340, CL = 95% $[\widetilde{t} R$ -hadrons] Mass m > 1250, CL = 95% $[\widetilde{b} R$ -hadrons]

$$\begin{split} \widetilde{b} & -- \text{ scalar bottom (sbottom)} \\ & \text{Mass } m > \ 1230 \text{ GeV, CL} = 95\% \\ & \text{ [jets+} \cancel{E}_T, \text{ Tsbot1, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 307 \text{ GeV, CL} = 95\% \quad \text{ [R-Parity Violating]} \\ & \widetilde{b} \rightarrow \ t \, d \text{ or } t \, s, \ \lambda_{332}'' \text{ or } \lambda_{331}'' \text{ coupling]} \\ \widetilde{t} & -- \text{ scalar top (stop)} \\ & \text{Mass } m > \ 1190 \text{ GeV, CL} = 95\% \\ & \text{ [jets+} \cancel{E}_T, \text{ Tstop1, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 1100 \text{ GeV, CL} = 95\% \quad \text{ [R-Parity Violating]} \\ & \widetilde{t} \rightarrow \ b \, e, \text{ Tstop2RPV, prompt]} \\ \widetilde{g} & -- \text{ gluino} \\ & \text{Mass } m > \ 2.000 \times 10^3 \text{ GeV, CL} = 95\% \\ & \text{ [jets+} \cancel{E}_T, \text{ Tglu1A, } m_{\widetilde{\chi}_1^0} = 0 \text{ GeV]} \\ & \text{Mass } m > \ 2.260 \times 10^3 \text{ GeV, CL} = 95\% \quad \text{ [R-Parity Violating]} \\ & [\ge 4\ell, \ \lambda_{12k} \ \neq \ 0, \ m_{\widetilde{\chi}_1^0} \ > \ 1000 \text{ GeV}] \\ \end{split}$$

Technicolor

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full *Review* (the data listings).

Quark and Lepton Compositeness, Searches for

Scale Limits Λ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \overline{\psi}_L \gamma_\mu \psi_L \overline{\psi}_L \gamma^\mu \psi_L$$

(with $g^2/4\pi$ set equal to 1), then we define $\Lambda \equiv \Lambda_{LL}^{\pm}$. For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

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$$\Lambda_{LL}^{+}(\textit{eeee})$$
 > 8.3 TeV, CL = 95% $\Lambda_{LL}^{-}(\textit{eeee})$ > 10.3 TeV, CL = 95%

$$\begin{array}{lll} \Lambda_{LL}^{+}(ee\mu\mu) &> 8.5 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{-}(ee\mu\mu) &> 9.5 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(ee\tau\tau) &> 7.9 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(ee\tau\tau) &> 7.2 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{-}(\ell\ell\ell\ell) &> 9.1 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\ell\ell\ell\ell) &> 9.1 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eeqq) &> 24 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eeqq) &> 23.3 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eeuu) &> 23.3 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eeuu) &> 12.5 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eedd) &> 11.1 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eedd) &> 26.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eecc) &> 9.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eebb) &> 9.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(eebb) &> 10.2 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\mu\mu qq) &> 32.7 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\mu\mu qq) &> 32.7 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\mu\mu qq) &> 32.7 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\mu qqq) &> 31.1 \ \mathrm{none} \ 17.4-29.5 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(qqqq) &> 13.1 \ \mathrm{none} \ 17.4-29.5 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(qqqq) &> 21.8 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\mu\nu qq) &> 5.0 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}(\nu\nu qq) &> 5.4 \ \mathrm{TeV}, \ \mathrm{CL} = 95\% \\ \Lambda_{LL}^{+}$$

Excited Leptons

The limits from $\ell^{*+}\ell^{*-}$ do not depend on λ (where λ is the $\ell\ell^{*}$ transition coupling). The λ -dependent limits assume chiral coupling.

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 $e^{*\pm}$ — excited electron

Mass
$$m>103.2$$
 GeV, CL $=95\%$ (from e^*e^*)
Mass $m>5.600\times10^3$ GeV, CL $=95\%$ (from $e\,e^*$)
Mass $m>356$ GeV, CL $=95\%$ (if $\lambda_{\gamma}=1$)

```
\mu^{*\pm} — excited muon
     Mass m > 103.2 \text{ GeV}, CL = 95\% (from \mu^* \mu^*)
    Mass m > 5.700 \times 10^3 GeV, CL = 95\% (from \mu \mu^*)
	au^{*\pm} — excited tau
    Mass m > 103.2 GeV, CL = 95\% (from \tau^* \tau^*)
    Mass m > 2.500 \times 10^3 GeV, CL = 95\% (from \tau \tau^*)
\nu^* — excited neutrino
    Mass m > 1.600 \times 10^3 GeV, CL = 95\% (from \nu^* \nu^*)
     Mass m > 213 GeV, CL = 95\% (from \nu^* X)
q^* — excited quark
    Mass m > 338 GeV, CL = 95\% (from q^* q^*)
    Mass m > 6700 \text{ GeV}, CL = 95\% (from g^*X)
Color Sextet and Octet Particles
Color Sextet Quarks (q_6)
     Mass m > 84 GeV, CL = 95\% (Stable q_6)
Color Octet Charged Leptons (\ell_8)
     Mass m > 86 GeV, CL = 95\% (Stable \ell_8)
Color Octet Neutrinos (\nu_8)
```

Extra Dimensions

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

Mass m > 110 GeV, CL = 90% $(\nu_8 \rightarrow \nu g)$

Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

(direct tests of Newton's law) $R < 4.8~\mu \text{m}$, CL = 95% $(p\,p \to j\,G)$ R < 0.16–916 nm (astrophysics; limits depend on technique and assumptions)

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Constraints on the fundamental gravity scale

$$M_{TT}>9.02$$
 TeV, CL $=95\%$ ($p\,p
ightarrow$ dijet, angular distribution) $M_c>4.16$ TeV, CL $=95\%$ ($p\,p
ightarrow$ $\ell \overline{\ell}$)

Constraints on the Kaluza-Klein graviton in warped extra dimensions

$$M_G$$
 $>$ 4.25 TeV, CL $=$ 95% $(pp \rightarrow \gamma \gamma)$

Constraints on the Kaluza-Klein gluon in warped extra dimensions

$$M_{g_{KK}}~>~3.8$$
 TeV, CL $=95\%~~(g_{KK}
ightarrow~t\,\overline{t})$

WIMP and Dark Matter Searches

No confirmed evidence found for galactic WIMPs from the GeV to the TeV mass scales and down to 1×10^{-10} pb spin independent cross section at M = 100 GeV.

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