

Review of particle properties

Particle Data Group

Thomas G. Trippe, Angela Barbaro-Galtieri, Robert L. Kelly, Alan Rittenberg, Arthur H. Rosenfeld, and George P. Yost

Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720, USA*

Naomi Barash-Schmidt

Brandeis University, Waltham, Massachusetts 02154, USA

Claude Bricman, Richard J. Hemingway, and Michael J. Losty

CERN, CH-1211 Genève 23, Switzerland

Matts Roos

Department of Nuclear Physics, University of Helsinki, Helsinki 17, Finland

Vladimir Chaloupka

Stanford Linear Accelerator Center, Stanford, California 94305, USA

Betty Armstrong (Technical Associate)

Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720, USA*

This review of the properties of leptons, mesons, and baryons is an updating of Review of Particle Properties, Particle Data Group [Phys. Letters **50B**, No.1 (1974), and Supplement, Rev. Mod. Phys. 47 (1975) 535]. Data are evaluated, listed, averaged, and summarized in tables. Numerous tables, figures, and formulae of interest to particle physicists are also included. A data booklet is available.

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I. INTRODUCTION, CREDITS, CONSULTANTS

This review is an updating through December 1975 of our previous review (Particle Data Group, 1974, and Supplement, 1975). In this version, we have attempted to make the text as complete and self-contained as possible.

As usual, the results of our compilation are presented in two sections. The Tables of Particle Properties, usually referred to as simply the Tables, contain a summary of the properties of only those peaks or resonances which in our judgment have a large probability of standing the test of time. This is a conservative judgment, and surely means that some genuine resonances are (temporarily) omitted. See Sec. V below.

The Data Card Listings, on the other hand, are an attempt to give up-to-date information, along with references, on all the reported particles whose existence is either considered confirmed (in which case they are also included in the Tables) or not yet confirmed. The Listings also contain mini-reviews pertaining to questions of interest.

A short survey of the history of some of the constants we compile, as well as some of those compiled by others, has been added to the text as a new feature (Sec. VIII). In general, the reliability of the data has been pretty good; that is, the percentage of the time that our best estimates of various constants have changed by more than one standard deviation over the years is roughly what would be expected from statistical considerations. A history of the Particle Data Group, with a discussion of procedures and problems, has been given by Rosenfeld (1975).

This year we are experimenting with a new statistical procedure in our data averaging. Given in the Tables are our best estimates (and their errors), calculated in the old way. In the Listings, we give simultaneously the old (labeled "AVG") and new (labeled "STUDENT") values. In most cases there is little difference. Details may be found in Sec. VII. User comments are solicited.

A pocket-sized data booklet, containing the Tables and a reprint of the figures and formulae from the first part of the book, is available on request.

For North and South America, Australia, and the Far East, write to Technical Information Division, Lawrence Berkeley Laboratory, Berkeley, California 94720, USA.

For all other areas, write to CERN Scientific Information Service, CH-1211 Genève 23, Switzerland.

As usual, we wish to emphasize that we compile the experimental results of others. It is inappropriate to give us the credit for their countless hours of effort. We urge that references be given directly to the original data, and we provide complete references in the Data Card Listings for that purpose.

The responsibilities for the various sections can be broken down as follows:

(1) *Stable particles*: A. Barbaro-Galtieri, N. Barash-Schmidt, and T. G. Trippe.

(2) *Meson resonances*: V. Chaloupka, R. J. Hemingway, M. J. Losty, and M. Roos.

(3) *Baryon resonances*: A. Barbaro-Galtieri, C. Bricman, and R. L. Kelly.

(4) *General, including text*: All authors.

Consultants: To overcome unavoidable gaps in our coverage, both intellectual and geographical, we have solicited the help of consultants:

- Ugo Amaldi (CERN)
- Stanley J. Brodsky (Stanford Linear Accelerator Center)
- Denyse M. Chew (Lawrence Berkeley Laboratory)
- Ronald Crawford (University of Glasgow)
- J. Engler (CERN)
- Anatoli Kuznetsov (JINR, Dubna)
- Gerald R. Lynch (Lawrence Berkeley Laboratory)
- F. Mönning (CERN)
- R. Gordon Moorhouse (University of Glasgow)
- David R. Nygren (Lawrence Berkeley Laboratory)
- Oliver E. Overseth (University of Michigan)
- Sherwood I. Parker (Lawrence Berkeley Laboratory)
- Bernard Sadoulet (Lawrence Berkeley Laboratory)
- Paul Söding (DESY)
- Fumiyo Uchiyama (Lawrence Berkeley Laboratory)

The usefulness of this compilation depends in large part on the interaction between the users and the authors and consultants. We appreciate comments, criticisms, and suggestions for improvements of all stages of data retrieval, processing, and presentation.

II. SELECTION OF DATA

All particles are considered to fall into one of the three groups:

- (1) Stable particles, immune to decay via the strong interaction, including the η .
- (2) Meson resonances (including all the new high mass particles, whether "stable" or not).
- (3) Baryon resonances.

These groups are maintained within the two main parts of the compilation:

- (1) Tables of Particle Properties.
- (2) Data Card Listings.

The Data Card Listings contain the original information (data, references etc.), weighted averages, comments, and "mini-reviews". Immediately preceding the Data Card Listings is an Illustrative Key thereto. We attempt to give complete Data Card Listings up to our closing date (January 1, 1976) for all journals listed in the Illustrative Key. We also include preprints and unpublished conference reports which have come to our attention, but make no attempt at completeness.

Roughly 40% of our encoded results are not used for averaging. They are set off in parentheses; our reasoning is then often given in a footnote below the data. If the reason is not given, it is one of the following:

- The quantity was presented with no error stated.
- The result comes from a preprint or conference report. It is our experience that such results (and particularly the errors) often change before final publication. Accordingly we keep these new results in parentheses until they are published (or explicitly verified to us by the authors).
- It involves some assumptions that we do not wish to incorporate.

- It is of poor quality, e.g., bad signal-to-noise-ratio.
- The result is inconsistent with others, e.g., because of different methods employed, rendering averaging meaningless.
- It is not independent of other results, e.g., a result from one of several partial-wave analyses all using the same data, again rendering averaging meaningless.

When the data for a particle have received special treatment or when they present special problems, this is noted in a mini-review in the Data Card Listings.

The Tables of Particle Properties represent the output of weighted averages and some critical judgment. The extent to which "blind" averaging has been tempered with judgment is explained in footnotes to the Tables. In general, however, the footnotes are less complete than is the collection of notes and mini-reviews in the Data Card Listings. The reader is therefore encouraged to familiarize himself with the Data Card Listings and, ultimately, with the original experiments.

III. NOMENCLATURE

A. Quantum numbers

The symbols $I^G(J^P)C$ represent:

I = isospin

G = G -parity

J = spin

P = space parity

C = charge conjugation parity .

1. Mesons

The charge conjugation operator C turns particle into antiparticle and has eigenvalues ± 1 only for neutral states; so it is useful to define an extension G which has eigenvalues for charged states too. It is usually¹ defined by

$$G = C \exp(i\pi I_y) . \quad (1)$$

A neutral nonstrange state is an eigenstate of $\exp(i\pi I_y)$ with eigenvalue $(-1)^I$. Then we can write the eigenvalue equation for the whole multiplet as

$$G = C_n (-1)^I , \quad (2)$$

where C_n (n for neutral) is the eigenvalue C would have if applied to the neutral member of the multiplet. Thus, for a π^0 , C has the eigenvalue $+1$, and since $I=1$, $G = -1$. For the charged pion there are no eigenvalues corresponding to C and to the isospin rotation, but Eqs. (1) and (2) still give $G = -1$.

Consider a meson as a bound state of fermion-antifermion, e.g., $\bar{q}q$, with orbital angular momentum l , and with the two fermion spins coupling to give a spin S . Then one can show that the charge-conjugation eigenvalue [defined in Eq. (2)] is

$$C_n = (-1)^{l+S} . \quad (3)$$

Equations (2) and (3) combine to give

$$G = (-1)^{l+S+I} . \quad (4)$$

The parity is

$$P = -(-1)^l . \quad (5)$$

Equations (3) and (5) combine to give

$$C_n P = -(-1)^S \quad (6)$$

so all singlet ($^1S_0, ^1P_1, \dots$) have $C_n P = -1$, and all triplet ($^3S_1, \dots$) have $C_n P = +1$. For proofs of the above, see our 1969 text (Particle Data Group, 1969) and Appendix by C. Zemach.

If, instead of $\bar{q}q$, we consider the meson as a state of *boson-antiboson* (e.g., $A_2 - \bar{K}K$), it turns out that some signs cancel, and Eqs. (3) and (4) [not (5)!] apply *unchanged*. Of course the mesons are often spinless so S is zero, but the equations are more general. Eqs. (3) and (4) can be considered as selection rules forbidding many decays.

We now use Eqs. (3) and (4) to introduce the concept of "Abnormal- C " mesons, i.e., mesons that cannot be composed of $\bar{q}q$.

The unitary triplet of quarks is of course defined to have isospin and hypercharge properties such that $\bar{q}q$ can combine (according to the SU(3) relations $\{3\} \otimes \{3\} = \{8\} \oplus \{1\}$) so as to form only unitary octets and singlets. The non-observation of "exotic" mesons (i.e., mesons in more complicated supermultiplets) is of course one of the bases of the naive quark model. But it is slightly less obvious that even some *octets* are forbidden by the model, namely those with $(J^P)C_n = (0^*)-, (1^-)+, (2^*)-, \dots$. Such states are also not observed, and this is an additional success of the naive quark model classification scheme.

In what follows, do not confuse "Abnormal- C " with "Normal" or "Abnormal" J^P , both of which are allowed by the quark model. The series $J^P = 0^*, 1^-, 2^*, \dots$ is called Normal because $P = (-1)^J$ as for normal spherical harmonics, and $J^P = 0^-, 1^+, \dots$ is called Abnormal.

The top part of Table I shows all the low angular momentum states that can be formed from $\bar{q}q$. Note that half of the J^P states can be formed by both a triplet and a singlet $\bar{q}q$ state, e.g., $^3P_1, ^1P_1$ or $^3D_2, ^1D_2$. Eq. (3) shows that 3P_1 and 1P_1 have opposite C_n , so the $\bar{q}q$ model allows both. But the states 3P_0 and 3P_2 have no 1P counterparts. According to Eq. (6) they have $C_n P = +1$, and with the $\bar{q}q$ model there is no way to form a state with a J^P of $^3P_{0,2}$ (i.e., $J^P = \text{Normal}$) and with $C_n P = -1$. As mentioned, such octets have not shown up. With the help of Table I one can also see that the special state $^1S_0, C_n P = +1$, cannot be formed, so has Abnormal C .

2. General remarks

Well-established quantum numbers are underlined in the Tables of Particle Properties (except for stable particles, where most of the quantum numbers are established). We have used flimsy evidence to guess many of the remaining ones, and we have indicated with "?" ones (in the baryon table) for which there is almost no evi-

¹Most texts define it as in Eq. (1); see, e.g., Gasiorowicz (1966); however, sometimes the rotation is taken about I_x . The difference between the two conventions is mentioned in a footnote in Källén (1964).

TABLE I. $I^G(J^P)$ of non-strange mesons from $\bar{q}q$ model. For the distinction between abnormal J^P and abnormal C , see text following Eq. (6). K mesons share the same values of J^P as the $I=0$ and 1 states shown, but are not eigenstates of G . The middle column, which gathers together $(J^P)_{\text{Normal or } A} CP$, is a redundant intermediate step intended to make the table easier to read.

Parity	$\bar{q}q$ State		$(J^P)_{\text{Normal or abnormal}} \uparrow CP$	$I^G(J^P)C_n$	Examples and comments
	CP	CP			
	-	+			
Parity -	$1S_0$		$(0^-)_{A^-}$	$\begin{cases} 0^+(0^-)+ \\ 1^-(0^-)+ \end{cases}$	$\eta, \eta', E ?$ π
	$3S_1$		$(1^-)_{N^+}$	$\begin{cases} 0^-(1^-)- \\ 1^+(1^-)- \end{cases}$	ω, ϕ ρ
Parity +	$1P_1$		$(1^+)_{A^-}$	$\begin{cases} 0^-(1^+)- \\ 1^+(1^+)- \end{cases}$	B
	$3P_0$		$(0^+)_{N^+}$	$\begin{cases} 0^+(0^+)+ \\ 1^-(0^+)+ \end{cases}$	ϵ, S^* $\delta ?$
	$3P_1$		$(1^+)_{A^+}$	$\begin{cases} 0^+(1^+)+ \\ 1^-(1^+)+ \end{cases}$	D ? A1
	$3P_2$		$(2^+)_{N^+}$	$\begin{cases} 0^+(2^+)+ \\ 1^-(2^+)+ \end{cases}$	f, f' A2
Parity -	$1D_2$		$(2^-)_{A^-}$	$\begin{cases} 0^+(2^-)+ \\ 1^-(2^-)+ \end{cases}$	A3
	$3D_1$		$(1^-)_{N^+}$	same as $3S_1$	ρ'
	$3D_2$		$(2^-)_{A^+}$	$\begin{cases} 0^-(2^-)- \\ 1^+(2^-)- \end{cases}$	Regge recurrence of the abnormal-C state $(J^P)_{C_n} = (0^-)-$
	$3D_3$		$(3^-)_{N^+}$	$\begin{cases} 0^-(3^-)- \\ 1^+(3^-)- \end{cases}$	g
Parity +	$1F_3$		$(3^+)_{A^-}$	$\{ J > 2 \}$	
	$3F_2$		$(2^+)_{N^+}$	same as $3P_2$	
	$3F_3$		$(3^+)_{A^+}$	$\{ J > 2 \}$	
	$3F_4$		$(4^+)_{N^+}$	etc.	

ABNORMAL C STATES THAT CANNOT COME FROM $\bar{q}q$ MODEL

Abnormal C states Have no $\bar{q}q$ model	$(0^-)_{A^+}$	$\begin{cases} 0^-(0^-)- \\ 1^+(0^-)- \end{cases}$	All except $J^P = 0^-$ are $J^P = \text{normal}, CP = -1$
	$(1^-)_{N^-}$	$\begin{cases} 0^+(1^-)+ \\ 1^-(1^-)+ \end{cases}$	
	$(0^+)_{N^-}$	$\begin{cases} 0^-(0^+)- \\ 1^+(0^+)- \end{cases}$	
	$(2^+)_{N^-}$	$\begin{cases} 0^-(2^+)- \\ 1^+(2^+)- \end{cases}$	
	$(3^-)_{N^-}$	$\begin{cases} 0^+(3^-)+ \\ 1^-(3^-)+ \end{cases}$	

dence.

As is customary, we define antiparticles as the result of operating with CPT on particles, so both share the same spins, masses, and mean lives. Whenever there is a particularly interesting test of CPT invariance we include it in the Stable Particles Table.

B. Particle names

If a *meson* has a well-accepted colloquial name, we use it. If not, we name it by a single symbol which specifies its baryon number B ($=0$ for mesons), its isospin I , its hypercharge Y , and, for a nonstrange meson its G parity. For convenience, we also list the strangeness S , which is related to Y and B by

$$S = Y - B.$$

The name conventions for mesons are given in the first part of Table II.

To crowd even more information onto the symbol, we sometimes add a subscript giving J^P . If J^P is not known, but must be "Normal" ($0^+, 1^-, 2^+, \dots$), e.g., because $K\pi$ decays are seen, we use the subscript N . If such modes are *not* seen (and are not otherwise forbidden), we *guess* that it is because J is "Abnormal", and we use the subscript A .

For *some* pairs of *mesons* with supposedly identical quantum numbers, we also use primes; e.g., η, η' ; f, f' .

For *baryons* no attempt has been made to attach a subscript about J and P . The name conventions are given in the second part of Table II. For stable baryons of each I and Y we use the symbol standing alone; for resonances, the mass is in parentheses [i.e., $N(1688)$, $\Lambda(1405)$, $\Sigma(1765)$, etc.]. The J^P assignments are reported in the Baryon Table as $\frac{1}{2}^+, \frac{3}{2}^-, \frac{5}{2}^+$, etc., and also by the symbols P_{11}, D_{13}, F_{15} , which refer to the πp or Kp partial-wave amplitude in which the resonant state occurs (the first subscript refers to the isospin state:

TABLE II. Particle name conventions.

Name	I	Y	S	G
Mesons				
η	0	0	0	+
ω or ϕ^a	0	0	0	-
ρ	1	0	0	+
π	1	0	0	-
K^+, K^0	$\frac{1}{2}$	+1	+1	
K^-, \bar{K}^0	$\frac{1}{2}$	-1	-1	
Baryons				
N	$\frac{1}{2}$	+1	0	
Δ	$\frac{3}{2}$	+1	0	
Z_0, Z_1	0, 1	+2	+1	
Λ	0	0	-1	
Σ	1	0	-1	
Ξ	$\frac{1}{2}$	-1	-2	
Ω	0	-2	-3	

^a Since 1973, we have used the symbol ω for those $I^G=0^-$ mesons that decay mainly into 3π [$\omega(783)$, $\omega(1670)$]; we reserve the symbol ϕ for $\phi(1020)$ and possible future higher-mass $I^G=0^-$ mesons that decay mainly into $K\bar{K}$.

$2 \times I$ for N and Δ and just I for Z , Λ , and Σ). When two or more baryons have identical quantum numbers we warn the reader by adding primes to the spectroscopic symbol as explained in footnote (a) of the Baryon Table.

IV. CONVENTIONS AND PARAMETERS FOR STRONG INTERACTIONS

A. Partial-wave amplitudes and resonance parameters

The vast majority of information concerning baryon resonances comes in the form of partial-wave analyses. In addition data concerning meson resonances ($\pi\pi, K\pi, \pi\pi\pi$) are, with increasing frequency, being subjected to partial-wave analyses. We thus find it natural to introduce the resonance parameters which we compile in terms of a Breit-Wigner approximation for the partial-wave amplitude.

In general the elastic amplitude for a given angular momentum l may be written as

$$T_{11} = \frac{\eta \exp(2i\delta) - 1}{2i}, \quad (1)$$

where η is referred to as the absorption parameter ($0 \leq \eta \leq 1$) and δ , as the phase shift. The subscripts 11 on T denote scattering from channel 1 to channel 1 (i.e., $\pi\pi \rightarrow \pi\pi$ or $\bar{K}N \rightarrow \bar{K}N$).

In Fig. 1 we show an Argand plot of the elastic partial wave amplitude T_{11} . It illustrates geometrically how the real parameters η and δ are related to the real and imaginary parts of T_{11} . Many examples of such Argand plots may be found in the Baryon Data Card Listings.

Consider the so-called non-relativistic Breit-Wigner approximation for T_{11} :

$$T_{11} = \frac{1}{2} \Gamma_1 / (M - E - \frac{1}{2} i\Gamma) \quad (2)$$

where E is the c.m. energy of invariant mass, Γ_1 and Γ are the *elastic* and *total* widths, and M is called the *resonance mass*. Equation (2) is, of course, not the only possible description of a resonant amplitude; it suffices to illustrate the properties of partial-wave amplitudes which we associate with resonance behavior in the absence of any background in the same partial wave (see, e.g., the $\pi N D_{15}$ and F_{15} waves in the Baryon Data Card Listings). Usually the widths contain barrier-

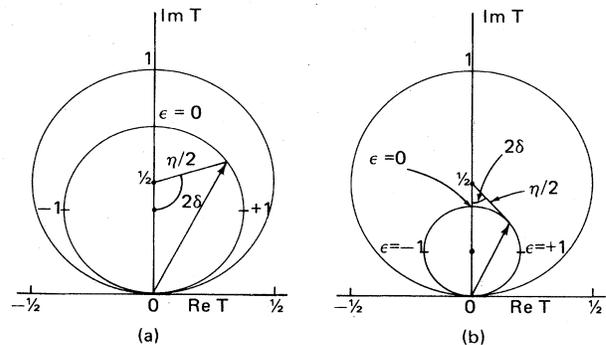


FIG. 1. Argand plots for the partial wave amplitude T_{11} . The outer circles are the unitarity bound ($\eta=1$). The inner circles correspond to the Breit-Wigner approximation of Eq. (2) for (a) $x_1 = \Gamma_1/\Gamma = 0.75$ and (b) $x_1 = 0.4$. Note: $\epsilon = 2(M - E)/\Gamma$.

penetration factors which can vary rapidly with energy. Near threshold, $\Gamma_1(E)$ should start up as q^{2l+1} (also true for the inelastic width Γ_β). Various E dependences are then used for Γ_1 , mostly of the form

$$\Gamma_1(E) \propto \frac{(qR)^{2l+1}}{\text{const.} + \dots + (qR)^{2l}}, \quad (3)$$

see Jackson (1964), Pišút and Roos (1968), and Barbaro-Galtieri (1968).

The BW-approximation to the amplitude for an inelastic process leading from channel 1 to channel β ($\pi\pi \rightarrow \bar{K}K$ or $KN \rightarrow \Sigma\pi$, for example), is

$$T_{1\beta} = \frac{1}{2}(\Gamma_1\Gamma_\beta)^{1/2}/(M-E - \frac{1}{2}i\Gamma) \\ = (x_1x_\beta)^{1/2}[\frac{1}{2}\Gamma/(M-E - \frac{1}{2}i\Gamma)], \quad (4)$$

where

$$\Gamma = \sum_1^N \Gamma_\beta, \quad x_\beta = \Gamma_\beta/\Gamma, \quad (5)$$

and x_1 (called the elasticity) is often written x_e . (Note that in the Data Card Listings we use the symbol P_β to denote x_β .) The channel cross-section $\sigma_{1\beta}$ for the reaction $1 \rightarrow \beta$, for spin $0 \rightarrow$ spin $\frac{1}{2}$ scattering, is

$$\sigma_{1\beta} = 4\pi\lambda^2(J + \frac{1}{2}) |T_{1\beta}|^2, \quad (6)$$

where $J = l \pm \frac{1}{2}$.

The important features of Eq. (4) which characterize resonant behavior in the Argand diagram ($\text{Im}T_{1\beta}$ versus $\text{Re}T_{1\beta}$) are:

(1) Energy variation given by circles with diameter $(x_1x_\beta)^{1/2}$ and maximum amplitude at $E=M$ of

$$T_{1\beta}^{\text{max}} = i(x_1x_\beta)^{1/2}. \quad (7)$$

(2) A maximum in the speed near resonance, given approximately by

$$\text{“Speed” (res)} = |dT_{1\beta}/dE|_{E=M} = \frac{2(x_1x_\beta)^{1/2}}{\Gamma(E)}, \quad (8)$$

for slowly varying $\Gamma(E)$. These features may be related to the η, δ representation of T_{11} . Thus when $E=M$, δ is either $90^\circ (x_1 > \frac{1}{2})$ or $0^\circ (x_1 < \frac{1}{2})$ and η dips to its minimum

value.

These simple properties can be used to judge the presence or absence of resonance behavior in an Argand plot. However, it must be kept in mind that Eqs. (2) and (4) are only approximations to the “true” amplitude. The simple picture given above can be distorted by various effects:

- the presence of “background” in the same partial wave as the resonance
- two resonances in the same partial wave overlapping in energy
- the resonant energy M being close to an inelastic channel threshold, in which case a K -matrix-like parametrization is more appropriate.

B. SU(3) sign conventions for Λ and Σ resonances

Consider the partial width Γ_β of a resonance decaying into the channel β . We can always define a coupling constant such that

$$\Gamma_\beta \propto G_\beta^2.$$

In this case the inelastic amplitude in the Breit-Wigner approximation, Eq. (4) will go as

$$T_{1\beta} \propto G_1G_\beta/(M-E - \frac{1}{2}i\Gamma),$$

where G_1 is the coupling constant for the elastic channel. In the context of exact SU(3) symmetry the relative signs of the product G_1G_β for different resonances are often useful as a consistency check on SU(3) assignment of Λ and Σ resonances. See Appendix II for further details.

In the Data Card Listings for Λ and Σ resonances, we tabulate measured values for $(x_1x_\beta)^{1/2} \propto G_1G_\beta$. Whenever there is an explicit sign, it will be according to the convention advocated by Levi-Setti (1969) and used in the table of SU(3) Isoscalar Factors presented in this review. Thus the signs multiplying the Breit-Wigner amplitudes for $\bar{K}N \rightarrow \Sigma(1385) \rightarrow \Sigma\pi$, $\Lambda\pi$ and $\bar{K}N \rightarrow \Lambda(1405) \rightarrow \Sigma\pi$ are simply the product of the phases of the appropriate isoscalar factors. This convention is shown in Fig. 2 from Levi-Setti (1969).

SU(3) RELATIVE SIGN OF RESONANT AMPLITUDES

$$T_{\text{RES}} \sim \alpha (G_{N\bar{K}Y} \cdot G_{Y\pi Y^*}) / (M-E - i\Gamma/2)$$

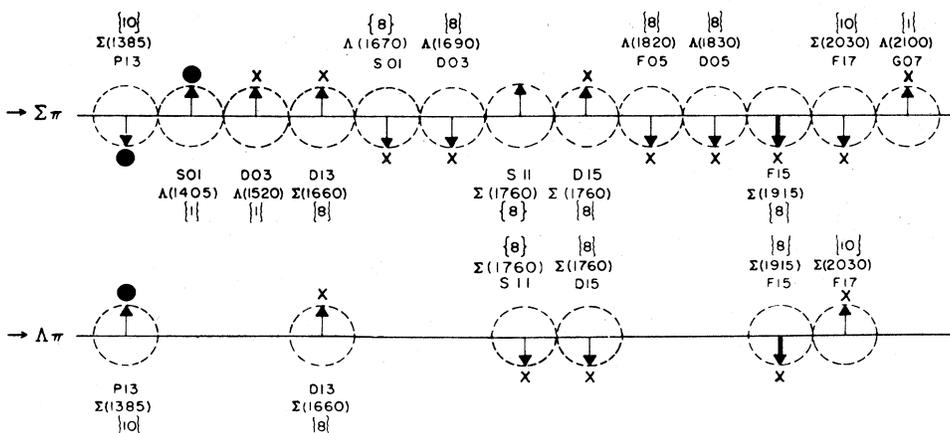


FIG. 2. Plot adapted from Levi-Setti (1969) showing the sign convention adopted here for the $\Sigma\pi$ and $\Lambda\pi$ amplitudes. Once the signs of one $I=0$ and one $I=1$ amplitude are fixed, the others can be measured relative to these two. Arrows here indicate signs predicted by SU(3); \times marks indicate the observed phases; \bullet indicates phase chosen according to sign convention described in text. The $\Sigma(1915)$ predictions have been changed from Levi-Setti's original figure.

C. Types of partial-wave analyses

Partial-wave analyses (PWA) are classified into three categories in the Data Card Listings: energy-independent partial-wave analyses (IPWA), energy-dependent partial-wave analyses (DPWA), and model-dependent partial-wave analyses (MPWA), in increasing order of the number of explicit supplementary hypotheses that are used to extract the amplitudes from experimental data.

In an IPWA, data at different energies are analyzed separately. Usually each partial wave included in the fit is allowed to vary freely (subject to unitarity constraints) over some large region, and waves whose angular momenta are above some cutoff value are assumed to be negligible. The sharp cutoff in angular momentum resolves continuum ambiguities in the solution (such as the overall phase ambiguity), but there remains a finite number of indistinguishable "best" solutions (i.e., solutions corresponding to identical physical observables) which have been codified by Barrelet (1972). In addition, there are generally some nearby solutions (and their associated Barrelet ambiguities) which have chi-squared values close to the minimum one.

At the end of the analysis a choice is made among these many solutions, usually on the basis of energy continuity. A popular criterion for making this choice is the shortest path technique in which the total "length" of the preferred solution is chosen to be a minimum. The definition of "length" used here is not universal but is usually closely related to the total geometrical length of the lines representing the various partial-wave amplitudes in Argand plots (see the Baryon section of the Data Card Listings for examples of Argand plots). Various other criteria which are also used in some analyses are, e.g., matching with known solutions at low energies, the presence of known resonances in the final results, and limited inelasticity in high partial waves.

In a DPWA, data at different energies are fit simultaneously by using an energy dependent parametrization of the partial-wave amplitudes. The parametrization is usually chosen to include both resonances and nonresonant background of some sort and an attempt is made to keep it as "model independent" as possible. Often the data are grouped into several energy bins which are fit separately rather than trying to fit the whole energy range under consideration simultaneously. One of the main advantages of DPWA over IPWA is that sparse data spread over many different energies can be analyzed, e.g., nearly all $S=-1$ analyses are DPWA. In addition, the built-in energy continuity helps to resolve the ambiguities that plague IPWA and eases the problems associated with resonance parameter extraction. The price one pays for these advantages lies in the danger of systematic error in the amplitudes and poor fits to the data if the parametrization is poorly chosen or insufficiently flexible.

An MPWA also uses an energy-dependent parametrization, but one based on explicit model-dependent theoretical assumptions such as Regge exchanges. This technique is usually applied to reactions where the data are incomplete. There is, of course, no sharp distinction between DPWA and MPWA, and a well chosen MPWA

parametrization may actually be less biased than a model-independent but poorly chosen DPWA parametrization.

D. Production of resonances

Hereby, we mean the observation of statistically significant peaks in invariant mass plots or, loosely, in integrated cross sections. Many meson resonances are of this type. We expect most of these peaks to be associated with Breit-Wigner behavior in appropriate Argand plots; thus the ρ meson peak in $\pi\pi$ mass plots is firmly related to the $I=1, l=1$ $\pi\pi$ phase shift passing through 90° .

From mass plots we can determine M , Γ , and the approximate branching ratios

$$x_\alpha/x_\beta = \Gamma_\alpha/\Gamma_\beta . \quad (9)$$

In the case of total cross sections, the peak above background gives us, using the optical theorem, the product $(J+\frac{1}{2})x_e$:

$$\sigma^{\text{tot}}(E=M) = 4\pi\chi^2(J+\frac{1}{2})x_e . \quad (10)$$

V. CRITERIA FOR RESONANCES

An experimentalist who sees indications of a resonance in some energy (or mass) region will of course want to know what has been seen in that region in the past; hence, we strive to have the Data Card Listings serve as an archive for all substantial claims or evidences for resonances.

For the Tables of Particle Properties, on the other hand, we wish to be more conservative and to include only those peaks or resonances which we feel have a large chance of survival. An arrow (→) at the left of the Tables of Particle Properties indicates that a questionable candidate has been omitted from the Table, but that it can be found in the corresponding part of the Data Card Listings. One's betting odds for survival are of course subjective; therefore no precise criteria can be defined. In what follows we shall attempt to specify some guide lines.

(a) When energy-independent partial-wave analyses are available (mostly for N^* 's), approximate Breit-Wigner behavior of the amplitude appears to us to be the most satisfactory test for a resonance. We can check that the Argand plot follows roughly a left-hand circle, and that the "speed" of the amplitude also shows a maximum near the resonance energy; further, there should be data well above the resonance, showing that the speed again decreases. Indeed proper behavior of the partial-wave amplitude could accredit a resonance even if its elasticity is too small to make a noticeable peak in the cross section.

Of course even if Argand plots are available, it may still be a matter of opinion as to what behavior constitutes a resonance. Such an example is the $Z_0(1780)$ state seen in KN total cross-section experiments and in partial-wave analysis. The recent partial-wave analyses of Giacomelli (1974) and Martin (1975) find preferred solutions which exhibit a resonance-like loop in the P_{01} wave near 1740 MeV (see Fig. 5 of the $S=+1$

mini-review in the Baryon Data Card Listings). However, Giacomelli *et al.* and Martin point out that, despite the resonant-like *appearance* of the loop, the evidence for resonant *energy dependence* is inconclusive. Thus we omit the $Z_0(1780)$ from the Baryon Table. A similar quandary has existed for some time concerning the $Z_1(1900)$, and it too has been omitted from the Tables.

(b) When there are insufficient data to perform energy-independent analyses, one often resorts to energy-dependent partial-wave analyses (mostly for Y^* 's). In this case Breit-Wigner behavior is an input. We therefore require that resonance solutions be found by several different analyses, preferably in different channels ($\bar{K}N \rightarrow \bar{K}N, \pi\Sigma$, etc.), before putting the claim in the table.

(c) Partial-wave analyses of three-body final states ($\pi N \rightarrow \pi\pi N$) are becoming available. While these analyses are based on the isobar model ($\pi N \rightarrow \rho N, \pi\Delta$, etc.) and are subject to theoretical objections of varying importance, they provide increasingly reliable information on inelastic decay modes of otherwise established resonances.

(d) Most mesons, Ξ^* peaks, and high mass N^* and Y^* peaks fall into a category for which no partial-wave analyses exist. In general we accept such peaks if they are experimentally reliable, of high statistical significance or if they are observed in several different production processes.

(e) A special category of "diffractive mesons" consists of statistically significant peaks like A_1, A_3 or Q , which are not far above the $\rho\pi, f\pi$, or $K^*\pi$ thresholds. The question of a resonance interpretation for these states is complicated, because the behavior near threshold in these channels may be described by the Deck effect. Modern partial-wave analyses can shed considerable light on these problems. See the mini-reviews for details.

Thus, we enter into the tables of Particle Properties only states for which there is experimentally convincing evidence, and we expect that most of these will be confirmed as resonances.

VI. CONVENTIONS AND PARAMETERS FOR WEAK AND ELECTROMAGNETIC DECAYS

A. Muon-decay parameters

The μ -decay parameters describe the momentum spectrum (ρ and η), the asymmetry (ξ and δ), and the helicity (h) of the electron in the process $\mu^+ \rightarrow e^+ + \nu + \bar{\nu}$. Assuming a local and lepton-conserving interaction, the matrix element may be written as

$$\sum_i \langle \bar{e} | \Gamma_i | \mu \rangle \langle \bar{\nu} | \Gamma_i (C_i + C'_i \gamma_5) | \nu \rangle,$$

where the summation is taken over $i=S, V, T, A, P$. Using the definitions and sign conventions of Kinoshita and Sirlin (1957), we have for the momentum parameters

$$\rho = [3g_A^2 + 3g_V^2 + 6g_T^2] / D,$$

$$\eta = [g_S^2 - g_P^2 + 2g_A^2 - 2g_V^2] / D,$$

for the asymmetry parameters:

$$\xi = \frac{6g_S g_P \cos\phi_{SP} - 8g_A g_V \cos\phi_{AV} + 14g_T^2 \cos\phi_{TT}}{D},$$

$$\delta = [-6g_A g_V \cos\phi_{AV} + 6g_T^2 \cos\phi_{TT}] / D\xi,$$

and for the parameter describing the helicity of the electron:

$$h = \frac{2g_S g_P \cos\phi_{SP} - 8g_A g_V \cos\phi_{AV} - 6g_T^2 \cos\phi_{TT}}{D}.$$

Here

$$D = g_S^2 + g_P^2 + 4g_V^2 = 6g_T^2 + 4g_A^2,$$

$$g_i^2 = |C_i|^2 + |C'_i|^2,$$

and

$$\cos\phi_{ij} = \text{Re}(C_i^* C'_j + C'_i C_j^*).$$

The quantities g_i are defined to be real non-negative numbers, and the ϕ_{ij} are phase angles between the i -type and j -type interactions. Under the assumption of two-component neutrinos $C'_i = -C_i$ and $C'_j = -C_j$, the S, P , and T terms vanish, and ϕ_{AV} is the phase angle between C_A and C_V in the complex plane.

By using the above equations and the experimental determinations of ρ, η, ξ, δ , and h , limits can be placed on $g_S/g_V, g_A/g_V, g_T/g_V, g_P/g_V$, and ϕ_{AV} . The results, given in the Data Card Listings assume neither two-component neutrinos nor time-reversal invariance. If, however, two-component neutrinos are assumed, then $\sin\phi_{AV}$ is the amplitude of time-reversal violation. Note that most experiments study only the upper end of the spectrum where ρ and η are highly correlated, so they can only report ρ for $\eta=0$ and η for $\rho=\frac{3}{4}$. The values for ρ and η we use here were obtained by combining measurements of both upper and lower ends of the spectrum and turn out to be nearly uncorrelated.

Note also that the radiative corrections are unambiguous only when $g_S = g_T = g_P = 0$. The same limits on g_A/g_V and ϕ_{AV} are obtained, however, as when g_S, g_T , and g_P are left free.

Current values for the asymmetry parameters as well as $|g_A/g_V|$ and ϕ_{AV} are given in the Addendum to the Stable Particle Table. In addition, upper limits on $|g_S/g_V|, |g_T/g_V|$ and $|g_P/g_V|$ are given in the μ section of the Stable Particle Data Card Listings.

B. K -decay parameters

1. Dalitz plot for $K \rightarrow 3\pi$ decays

The small deviation from uniformity of the Dalitz plot for the 3π decay of the K meson is usually described by a "slope parameter" (Dalitz, 1956). For the τ and τ' decays of the charged K 's, and the τ^0 decay mode of the K_L^0 , we parametrize the Dalitz plot distribution by the expression

$$|M|^2 \propto 1 + g \frac{S_3 - S_0}{m_{\pi^+}^2} + h \left(\frac{S_3 - S_0}{m_{\pi^+}^2} \right)^2 + j \frac{S_2 - S_1}{m_{\pi^+}^2} + k \left(\frac{S_2 - S_1}{m_{\pi^+}^2} \right)^2 + \dots, \quad (1)$$

where $m_{\pi^+}^2$ has been introduced so as to make the coeffi-

cients g , h , j , and k dimensionless, and

$$s_i = (P_K - P_i)^2 = (m_K - m_i)^2 - 2m_K T_i, \quad i = 1, 2, 3$$

$$s_0 = \frac{1}{3} \sum_i s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2).$$

Here the P_i are 4-vectors, m_i and T_i are the mass and kinetic energy of the i th pion, and the index 3 is used for the odd pion.

The coefficient g is a measure of the slope in the variable s_3 (or T_3) of the Dalitz plot, while h and k measure the quadratic dependence on s_3 and $(s_2 - s_1)$, respectively. The coefficient j is related to the asymmetry of the plot and must be zero if CP invariance holds. Note also that if CP is good, g must be the same for τ^+ and τ^- , and similarly for h and k .

At present there is no compelling experimental evidence for the h , k , or j term (for upper limits on the j term, see Sec. B.3(b) below). Thus we stop the above expansion at the first term and list only g . Since different experiments use different forms for $|M|^2$, in order to compare the experiments we have converted to g whatever coefficients have been measured. See the mini-review in the K^* section of the Stable Particle Data Card Listings for details on this point. The results are given in the Addendum to the Stable Particle Table and in the K^* and K_L^0 sections of the Stable Particle Data Card Listings.

Relations among τ^\pm , τ'^\pm , and τ^0 are predicted by the $\Delta I = \frac{1}{2}$ rule. See Appendix I for these relations and a discussion of this rule.

2. Form factors in K_{l3} leptonic decays

Assuming that only the vector current contributes to these decays, we write the matrix element as

$$M \propto f_+(t) [(P_K + P_\pi)_\mu \bar{u}_l \gamma_\mu (1 + \gamma_5) u_\nu] + f_-(t) [m_l \bar{u}_l (1 + \gamma_5) \mu_\nu], \quad (2)$$

where P_K and P_π are the four momenta of K and π mesons; m_l is the lepton mass; f_+ and f_- are dimensionless form factors which can depend only on $t = (P_K - P_\pi)^2$, the square of the four-momentum transfer to the leptons. f_+ and f_- are relatively real if time reversal invariance holds for these decays. $K_{\mu 3}$ experiments measure f_+ and f_- , while K_{e3} experiments are sensitive only to f_+ because the presence of the lepton mass makes the f_- term negligible.

(a) $K_{\mu 3}$ experiments

Analyses of $K_{\mu 3}$ data frequently assume a linear dependence of f_+ and f_- on t , i.e.

$$f_\pm(t) = f_\pm(0) [1 + \lambda_\pm (t/m_\pi^2)]. \quad (3)$$

Most $K_{\mu 3}$ data are adequately described by Eq. (3) for f_+ and a constant f_- (i.e. $\lambda_- = 0$). There are two equivalent parametrizations commonly used in these analyses:

(1) $\lambda, \xi(0)$ Parametrization. Analyses of $K_{\mu 3}$ data often introduce the ratio of the two form factors

$$\xi(t) = f_-(t)/f_+(t).$$

The $K_{\mu 3}$ decay distribution is then described by the two parameters λ_+ and $\xi(0)$ (assuming time reversal invariance and $\lambda_- = 0$). These parameters can be determined by three different methods:

Method A: By studying the Dalitz plot or the pion spectrum of $K_{\mu 3}$ decay. The Dalitz plot density is (See, e.g. Chounet *et al.*, 1972):

$$\rho(E_\pi, E_\mu) \propto f_+^2(t) [A + B\xi(t) + C\xi(t)^2],$$

where

$$A = m_K(2E_\mu E_\nu - m_K E'_\pi) + m_\mu^2 (\frac{1}{4} E'_\pi - E_\nu),$$

$$B = m_\mu^2 (E_\nu - \frac{1}{2} E'_\pi),$$

$$C = \frac{1}{4} m_\mu^2 E'_\pi,$$

$$E'_\pi = E_\pi^{\max} - E_\pi = \frac{m_K^2 + m_\pi^2 - m_\mu^2}{2m_K} - E_\pi.$$

Here E_π , E_μ , and E_ν are respectively the pion, muon, and neutrino energies in the kaon center of mass. The density ρ is fit to the data to determine the values of λ_+ , $\xi(0)$, and their correlation.

Method B: By measuring the $K_{\mu 3}/K_{e3}$ branching ratio and comparing it with the theoretical ratio (See, e.g., Fearing *et al.*, 1970) as given in terms of λ^+ and $\xi(0)$, assuming μ - e universality.

$$\Gamma(K_{\mu 3}^\pm)/\Gamma(K_{e3}^\pm) = 0.6457 + 1.4115\lambda_+ + 0.1264\xi(0) + 0.0192\xi(0)^2 + 0.0080\lambda_+\xi(0),$$

$$\Gamma(K_{\mu 3}^0)/\Gamma(K_{e3}^0) = 0.6452 + 1.3162\lambda_+ + 0.1246\xi(0) + 0.0186\xi(0)^2 + 0.0064\lambda_+\xi(0).$$

This cannot determine λ_+ and $\xi(0)$ simultaneously but simply fixes a relationship between them.

Method C: By measuring the muon polarization in $K_{\mu 3}$ decay. In the rest frame of the K , the μ is expected to be polarized in the direction \vec{A} with $\vec{P} = \vec{A}/|\vec{A}|$, where \vec{A} is given (Cabibbo and Maksymowicz, 1964) by

$$\vec{A} = a_1(\xi) \vec{p}_\mu - a_2(\xi) \left\{ \frac{\vec{p}_\mu}{m_\mu} \left[m_K - E_\pi + \frac{\vec{p}_\pi \cdot \vec{p}_\mu}{|\vec{p}_\mu|^2} (E_\mu - m_\mu) \right] + \vec{p}_\pi \right\} + m_K \text{Im} \xi(t) (\vec{p}_\pi \times \vec{p}_\mu).$$

If time-reversal invariance holds, ξ is real, and thus there is no polarization perpendicular to the K -decay plane. Polarization experiments measure the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$.

(2) λ, λ_0 Parametrization. Some of the more recent $K_{\mu 3}$ analyses have parametrized in terms of the form factors f_+ and f_0 which are associated with vector and scalar exchange respectively to the lepton pair. f_0 is related to f_+ and f_- by

$$f_0(t) = f_+(t) + [t/(m_K^2 - m_\pi^2)] f_-(t).$$

Here $f_0(0)$ must equal $f_+(0)$ unless $f_-(t)$ diverges at $t=0$. The earlier assumption that f_+ is linear in t and f_- is constant leads to f_0 linear in t

$$f_0(t) = f_0(0) [1 + \lambda_0 (t/m_\pi^2)].$$

With the assumption that $f_0(0) = f_+(0)$, the two parametrizations, $(\lambda_+, \xi(0))$ and (λ_+, λ_0) are equivalent as long as correlation information is retained. (λ_+, λ_0) correlations tend to be less strong than $(\lambda_+, \xi(0))$ correlations.

The experimental results for $\xi(0)$ and its correlation with λ_+ are listed in the K^+ and K_L^0 sections of the Stable Particle Data Card Listings in subsection XIA, XIB, or XIC depending on whether Method A, B or C discussed above was used. The corresponding values of λ^+ are listed in subsection $L+M$.

Because current experiments tend to use the (λ_+, λ_0) parametrization, we have added a subsection $L0$ for λ_0 results. Wherever possible we have converted $\xi(0)$ results into λ_0 results and vice versa.

(b) K_{e3} experiments

Analysis of K_{e3} data is simpler than that of $K_{\mu 3}$ because the second term of the matrix element assuming a pure vector current (Eq. 2 above) can be neglected. Here f_+ is usually assumed to be linear in t , and the linear coefficient λ_+ of Eq. (3) is determined.

If we remove the assumption of a pure vector current, then the matrix element for the decay, in addition to the terms in Eq. (2), would contain

$$+ 2m_K(f_S \bar{u}_1(1 + \gamma_5)u_\nu + (2f_T/m_K)(P_K)_\lambda (P_\pi)_\mu \bar{u}_1 \sigma_{\lambda\mu}(1 + \gamma_5)u_\nu),$$

where f_S is the scalar form factor, and f_T is the tensor form factor. In the case of the K_{e3} decays where the f_- term can be neglected, experiments have yielded limits on $|f_S/f_+|$ and $|f_T/f_+|$.

The K_{e3} results for λ_+ , $|f_S/f_+|$ and $|f_T/f_+|$ are listed in the subsections $L+M$, FS , and FT respectively of the K^+ and K_L^0 sections of the Stable Particle Data Card Listings.

See also the *Note on K_{13}^+ and K_{13}^0 Form Factors* in the K^+ section of the Stable Particle Data Card Listings for additional discussion of the $K_{\mu 3}$ parameters, correlations, and conversion between parametrization and also for a comparison of the experimental results.

3. CP violation in K^0 decays

We list parameters for four different reactions in which CP can be tested [for details, see Okun and Rubbia (1967), Steinberger (1969), and Wolfenstein (1969)].

(a) $K_S \rightarrow \pi^+ \pi^- \pi^0$. The quantity measured here is the ratio of amplitudes

$$A_S(K_S \rightarrow \pi^+ \pi^- \pi^0) / A_L(K_L \rightarrow \pi^+ \pi^- \pi^0) = x + iy. \quad (4)$$

If CPT invariance holds and there is no $I=3$ state present, then x can be neglected and CP violation would be observed as a nonzero y . We give the result for Eq. (4) in the K_L^0 section of the Stable Particle Table and under Branching Ratio $R4$ in the K_S^0 section of the Stable Particle Data Card Listings. Our procedure is to assume that $x=0$, and to list $(A_S/A_L)^2$ in the form of a branching ratio.

(b) *Charge asymmetry in $K_L \rightarrow 3\pi$ decays.* As mentioned above, the presence of a term in $(s_2 - s_1)$ in expression (1) describing the Dalitz plot distribution for τ^\pm , τ^0 decays of K mesons would be an indication of CP violation. Rather than listing values of the $(s_2 - s_1)$

coefficient j in Eq. (1), we choose to list σ_\pm from the equivalent expression

$$|M|^2 \propto 1 + \sigma_\pm (2/\sqrt{3})(T_+ - T_-)/T_{\pm \max} + (CP \text{ nonviolating terms}), \quad (5)$$

where T_\pm are the kinetic energies of the charged pions. We have momentarily abandoned the form involving the Mandelstam variables s_i in favor of Eq. (5) because the latter has been consistently used by experimenters searching for CP violation. We list σ_\pm among the CP -violating parameters at the back of the K_L^0 section of the Stable Particle Data Card Listings. Note that only upper limits have been reported for this quantity.

(c) *Asymmetry in the $K_L - \pi^\mp l^\pm \nu$ decays.* The quantity measured and compiled here is

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^- l^+ \nu) - \Gamma(K_L \rightarrow \pi^+ l^- \nu)}{\Gamma(K_L \rightarrow \pi^- l^+ \nu) + \Gamma(K_L \rightarrow \pi^+ l^- \nu)}.$$

This asymmetry violates CP invariance. If CPT is good, for a pure K_L^0 beam, δ can be written as

$$\delta = 2[(1 - |x|^2)/(1 + |x|^2)] \text{Re} \epsilon,$$

where x is the $\Delta S = \Delta Q$ -violating parameters defined in section B4, and ϵ is the parameter of the expansion

$$|K_L\rangle = [(1 + \epsilon)|K\rangle - (1 - \epsilon)|\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}, \quad (6a)$$

$$|K_S\rangle = [(1 + \epsilon)|K\rangle + (1 - \epsilon)|\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}. \quad (6b)$$

We give δ in the Addendum to the Stable Particle Table. In addition, in the K_L^0 CP -violation section of the Stable Particle Data Card Listings, we list δ separately for $K_L^0 \rightarrow \pi \mu \nu$ and $K_L^0 \rightarrow \pi e \nu$.

(d) $K_L \rightarrow 2\pi$ decay. The relevant parameters are

$$\eta_{+-} = A(K_L \rightarrow \pi^+ \pi^-) / A(K_S \rightarrow \pi^+ \pi^-) = |\eta_{+-}| \exp(i\phi_{+-}),$$

$$\eta_{00} = A(K_L \rightarrow \pi^0 \pi^0) / A(K_S \rightarrow \pi^0 \pi^0) = |\eta_{00}| \exp(i\phi_{00}),$$

ϵ , defined in Eqs. (6) above, and

$$\epsilon' = \frac{1}{2} i \sqrt{2} [\exp(i\delta_2 - \delta_0)] \text{Im}(A_2/A_0).$$

Here, A_i and δ_i are the amplitude and phase of $\pi\pi$ scattering at the K mass, defined by

$$\langle I=0 | T | K \rangle = \exp(i\delta_0) A_0,$$

$$\langle I=2 | T | K \rangle = \exp(i\delta_2) A_2.$$

Wu and Yang (1964) have derived the relationships

$$\eta_{+-} = \epsilon + \epsilon',$$

$$\eta_{00} = \epsilon - 2\epsilon'.$$

We give η_{+-} , η_{00} , ϕ_{+-} , and ϕ_{00} in the Addendum to the Stable Particle Table. The phases are measured directly, whereas the magnitudes η_{+-} and η_{00} are derived parameters. We use, as far as we can, the directly measured quantities as input and calculate η_{+-} and η_{00} from the values given by our constrained fits. Therefore, if one looks at the Data Card Listings, most of the $|\eta|$ measurements appear in the form of branching ratios, with appropriate comments. We then give the values of η_{+-} and $|\eta_{00}|^2$ in a separate list at the end of the CP -violating parameters section of the K_L^0 section of the Stable Particle Data Card Listings.

4. $\Delta S = \Delta Q$ rule in K^0 decays

The relative amount of $\Delta S \neq \Delta Q$ component present is measured by the parameter x , defined as

$$x = A(\bar{K}^0 \rightarrow \pi^- l^+ \nu) / A(K^0 \rightarrow \pi^- l^+ \nu).$$

We list $\text{Re}\{x\}$ and $\text{Im}\{x\}$ for both K_{e3} and $K_{\mu 3}$ at the end of the Stable Particle Data Card Listings and give values in the Addendum to the Stable Particle Table.

C. η -decay parameters

1. C -violation in η -decays

As a test of possible C -violation in electromagnetic interactions, a number of experiments have looked for possible charge asymmetries in the decays $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$. We list the following parameters:

(a) The left-right asymmetry

$$A = (N^+ - N^-) / (N^+ + N^-),$$

where $N^{(\pm)}$ means the number of events with the $\pi^{(\pm)}$ energy greater than the $\pi^{(\mp)}$ energy in the η rest frame.

(b) The sextant asymmetry

$$A_s = \frac{N_1 + N_3 + N_5 - N_2 - N_4 - N_6}{N_1 + N_2 + N_3 + N_4 + N_5 + N_6}$$

for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$. The numbers refer to the sextants of the Dalitz plot [see, for example, Layter (1972)]. A_s is sensitive to an $I=0$ C -violating asymmetry.

(c) The quadrant asymmetry A_q , defined in a similar way as A_s , but with each sector of the Dalitz plot now containing $\pi/2$ rather than $\pi/3$ radians. A_q is sensitive to an $I=2$ C -violating final state.

(d) The d -wave contribution to the C -violating amplitude in the decay $\eta \rightarrow \pi^+ \pi^- \gamma$. The upper limit for this contribution is measured by the parameter β , defined by

$$dN/d|\cos\theta| \propto \sin^2\theta(1 + \beta \cos^2\theta),$$

where θ is the angle between the π^+ and the γ in the dipion center of mass. A term proportional to $\cos^2\theta$ could also be due to p - and f -wave interference.

We list A for the decay modes $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$, A_s and A_q for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$, and β for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ in the η section of the Stable Particle Data Card Listings.

2. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \pi^0$

The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$ may be fit by the distribution

$$|M(x, y)|^2 \propto 1 + ay + by^2 + cx + dx^2 + exy.$$

Here,

$$x = \sqrt{3} (T_+ - T_-) / Q,$$

$$y = (3T_0 / Q) - 1,$$

T_+ , T_- , T_0 are the kinetic energies of the π^+ , π^- , and π^0 in the η rest system, and $Q = m_\eta - m_{\pi^+} - m_{\pi^-} - m_{\pi^0}$. The coefficient of the term linear in x is sensitive to C -violation due to an $I=0$ or $I=2$ final state. We list papers presenting determinations of the parameters a , b , c , and d in the η section of the Stable Particle

Data Card Listings. However, we do not tabulate values of these parameters because the assumptions made by different authors are not compatible and do not allow comparison of the numerical values.

3. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \gamma$

The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ may be fit to the expression

$$|M|^2 \propto 1 + 2\alpha z,$$

where

$$z = \frac{2}{3} \sum_{i=1}^3 \left[\frac{3}{m_\eta - 3m_\pi} \left(E_i - \frac{1}{3} m_\eta \right) \right]^2 = \frac{\rho^2}{\rho_{\max}^2}.$$

Here E_i is the energy of the i th pion in the η rest frame, and ρ is the distance to the center of the Dalitz plot. We list the parameter α in the η section of the Stable Particle Data Card Listings.

D. Baryon-decay parameters

1. A/V ratio for baryon leptonic decays

Consider the decay

$$B_i \rightarrow B_f + l + \nu.$$

Assuming V, A theory, neglecting "induced" scalar, "induced" pseudoscalar, and axial weak-magnetism terms, and neglecting the q^2 dependence of the form factors, the baryon part of the matrix element for these decays may be written (Goldberger and Treiman, 1958) as

$$\langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) + (g_W / m_{B_i}) \sigma^{\lambda\nu} q_\nu | B_i \rangle,$$

where B_i and B_f represent initial and final baryons, g_A and g_V the axial and vector coupling constant, g_W the weak magnetism coupling constant, and q_ν the sum of the lepton momenta. Here the Pauli representation is used for the γ matrices. The definition of g_A/g_V is

$$g_A/g_V = |g_A/g_V| \exp(i\delta),$$

where δ is $0 + n\pi$ if time-reversal invariance holds (see Jackson *et al.*, 1957).

In neutron beta decay the measurements are consistent with time reversal, so g_A/g_V is nearly real and has been considered to be such in all the baryon leptonic decays. Notice that by using the above definition of the matrix element with the Pauli representations, the value of g_A/g_V in neutron beta decay is negative.

Due to statistical limitation the weak magnetism form factor g_W is usually assumed from CVC and SU(3), so only g_A and g_V are determined experimentally. This determination is accomplished in a variety of ways:

(a) The lepton-neutrino angular correlation provides a measure of the absolute value of g_A/g_V (for relevant formulas, see, e.g., Albright, 1959).

(b) The up-down asymmetry of the lepton from polarized baryon decays provides a measure of g_A/g_V with its sign (for relevant formulas, see, e.g. Albright, 1959).

(c) The lepton spectrum, given enough statistics, provides a measure of g_A/g_V with its sign (for relevant formulas, see, e.g., Bender, 1968).

(d) The polarization of the decay baryon, from polar-

ized or unpolarized initial baryon, also provides g_A/g_V with its sign (for formulas, see, e.g., Willis and Thompson, 1968).

We compile the ratio g_A/g_V with its sign, for those decays for which it has been measured. For the neutron beta decay we compile also the phase δ .

All the coupling constants and decay rates for baryon leptonic decays are related by Cabbibo's theory (Cabbibo, 1964). A recent fit to this theory has been done by Roos (1974).

2. Asymmetry parameters in nonleptonic hyperon decays

The transition matrix for the hyperon decay may be written as

$$M = s + p(\vec{\sigma} \cdot \vec{q}), \quad (7)$$

where s and p are the parity-changing and the parity conserving amplitudes, respectively; $\vec{\sigma}$ is the Pauli spin operator, and \vec{q} is a unit vector along the direction of the decay baryon in the hyperon rest frame.

The asymmetry parameters are defined by the relations

$$\begin{aligned} \alpha &= 2 \operatorname{Re}(s^*p) / (|s|^2 + |p|^2), \\ \beta &= 2 \operatorname{Im}(s^*p) / (|s|^2 + |p|^2), \\ \gamma &= (|s|^2 - |p|^2) / (|s|^2 + |p|^2). \end{aligned}$$

With the transition matrix (7), the angular distribution of the decay baryon, in the hyperon rest system, is of the form

$$I = 1 + \alpha \vec{P}_Y \cdot \vec{q},$$

where $\vec{P}_Y = \langle Y | \sigma | Y \rangle$ is the hyperon polarization.

In the notation of Lee and Yang (1957) the polarization \vec{P}_B of the decay baryon is²

$$\vec{P}_B = \frac{(\alpha + \vec{P}_Y \cdot \vec{q})\vec{q} + \beta(\vec{P}_Y \times \vec{q}) + \gamma\vec{q} \times (\vec{P}_Y \times \vec{q})}{1 + \alpha \vec{P}_Y \cdot \vec{q}}$$

where \vec{P}_B is defined in that rest system of the baryon obtained by a Lorentz transformation along \vec{q} from the hyperon rest system in which \vec{q} and \vec{P}_Y are defined. Note that α is the helicity of the decay baryon for unpolarized hyperons.

The three parameters α , β , and γ satisfy the relation

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

It is then convenient to describe hyperon nonleptonic decays in terms of the two independent parameters α and the angle ϕ defined by

$$\begin{aligned} \beta &= (1 - \alpha^2)^{1/2} \sin \phi, \\ \gamma &= (1 - \alpha^2)^{1/2} \cos \phi, \end{aligned}$$

which has a more nearly Gaussian distribution than β or γ . Evidently

²Note that Lee and Yang (1957) contains a misprint. The minus sign in the definition of β should be replaced by a 2. In addition, our unit vector \vec{q} is the direction of the baryon, whereas their unit vector \vec{p} is the direction of the pion.

$$-\frac{1}{2}\pi \leq \phi \leq \frac{1}{2}\pi \quad \text{for } \gamma > 0.$$

$$+\frac{1}{2}\pi \leq \phi \leq \frac{3}{2}\pi \quad \text{for } \gamma < 0.$$

In discussing time-reversal invariance, the quantity of interest is Δ , defined by

$$\begin{aligned} \alpha &= 2 |s| |p| \cos \Delta / (|s|^2 + |p|^2), \\ \beta &= -2 |s| |p| \sin \Delta / (|s|^2 + |p|^2); \end{aligned}$$

that is, Δ is the phase angle of s relative to p . Evidently

$$-\frac{1}{2}\pi \leq \Delta \leq \frac{1}{2}\pi \quad \text{for } \alpha > 0,$$

$$+\frac{1}{2}\pi \leq \Delta \leq \frac{3}{2}\pi \quad \text{for } \alpha < 0.$$

Under the assumption of time-reversal invariance, the angle Δ must satisfy the relation

$$\Delta = \delta_s - \delta_p,$$

modulo π , where δ_s and δ_p are the pion-baryon scattering phase shifts at the appropriate energy and for the appropriate isospin state. For Λ decay, assuming the validity of the $|\Delta I| = \frac{1}{2}$ rule,

$$\Delta = \delta_s - \delta_p = (6.8 \pm 2.0) \text{ deg.}^3$$

In the Stable Particle Data Card Listings we give α and ϕ for each decay since they are the most closely related to the experiments and are essentially uncorrelated. Whenever necessary we have changed the signs of the reported values, so as to agree with our conventions. In the Stable Particle Table we give α , ϕ , and Δ with errors; and for convenience we also give the central value of γ , without an error.

VII. STATISTICAL PROCEDURES

We divide this discussion on obtaining averages and errors into two sections:

- A. the unconstrained case, or "simple averaging", and
- B. the constrained case.

In what follows, the term "error" means one standard deviation (1σ); that is, for central value \bar{x} and error $\delta\bar{x}$, the range $\bar{x} \pm \delta\bar{x}$ constitutes a 68.3% confidence interval.

A. Unconstrained averaging

We first describe the standard procedure which we have used for several years to determine averages and errors. We will then discuss a second method, newly proposed, which we feel offers a less conservative, and possibly more accurate, estimate of errors.

1. Standard procedure—Gaussian distribution with scale factor

We begin by assuming that measurements of a given quantity obey a Gaussian distribution, and thus we calculate a weighted average and error

³This value for $\delta_s - \delta_p$ is derived from the phase-shift analyses by Roper *et al.* (1965). The error is our estimation of the uncertainty.

$$\bar{x} \pm \delta\bar{x} = \left(\frac{\sum_i w_i x_i}{\sum_i w_i} \right) \pm \left[1 / \left(\sum_i w_i \right) \right]^{-1/2}, \quad (1)$$

$$w_i = [1/(\delta x_i)^2],$$

where x_i and δx_i are the value and error, respectively, reported by the i th experiment, and the sums run over N experiments. We also calculate χ^2 and compare it with its expectation value of $N-1$.

If $\chi^2/(N-1)$ is less than or equal to 1, and there are no known problems with the data, we accept the above results.

If $\chi^2/(N-1)$ is ridiculously large, or if there is prior knowledge of extremely large inconsistencies between experiments, we may choose not to average the data at all. Or, in some cases, we may quote the calculated average, but then give an "educated guess" as to the error; such a guess is a generally quite conservative estimate designed to take into account known problems with the data.

Finally, if $\chi^2/(N-1)$ is greater than 1, but not to such a large extent, we still average the data, but then try to make up for this fact in two ways:

(i) We plot an ideogram to guide the reader in deciding which data he might reject before making his own selected average. An example of such an ideogram is given in Fig. 3 below. Each experiment appearing in the plot is represented by a Gaussian with central value x_i , error δx_i , and area proportional to $1/\delta x_i$. The choice of area is a somewhat arbitrary one; it is based on the assumption that an experimenter will work to reduce his systematic errors until they are slightly smaller (but seldom much smaller) than his statistical

errors. Thus as a bubble-chamber physicist gets more events, he will use them both to reduce his statistical errors and to study his biases. Our confidence that a significant systematic error has not been made in his experiment, as compared with other contradictory experiments, then tends to go up as $1/\delta x_i$.

But why not assign a weight $1/\delta x_i^2$, as is done when computing a weighted average? We feel that this is equivalent to assuming that large systematic errors are as infrequent as large statistical fluctuations, and that this is unrealistic.

We want to emphasize the difference between least-squares averaging (where the weighting factor is the inverse square of the error) and the ideograms prepared for visual display. The former arithmetic is of course best if one has statistically distributed input, and yields a narrow Gaussian distribution centered at the weighted mean. The ideogram (often multi-peaked and certainly not Gaussian) is based on the opposite hypothesis that some of the input is systematically in error. The idea behind least-squares averaging is that experiments 1, 2, 3, etc., are *all* valid (so we should multiply their probabilities); our *ideograms* are based on the assumption that 1 *or* 2 *or* 3, etc., is valid, "hedged" with $1/\delta x_i$ betting odds; we then add their probabilities. Both approaches cannot simultaneously be right; we leave it to the reader to choose. A glance at the ideogram will show, however, that the discrepancy is often not severe for reasonably distributed input.

(ii) The second way in which we try to take account of $\chi^2/(N-1)$ being greater than 1 is to scale up our quoted error $\delta\bar{x}$ in Eq. (1) by a factor

$$\text{SCALE} = [\chi^2/(N-1)]^{1/2}. \quad (2)$$

Our reasoning is as follows. Since we don't know which one or more of the experiments are wrong, we assume that all experimentalists underestimated their errors by the same scale factor (2). If we scaled up all input errors by this factor, χ^2 returns to $N-1$, and of course the output error scales up by the same factor.

If all the experiments have errors of about the same size, the above (straightforward) procedure for calculating SCALE is carried out. If, however, we are to combine experiments with widely varying errors, we must modify the procedure slightly. This is because it is the more precise experiments that most influence not only the average value \bar{x} , but also the error $\delta\bar{x}$. Now, on the average, the low-precision experiments each contribute about unity to *both* the numerator and the denominator of SCALE, hence the χ^2 contribution of the sensitive experiments is diluted, i.e., reduced. Therefore, we evaluate SCALE by using *only* experiments for which the error are not much greater than those of the more precise experiments. Explicitly, to calculate SCALE we use only the most sensitive experiments, i.e., those with errors less than δ_0 , where the ceiling δ_0 is (arbitrarily) chosen to be

$$\delta_0 = 3N^{1/2}\delta\bar{x}.$$

Here $\delta\bar{x}$ is the unscaled error of the mean of all the experiments. Note that if each experiment had the same error δx_i , then $\delta\bar{x}$ would be $\delta x_i/N^{1/2}$, so each individual

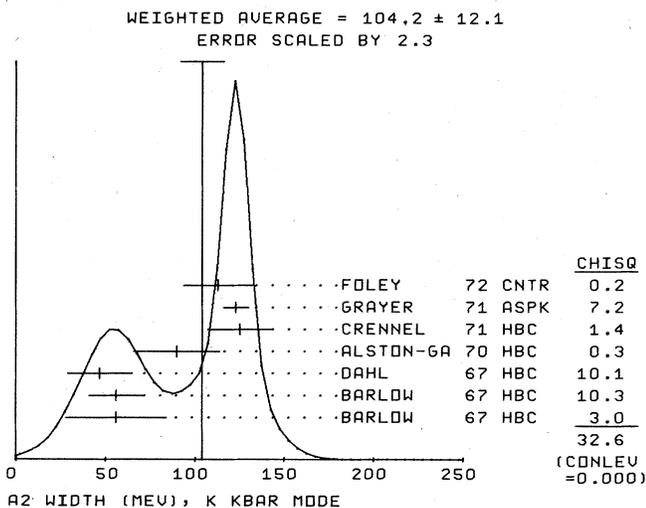


FIG. 3. Ideogram of measurements of the A_2 width, as determined from the $K\bar{K}$ mode. The vertical line indicates the position of the weighted average, while the horizontal bar atop the line gives the error in the average after scaling by the SCALE factor. Only those experiments indicated by + error flags were precise enough to be accepted in the calculation of the SCALE factor; the column on the far right gives the χ^2 contribution of each of these experiments. The less precise experiments were included in the calculation of the weighted average, but not SCALE; they have \pm error flags.

experiment would be well under the ceiling on SCALE.

This scaling approach has the property that if there are two values with comparable errors separated by much more than their stated errors (with or without a number of other experiments of lower accuracy), the error on the mean value $\delta\bar{x}$ is increased so that it is approximately half the interval between the two discrepant values.

We wish to emphasize the fact that our scaling procedures for errors in no way affect central values. In addition, if one wishes to recover the unscaled error $\delta\bar{x}$, he need only divide the given error by the SCALE factor for that error.

2. A second procedure—Student's distribution

The newly proposed method of averaging data, described in detail in Roos *et al.* (1975), relies upon an empirical determination of the distribution of the residuals for the ensemble of data appearing in the Review. The residual for the i th measurement of a quantity with average value \bar{x} is defined as

$$h_i = (x_i - \bar{x})/\delta x_i.$$

Roos *et al.* select several different subsamples of the data, and show that the residuals for each subsample have approximately the same properties; in particular, their first few even moments are similar. Since the distributions have longer tails than a Gaussian, the authors choose to represent them by a distribution function having such a property, namely the Student distribution

$$S_n(h/c) = K \left[1 + \frac{(h/c)^2}{n} \right]^{-(n+1)/2}. \quad (3)$$

Here K is a normalization constant, and n and c are parameters which the authors then fit to the combined sample of data. The resulting empirical distribution is

$$S_{10}(h/1.11) = 0.351 \left[1 + \frac{(h/1.11)^2}{10} \right]^{-11/2}. \quad (4)$$

Note that the shape of S_{10} is somewhere between that of a Gaussian ($=S_\infty$) and that of a Breit-Wigner ($=S_1$).

The proposed method of averaging the data for a given quantity then consists of finding the value of \bar{x} which maximizes the log-likelihood function

$$\text{Log} \mathcal{L}(\{x_i\} | \bar{x}) = \sum_i S_{10} \left(\frac{x_i - \bar{x}}{1.11 \delta x_i} \right); \quad (5)$$

the sum here is again taken over all N measurements of x . The error $\delta\bar{x}$ is determined by finding the variation in \bar{x} needed to decrease the log-likelihood by 1/2:

$$\log \mathcal{L}(\{x_i\} | \bar{x}) - \log \mathcal{L}(\{x_i\} | \bar{x} \pm \delta\bar{x}) = \frac{1}{2}. \quad (6)$$

3. Comparison of procedures

Both of the procedures described above adopt a partially empirical approach to the problem that measured values for the quantities tabulated in this Review do not exhibit the Gaussian behavior naively expected. (This problem, it should be noted, persists even when careful attempts are made to resolve difficulties and inconsistencies in the data prior to averaging.)

The first approach operates on a quantity-by-quantity basis and adjusts the error in each case so that no scaled $\chi^2/(N-1)$ is greater than 1. This is obviously rather conservative, since even if the data obeyed a Gaussian distribution, about half of the quantities would be expected to have $\chi^2/(N-1) > 1$.

The second approach, on the other hand, assumes that (provided we first eliminate quantities with obvious, known problems) all quantities have the same theoretical distribution function, namely the fairly long-tailed $S_{10}(h/1.11)$. With this supposition, if a particular quantity has a large χ^2 , it is assumed to be just a happenstance, occasioned by a random fluctuation into the long tails, and no special scaling for this quantity is done. This procedure thus results in generally smaller, or less conservative, error estimates for quantities having $\chi^2/(N-1) > 1$. (However, it should be noted that, because of the overall scale of 1.11 appearing in the empirical Student's distribution, the errors for quantities with $\chi^2/(N-1) \leq 1$ are actually increased by about 10%.) Table III shows some comparisons of sample results from the two procedures, using data from the 1974 edition of the Review. Shifts in both \bar{x} and $\delta\bar{x}$ can be observed, especially where $\text{SCALE} > 1$.

Since the new procedure is a significant departure from the past, we have adopted the following approach for this year: in the Data Card Listings we give the average-and-error for each quantity calculated both ways; the standard way is labeled at the left with the code "AVG," while the proposed newer way is labeled

TABLE III. Comparison of procedures.

Particle property	Pure Gaussian $\bar{x} \pm \delta\bar{x}$	Standard method: Gaussian + scale factor $\bar{x} \pm \delta\bar{x}$	SCALE	Proposed method: Student's distribution $\bar{x} \pm \delta\bar{x}$
ρ^0 mass (MeV)	770.32 ± 0.65	770.32 ± 0.91	1.4	770.37 ± 0.82
η' mass (MeV)	957.59 ± 0.24	957.59 ± 0.24	1.0	957.58 ± 0.28
ϕ mass (MeV)	1019.69 ± 0.15	1019.69 ± 0.28	1.9	1019.83 ± 0.20
$K^*(1420)$ mass (MeV)	1421.3 ± 2.3	1421.3 ± 2.3	1.0	1421.3 ± 2.6
K_S^0 mean life (10^{-8} sec)	5.158 ± 0.042	5.158 ± 0.042	1.0	5.158 ± 0.046
Σ^+ mean life (10^{-10} sec)	0.8004 ± 0.0058	0.8004 ± 0.0058	1.0	0.8004 ± 0.0064
Σ^- mean life (10^{-10} sec)	1.482 ± 0.011	1.482 ± 0.017	1.5	1.479 ± 0.013
$K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ (%)	5.521 ± 0.075	5.521 ± 0.098	1.3	5.533 ± 0.089
$\Lambda \rightarrow p \pi^-$ (%)	63.99 ± 0.49	63.99 ± 0.49	1.0	63.98 ± 0.55

“STUDENT.” In the Tables of Particle Properties, we continue to use the standard procedure—Gaussian with SCALE factor. As in the past, a SCALE factor greater than 1 is indicated by the appearance of “S=...” next to the value and error.

We heartily invite your comments on the proposed Student’s distribution method. They will assist us in deciding on procedures for future editions.

B. Constrained fits

Except for trivial cases, all branching ratios and rate measurements are analyzed by computer program AHR. This program makes a simultaneous least-squares fit to all the data, and outputs the partial-decay fractions \bar{P}_i , width Γ , partial widths Γ_i , and their error matrix.

The original version of AHR was written by J. Peter Berge. It is documented separately, and we wish here only to give the simplest nontrivial example that permits us to comment on the error matrix and the scale factor.

Assume that a state has only three partial-decay fractions, P_1 , P_2 , and P_3 ($\sum P_i = 1$), which have been measured in four different ratios, R_1, \dots, R_4 , where, e.g., $R_1 = P_1/P_2, R_2 = P_1/P_3$, etc.⁴ Further assume that *each* ratio has been measured by N experiments (we designate each experiment with a subscript x , e.g., R_{1x}). Then AHR finds the best values of P_1 , P_2 , and P_3 by minimizing χ^2 , namely

$$\chi^2 = \sum_{r=1}^4 \left[\sum_{x=1}^N \left(\frac{R_{rx} - R_r(P_1, P_2, P_3)}{\delta R_{rx}} \right)^2 \right]. \quad (7)$$

In addition to the fitted values \bar{P}_i , the program calculates an error matrix $\langle \delta \bar{P}_i \delta \bar{P}_j \rangle$. We tabulate the diagonal elements $\delta \bar{P}_i = \langle \delta \bar{P}_i \delta \bar{P}_i \rangle^{1/2}$ [except that some errors are scaled according to Eq. (2) as discussed below]. In the listings we give the complete error matrix; we also calculate the fitted value of each ratio, for comparison with the input data, and list it below the relevant input, along with a simple unconstrained average of the same input.

Two further comments on the example above:

(1) There was no connection between measurements of the width and the branching ratios. But often we also have information on partial widths Γ_i as well as total width Γ . In this case AHR must introduce Γ as a parameter into the fit, along with the relations $\Gamma_i = \Gamma P_i$, $\sum \Gamma_i = \Gamma$. When appropriate, we tabulate the Γ_i along with the P_i , and give error matrices in the listings.

(2) Note that we do *not* allow for correlations between input data. We *do* try to pick those ratios and widths which are as independent and as close to the original data as possible.

In *asymmetric* errors, we use a continuous function of $\delta(P)^+$ and $\delta(P)^-$ in the fitting. When no errors are reported, we merely list the data for inspection.

Hyperon-decay parameters

The program AHR handles any type of input, α , Φ , Δ , β , or γ , according to the definitions of Sec. VI. If for a

⁴We can handle any R of the form $R = \sum \alpha_i P_i / \sum \beta_i P_i$, where α_i and β_i are constants, usually 1 or 0.

particular hyperon decay there are data for more than two of the decay parameters, they are analyzed by using the constraint

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

Inconsistent constrained data

According to our simple example, which led to Eq. (7), the double sum for χ^2 is summed over experiments $x = 1$ to N , leaving a single sum over ratios

$$\chi^2 = \sum_r \chi_r^2.$$

Even before fitting, some of the χ_r^2 may be too large. But if we scaled them before fitting, then the scaling would move the central value, contrary to our policy. So we do not scale until after the first fit; then, knowing the fitted χ_r^2 and its expectation value $\langle \chi_r^2 \rangle$ we form SCALE factors (just as before), i.e.,

$$(\text{SCALE})_r^2 = \chi_r^2 / \langle \chi_r^2 \rangle,$$

and if any $(\text{SCALE})_r$ is greater than 1, all N of the measurements of that particular ratio are equally penalized by having their errors increased by SCALE. Program AHR then recycles on all the data, those with errors unchanged as well as those with errors increased. We then get new values, $\delta \bar{P}'_i$ for the errors in the partial decay modes.

Because of the constraint ($\sum P_i = 1$) some SCALE factors may still be greater than 1 even after this second pass. If this is so, the whole procedure (i.e., increasing errors by the new SCALE factors and recycling through AHR) is repeated.

At the end of AHR’s final pass we have *two* measures of the errors for the \bar{P}_i . One is, of course, the $\delta \bar{P}'_i$, i.e., the errors in the final fitted values \bar{P}'_i which include the effects of scaling the input errors. The other measure of the errors is $(\bar{P}_i - \bar{P}'_i)$, i.e., the *shift* in the central values of the i th mode between the first (unscaled) fit and the final (scaled) fit. In practice we find that on the average these two measures of the uncertainty are about equal. Rather than selecting just one or the other, our tabulated errors are given by the combination

$$(\delta \bar{P}_i)_{\text{tab}} = [\delta \bar{P}'_i{}^2 + (\bar{P}_i - \bar{P}'_i)^2]^{1/2},$$

where \bar{P}_i is the fitted value of the i th partial-decay mode before scaling, \bar{P}'_i is its value after scaling, and $\delta \bar{P}'_i$ is the error in \bar{P}'_i . The SCALE factors we finally list in such cases are defined by

$$(\text{SCALE})_i = (\delta \bar{P}_i)_{\text{tab}} / \delta \bar{P}'_i.$$

However, in line with our policy of not letting SCALE affect the central values, we give the values of \bar{P}_i obtained from the original (unscaled) fits. [The differences between the \bar{P}_i calculated with either the scaled or the unscaled errors are, of course, always within the tabulated errors, $(\delta \bar{P}_i)_{\text{tab}}$.]

VIII. A LOOK AT HISTORY

It may be said that one can estimate the age of a high energy physicist by asking him or her the mass of the

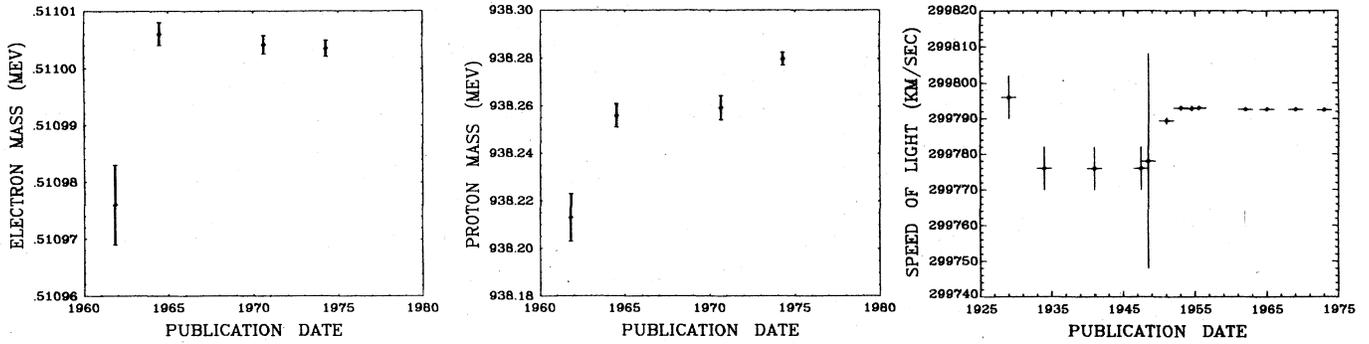


FIG. 4. The "generally accepted values" of the proton mass, the electron mass, and the speed of light, as a function of the publication data of the compilation used (not done by Particle Data Group). Data for the speed of light plot courtesy of E. R. Cohen, Rockwell International Science Center. See the Stable Particle Data Card Listings for references on proton and electron masses.

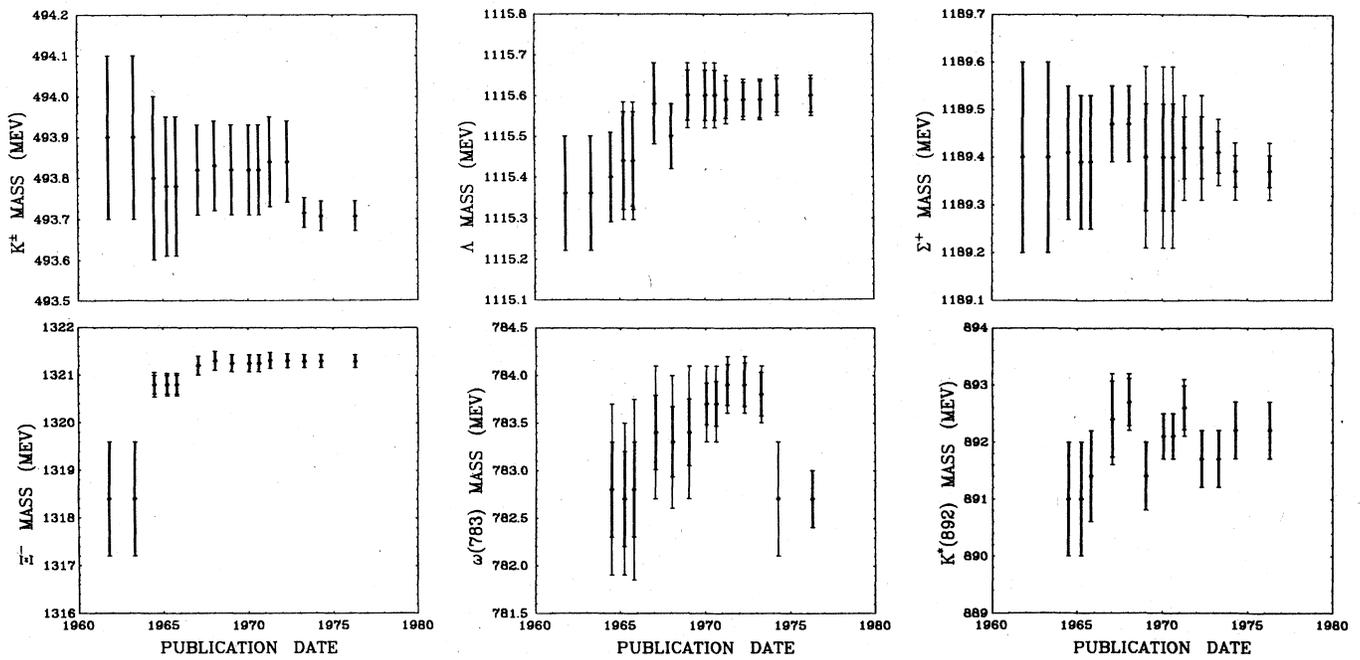


FIG. 5. Particle Data Group averages of the masses of various particles, as a function of date of publication of Review of Particle Properties (Adapted, with permission, from *Annual Review of Nuclear Science*, Volume 25. Copyright © 1975 by Annual Reviews, Inc. All rights reserved).

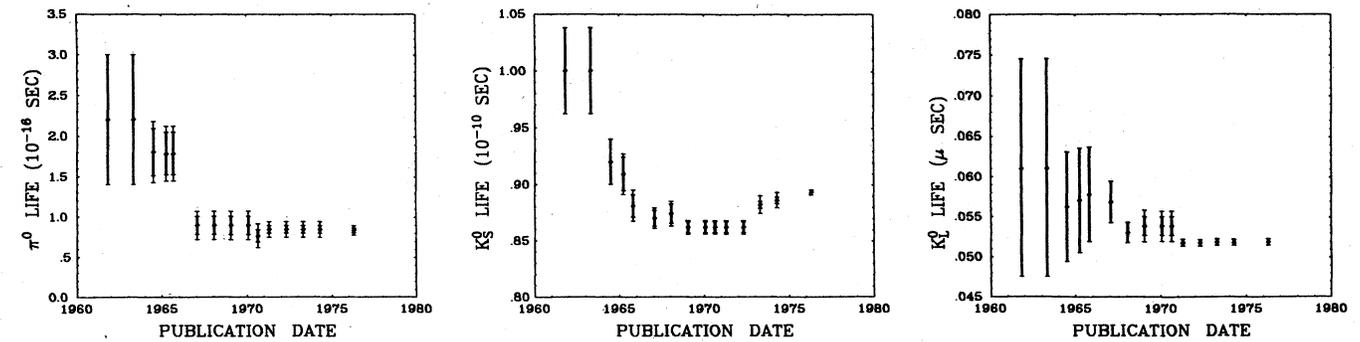


FIG. 6. Particle Data Group averages of the lifetimes of various particles, as a function of publication date of RPP.

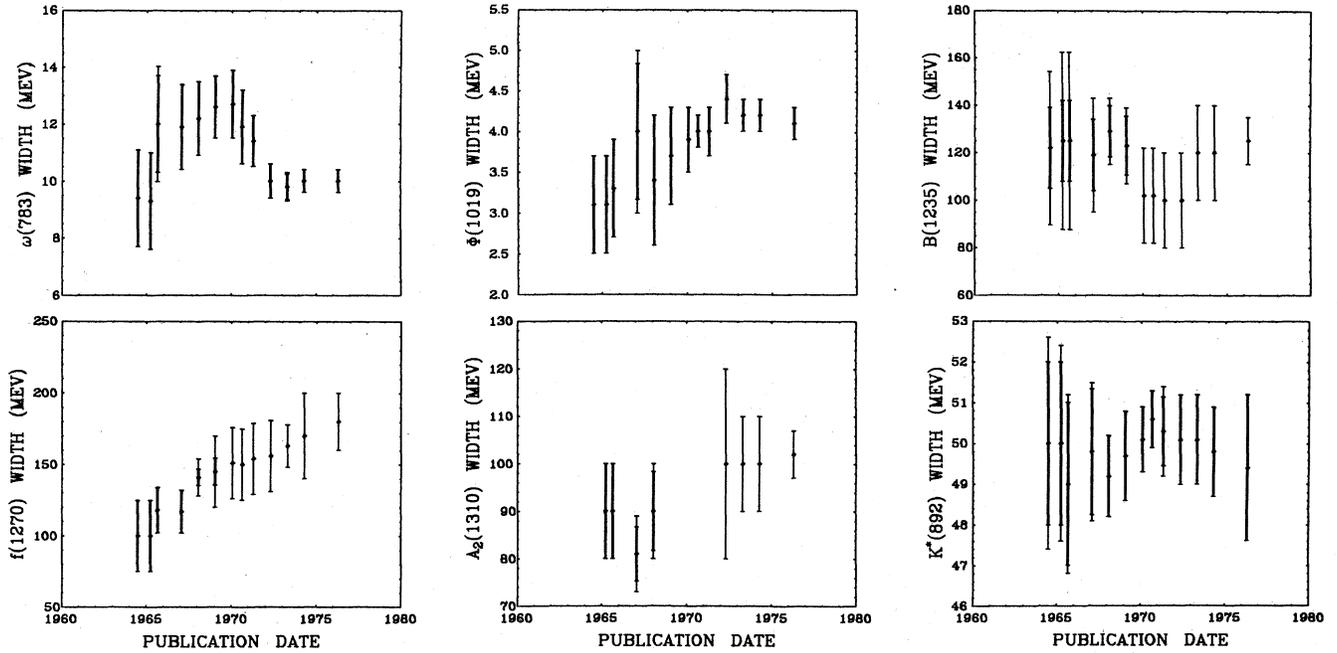


FIG. 7. Particle Data Group averages of the widths of various resonances, as a function of date of publication of RPP. The gap in the A_2 data indicates the years when the A_2 was thought to be split.

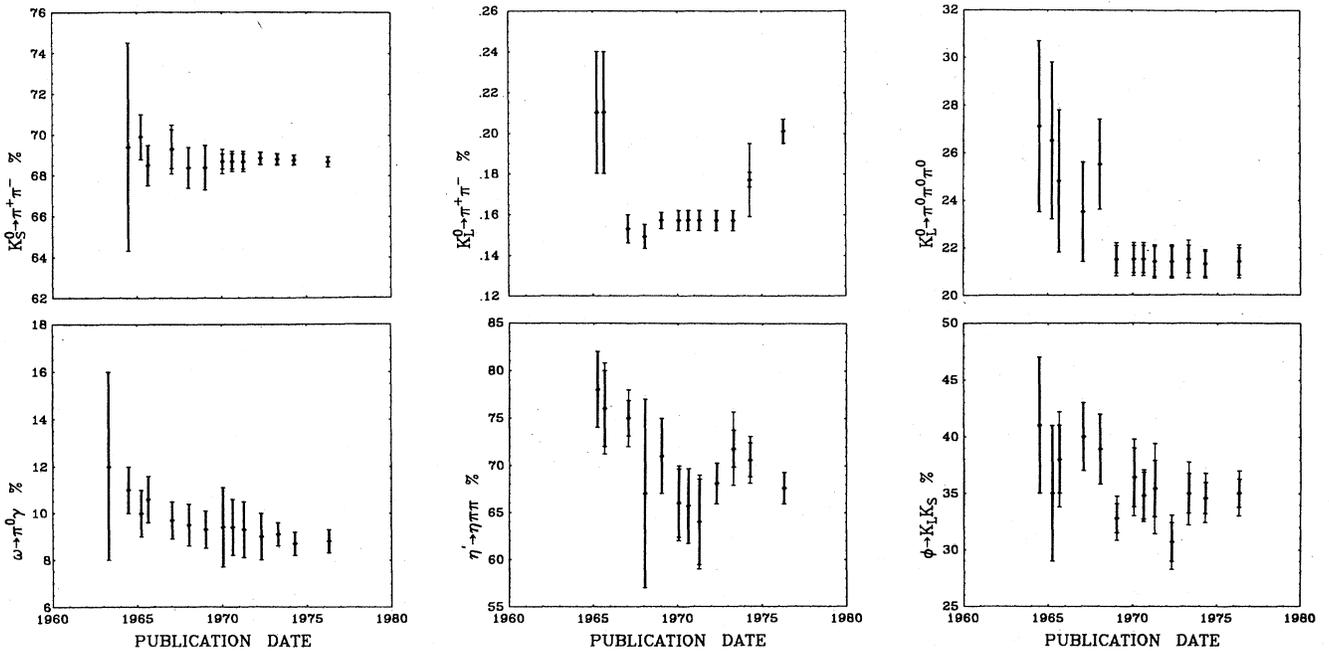


FIG. 8. Particle Data Group averages of various branching fractions, as a function of date of publication of RPP.

Λ^0 . If the answer is 1115.44 MeV, he probably was deep into his graduate training in 1965.

A history of the values of physical constants has more than whimsical value. In Fig. 4 we show how the generally accepted values for the speed of light and a couple of other constants have changed with time. The "generally accepted value" is usually an average over several experiments, performed by a compiler (in Fig. 4, the compiler is other than the Particle Data Group in all cases, although we do of course quote the compiled results). The x axis on all these figures is the date of publication of the value shown. Clearly there is a general progression toward better understanding—at least as measured by the size of the error bars. However, the size of the error bars do not tell the full story, as we can see by the frequency with which the "best" value has changed by more than one standard deviation. Changes in these values can come from several sources—a new experimental measurement, re-evaluation of an old measurement (which can come about if a previously unrecognized source of bias is discovered and corrected, or if a new value for one of the input constants, e.g. the electric charge, is available), or a change in the averaging procedure.

In Fig. 5 we show the history of some masses (including the Λ^0 , for radioactive Λ^0 dating of your colleagues), based on averages which we ourselves have performed. All of these were originally presented in Rosenfeld (1975). The publication date refers to the publication of the Review of Particle Properties.

In Fig. 6 we show the best estimates for the lifetimes of some of the particles stable against strong decay. These and subsequent figures have been compiled since publication of Rosenfeld (1975). In Fig. 7 we show the widths of some of the resonances, and in Fig. 8, the values of some of the branching fractions. All values are taken from the Tables. Before 1964, very few branching fractions were listed in the Tables. In all cases, a representative sample is chosen. In each figure, the heavy inner error bar represents the statistical error computed in the averaging procedure, and the thin outer error bars, when present, indicate the increase in the error due to the scale factor. The scale factor is described in the preceding section. It represents an attempt to quantify the increase in the uncertainty which is present in the case of experiments which disagree by more than a certain amount. In the case where the error represents an "educated guess," rather than a calculation, the inner error bar is absent.

On the whole, the number of times the values have changed by more than one standard deviation over the years is remarkably few. Even those branching fractions which involve rare decays and which are therefore presumably difficult to measure (Fig. 8) are, for the most part, within one or two standard deviations in 1974 of their value in any year since 1960. This is in spite of the vast amount of new experimental input, and indicates the general reliability of the results.

Of course, the data points for the different years are hardly independent of each other, but those differing by several years frequently have quite different experimental input. The relative lack of change is a comment both

on the experiments and on the averaging procedures. We, of course, are responsible only for the averages (but not on Fig. 4). These averages entail considerable exercise of judgment: there are conflicting experiments, experiments with impossibly small errors, "preliminary" results, and so forth. Statistical procedures will tell us that two experiments do not agree; they do not give a clue as to which (if either) is a good representation of the truth. Major decisions, and their motivations, are usually discussed on a case-by-case basis in the Data Card Listings; general comments may be found in Sec. II and in Rosenfeld (1975). Note that, occasionally, the error bars increase from one publication to the next, in these figures. This is usually the result of decision making by the compiler, e.g., to cease using a particular result.

We show these figures not only to demonstrate that there is not much change in these averages in the usual case, but also to show that there exist cases with relatively large changes. There is a psychological danger in preparing tables of "right" answers. The old joke about the experimenter who fights the systematics until he or she get the "right" answer (read "agrees with previous experiments"), and then publishes, contains a germ of truth (presumably, those who compile and average experimental results are also not immune to this disease). A result can disagree with the average of all previous experiments by five standard deviations, and still be right! Hence, perhaps it is of value to show that large changes can (and do) sometimes occur.

In summary, with the addition of Figs. 7–8, not available at the time of publication of Rosenfeld (1975), we find we can reiterate his conclusions. Namely, that the combination of careful work by experimenters and by compilers (which involved excluding around 40% of the data from the averages, adjusting impossibly small errors, etc.), and the frequent use of the conservative scale factor, has produced averages whose reliability is remarkably good.

ACKNOWLEDGMENTS

We would like to take this opportunity to thank all those who have assisted in the many phases of preparing this Review. In particular, we want to acknowledge the usefulness of the feed-back from the physics community, especially those who have taken the trouble to make suggestions or point out errors. The European members of the Particle Data Group wish to acknowledge the generous support of CERN; in particular through Dr. A. Günther and his services.

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TABLES OF PARTICLE PROPERTIES

April 1976

N. Barash-Schmidt, A. Barbaro-Galtieri, C. Bricman, V. Chaloupka,
 R. J. Hemingway, R. L. Kelly, M. J. Losty, A. Rittenberg,
 M. Roos, A. H. Rosenfeld, T. G. Trippe, G. P. Yost

(Closing date for data: Jan. 1, 1976)

Stable Particle Table

For additional parameters, see Addendum to this table.

Quantities in italics have changed by more than one (old) standard deviation since April 1974.

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean Life (sec) $c\tau$ (cm)	Partial decay mode		
				Mode	Fraction ^a	p or P _{max} ^b (MeV/c)
γ	0,1(1 ⁻) ⁻	0(<7×10 ⁻²²)	stable	stable		
ν	J=1/2 ν_e : 0(<0.00006) ν_μ : 0(<0.65)		stable	stable		
e	J=1/2	0.5110034 ±.0000014	stable (>5×10 ²¹ y)	stable		
μ	J=1/2 $m_\mu - m_\pi = -33.909$ ±.006	105.65948 ±.00035 $m^2 = 0.01116$	<i>2.197134</i> ×10 ⁻⁶ <i>±.000077</i> $c\tau = 6.5868 \times 10^4$	$e\nu\bar{\nu}$ $e\gamma\gamma$ 3e $e\gamma$ $e^+\nu_e\nu_\mu$	100 (<4)×10 ⁻⁶ (<6)×10 ⁻⁹ (<2.2)×10 ⁻⁸ (<25)%	53 53 53 53 53
π^\pm	1 ⁻ (0 ⁻)	139.5688 ±.0064 $m^2 = 0.0195$	2.6030×10 ⁻⁸ ±.0023 $c\tau = 780.4$ $(\tau^+ - \tau^-)/\tau =$ <i>(0.05±0.07)%</i> (test of CPT)	$\mu\nu$ $e\nu$ $\mu\nu\gamma$ $\pi^0 e\nu$ $e\nu\gamma$ $e\nu e^+e^-$	100 (1.267±0.023)×10 ⁻⁴ ^c (1.24±0.25)×10 ⁻⁴ (1.02±0.07)×10 ⁻⁸ ^c (3.0±0.5)×10 ⁻⁸ (<3.4)×10 ⁻⁸	30 70 30 5 70 70
π^0	1 ⁻ (0 ⁻) ⁺	134.9645 ±.0074 $m^2 = 0.182$ $m_\pi^\pm - m_\pi^0 = 4.6043$ ±.0037	0.828×10 ⁻¹⁶ ±.057 S=1.8* $c\tau = 2.5 \times 10^{-6}$	$\gamma\gamma$ γe^+e^- $\gamma\gamma\gamma$ $e^+e^-e^+e^-$ ^d $\gamma\gamma\gamma$ e^+e^-	(98.85±0.05)% (1.15±0.05)% (<5)×10 ⁻⁶ (3.32)×10 ⁻⁵ (<6)×10 ⁻⁵ (<2)×10 ⁻⁶	67 67 67 67 67 67

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean life (sec) $c\tau$ (cm)	Partial decay mode		p or Pmax ^b (MeV/c)
				Mode	Fraction ^a	
K[±]	$\frac{1}{2}(0^-)$	493.707 ±0.037 $m^2=0.244$ $m_{K^\pm}-m_{K^0}=-3.99$ ±0.13 $S=1.1^*$	1.2371×10 ⁻⁸ ±0.0026 $S=1.9^*$ $c\tau=370.9$ ($\tau^+-\tau^-$)/ $\bar{\tau}^-$ (±.11±.09)% (test of CPT) $S=1.2^*$	$\mu\nu$	(63.61±0.16)%	236
				$\pi\pi^0$	(21.05±0.14)%	205
				$\pi\pi^-\pi^+$	(5.59±0.03)% $S=1.1^*$	125
				$\pi\pi^0\pi^0$	(1.73±0.05)% $S=1.4^*$	133
				$\mu\pi^0\nu$	(3.20±0.09)% $S=1.7^*$	215
				$e\pi^0\nu$	(4.82±0.05)% $S=1.1^*$	228
				$\mu\nu\gamma$	$c(5.8 \pm 3.5) \times 10^{-3}$	236
				$e\pi^0\mu^0\nu$	(1.8 ± 2.4) × 10 ⁻⁵	207
				$\pi\pi^+e^\pm\nu$	(3.7 ± 0.5) × 10 ⁻⁵	203
				$\pi\pi^+e^\mp\nu$	(<5) × 10 ⁻⁷	203
				$\pi\pi^+\mu^\pm\nu$	(0.9 ± 0.4) × 10 ⁻⁵	151
				$\pi\pi^+\mu^\mp\nu$	(<3.0) × 10 ⁻⁶	151
				$e\nu$	(1.54±0.09) × 10 ⁻⁵	247
				$e\nu\gamma$	$c(1.62 \pm 0.47) \times 10^{-5}$	247
				$\pi\pi^0\gamma$	$e,c(2.71 \pm 0.19) \times 10^{-4}$	205
				$\pi\pi^+\pi^-\gamma$	$c(1.0 \pm 0.4) \times 10^{-4}$	125
				$\mu\pi^0\nu\gamma$	$c(<6) \times 10^{-5}$	215
				$e\pi^0\nu\gamma$	$c(3.7 \pm 1.4) \times 10^{-4}$	228
				πe^+e^-	(2.6 ± 0.5) × 10 ⁻⁷	227
				$\pi^+e^+e^\pm$	(<1.5) × 10 ⁻⁵	227
				$\pi\mu^+\mu^-$	(<2.4) × 10 ⁻⁶	172
				$\pi\gamma\gamma$	$c(<3.5) \times 10^{-5}$	227
				$\pi\gamma\gamma\gamma$	$c(<3.0) \times 10^{-4}$	227
				$\pi\nu\nu$	(<0.6) × 10 ⁻⁶	227
				$\pi\gamma$	(<4) × 10 ⁻⁶	227
				$e\pi^+e^\pm$	(<2.8) × 10 ⁻⁸	214
$e\pi^+e^\mp$	(<1.4) × 10 ⁻⁸	214				
$\mu\nu\nu$	(<6) × 10 ⁻⁶	236				
K⁰	$\frac{1}{2}(0^-)$	497.70 ±0.13 $S=1.1^*$ $m^2=0.248$	50% K _{Short} , 50% K _{Long}			
K_S⁰	$\frac{1}{2}(0^-)$	0.8930×10 ⁻¹⁰ (f) ±0.0023 $c\tau=2.68$	$\pi^+\pi^-$	(68.67 ± 0.25)% $S=1.1^*$	206	
			$\pi^0\pi^0$	(31.33 ± 0.25)%	209	
			$\mu^+\mu^-$	(<3.2) × 10 ⁻⁷	225	
			e^+e^-	(<3.4) × 10 ⁻⁴	249	
			$\pi^+\pi^-\gamma$	$c(2.0 \pm 0.4) \times 10^{-3}$	206	
			$\gamma\gamma$	(<0.4) × 10 ⁻³	249	
K_L⁰	$\frac{1}{2}(0^-)$	5.181×10 ⁻⁸ ±0.040 $c\tau=1553$ $m_{K_L}-m_{K_S}=0.5349 \times 10^{10} \hbar \text{ sec}^{-1}$ ±0.0022	$\pi^0\pi^0\pi^0$	(21.4 ± 0.7)% $S=1.2^*$	139	
			$\pi^+\pi^-\pi^0$	(12.25±0.18)% $S=1.1^*$	133	
			$\pi\mu\nu$	(27.1 ± 0.5)%	216	
			$\pi e\nu$	$E(39.0 \pm 0.5)\%$ $S=1.1^*$	229	
			$\pi e\nu\gamma$	$E,c(1.3 \pm 0.8)\%$	229	
			$\pi^+\pi^-$	$f(0.201 \pm 0.006)\%$	206	
			$\pi^0\pi^0$	$c(0.094 \pm 0.019)\%$ $S=1.5^*$	209	
			$\pi^+\pi^-\gamma$	$c(6.0 \pm 2.0) \times 10^{-5}$	206	
			$\pi^0\gamma\gamma$	(<2.4) × 10 ⁻⁴	231	
			$\gamma\gamma$	(4.9 ± 0.5) × 10 ⁻⁴	249	
			$e\mu$	(<2.0) × 10 ⁻⁹	238	
			$\mu^+\mu^-$	$h(1.0 \pm 0.3) \times 10^{-8}$	225	
			$\mu^+\mu^-\gamma$	(<7.8) × 10 ⁻⁶	225	
			$\mu^+\mu^-\pi^0$	(<5.7) × 10 ⁻⁵	177	
			e^+e^-	(<2.0) × 10 ⁻⁹	249	
			$e^+e^-\gamma$	(<2.8) × 10 ⁻⁵	249	
$\pi^+\pi^-e^+e^-$	(<7.2) × 10 ⁻⁶	206				
$\pi^0\pi^+e^\mp\nu$	(<2.2) × 10 ⁻³	207				
η	$0^+(0^-)+$	548.8 ±0.6 $S=1.4^*$ $m^2=0.301$ $e\Gamma=(0.85 \pm 0.12)\text{keV}^{(1)}$ Neutral decays (71.0±0.7)% $S=1.1^*$ Charged decays (29.0±0.7)% $S=1.1^*$	$\gamma\gamma$	(38.0 ± 1.0)% $S=1.2^*$	274	
			$\pi^0\gamma\gamma$	$l(3.1 \pm 1.1)\%$ $S=1.2^*$	258	
			$3\pi^0$	(29.9 ± 1.1)% $S=1.1^*$	180	
			$\pi^+\pi^-\pi^0$	(23.6 ± 0.6)% $S=1.1^*$	175	
			$\pi^+\pi^-\gamma$	(4.89±0.13)% $S=1.1^*$	236	
			$e^+e^-\gamma$	(0.50±0.12)%	274	
			$\pi^0e^+e^-$	(<0.04)%	258	
			$\pi^+\pi^-$	(<0.15)%	236	
			$\pi^+\pi^-e^+e^-$	(0.1 ± 0.1)%	236	
			$\pi^+\pi^-\pi^0\gamma$	(<6) × 10 ⁻⁴	175	
			$\pi^+\pi^-\gamma\gamma$	(<0.2)%	236	
			$\mu^+\mu^-\pi^0$	(2.2 ± 0.8) × 10 ⁻⁵	253	
$\mu^+\mu^-\pi^0$	(<5) × 10 ⁻⁴	211				

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean Life (sec) $c\tau$ (cm)	Partial decay mode		
				Mode	Fraction ^a	p or p_{max}^b (MeV/c)
p	$\frac{1}{2}(\frac{1}{2}^+)$	938.2796 ± 0.0027 $m^2 = 0.8804$	stable ($>2 \times 10^{30}y$)			
n	$\frac{1}{2}(\frac{1}{2}^+)$	939.5731 ± 0.0027 $m^2 = 0.8828$ $m_p - m_n = -1.29343$ ± 0.00004	918 \pm 14 $c\tau = 2.75 \times 10^{13}$	$pe^- \nu$	100 %	1
Λ	$0(\frac{1}{2}^+)$	1115.60 ± 0.05 $S = 1.2^*$ $m^2 = 1.245$	2.578×10^{-10} ± 0.021 $S = 1.6^*$ $c\tau = 7.73$	$p\pi^-$ $n\pi^0$ $pe^- \nu$ $p\mu^- \nu$ $p\pi^- \gamma$	(64.2 \pm 0.5)% (35.8 \pm 0.5)% (8.13 \pm 0.29) $\times 10^{-4}$ (1.57 \pm 0.35) $\times 10^{-4}$ ^c (0.85 \pm 0.14) $\times 10^{-3}$	100 104 163 131 100
Σ^+	$1(\frac{1}{2}^+)$	1189.37 ± 0.06 $S = 1.8^*$ $m^2 = 1.415$ $m_{\Sigma^+} - m_{\Sigma^-} = -7.98$ ± 0.08 $S = 1.2^*$	0.800×10^{-10} ± 0.006 $c\tau = 2.40$ $\frac{\Gamma(\Sigma^+ \rightarrow l^+ \nu)}{\Gamma(\Sigma^- \rightarrow l^- \nu)} < 0.043$	$p\pi^0$ $n\pi^+$ $p\gamma$ $n\pi^+ \gamma$ $\Lambda e^+ \nu$ $\Lambda \mu^+ \nu$ $pe^+ e^-$	(51.6 \pm 0.7)% (48.4)% (1.24 \pm 0.18) $\times 10^{-3}$ $S = 1.4^*$ ^c (0.93 \pm 0.10) $\times 10^{-3}$ (2.02 \pm 0.47) $\times 10^{-5}$ (<3.0) $\times 10^{-5}$ (<0.5) $\times 10^{-5}$ (<7) $\times 10^{-6}$	189 185 225 185 71 202 224 225
Σ^0	$1(\frac{1}{2}^+)$	1192.47 ± 0.08 $m^2 = 1.422$	$< 1.0 \times 10^{-14}$ $c\tau < 3 \times 10^{-4}$	$\Lambda\gamma$ $\Lambda e^+ e^-$ $\Lambda\gamma\gamma$	100 % ^d (5.45) $\times 10^{-3}$ (<3)%	74 74 74
Σ^-	$1(\frac{1}{2}^+)$	1197.35 ± 0.06 $m^2 = 1.434$ $m_{\Sigma^0} - m_{\Sigma^-} = -4.88$ ± 0.06	1.482×10^{-10} ± 0.017 $S = 1.5^*$ $c\tau = 4.44$	$n\pi^-$ $ne^- \nu$ $n\mu^- \nu$ $\Lambda e^- \nu$ $n\pi^- \gamma$	100 % (1.08 \pm 0.04) $\times 10^{-3}$ (0.45 \pm 0.04) $\times 10^{-3}$ (0.60 \pm 0.06) $\times 10^{-4}$ ^c (4.6 \pm 0.6) $\times 10^{-4}$	193 230 210 79 193
Ξ^0	$\frac{1}{2}(\frac{1}{2}^+)(j)$	1314.9 ± 0.6 $m^2 = 1.729$ $m_{\Xi^0} - m_{\Xi^-} = -6.4$ ± 0.6	2.96×10^{-10} ± 0.12 $c\tau = 8.87$	$\Lambda\pi^0$ $\Lambda\gamma$ $\Sigma^0\gamma$ $p\pi^-$ $pe^- \nu$ $\Sigma^+ e^- \nu$ $\Sigma^- e^+ \nu$ $\Sigma^+ \mu^- \nu$ $\Sigma^- \mu^+ \nu$ $p\mu^- \nu$	100 % (0.5 \pm 0.5)% (<7)% (<3.6) $\times 10^{-5}$ (<1.3) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<0.9) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<0.9) $\times 10^{-3}$ (<1.3) $\times 10^{-3}$	135 184 117 299 323 120 112 64 49 309
Ξ^-	$\frac{1}{2}(\frac{1}{2}^+)(j)$	1321.29 ± 0.14 $m^2 = 1.746$	1.652×10^{-10} ± 0.023 $S = 1.1^*$ $c\tau = 4.95$	$\Lambda\pi^-$ $\Lambda e^- \nu$ $\Sigma^0 e^- \nu$ $\Lambda\mu^- \nu$ $\Sigma^0 \mu^- \nu$ $n\pi^-$ $ne^- \nu$ $n\mu^- \nu$ $\Sigma^- \gamma$ $p\pi^- \pi^-$ $p\pi^- e^- \nu$ $p\pi^- \mu^- \nu$ $\Xi^0 e^- \nu$	100 % ^k (0.69 \pm 0.18) $\times 10^{-3}$ (<0.5) $\times 10^{-3}$ (3.5 \pm 3.5) $\times 10^{-4}$ (<0.8) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<3.2) $\times 10^{-3}$ (<1.5)% (<1.2) $\times 10^{-3}$ (<4) $\times 10^{-4}$ (<4) $\times 10^{-4}$ (<4) $\times 10^{-4}$ (<2.3) $\times 10^{-4}$	139 190 123 163 70 303 327 313 118 223 304 250 6
Ω^-	$0(\frac{3}{2}^+)(j)$	1672.2 ± 0.4 $m^2 = 2.796$	$1.3_{-0.2}^{+0.3} \times 10^{-10}$ $c\tau = 4.0$	$\Xi^0 \pi^-$ $\Xi^- \pi^0$ ΛK^-	Total of 43 events seen	293 290 211

ADDENDUM TO
Stable Particle Table

e $1.001\ 159\ 6567 \frac{e\hbar}{2m_e c}$ $\pm .000\ 000\ 0035$		μ Decay parameters (4)			
μ $1.001\ 166\ 897 \frac{e\hbar}{2m_\mu c}$ $\pm .000\ 000\ 027$		$\rho = 0.752 \pm 0.003$ $\xi = 0.972 \pm 0.013$ $ g_A/g_V = 0.86^{+0.33}_{-0.11}$	$\eta = -0.12 \pm 0.21$ $\delta = 0.755 \pm 0.009$ $\phi = 180^\circ \pm 15^\circ$	$h = 1.00 \pm 0.13$	
K^\pm Mode	Partial rate (sec ⁻¹)		$\Delta I = \frac{1}{2}$ rule for $K \rightarrow 3\pi$ (m)		
	$\mu\nu$ $\pi\pi^0$ $\pi\pi^+\pi^-$ $\pi\pi^0\pi^0$ $\mu\pi^0\nu$ $e\pi^0\nu$	(51.42±0.17)×10 ⁶ (17.02±0.12)×10 ⁶ (4.52±0.02)×10 ⁶ (1.40±0.04)×10 ⁶ (2.58±0.07)×10 ⁶ (3.90±0.04)×10 ⁶	S=1.2* S=1.1* S=1.1* S=1.4* S=1.7* S=1.1*	$K^+ \rightarrow \pi^+\pi^+\pi^-$ $g = -0.214 \pm 0.005$ $S = 1.7^*$ $K^- \rightarrow \pi^-\pi^-\pi^+$ $g = -0.214 \pm 0.007$ $S = 2.7^*$ $K^\pm \rightarrow \pi^0\pi^0\pi^\pm$ $g = 0.550 \pm 0.020$ $S = 1.6^*$ $K_L^0 \rightarrow \pi^+\pi^-\pi^0$ $g = 0.646 \pm 0.014$ $S = 2.5^*$	
K_S^0	$\pi^+\pi^-\pi^0$ $\pi^0\pi^0$	$n(0.7689 \pm 0.0035) \times 10^{10}$ $n(0.3509 \pm 0.0030) \times 10^{10}$	$S = 1.1^*$	Form factors for $K_{\ell 3}$ decays (For ξ , λ_+^μ , and λ_0^μ see Data Card Listings, especially note in K^\pm section.) K_{e3}^+ $\lambda_+^e = 0.029 \pm 0.004$ K_{e3}^0 $\lambda_+^e = 0.0288 \pm 0.0028$ $S = 1.4^*$	
K_L^0	$\pi^0\pi^0\pi^0$ $\pi^+\pi^-\pi^0$ $\pi\mu\nu$ $\pi e\nu$ $\pi^+\pi^0$ f,n $\pi^0\pi^0$ n	(4.13 ± 0.14) × 10 ⁶ (2.36 ± 0.04) × 10 ⁶ (5.23 ± 0.10) × 10 ⁶ (7.52 ± 0.11) × 10 ⁶ (3.88 ± 0.11) × 10 ⁴ (1.81 ± 0.37) × 10 ⁴	S=1.2* S=1.1* S=1.5*	CP violation parameters (o,n,t) $ \eta_{+-} = (2.272 \pm 0.023) \times 10^{-3}$ $ \eta_{00} = (2.32 \pm 0.09) \times 10^{-3}$ $S = 1.1^*$ $\phi_{+-} = (45.0 \pm 1.2)^\circ$ $\phi_{00} = (48 \pm 13)^\circ$ $ \eta_{+-} ^2 < 0.12$ $ \eta_{00} ^2 < 0.28$ $\delta = (0.330 \pm 0.012) \times 10^{-2}$ $\Delta S = -\Delta Q$ $\text{Re } x = 0.008 \pm 0.020$ $S = 1.4^*$ $\text{Im } x = -0.003 \pm 0.027$ $S = 1.2^*$	
	η Mode	Left-right asymmetry $\pi^+\pi^-\pi^0$ (0.12 ± 0.17)% $\pi^+\pi^-\gamma$ (0.88 ± 0.40)%	Sextant asymmetry (0.19 ± 0.16)%	Quadrant asymmetry (-0.17 ± 0.17)% $\beta = 0.047 \pm 0.062$	
P	Magnetic moment (eħ/2m _p c)	Decay parameters (p)			
	2.7928456 ± 0.000011	Measured α	Derived ϕ (degree)	γ	Δ (degree)
π	-1.913148 ± 0.000066	$p e^- \nu$	-1.250 ± 0.009 $\delta = (181.1 \pm 1.3)^\circ$		
Λ	-0.67 ± 0.06	$p\pi^-$ 0.647 ± 0.013 (-6.5 ± 3.5)°	0.76	$(7.6^{+4.0}_{-4.1})^\circ$	
		$n\pi^0$ 0.651 ± 0.045 $p e^- \nu$	-0.66 ± 0.05 $S = 1.2^*$		
Σ^+	2.62 ± 0.41	$p\pi^0$ -0.979 ± 0.016 (36 ± 34)°	0.17	$(187 \pm 6)^\circ$	
		$n\pi^+$ +0.066 ± 0.016 (167 ± 20)° $p\gamma$ -1.03 ⁺⁵² ₋₄₂ $S = 1.1^*$	-0.97	$(-73^{+136}_{-10})^\circ$	
Σ^-	-1.48 ± 0.37	$n\pi^-$ -0.069 ± 0.008 (10 ± 15)°	0.98	$(249^{+12}_{-115})^\circ$	
		$n e^- \nu$ $\Lambda e^- \nu$	$\pm(0.435 \pm 0.035)$ 0.24 ± 0.23 $S = 1.3^*$		
Ξ^0	-0.44 ± 0.08 $S = 1.3^*$	$\Lambda\pi^0$ (21 ± 12)°	0.84	$(216^{+13}_{-19})^\circ$	
Ξ^-	-1.85 ± 0.75	$\Lambda\pi^-$ -0.392 ± 0.021 (2 ± 6)° $S = 1.1^*$	0.92	$(185 \pm 13)^\circ$	
Ω^-	$-0.66^{+0.36}_{-0.30}$	ΛK^-			

Stable Particle Table (cont'd)

*S = Scale factor = $\sqrt{\chi^2/(N-1)}$, where $N \approx$ number of experiments. S should be ≈ 1 . If $S > 1$, we have enlarged the error of the mean, $\delta\bar{x}$; i.e., $\delta\bar{x} \rightarrow S\delta\bar{x}$. This convention is still inadequate, since if $S \gg 1$ the experiments are probably inconsistent, and therefore the real uncertainty is probably even greater than $S\delta\bar{x}$. See text, and ideograms in Stable Particle Data Card Listings.

- a. Quoted upper limits correspond to a 90% confidence level.
- b. In decays with more than two bodies, p_{\max} is the maximum momentum that any particle can have.
- c. See Stable Particle Data Card Listings for energy limits used in this measurement.
- d. Theoretical value; see also Stable Particle Data Card Listings.
- e. The direct emission branching fraction is $(1.56 \pm .35) \times 10^{-5}$.
- f. The $\tau(K_S^0)$ and $|\eta_{+-}|$ averages (and the related $K_L^0 \rightarrow \pi^+\pi^-$ branching fraction and rate averages) contain only post-1971 results. The pre-1971 averages were $|\eta_{+-}| = (1.95 \pm 0.03) \times 10^{-3}$ and $\tau(K_S^0) = (0.862 \pm 0.006) \times 10^{-10}$ sec. See notes on $|\eta_{+-}|$ and $\tau(K_S^0)$ discrepancies in Stable Particle Data Card Listings.
- g. The branching fraction for $K_L^0 \rightarrow \pi\nu$ includes the radiative events $K_L^0 \rightarrow \pi\nu\gamma$.
- h. This is above the contradictory result of Clark et al. ($< 0.3 \times 10^{-8}$). See note in Stable Particle Data Card Listings.
- i. See note in Stable Particle Data Card Listings.
- j. P for Ξ and J^P for Ω^- not yet measured. Values reported are SU(3) predictions.
- k. Assumes rate for $\Xi^- \rightarrow \Sigma^0 e^- \nu$ small compared with $\Xi^- \rightarrow \Lambda e^- \nu$.
- l. $|g_A/g_V|$ defined by $g_A^2 = |C_A|^2 + |C'_A|^2$, $g_V^2 = |C_V|^2 + |C'_V|^2$, and $\Sigma(\bar{e}|\Gamma_1|\mu)(\bar{\nu}|\Gamma_1(C_1+C'_1\gamma_5)|\nu)$; ϕ defined by $\cos \phi = -\text{Re}(C_A^* C'_V + C'_A C_V)/g_A g_V$ [for more details, see text Section VI A].
- m. The definition of the slope parameter of the Dalitz plot is as follows [see also text Section VI B.1]:

$$|M|^2 = 1 + g \left(\frac{s_3 - s_0}{m_{\pi^+}^2} \right).$$

- n. The $K_S^0 \rightarrow \pi\pi$ and $K_L^0 \rightarrow \pi\pi$ rates (and branching fractions) are from independent fits and do not include results of K_L^0 - K_S^0 interference experiments. The $|\eta_{+-}|$ and $|\eta_{00}|$ values given in the addendum are these rates combined with the $|\eta_{+-}|$ and $|\eta_{00}|$ results from interference experiments.
- o. The definition for the CP violation parameters is as follows [see also text Section VI B.3]:

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L^0 \rightarrow \pi^+\pi^-)}{A(K_S^0 \rightarrow \pi^+\pi^-)} \quad \eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L^0 \rightarrow \pi^0\pi^0)}{A(K_S^0 \rightarrow \pi^0\pi^0)}$$

$$\delta = \frac{\Gamma(K_L^0 \rightarrow \pi^+\pi^-) - \Gamma(K_S^0 \rightarrow \pi^+\pi^-)}{\Gamma(K_L^0 \rightarrow \pi^+\pi^-) + \Gamma(K_S^0 \rightarrow \pi^+\pi^-)}, \quad |\eta_{+-}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^+\pi^-)}{\Gamma(K_L^0 \rightarrow \pi^+\pi^-)}, \quad |\eta_{00}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^0\pi^0)}{\Gamma(K_L^0 \rightarrow \pi^0\pi^0)}.$$

- p. The definition of these quantities is as follows [for more details on sign convention, see text Section VI B]:

$$\alpha = \frac{2|s||p|\cos\Delta}{|s|^2 + |p|^2} \quad \left| \quad \beta = \sqrt{1 - \alpha^2} \sin\phi \quad \left| \quad g_A/g_V \text{ defined by } \langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) | B_i \rangle \right.$$

$$\beta = \frac{-2|s||p|\sin\Delta}{|s|^2 + |p|^2} \quad \left| \quad \gamma = \sqrt{1 - \alpha^2} \cos\phi \quad \left| \quad \delta \text{ defined by } g_A'/g_V = |g_A/g_V| e^{i\delta} \right.$$

Meson Table

April 1976

In addition to the entries in the Meson Table, the Meson Data Card Listings contain all substantial claims for meson resonances. See Contents of Meson Data Card Listings⁽¹⁾.

Quantities in italics have changed by more than one (old) standard deviation since April 1974.

Name	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	$M^2 \pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode			
					Mode	Fraction (%) [Upper limits are 1σ (%)]	p or Pmax (b) (MeV/c)	
π^\pm π^0	$1^-(0^-)_+$	139.57 134.96	0.0 7.8 eV ± 1.9 eV	0.019479 0.018215	See Stable Particle Table			
η	$0^+(0^-)_+$	548.8 ± 0.6	2.63 keV ± 0.58 keV	0.301 ± 0.000	Neutral Charged	71.1 28.9	See Stable Particle Table	
$\rho(770)$	$1^+(1^-)_-$	773 ± 3 _s	152 ± 3 _s	0.598 ± 0.117	$\pi\pi$ $\pi\gamma$ e^+e^- $\mu^+\mu^-$	≈ 100 0.024 ± 0.007 0.0043 ± 0.0005 (d) 0.0067 ± 0.0012 (d)	360 374 386 372	
M and Γ from neutral mode.								
$\omega(783)$	$0^-(1^-)_-$	782.7 ± 0.3	10.0 ± 0.4	0.613 ± 0.008	$\pi^+\pi^-\pi^0$ $\pi^0\pi^-$ $\pi^0\gamma$ e^+e^-	89.9 ± 0.6 1.3 ± 0.3 8.8 ± 0.5 0.0076 ± 0.0017	S=1.2* S=1.5* S=1.9*	327 366 380 391
η' (958)	$0^+(0^-)_+$	957.6 ± 0.3	< 1	0.917 < .001	$\eta\pi\pi$ $\rho^0\gamma$ $\gamma\gamma$	67.6 ± 1.7 30.4 ± 1.7 2.0 ± 0.3	S=1.1*	231 167 479
For upper limits, see footnote (g)								
$\delta(970)$	$1^-(0^+)_+$	976 ^(h) ± 10 _s	50 ^(h) ± 20 _s	0.953 ± 0.049	$\eta\pi$	seen	315	
Possibly coupled to the $I = 1$ $K\bar{K}$ system ^(f) .								
S^* (993)	$0^+(0^+)_+$	~ 993 ^(c) ± 5	40 ^(c) ± 8	0.986 ± 0.040	$K\bar{K}$ $\pi\pi$	near threshold	53 476	
See note on $\pi\pi$ S wave ^(f) .								
$\Phi(1020)$	$0^-(1^-)_-$	1019.7 ± 0.3 S=1.6*	4.1 ± 0.2	1.040 ± 0.004	K^+K^- $K_L^0K_S^0$ $\pi^+\pi^-\pi^0$ (incl. $\rho\pi$) $\eta\gamma$ $\pi^0\gamma$ e^+e^- $\mu^+\mu^-$	46.6 ± 2.3 35.0 ± 2.0 16.4 ± 1.5 2.0 ± 0.4 0.14 ± 0.05 .032 ± 0.002 .025 ± 0.003	S=1.6* S=1.6* S=1.1* S=1.4*	128 111 462 362 501 510 499
For upper limits, see footnote (i)								
$A_1(1100)$	$1^-(1^+)_+$	~ 1100	~ 300	1.21 ± 0.33	$\rho\pi$	~ 100	251	
Broad enhancement in the $J^P=1^+$ $\rho\pi$ partial wave; not an established resonance ^(f) .								
$\epsilon(1200)$	$0^+(0^+)_+$	1100 to 1300	~ 600		$\pi\pi$			
Existence of pole not established. See note on $\pi\pi$ S wave ^(f) .								
$B(1235)$	$1^+(1^+)_-$	1228 _s ± 10 _s	125 _s ± 10 _s	1.51 ± 0.15	$\omega\pi$	only mode seen [D/S amplitude ratio = .25 ± 0.06] For upper limits, see footnote (j)	345	
$f(1270)$	$0^+(2^+)_+$	1271 _s ± 5 _s	180 _s ± 20 _s	1.62 ± 0.23	$\pi\pi$ $2\pi^+2\pi^-$ $K\bar{K}$ $\pi^+\pi^-2\pi^0$	81 ± 15 2.8 ± 0.3 2.7 ± 0.6 seen	S=1.1* 620 557 395 560	
For upper limits, see footnote (l)								
$D(1285)$	$0^+(A)_+$	1286 _s ± 10 _s	30 _s ± 20 _s	1.65 ± 0.04	$K\bar{K}\pi$ $\eta\pi\pi$ $[\delta\pi]$ $2\pi^+2\pi^-$ (prob. $\rho^0\pi^+\pi^-$)	seen seen seen seen	305 484 245 565	
$J^P = 0^-, 1^+, 2^-,$ with 1^+ favoured								

Meson Table (cont'd)

Name	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Mode	Partial decay mode		p or P _{max} (MeV/c) ^(b)
						Fraction (%)	p or P _{max} (%)	
						[Upper limits are 1 σ (%)]		
A ₂ (1310)	$1^-(2^+)_{+}$	1310 _{±5}	102 _{±5}	1.72 ±.13	$\rho\pi$ $\eta\pi$ $\omega\pi\pi$ $K\bar{K}$ $\eta'\pi$	70.9±1.8 15.0±1.2 9.3±1.9 4.7±0.5 <1	S=1.1* S=1.2*	411 529 354 428 279
E(1420)	$0^+(A)_{+}$	1416 _{±10}	60 _{±20}	2.01 ±.08	$K\bar{K}\pi$ $[K^*\bar{K} + \bar{K}^*K]$ $\eta\pi\pi$ $[\delta\pi]$	~ 40 ~ 20 ~ 60 possibly seen]		421 150 564 352
Not a well established resonance.								
f'(1514)	$0^+(2^+)_{+}$	1516 ±3	40 ±10	2.30 ±.06	$K\bar{K}$	only mode seen		572
For upper limits, see footnote (k)								
F ₁ (1540)	$1(A)_{-}$	1540 ±5	40 ±15	2.37 ±.06	$K^*\bar{K} + \bar{K}^*K$ 3π	seen possibly seen		321 737
Not a well established resonance.								
ρ' (1600)	$1^+(1^-)_{-}$	~ 1600	200-800	2.56	4π $[0\pi^+\pi^-]$ $\pi\pi\pi$ $K\bar{K}$	seen with $\pi^+\pi^-$ in S-wave] possibly seen < 8		738 573 788 629
Not a well established resonance. ^{fl}								
A ₃ (1640)	$1^-(2^-)_{+}$	~ 1640	~ 300	2.69 ±.49	$f\pi$			304
Broad enhancement in the $J^P = 2^- f\pi$ partial wave; not a well established resonance. ^{fl}								
ω (1675)	$0^-(3^-)_{-}$	1667 _{±10}	150 _{±20}	2.78 ±.25	$\rho\pi$ 3π 5π $[\omega\pi\pi]$	seen possibly seen possibly seen possibly seen]		646 806 778 615
g(1680) ^{fl}	$1^+(3^-)_{-}$	1690 _{±20}	180 _{±30}	2.86 ±.30	2π 4π (incl. $\pi\pi\rho, \rho\rho, A_2\pi, \omega\pi$) $K\bar{K}$ $K\bar{K}\pi$ (incl. $K^*\bar{K}$)	24±1 large small small		833 787 683 624
J^P, M and Γ from the 2π mode.								
h(2040)	$0^+(4^+)_{+}$	2040 ±20	193 ±50	4.16 ±.39	$\pi\pi$ $K\bar{K}$	seen seen		1010 890
See note (1) for possible heavier states.								
K^+ K^0	$1/2(0^-)$	493.71 497.70		0.244 0.248	See Stable Particle Table			
K^* (892)	$1/2(1^-)$	892.2 ±0.5	49.4 ±1.8	0.796 ±.044	$K\pi$ $K\pi\pi$ $K\gamma$	≈ 100 < 0.2 0.15±0.07		288 216 309
M and Γ from charged mode; $m^0 - m^\pm = 4.1 \pm 0.6$ MeV.								
κ (1250)	$1/2(0^+)_{+}$	1250 _{±100}	~ 450	1.56 ±.56	$K\pi$			
See note on $K\pi$ S wave ^{fl} .								
Q region	$1/2(A)_{-}$	1200 to 1400			$K\pi\pi$ $+[K^*\pi]$ $+[K\rho]$ $+[K(\pi\pi)]_{\ell=0}$	only mode seen large] seen] possibly seen]		
$J^P = 1^+$ is dominant contribution; not a well established resonance ^{fl} .								
K^* (1420)	$1/2(2^+)_{+}$	1421 _{±3}	108 _{±10}	2.02 ±.15	$K\pi$ $K^*\pi$ $K\rho$ $K\omega$ $K\eta$	56.1±2.6 30.9±2.1 6.6±1.7 4.5±1.7 2.0±2.0		616 415 316 305 482
See note (m).								

Meson Table (cont'd)

Name	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		p or Pmax (MeV/c) ^(b)
					Mode	Fraction (%) [Upper limits are 1 σ (%)]	
L(1770)	$1/2(A)$	1765 _s ± 10	140 _s ± 50	3.11 $\pm .25$	$K\pi\pi$ $K\pi\pi\pi$	dominant seen	788 757
Not a well established resonance [¶] . †[K*(1420) π and other subreactions [¶]]							
See note (1) for possible heavier states.							
J/ ψ (3100)	$0^-(1^-)$	3098 ± 3	0.067 $\pm .012$	9.6 $\pm .0$	e^+e^- $\mu^+\mu^-$ hadrons	7 \pm 1 7 \pm 1 86 \pm 2	1549 1545
†[identified hadron modes + [X(2750) possibly seen] ^{¶¶} ~ 12] ^{¶¶}							
ψ (3700)	$0^-(1^-)$	3684 ± 4	0.228 $\pm .056$	13.6 $\pm .0$	e^+e^- $\mu^+\mu^-$ hadrons	0.9 \pm .2 0.9 \pm .2 98.1 \pm .3	1842 1839
†[J/ ψ $\pi^+\pi^-$ 33 \pm 3] 474 †[J/ ψ $\pi^0\pi^0$ 17 \pm 2] 478 †[J/ ψ η 4.2 \pm .7] 189 †[$\gamma P_C, P_C \rightarrow J/\psi \gamma$ 3.6 \pm .7] ^{¶¶} †[$\gamma\chi$ (3410) seen] ^{¶¶} 264 †[$\gamma\chi$ (3530) seen] ^{¶¶} 151 †[other identified hadron modes ~ 0.5] ^{¶¶}							
ψ (4100)	(1^-)	~ 4100	~ 200	16.8 $\pm .8$	Broad enhancement in the e^+e^- total cross section; probably not a single resonance. ^{¶¶}		
ψ (4400)	(1^-)	4414 ± 7	33 ± 10	19.5 $\pm .1$	e^+e^-	.0013 \pm .0003	2207
†X(2750) †P _C (3300 or 3500) †X(3410) †X(3530)							
States observed in radiative decays of J/ ψ (3100) and ψ (3700). See Meson Data Card Listings for a compilation and discussion of the experimental data.							

(1) Contents of Meson Data Card Listings

Non-strange (Y = 0)				Strange (Y = 1)	
entry	$I^G(J^P)C_n$	entry	$I^G(J^P)C_n$	entry	I (J ^P)
π	$1^-(0^-)+$	A_1 (1100)	$1^-(1^+)+$	ρ' (1600)	$1^+(1^-)-$
η	$0^+(0^-)+$	$\rightarrow M$ (1150)		A_3 (1640)	$1^-(2^-)+$
ρ (770)	$1^+(1^-)-$	$\rightarrow A_{1,s}$ (1170)		ω (1675)	$0^-(3^-)-$
ω (783)	$0^-(1^-)-$	ϵ (1200)	$0^+(0^+)+$	g (1680)	$1^+(3^-)-$
$\rightarrow M$ (940)		B (1235)	$1^+(1^+)-$	$\rightarrow X$ (1690)	-
$\rightarrow M$ (953)		$\rightarrow \rho'$ (1250)	$1^+(1^-)-$	$\rightarrow X$ (1795)	1
η' (958)	$0^+(0^-)+$	f (1270)	$0^+(2^+)+$	$\rightarrow A_4$ (1900)	1^-
δ (970)	$1^-(0^+)+$	D (1285)	$0^+(A)+$	$\rightarrow S$ (1930)	1
$\rightarrow H$ (990)		A_2 (1310)	$1^-(2^+)+$	h (2040)	$0^+(4^+)+$
S^* (993)	$0^+(0^+)+$	E (1420)	$0^+(A)+$	$\rightarrow \rho$ (2100)	
ϕ (1020)	$0^-(1^-)-$	$\rightarrow X$ (1430)	0	$\rightarrow T$ (2200)	1
$\rightarrow M$ (1033)		$\rightarrow X$ (1440)	1	$\rightarrow U$ (2360)	1
$\rightarrow B_1$ (1040)		f' (1514)	$0^+(2^+)+$	$\rightarrow N\bar{N}$ (2375)	0
$\rightarrow \eta_N$ (1080)	$0^+(N)+$	F_1 (1540)	1 (A)	$\rightarrow X$ (2500-3600)	

New heavy mesons

J/ ψ (3100) ψ (3700) ψ (4100) ψ (4400) †X(2750) †P_C(3300 or 3500) †X(3410) †X(3530)

Meson Table (cont'd)

- + Indicates an entry in Meson Data Card Listings not entered in the Meson Table. We do not regard these as established resonances.
- ¶ See Meson Data Card Listings.
- * Quoted error includes scale factor $S = \sqrt{\chi^2/(N-1)}$. See footnote to Stable Particle Table.
- † Square brackets indicate a subreaction of the previous (unbracketed) decay mode(s).
- § This is only an educated guess; the error given is larger than the error of the average of the published values. (See Meson Data Card Listings for the latter.)
- (a) ΓM is approximately the half-width of the resonance when plotted against M^2 .
- (b) For decay modes into ≥ 3 particles, p_{\max} is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated by using the averaged central mass values, without taking into account the widths of the resonances.
- (c) From pole position $(M - i\Gamma/2)$.
- (d) The e^+e^- branching ratio is from $e^+e^- \rightarrow \pi^+\pi^-$ experiments only. The $\omega\rho$ interference is then due to $\omega\rho$ mixing only, and is expected to be small. See note in Meson Data Card Listings. The $\mu^+\mu^-$ branching ratio is compiled from 3 experiments; each possibly with substantial $\omega\rho$ interference. The error reflects this uncertainty; see notes in Meson Data Card Listings. If $e\mu$ universality holds, $\Gamma(\rho^0 \rightarrow \mu^+\mu^-) = \Gamma(\rho^0 \rightarrow e^+e^-) \times 0.99785$.
- (e) Empirical limits on fractions for other decay modes of $\rho(770)$ are $\pi^+\eta < 0.8\%$, $\pi^+\pi^+\pi^-\pi^- < 0.15\%$, $\pi^+\pi^+\pi^-\pi^0 < 0.2\%$.
- (f) Empirical limits on fractions for other decay modes of $\omega(783)$ are $\pi^+\pi^-\gamma < 5\%$, $\pi^0\pi^0\gamma < 1\%$, $\eta + \text{neutral}(s) < 1.5\%$, $\mu^+\mu^- < 0.02\%$, $\pi^0\mu^+\mu^- < 0.2\%$, $\eta\gamma < 0.5\%$.
- (g) Empirical limits on fractions for other decay modes of $\eta'(958)$: $\pi^+\pi^- < 2\%$, $\pi^+\pi^-\pi^0 < 5\%$, $\pi^+\pi^+\pi^-\pi^- < 1\%$, $\pi^+\pi^+\pi^-\pi^0 < 1\%$, $6\pi < 1\%$, $\pi^+\pi^+e^+e^- < 0.6\%$, $\pi^0e^+e^- < 1.3\%$, $\eta e^+e^- < 1.1\%$, $\pi^0\rho^0 < 4\%$, $\gamma\omega < 5\%$.
- (h) The mass and width are from the $\eta\pi$ mode only. If the $K\bar{K}$ channel is strongly coupled, the width may be 120 MeV or more.
- (i) Empirical limits on fractions for other decay modes of $\phi(1020)$ are $\pi^+\pi^- < 0.03\%$, $\pi^+\pi^-\gamma < 0.7\%$, $\omega\gamma < 5\%$, $\rho\gamma < 2\%$, $2\pi^+2\pi^-\pi^0 < 1\%$.
- (j) Empirical limits on fractions for other decay modes of $B(1235)$: $\pi\pi < 15\%$, $K\bar{K} < 2\%$, $4\pi < 50\%$, $\phi\pi < 1.5\%$, $\eta\pi < 25\%$, $(\bar{K}K)\pi^0 < 8\%$, $K_S K_S \pi^\pm < 2\%$, $K_S K_L \pi^\pm < 6\%$.
- (k) Empirical limits on fractions for other decay modes of $f'(1514)$ are $\pi^+\pi^- < 20\%$, $\eta\pi < 50\%$, $\eta\pi\pi < 30\%$, $K\bar{K}\pi + K^*\bar{K} < 35\%$, $2\pi^+2\pi^- < 32\%$.
- (l) Empirical limits on fractions for other decay modes of $f(1270)$ are $\eta\pi\pi < 1\%$, $K^0K^-\pi^+ + \text{c.c.} < 1\%$, $\eta\eta < 2\%$.
- (m) The tabulated mass of 1421 MeV comes from the $K\pi\pi$ mode; the $K\pi\pi$ mode can be contaminated with diffractively produced Q^\pm .

Established Nonets, and octet-singlet mixing angles from Appendix IIB, Eq. (2'). Of the two isosinglets, the "mainly octet" one is written first, followed by a semicolon.

$(J^P)C_n$	Nonet members	$\theta_{\text{lin.}}$	$\theta_{\text{quadr.}}$
$(0^-)^+$	$\pi, K, \eta; \eta'$	$-24 \pm 1^\circ$	$-11 \pm 1^\circ$
$(1^-)^-$	$\rho, K^*, \phi; \omega$	$37 \pm 1^\circ$	$40 \pm 1^\circ$
$(2^+)^+$	$A_2, K^*(1420), f'; f$	$29 \pm 2^\circ$	$31 \pm 2^\circ$

Baryon Table

April 1976

The following short list gives the status of all the Baryon States in the Data Card Listings. In addition to the status, the name, the nominal mass, and the quantum numbers (where known) are shown. States with three- or four-star status are included in the main Baryon Table; the others have been omitted because the evidence for the existence of the effect and/or for its interpretation as a resonance is open to considerable question.

N(939)	P11	****	$\Delta(1232)$	P33	****	$\Lambda(1116)$	P01	****	$\Sigma(1193)$	P11	****	$\Xi(1317)$	P11	****
N(1470)	P11	****	$\Delta(1650)$	S31	****	$\Lambda(1330)$	Dead		$\Sigma(1385)$	P13	****	$\Xi(1530)$	P13	****
N(1520)	D13	****	$\Delta(1670)$	D33	****	$\Lambda(1405)$	S01	****	$\Sigma(1440)$	Dead		$\Xi(1630)$		**
N(1535)	S11	****	$\Delta(1690)$	P33	*	$\Lambda(1520)$	D03	****	$\Sigma(1480)$	*		$\Xi(1820)$		***
N(1670)	D45	****	$\Delta(1890)$	F35	**	$\Lambda(1600)$	F01	*	$\Sigma(1580)$	D13	**	$\Xi(1940)$		***
N(1688)	F15	****	$\Delta(1900)$	S31	*	$\Lambda(1670)$	S01	****	$\Sigma(1620)$	S11	**	$\Xi(2030)$		**
N(1700)	S11	****	$\Delta(1910)$	P31	****	$\Lambda(1690)$	D03	****	$\Sigma(1660)$	P11	**	$\Xi(2250)$		*
N(1700)	D13	**	$\Delta(1950)$	F37	****	$\Lambda(1800)$	P01	**	$\Sigma(1670)$	D13	****	$\Xi(2500)$		**
N(1780)	P11	***	$\Delta(1960)$	D35	**	$\Lambda(1800)$	G09	*	$\Sigma(1670)$	**				
N(1810)	P13	***	$\Delta(2160)$	**	*	$\Lambda(1815)$	F05	****	$\Sigma(1690)$	**		$\Omega(1672)$	P03	****
N(1990)	F17	**	$\Delta(2420)$	H311	***	$\Lambda(1830)$	D05	****	$\Sigma(1750)$	S11	***			
N(2000)	F15	*	$\Delta(2850)$	***	*	$\Lambda(1860)$	P03	***	$\Sigma(1765)$	D15	****			
N(2040)	D13	**	$\Delta(3230)$	***	*	$\Lambda(1870)$	S01	**	$\Sigma(1770)$	P11	*			
N(2100)	S11	*			*	$\Lambda(2010)$	**	*	$\Sigma(1840)$	P13	*			
N(2100)	D45	*	Z0(1780)	P01	*	$\Lambda(2020)$	F07	*	$\Sigma(1880)$	P11	**			
N(2190)	G17	***	Z0(1865)	D03	*	$\Lambda(2100)$	G07	****	$\Sigma(1915)$	F15	****			
N(2220)	H19	****	Z1(1900)	P13	*	$\Lambda(2110)$	F05	**	$\Sigma(1940)$	D13	***			
N(2650)		****	Z1(2150)	*	*	$\Lambda(2350)$	****	*	$\Sigma(2000)$	S11	*			
N(3030)		****	Z1(2500)	*	*	$\Lambda(2585)$	****	*	$\Sigma(2030)$	F17	****			
N(3245)		*							$\Sigma(2070)$	F15	*			
N(3690)		*							$\Sigma(2080)$	P13	**			
N(3755)		*							$\Sigma(2100)$	G17	*			
									$\Sigma(2250)$	****				
									$\Sigma(2455)$	***				
									$\Sigma(2620)$	***				
									$\Sigma(3000)$	**				

 **** Good, clear, and unmistakable. *** Good, but in need of clarification or not absolutely certain.
 ** Needs confirmation. * Weak.

[See notes on N's and Δ 's, on possible Z^{*}'s, and on Y^{*}'s and Ξ ^{*}'s at the beginning of those sections in the Baryon Data Card Listings; also see notes on individual resonances in the Baryon Data Card Listings.]

Particle ^a	I (J ^P) ^a estab.	π or K Beam ^b p _{beam} (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass ^c M (MeV)	Full Width ^c Γ (MeV)	M ² $\pm\Gamma M$ ^b (GeV ²)	Partial decay mode		
						Mode	Fraction %	p or d P _{max} (MeV/c)
p	$1/2(1/2^+)$		938.3		0.880	See Stable Particle Table		
n	$1/2(1/2^+)$		939.6		0.883			
N(1470) ^g	$1/2(1/2^+)$ P' ₁₁	p = 0.66 $\sigma = 27.8$	1390 to 1470	180 to 220 (200)	2.16 ± 0.29	N π N η N $\pi\pi$ [N ϵ [$\Delta\pi$ [N ρ p γ ^f n γ ^f	~60 ~18 ~25 ~ 7] ^e ~19] ^e < 9] ^e 0.07-0.14 < 0.05	420 d 368 d 177 d 435 435
N(1520) ^g	$1/2(3/2^-)$ D' ₁₃	p = 0.74 $\sigma = 23.5$	1510 to 1530	110 to 150 (125)	2.31 ± 0.19	N π N $\pi\pi$ [N ϵ [N ρ [$\Delta\pi$ N η p γ ^f n γ ^f	~55 ~45 < 5] ^e ~15] ^e ~25] ^e < 1 0.4-0.7 0.3-0.6	456 410 d d 228 d 471 471
N(1535) ^g	$1/2(1/2^-)$ S' ₁₁	p = 0.76 $\sigma = 22.5$	1500 to 1530	50 to 120 (100)	2.36 ± 0.15	N π N η N $\pi\pi$ [N ρ [N ϵ [$\Delta\pi$ p γ ^f n γ ^f	~30 ~65 ~ 5 ~ 3] ^e ~ 2] ^e ~ 1] ^e < 0.4 < 0.4	467 182 422 d d 243 481 481

Baryon Table (cont'd)

Particle ^a	I (J ^P) ^a estab.	$\frac{\pi \text{ or } K \text{ Beam}^b}{p_{\text{beam}}(\text{GeV}/c)}$ $\sigma = 4\pi\lambda^2$ (mb)	Mass M ^c (MeV)	Full Width Γ^c (MeV)	M ² $\pm\Gamma M^b$ (GeV ²)	Partial decay mode		
						Mode	Fraction %	p or P _{max} ^d (MeV/c)
N(1670) ^g	<u>1/2(5/2⁻)D₁₅</u>	p = 1.00 $\sigma = 15.6$	1660 to 1685	145 to 165 (155)	2.79 ± 0.26	N π N $\pi\pi$ [$\Delta\pi$ ΛK N η p γ^f n γ^f	~45 ~55 ~50] ^e <0.3 <0.5 <0.03 <0.14	560 525 360 200 368 572 572
N(1688) ^g	<u>1/2(5/2⁺)F₁₅</u>	p = 1.03 $\sigma = 14.9$	1670 to 1690	120 to 145 (140)	2.85 ± 0.24	N π N $\pi\pi$ [N ϵ [N ρ [$\Delta\pi$ N η p γ^f n γ^f	~60 ~40 ~14] ^e ~14] ^e ~11] ^e <0.3 0.1-0.4 <0.03	572 538 340 d 375 388 583 583
N(1700) ^g	<u>1/2(1/2⁻)S₁₁</u>	p = 1.05 $\sigma = 14.3$	1660 to 1690	100 to 200 (150)	2.89 ± 0.26	N π N $\pi\pi$ [N ϵ [N ρ [$\Delta\pi$ ΛK ΣK p γ^f n γ^f	~55 ~30 ~10] ^e ~7] ^e ~4] ^e ~4 ~2 <0.1 <0.15	580 547 355 d 385 250 109 591 591
N(1780)	<u>1/2(1/2⁺)P₁₁</u>	p = 1.20 $\sigma = 12.2$	1700 to 1800	100 to 250 (200)	3.17 ± 0.36	N π N $\pi\pi$ [N ϵ [N ρ [$\Delta\pi$ ΛK ΣK N η p γ^f n γ^f	~20 >40 15-40] ^e 20-50] ^e 10-20] ^e ~7 ~10 2-20 ^h <0.15 <0.13	633 603 440 249 448 353 267 476 643 643
N(1810)	<u>1/2(3/2⁺)P₁₃</u>	p = 1.26 $\sigma = 11.5$	1700 to 1850	100 to 300 (200)	3.28 ± 0.36	N π N $\pi\pi$ [N ρ ΛK ΣK N η p γ^f n γ^f	~20 ~70 ~70] ^e ~5 ~2 <5 <0.2 <0.2	652 624 297 386 307 503 661 661
N(2190)	<u>1/2(7/2⁻)G₁₇</u>	p = 2.07 $\sigma = 6.21$	2100 to 2250	150 to 300 (250)	4.80 ± 0.55	N π ΛK ΣK	15-35 <0.2 <0.2	888 710 664
N(2220)	<u>1/2(9/2⁺)H₁₉</u>	p = 2.14 $\sigma = 5.97$	2200 to 2250	250 to 350 (300)	4.93 ± 0.67	N π	~20	905
N(2650)	<u>1/2(?)</u>	p = 3.26 $\sigma = 3.67$	~2650	~350 (350)	7.02 ± 0.93	N π	(J+1/2) _x <0.4] ^j	1154
N(3030)	<u>1/2(?)</u>	p = 4.41 $\sigma = 2.62$	~3030	~400 (400)	9.18 ± 1.21	N π	(J+1/2) _x <0.1] ^j	1366
$\Delta(1232)^g$	<u>3/2(3/2⁺)P₃₃</u>	p = 0.30 $\sigma = 94.3$	1230 to 1234	110 to 120 (115)	1.52 ± 0.14	N π N $\pi^+\pi^-$ p γ^f	~99.4 ~0 0.58-0.66	227 80 259
		$\Delta(++)$ Pole position: ^k M-i $\Gamma/2 = (1211.0\pm 0.8) - i(49.9\pm 0.6)$						
		$\Delta(0)$ Pole position: ^k M-i $\Gamma/2 = (1210.9\pm 1.0) - i(53.1\pm 1.0)$						
$\Delta(1650)^g$	<u>3/2(1/2⁻)S₃₁</u>	p=0.96 $\sigma = 16.4$	1615 to 1695	140 to 200 (140)	2.72 ± 0.23	N π N $\pi\pi$ [N ρ [$\Delta\pi$ p γ^f	~35 ~65 10-25] ^e ~50] ^e <0.25	547 511 d 344 558

Baryon Table (cont'd)

Particle ^a	I	(J ^P) ^a estab.	π or K Beam ^b $\frac{p_{beam}}{\sigma} \frac{(GeV/c)}{\lambda^2}$ (mb)	Mass M ^c (MeV)	Full Width Γ^c (MeV)	M ² $\pm \Gamma M^b$ (GeV ²)	Partial decay mode		
							Mode	Fraction %	p or P _{max} ^d (MeV/c)
$\Delta(1670)^g$		$3/2(3/2^-)D_{33}$	p = 1.00 $\sigma = 15.6$	1650 to 1720	190 to 260 (200)	2.79 ± 0.33	N π N $\pi\pi$ [N ρ [$\Delta\pi^f$ p γ^f	~ 15 ~ 85 30-60] ^e ~ 45] ^e 0.05-0.3	560 525 d 361 572
$\Delta(1890)^g$		$3/2(5/2^+)F_{35}$	p = 1.42 $\sigma = 9.88$	1860 to 1900	150 to 300 (250)	3.57 ± 0.47	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ΣK p γ^f	~ 15 ~ 80 ~ 60] ^e 10-30] ^e <3 <0.1	704 677 403 531 400 712
$\Delta(1910)^g$		$3/2(1/2^+)P_{31}$	p = 1.46 $\sigma = 9.54$	1780 to 1950	160 to 230 (200)	3.65 ± 0.38	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ΣK p γ^f	15-35 ? small] ^e small] ^e 2-20 <0.1	716 691 429 545 420 725
$\Delta(1950)^g$		$3/2(7/2^+)F_{37}$	p = 1.54 $\sigma = 8.90$	1910 to 1940	200 to 240 (220)	3.80 ± 0.43	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ΣK p γ^f	~ 40 >25 ~ 10] ^e ~ 20] ^e <1 0.09-0.15	741 716 471 574 460 749
$\Delta(2420)^g$		$3/2(11/2^+)H_{311}$	p = 2.64 $\sigma = 4.68$	2380 to 2450	300 to 500 (300)	5.86 ± 0.73	N π	10-15	1023
$\Delta(2850)$		$3/2(?^+)$	p = 3.85 $\sigma = 3.05$	2800 to 2900	~ 400 (400)	8.12 ± 1.14	N π	(J+1/2) _x ~ 0.25]	1266
$\Delta(3230)$		$3/2(?)$	p = 5.08 $\sigma = 2.25$	3200 to 3350	~ 440 (440)	10.43 ± 1.42	N π	(J+1/2) _x ~ 0.05]	1475
Z* Evidence for states with strangeness +1 is controversial. See the Baryon Data Card listings for discussion and display of data.									
Λ		$0(1/2^+)$		1115.6		1.245	See Stable Particle Table		
$\Lambda(1405)$		$0(1/2^-)S'_{01}$	below K ⁻ p threshold	1405 $\pm 5^l$	40 $\pm 10^l$ (40)	1.97 ± 0.06	$\Sigma\pi$	100	142
$\Lambda(1520)$		$0(3/2^-)D'_{03}$	p = 0.389 $\sigma = 84.5$	1519 $\pm 2^l$	15 $\pm 2^l$ (15)	2.31 ± 0.02	N \bar{K} $\Sigma\pi$ $\Delta\pi\pi$ $\Sigma\pi\pi$	46 ± 1 42 ± 1 10 ± 1 0.9 ± 0.1	234 258 250 140
$\Lambda(1670)$		$0(1/2^-)S''_{01}$	p = 0.74 $\sigma = 28.5$	1660 to 1680	20 to 60 (40)	2.79 ± 0.07	N \bar{K} $\Delta\eta$ $\Sigma\pi$	15-35 15-35 20-60	410 64 393
$\Lambda(1690)$		$0(3/2^-)D''_{03}$	p = 0.78 $\sigma = 26.1$	1690 $\pm 10^l$	30 to 80 (60)	2.86 ± 0.10	N \bar{K} $\Sigma\pi$ $\Delta\pi\pi$ $\Sigma\pi\pi$	20-30 15-40 ~ 25 ~ 20	429 409 415 352
$\Lambda(1815)$		$0(5/2^+)F'_{05}$	p = 1.05 $\sigma = 16.7$	1820 $\pm 5^l$	70 to 100 (85)	3.29 ± 0.15	N \bar{K} $\Sigma\pi$ $\Sigma(1385)\pi$	~ 60 ~ 12 15-20	542 508 362
$\Lambda(1830)$		$0(5/2^-)D_{05}$	p = 1.09 $\sigma = 15.8$	1810 to 1840	60 to 110 (95)	3.35 ± 0.17	N \bar{K} $\Sigma\pi$ $\Delta\eta$	<10 35-75 <4	554 519 367
$\Lambda(1860)$		$0(1/2^+)P_{03}$	p = 1.14 $\sigma = 14.7$	1860 to 1910	40 to 110 (80)	3.46 ± 0.15	N \bar{K} $\Sigma\pi$	15-35 5-10	576 534
$\Lambda(2100)$		$0(7/2^-)G_{07}$	p = 1.68 $\sigma = 8.68$	2100 to 2120	150 to 300 (250)	4.41 ± 0.53	N \bar{K} $\Sigma\pi$ $\Delta\eta$ ΞK $\Lambda\omega$	~ 30 ~ 5 <3 <3 <8	748 699 617 483 443

Baryon Table (cont'd)

Particle ^a	I	(J ^P) ^a estab.	π or K Beam ^b		Mass M ^c (MeV)	Full Width Γ ^c (MeV)	M ² ±ΓM ^b (GeV ²)	Partial decay mode		
			p _{beam} (GeV/c)	σ = 4πλ ² (mb)				Mode	Fraction %	p or d p _{max} (MeV/c)
Λ(2350)	0	(?)	p = 2.29 σ = 5.85		2340 to 2360	100 to 200 (120)	5.52 ±0.28	N \bar{K} Σπ	(J+1/2) _x ~0.9 _j seen	913 865
Λ(2585)	0	(?)	p = 2.91 σ = 4.37		~2585	~300 (300)	6.68 ±0.78	N \bar{K}	(J+1/2) _x ~1.0 _j	1058
Σ		1(1/2 ⁺)			(+)1189.4 (0)1192.5 (-)1197.4		1.415 1.422 1.434	See Stable Particle Table		
Σ(1385)		1(3/2 ⁺)P ₁₃ ⁱ	below K ⁻ p threshold		(+)1382.5±0.5 S=1.2 ^m (-)1386.6±1.2 S=2.3 ^m	(+)35±2 S=1.9 ^m (-)42±4 S=3.2 ^m (35)	1.92 ±0.05	Λπ Σπ	88±2 12±2	208 117
Σ(1670) ⁿ		1(3/2 ⁻)D ₁₃ ⁱⁱ	p = 0.74 σ = 28.5		1670 ±10 ^l	35 to 70 (50)	2.79 ±0.08	N \bar{K} Σπ Λπ	10-25 20-60 <20	410 387 447
Σ(1750)		1(1/2 ⁻)S ₁₁ ⁱⁱⁱ	p = 0.91 σ = 20.7		1700 to 1790	50 to 120 (75)	3.06 ±0.13	N \bar{K} Λπ Σπ Ση	10-40 5-20 <18 15-55	483 507 450 54
Σ(1765)		1(5/2 ⁻)D ₁₅ ^{iv}	p = 0.94 σ = 19.6		1723 ±7 ^l	110 to 150 (130)	3.12 ±0.23	N \bar{K} Λπ Λ(1520)π Σ(1385)π Σπ	~41 ~14 ~16 ~10 ~1	496 518 187 315 461
Σ(1915) ^g		1(5/2 ⁺)F ₁₅ ^v	p = 1.25 σ = 13.0		1905 to 1930	70 to 140 (100)	3.67 ±0.19	N \bar{K} Λπ Σπ	5-15 20 ?	612 619 568
Σ(1940) ⁱ		1(3/2 ⁻)D ₁₃ ^{vi}	p = 1.32 σ = 12.0		1900 to 1960	110 to 280 (220)	3.76 ±0.43	N \bar{K} Λπ Σπ	<20 ~4 ~7	678 680 589
Σ(2030) ^g		1(7/2 ⁺)F ₁₇ ^{vii}	p = 1.52 σ = 9.93		2020 to 2040	120 to 200 (180)	4.12 ±0.37	N \bar{K} Λπ Σπ ΞK	~20 ~20 5-10 <2	700 700 652 412
Σ(2250)		1(?)	p = 2.04 σ = 6.76		2200 to 2300	50 to 200 (150)	5.06 ±0.34	N \bar{K} Λπ Σπ	(J+1/2) _x ~0.3 _j seen seen	849 841 801
Σ(2455)		1(?)	p = 2.57 σ = 5.09		~2455	~120 (120)	6.03 ±0.29	N \bar{K}	(J+1/2) _x ~0.2 _j	979
Σ(2620)		1(?)	p = 2.95 σ = 4.30		~2600	~200 (200)	6.86 ±0.52	N \bar{K}	(J+1/2) _x ~0.3 _j	1064
Ξ		1/2(1/2 ⁺)			(0)1314.9 (-)1321.3		1.729 1.746	See Stable Particle Table		
Ξ(1530) ^o		1/2(3/2 ⁺)P ₁₃			(0)1531.8±0.3 S=1.3 ^m (-)1535.1±0.6	(0) 9.1±0.5 (-) 10.1±1.9 (10)	2.34 ±0.02	Ξπ	100	144
Ξ(1820) ^{o,P}		1/2(?)			1800 to 1850	12 to 100 (60)	3.31 ±0.11	Λ \bar{K} Σ \bar{K} Ξπ Ξ(1530)π	seen seen seen seen	396 306 413 234
Ξ(1940) ^{o,q}		1/2(?)			1900 to 1970	30 to 140 (90)	3.76 ±0.17	Ξπ Ξ(1530)π	seen seen	499 336
Ω ⁻		0(3/2 ⁺)			1672.2		2.796	See Stable Particle Table		

Baryon Table (*cont'd*)

- For convenience all Baryon States for which information exists in the Baryon Data Card Listings are listed at the beginning of the Baryon Table. States with only a one or two star (*) rating in that list have been omitted from the main Baryon Table; each omitted state is indicated by an arrow in the left-hand margin of the Table. In the Listings there is an arrow under the name of each state omitted from the Table.
- a. The names of the Baryon States in Col. 1 [such as N(1470)] contain a nominal mass which is a rounded average of the reported values in the Data Card Listings. The convention for using primes in the spectroscopic notation for the quantum numbers in Col. 2 [such as P_{11}^{\prime}] is as follows: no prime is attached when the Data Card Listings include only one resonance in the given partial wave; when there is more than one resonance the first has been designated with a prime, the second with a double prime, etc. The name and the quantum numbers for each state are also given in large print at the beginning of the Data Card Listings for that state.
- b. The numbers in Col. 3 and Col. 6 are calculated using the nominal mass (see a. above) for M and the nominal width (see c. below) for Γ .
- c. For M and Γ of most baryons we report here an interval instead of an average. Averages are appropriate if each result is based on independent measurements, but inappropriate where the spread in parameters arises because different models or procedures have been applied to a common set of data. A single value with an approximation sign (~) indicates that there is not enough data to give a meaningful interval. A nominal width is included in parentheses in Col. 5; this nominal width is used to calculate the value of ΓM given in Col. 6.
- d. For two body decay modes we give the momentum, p, of the decay products in the decaying baryon rest frame. For decay modes into ≥ 3 particles we give the maximum momentum, p_{\max} , that any of the particles in the final state can have in this frame. The momenta are calculated using the nominal mass (see a. above) of the decaying baryon, and of any isobars in the final state. Some decays which would be energetically forbidden for the nominal masses actually occur because of the finite widths of the decaying Baryon and/or isobars in the final state. In these cases, the decay momentum is omitted from Col. 9 and replaced with a reference to this footnote.
- e. Square brackets around an isobar decay mode indicate that it is a sub-reaction of the previous unbracketed decay mode. In the case of N^* and Δ decays into isobar modes we have used the isobar model results of LONGACRE 75 in addition to other data from the listings (where available) to estimate the branching fractions.
- f. The tabulated radiative fractions involve a sum over two helicities (1/2, 3/2). In the case of $I = 1/2$ resonances, there are two distinct isospin couplings, whence γ_p and γ_n . For conventions and further details, see the Mini-Review preceding the Baryon Data Card Listings.
- g. Only information coming from partial-wave analyses has been used here. For the production experiments results see the Baryon Data Card Listings.
- h. The range given here does not include the branching ratio of approximately 80% reported by FELTESSE 75.
- i. There may be more than one state in this region. The only analysis which reports an elastic coupling (LEA 73) also finds unusually low mass and width values. The inelastic branching fractions quoted here are based on an elasticity of 10%, which is a compromise between LEA 73 and RLIC 76.
- j. This state has been seen only in an energy-dependent fit to total, channel, or fixed angle cross-section data. J is not known; x is Γ_{e1}/Γ .
- k. See note on determination of resonance parameters in the Baryon Data Card Listings. Values of mass and width are dependent upon resonance shape used to fit the data. The pole position is much less dependent upon the parametrization used. The pole positions given here are taken from results (in the Data Card Listings) of fits to the phase shifts of CARTER 73 without Coulomb corrections.
- l. The error given here is only an educated guess; it is larger than the error of the average of the published values (see the Baryon Data Card Listings for the latter).
- m. Quoted error includes an S (scale) factor. See first footnote to Stable Particle Table.
- n. In this energy region the situation is still confused. In addition to the effect at ~ 1670 MeV seen in both production and formation experiments, recent formation experiments have found evidence for fairly narrow S_{11} and/or P_{11} states at 1620-1660 MeV. A narrow bump in the $I = 1$ $\bar{K}N$ total cross section has also been seen recently at ~ 1590 MeV. It is not clear how many states really exist here. No one has reported a strong coupling of any of these states to $\bar{K}N$ but there is much disagreement about branching ratios into $\pi\Lambda$ and $\pi\Sigma$. See the mini-reviews preceding the $\Sigma(1620)$ and $\Sigma(1670)$ Data Card Listings for more information.
- o. Only $\Xi(1530)$ is firmly established; information on the other states comes from experiments that have poor statistics due to the fact that the cross sections for $S = -2$ states are very low. For Ξ states, because of the meager statistics, we lower our standards and tabulate resonant effects if they have at least a four-standard-deviation statistical significance and if they are seen by more than one group. See the Baryon Data Card Listings for the other states.
- p. All four decay modes shown have been seen. Branching ratios are not quoted because there may be more than one state here.
- q. This bump has been seen in both final states shown; it is not clear if one, or more, states are present.

PHYSICAL AND NUMERICAL CONSTANTS*

PHYSICAL CONSTANTS

		Uncert. (ppm)
N	= 6.0220943(63)×10 ²³ mole ⁻¹	1.05
V _m	= 22413.83(70) cm ³ mole ⁻¹ = molar volume of ideal gas at STP	31
c	= 2.99792458(1.2)×10 ¹⁰ cm sec ⁻¹	0.004
e	= 4.803242(14)×10 ⁻¹⁰ esu = 1.6021892(46)×10 ⁻¹⁹ coulomb	2.9; 2.9
1 MeV	= 1.6021892(46)×10 ⁻⁶ erg	2.9
ħ=h/2π	= 6.582173(17)×10 ⁻²² MeV sec = 1.0545887(57)×10 ⁻²⁷ erg sec	2.6; 5.4
ħc	= 1.9732858(51)×10 ⁻¹¹ MeV cm = 197.32858(51) MeV Fermi	2.6; 2.6
	= 0.6240078(16) GeV mb ^{1/2}	2.6
α	= e ² /ħc = 1/137.035982(30)	0.22
k _{Boltzmann}	= 1.380662(44)×10 ⁻¹⁶ erg °K ⁻¹	32
	= 8.61735(28)×10 ⁻¹¹ MeV °K ⁻¹ = 1 eV/11604.50(36) °K	32; 32
m _e	= 0.5110034(14) MeV = 9.109534(47)×10 ⁻³¹ kg	2.8; 5.1
m _p	= 938.2796(27) MeV = 1836.15152(70) m _e = 6.72270(31) m _{π±}	2.8; 0.38; 46
	= 1.007276470(11) amu	0.011
1 amu	= 1/12 m _{C12} = 931.5016(26) MeV	2.8
m _d	= 1875.628(5) MeV	3
r _e	= e ² /m _e c ² = 2.8179380(70) fermi (1 fermi = 10 ⁻¹³ cm)	2.5
λ _e	= ħ/m _e c = r _e α ⁻¹ = 3.8615905(64)×10 ⁻¹¹ cm	1.6
a _{∞Bohr}	= ħ ² /m _e e ² = r _e α ⁻² = 0.52917706(44)Å (1Å = 10 ⁻⁸ cm)	0.82
σ _{Thomson}	= (8/3)πr _e ² = 0.6652448(33)×10 ⁻²⁴ cm ² (10 ⁻²⁴ cm ² = 1 barn)	4.9
μ _{Bohr}	= eħ/2m _e c = 0.57883785(95)×10 ⁻¹⁴ MeV gauss ⁻¹	1.6
μ _p	= eħ/2m _p c = 3.1524515(53)×10 ⁻¹⁸ MeV gauss ⁻¹	1.7
μ _p /μ _{Bohr}	= 1.520993136(21)	0.014
1/2ω _{cyclotron}	= e/2m _e c = 8.794023(25)×10 ⁶ rad sec ⁻¹ gauss ⁻¹	2.8
1/2ω _p _{cyclotron}	= e/2m _p c = 4.789378(13)×10 ³ rad sec ⁻¹ gauss ⁻¹	2.8
Hydrogen-like atom (nonrelativistic, μ = reduced mass):		
	$\frac{v}{c}$ _{rms} = $\frac{ze^2}{n\hbar c}$, E _n = $\frac{\mu v^2}{2} = \frac{\mu z^2 e^4}{2(n\hbar)^2}$, a _n = $\frac{n^2 \hbar^2}{\mu z e^2}$	
R _∞ = m _e e ⁴ /2ħ ²	= m _e c ² α ² /2 = 13.605804(36) eV (Rydberg)	2.6
	= m _e cα ² /2ħ = 109737.3143(10) cm ⁻¹	0.009
pc = 0.3 H _p (MeV, kilogauss, cm)		
1 year (sidereal)	= 365.256 days = 3.1558×10 ⁷ sec (≈π×10 ⁷ sec)	
density of dry air	= 1.205 mg cm ⁻³ (at 20°C, 760 mm)	
acceleration by gravity	= 980.62 cm sec ⁻² (sea level, 45°)	
gravitational constant	= 6.6732(31)×10 ⁻⁸ cm ³ g ⁻¹ sec ⁻²	
1 calorie (thermochemical)	= 4.184 joules	
1 atmosphere	= 1033.2275 dynes cm ⁻² = 1.01325 bar	
1 eV per particle	= 11604.50(36) °K (from E = kT)	

NUMERICAL CONSTANTS

π	= 3.1415927	1 rad	= 57.2957795 deg	√π	= 1.7724539
e	= 2.7182818	1/e	= 0.3678794	√2	= 1.4142136
ln2	= 0.6931472	ln10	= 2.3025851	√3	= 1.7320508
log ₁₀ 2	= 0.3010300	log ₁₀ e	= 0.4342945	√10	= 3.1622777

*Prepared by Stanley J. Brodsky, based mainly on the adjustment of the fundamental physical constants by E. R. Cohen and B. N. Taylor, *J. Phys. Chem. Ref. Data* **2**, 663 (1973), plus current values for N, α, μ_p/μ_{Bohr}, R_∞ [see B. N. Taylor and E. R. Cohen, *Proceedings of the Fifth International Conference on Atom Masses and Fundamental Constants (AMCO-5), Paris, 1975*]. The figures in parentheses correspond to the one-standard-deviation uncertainty in the last digits of the main number. (Updated April 1976.)

CLEBSCH-GORDAN COEFFICIENTS, SPHERICAL HARMONICS, AND d FUNCTIONS

Note: A $\sqrt{\quad}$ is to be understood over every coefficient; e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation:

J	J	...
M	M	...
m_1	m_2	
m_1	m_2	Coefficients
.	.	.
.	.	.

$1/2 \times 1/2$	1	0	0
$+1/2 + 1/2$	1	0	0
$+1/2 - 1/2$	$1/2$	$1/2$	1
$-1/2 + 1/2$	$1/2$	$-1/2$	-1
$-1/2 - 1/2$	1	0	0

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

$$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$$

$$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$$

$2 \times 1/2$	$5/2$	$3/2$
$+2$	$1/2$	1
$+2 - 1/2$	$1/5$	$4/5$
$+1 + 1/2$	$4/5 - 1/5$	$5/2$
$0 + 1/2$	$2/5$	$3/5$
$0 + 1/2$	$3/5 - 2/5$	$5/2$
$0 - 1/2$	$3/5$	$2/5$
$-1 + 1/2$	$2/5 - 3/5$	$5/2$
$-1 - 1/2$	$3/5$	$3/2$

$1 \times 1/2$	$3/2$	$1/2$
$+1 + 1/2$	$3/2$	$1/2$
$+1 - 1/2$	$1/3$	$2/3$
$0 + 1/2$	$2/3 - 1/3$	$3/2$
$0 + 1/2$	$1/3$	$1/3$
$0 - 1/2$	$2/3$	$1/3$
$-1 + 1/2$	$1/3 - 2/3$	$-3/2$
$-1 - 1/2$	1	1

$3/2 \times 1/2$	2	1
$+3/2$	$1/2$	1
$+3/2 - 1/2$	$1/4$	$3/4$
$+1/2 + 1/2$	$3/4 - 1/4$	2
$+1/2 + 1/2$	$1/2$	$1/2$
$-1/2 + 1/2$	$1/2$	$-1/2$
$-1/2 - 1/2$	$3/4$	$1/4$
$-3/2 + 1/2$	$1/4 - 3/4$	-2
$-3/2 - 1/2$	1	1

2×1	3	2
$+2 + 1$	1	2
$+2 + 1$	$1/3$	$2/3$
$+1 + 1$	$2/3 - 1/3$	3
$+1 + 1$	$1/3$	$1/3$
$0 - 1/2$	$2/3$	$1/3$
$-1 + 1/2$	$1/3 - 2/3$	$-3/2$
$-1 - 1/2$	1	1

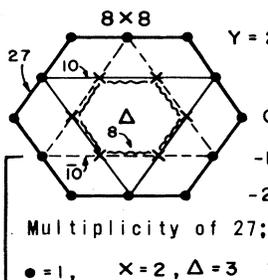
$3/2 \times 1$	$5/2$	$3/2$
$+3/2$	1	1
$+3/2$	0	$2/5$
$+1/2 + 1$	$3/5$	$3/5$
$+1/2 + 1$	$3/5 - 2/5$	$5/2$
$+1/2 + 1$	$1/5$	$1/2$
$0 - 1/2$	$2/5$	$3/2$
$-1/2 + 1/2$	$3/5$	$1/2$
$-1/2 + 1/2$	$1/5 - 1/3$	$5/2$
$-1/2 + 1/2$	$1/10$	$2/5$
$-1/2 + 1$	$3/5$	$1/15$
$-1/2 + 1$	$3/10$	$-8/15$
$-1/2 + 1$	$3/10$	$1/6$
$-1/2 + 1$	$1/5$	$1/2$
$-1/2 + 1$	$3/5$	$1/15$
$-1/2 + 1$	$3/10$	$-8/15$
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$-1/2 + 1$	$1/5$	$1/2$

SU(3) ISOSCALAR FACTORS

Adapted from J. J. de Swart, Rev. Mod. Phys. 35, 916 (1963)

The convention used here is: baryon first, meson second.

$$\{8\} \otimes \{8\} = \{27\} \oplus \{10\} \oplus \{10^*\} \oplus \{8\}_1 \oplus \{8\}_2 \oplus \{1\}.$$



* Five single-coefficient tables are omitted. The one involving a $\{10^*\}$ has a negative coefficient, i.e. $(NK|10^*) = -1$. The others, involving $\{27\}$ and $\{10\}$, are all $+1$.

		$Y=1, I=1/2, N$				$Y=1, I=3/2, \Delta$	
$\xi_1 \downarrow$		27	8 _D	8 _F	10*	27	10
N π		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	-1/2	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	-1/2	$\sqrt{2}/2$	$\sqrt{2}/2$
N η		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	1/2		
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	-1/2		

		$Y=0, I=0, \Lambda$				$Y=0, I=1, \Sigma$				
$\xi_1 \downarrow$		27	8 _D	1	8 _F	27	8 _D	8 _F	10	10*
N \bar{K}		$\sqrt{15}/10$	$\sqrt{10}/10$	1/2	$\sqrt{2}/2$	$\sqrt{5}/5$	$-\sqrt{30}/10$	$\sqrt{6}/6$	$-\sqrt{6}/6$	$-\sqrt{6}/6$
ΣK		$-\sqrt{15}/10$	$-\sqrt{10}/10$	-1/2	$\sqrt{2}/2$	$\sqrt{5}/5$	$-\sqrt{30}/10$	$-\sqrt{6}/6$	$\sqrt{6}/6$	$-\sqrt{6}/6$
$\Sigma \pi$		$-\sqrt{10}/20$	$-\sqrt{15}/5$	$\sqrt{6}/4$	0	0	0	$\sqrt{6}/3$	$\sqrt{6}/6$	$-\sqrt{6}/6$
$\Delta \eta$		$3\sqrt{30}/20$	$-\sqrt{5}/5$	$-\sqrt{2}/4$	0	$\sqrt{30}/10$	$\sqrt{5}/5$	0	1/2	1/2
$\Lambda \pi$						$\sqrt{30}/10$	$\sqrt{5}/5$	0	-1/2	-1/2

		$Y=-1, I=1/2, \Xi$				$Y=-1, I=3/2$	
$\xi_1 \downarrow$		27	8 _D	8 _F	10	27	10*
$\Xi \pi$		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	1/2	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΞK		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	1/2	$\sqrt{2}/2$	$\sqrt{2}/2$
$\Xi \eta$		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	1/2		
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	-1/2		

The phase factor $\xi_1 = \pm 1$, from de Swart's Table I, enters in his symmetry formula (14. 3):

$$(\mu_1 \mu_2 | \mu) = \xi_1 (-1)^{I_1 + I_2 - I} (\mu_2 \mu_1 | \mu).$$

This factor is irrelevant if you are doing your own self-consistent calculations; it enters when you try to check someone else who chose $\mu_2 \otimes \mu_1$ instead of $\mu_1 \otimes \mu_2$.

$$\{10\} \otimes \{8\} = \{35\} \oplus \{27\} \oplus \{10\} \oplus \{8\}.$$

* Four single coefficient tables are omitted; only the $\{27\}$ is -1; the three with $\{35\}$ are +1.

		$Y=1, I=1/2, N$		$Y=1, I=3/2, \Delta$		
$\xi_1 \downarrow$		27	8	35	27	10
$\Delta \pi$		$-\sqrt{5}/5$	$-2\sqrt{5}/5$	1/4	$-\sqrt{5}/4$	$\sqrt{10}/4$
ΣK		$-2\sqrt{5}/5$	$\sqrt{5}/5$	$\sqrt{5}/4$	3/4	$\sqrt{2}/4$
				$\sqrt{10}/4$	$-\sqrt{2}/4$	-1/2

		$Y=0, I=0, \Lambda$		$Y=0, I=1, \Sigma$				$Y=0, I=2$	
$\xi_1 \downarrow$		27	8	35	27	10	8	35	27
$\Sigma \pi$		$-\sqrt{10}/5$	$-\sqrt{15}/5$	$\sqrt{3}/6$	$-3\sqrt{5}/10$	$\sqrt{3}/3$	$-\sqrt{30}/15$	$\sqrt{3}/2$	-1/2
ΞK		$-\sqrt{15}/5$	$\sqrt{10}/5$	$\sqrt{2}/2$	$\sqrt{30}/10$	0	$-\sqrt{5}/5$	-1/2	$\sqrt{3}/2$
$\Sigma \eta$				$\sqrt{3}/3$	$-\sqrt{5}/5$	$-\sqrt{3}/3$	$\sqrt{30}/15$		
ΔK				$\sqrt{3}/6$	$\sqrt{5}/10$	$\sqrt{3}/3$	$2\sqrt{30}/15$		

		$Y=-1, I=1/2, \Xi$				$Y=-1, I=3/2$	
$\xi_1 \downarrow$		35	27	10	8	35	27
$\Xi \pi$		1/4	$-7\sqrt{5}/20$	$\sqrt{2}/4$	$-\sqrt{5}/5$	$\sqrt{2}/2$	$-\sqrt{2}/2$
$\Xi \eta$		3/4	$3\sqrt{5}/20$	$-\sqrt{2}/4$	$-\sqrt{5}/5$	$\sqrt{2}/2$	$\sqrt{2}/2$
ΩK		$\sqrt{2}/4$	$-3\sqrt{10}/20$	-1/2	$\sqrt{10}/5$	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		1/2	$\sqrt{5}/10$	$\sqrt{2}/2$	$\sqrt{5}/5$	$\sqrt{2}/2$	$\sqrt{2}/2$

		$Y=-2, I=0, \Omega^-$		$Y=-2, I=1$	
$\xi_1 \downarrow$		35	10	35	27
$\Omega \eta$		$\sqrt{2}/2$	$-\sqrt{2}/2$	1/2	$-\sqrt{3}/2$
ΞK		$\sqrt{2}/2$	$\sqrt{2}/2$	$\sqrt{3}/2$	1/2

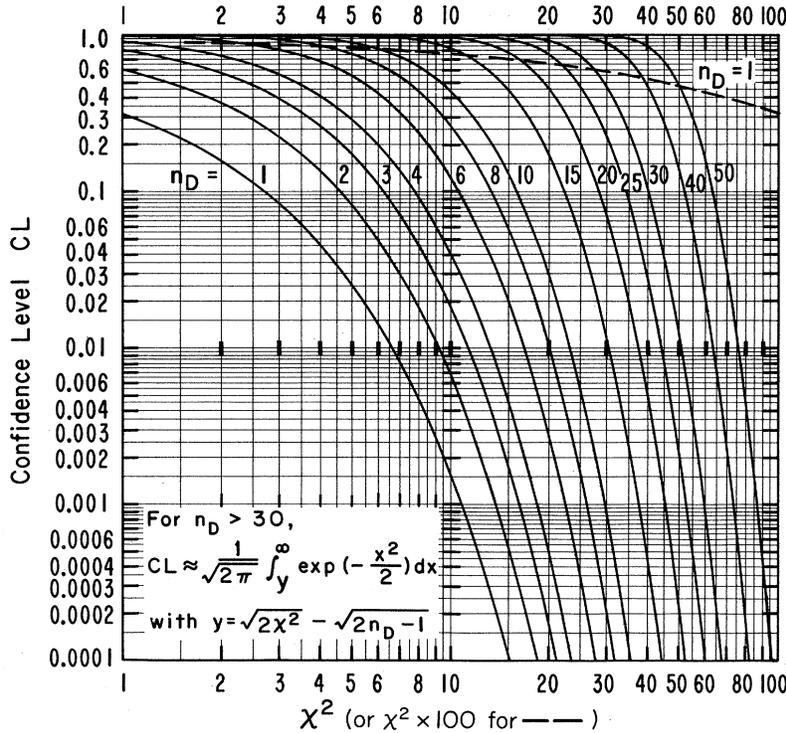
PROBABILITY AND STATISTICS

A. PROBABILITY DISTRIBUTIONS AND CONFIDENCE LEVELS

We give here properties of the three probability distributions most commonly used in high energy physics: Normal (or Gaussian), Chi-squared, and Poisson. We warn the reader that there is no universal convention for the term "confidence level"

as used by physicists; thus, explicit definitions are given for each distribution, and we have attempted to choose definitions that correspond to common usage. It is explained below how confidence levels for all three distributions can be extracted from the following figure.

χ^2 Confidence Level vs. χ^2 for n_D Degrees of Freedom

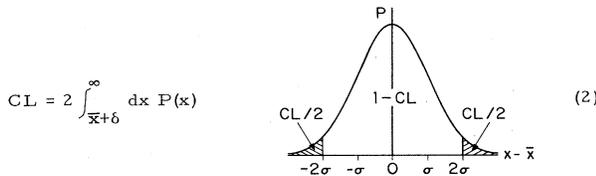


A.1. Normal Distribution

The normal distribution with mean \bar{x} and standard deviation σ (variance σ^2) is:

$$P(x)dx = \frac{1}{\sqrt{2\pi}\sigma} e^{-(x-\bar{x})^2/2\sigma^2} dx. \tag{1}$$

The confidence level associated with an observed deviation from the mean, δ , is the probability that $|x-\bar{x}| > \delta$, i.e.,



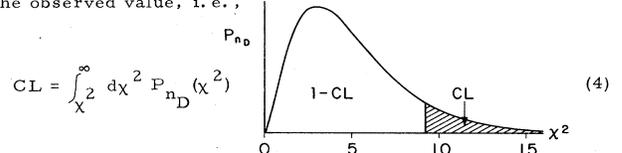
[The small figure in Eq. (2) is drawn with $\delta = 2\sigma$.] CL is given by the ordinate of the $n_D = 1$ curve in the figure at $\chi^2 = (\delta/\sigma)^2$. The confidence level for $\delta = 1\sigma$ is 31.7%; 2σ , 4.6%; 3σ , 0.3%. The central confidence interval, $1-CL$, (which is also sometimes called confidence level) for $\delta = 1\sigma$ is 68.3%; 2σ , 95.4%; 3σ , 99.7%. The odds against exceeding δ , $(1-CL)/CL$, for $\delta = 1\sigma$ are 2.15:1; 2σ , 21:1; 3σ , 370:1; 4σ , 16,000:1; 5σ , 1,700,000:1. Relations between σ and other measures of the width: probable error (CL = 0.5 deviation) = 0.67σ ; mean absolute deviation = 0.80σ ; RMS deviation = σ ; half width at half maximum = 1.18σ .

A.2. Chi-squared Distribution

The chi-squared distribution for n_D degrees of freedom is:

$$P_{n_D}(\chi^2) d\chi^2 = \frac{1}{2^h \Gamma(h)} (\chi^2)^{h-1} e^{-\chi^2/2} d\chi^2 \quad (\chi^2 \geq 0), \tag{3}$$

where h (for "half") = $n_D/2$. The mean and variance are n_D and $2n_D$, respectively. In evaluating Eq. (3) one may use Stirling's approximation: $\Gamma(h) = (h-1)! \approx 2.507 e^{-h} h^{h-1/2} \times (1 + 0.0833/h)$ which is accurate to $\pm 0.1\%$ for all $h \geq 1/2$. The confidence level associated with a given value of n_D and an observed value of χ^2 is the probability of chi-squared exceeding the observed value, i.e.,



[The small figure in Eq. (4) is drawn with $n_D = 5$ and $CL = 10\%$.] CL is plotted as a function of χ^2 for several values of n_D in the above figure. For large n_D , χ^2 becomes normally distributed about n_D . Thus,

$$y_1 = (\chi^2 - n_D) / \sqrt{2n_D} \tag{5}$$

becomes normally distributed with unit standard deviation. A better approximation, due to Fisher,¹ is that χ , not χ^2 , becomes normally distributed, specifically

$$y_2 = \sqrt{2\chi^2} - \sqrt{2n_D - 1} \tag{6}$$

approaches normality with unit standard deviation. For small CL's in particular, y_2 is much more accurate than y_1 . Thus, for $n_D = 50$ and $\chi^2 = 80$, the true CL = 0.45%, but y_1 is 3.0 corresponding to a CL of 0.13%, while y_2 is 2.7 corresponding to a CL of 0.35%.

PROBABILITY AND STATISTICS (Cont'd)

A.3. Poisson Distribution

The Poisson distribution with mean \bar{n} is:

$$P_{\bar{n}}(n) = \frac{e^{-\bar{n}} \bar{n}^n}{n!} \quad (n = 0, 1, 2, \dots) \quad (7)$$

The variance is equal to the mean. Confidence levels for Poisson distributions are usually defined in terms of quantities called "upper limits" as follows: The confidence level associated with a given upper limit N and an observed value n of n is the probability that $n > n_0$ if $\bar{n} = N$, i. e.,

$$CL = \sum_{n=n_0+1}^{\infty} P_N(n) = 1 - \sum_{n=0}^{n_0} P_N(n) \quad (8)$$

[The small figure in Eq. (8) is drawn with $n_0 = 2$ and $CL = 90\%$.] A useful relation between Poisson and chi-squared confidence levels allows one to look up this quantity on the above figure. Specifically, the quantity $1-CL$ is given by the ordinate of the $n_D = 2(n_0+1)$ curve at $\chi^2 = 2N$. Thus, 90% confidence level upper limits for $n_0 = 0, 1, \text{ and } 2$ are given by half the χ^2 value corresponding to an ordinate of 0.1 on the $n_D = 2, 4, \text{ and } 6$ curves, respectively; the values are $N = 2.3, 3.9, \text{ and } 5.3$.

Tables of confidence levels for all three of these distributions, the relation between Poisson and chi-squared confidence levels, and numerous other useful tables and relations may be found in Ref. 2.

B. STATISTICS

We consider here the situation in which one is presented with N independent data, $y_n \pm \sigma_n$, and it is desired to make some inference about the "true" value of the quantity represented by these data. For this purpose we interpret each datum y_n as a single sample point drawn randomly (and independently of the other data) from a distribution having mean \bar{y}_n (which we wish to estimate) and variance σ_n^2 . (Identification of the true σ_n with the σ_n datum is an approximation which may become seriously inaccurate when σ_n is an appreciable fraction of y_n .) Some methods of estimation commonly used in high energy physics are given below; see Ref. 3 for numerous applications. Section B.1. deals with the case in which all \bar{y}_n are the same, e. g., several different measurements of the same quantity; Sec. B.2. deals with the case in which $\bar{y}_n = \bar{y}(x_n)$, where x_n represents some set of independent variables, e. g., cross-section measurements at various values of energy and angle, $x_n = \{\theta_n, \theta_n\}$.

B.1. Single Mean and Variance Estimates

(1) If the y_n represent a set of values all supposedly drawn from a single distribution with mean \bar{y} and variance σ^2 (i. e., the σ_n are all the same, but their common value is unknown) then

$$\bar{y}_e = \frac{1}{N} \sum y_n \quad \text{and} \quad (9)$$

$$\sigma_e^2 = \frac{1}{N-1} \sum (y_n - \bar{y}_e)^2 \quad (10)$$

are unbiased estimates of \bar{y} and σ^2 . The variance of \bar{y}_e is σ^2/N . If the parent distribution is normal and N is large, the variance of σ_e^2 is $2\sigma^4/N$.

(2) If the \bar{y}_n all have the common value \bar{y} and the σ_n are known, then the weighted average

$$\bar{y}_e = \frac{1}{w} \sum w_n y_n \quad (11)$$

where $w_n = 1/\sigma_n^2$ and $w = \sum w_n$, is an appropriate unbiased estimate of \bar{y} . This choice of weighting factors in Eq. (11) minimizes the variance of the estimate; the variance is $1/w$.

B.2. Linear Least Squares Fit

A least squares fit of the function $y(x) = \sum_i a_i f_i(x)$ to independent data $y_n \pm \sigma_n$ at points x_n (e. g., a Legendre fit in which the f_i are Legendre polynomials and the a_i are Legendre coefficients) gives the following estimates of the parameters a_i :

$$a_{e,i} = \sum_{j=1}^N V_{ij} f_j(x_n) y_n / \sigma_n^2 \quad (12)$$

Here V is the covariance matrix of the fitted parameters

$$V_{ij} = \overline{(a_{e,i} - \bar{a}_{e,i})(a_{e,j} - \bar{a}_{e,j})} \quad (13)$$

which is given by

$$(V^{-1})_{ij} = \sum f_i(x_n) f_j(x_n) / \sigma_n^2 \quad (14)$$

The variance of an interpolated or extrapolated value of y at point x, $y_e = \sum a_{e,i} f_i(x)$, is:

$$(y_e - \bar{y}_e)^2 = \sum_{ij} V_{ij} f_i(x) f_j(x) \quad (15)$$

For the case of a straight line fit, $y(x) = a + bx$, one obtains the following estimates of a and b,

$$a_e = (S_y S_{xx} - S_x S_{xy}) / D, \quad (16)$$

$$b_e = (S_1 S_{xy} - S_x S_y) / D,$$

where

$$S_1, S_x, S_y, S_{xx}, S_{xy} = \sum (1, x_n, y_n, x_n^2, x_n y_n) / \sigma_n^2, \quad (17)$$

$$D = S_1 S_{xx} - S_x^2$$

The covariance matrix of the fitted parameters is:

$$\begin{pmatrix} V_{aa} & V_{ab} \\ V_{ab} & V_{bb} \end{pmatrix} = \frac{1}{D} \begin{pmatrix} S_{xx} & -S_x \\ -S_x & S_1 \end{pmatrix} \quad (18)$$

The variance of an interpolated or extrapolated value of y at point x is:

$$(y_e - \bar{y}_e)^2 = \frac{1}{S_1} + \frac{S_1}{D} \left(x - \frac{S_x}{S_1} \right)^2 \quad (19)$$

C. ERROR PROPAGATION

We consider here the situation in which one wishes to calculate the value and error of a function of some other quantities with errors, e. g., in a Monte Carlo program. Let {y} be a set of random variables with means $\{\bar{y}\}$ and covariance matrix V. Then the mean and variance of a function of these variables are approximately (to second order in $\{y - \bar{y}\}$):

$$\bar{f} \approx f(\{\bar{y}\}) + \frac{1}{2} \sum_{mn} V_{mn} \left(\frac{\partial^2 f}{\partial y_m \partial y_n} \right) \{\bar{y}\} = \{f\} \quad (20)$$

$$(f - \bar{f})^2 = \sum_{mn} V_{mn} \left(\frac{\partial f}{\partial y_m} \right) \{\bar{y}\} \left(\frac{\partial f}{\partial y_n} \right) \{\bar{y}\} = \{f\} \quad (21)$$

E. g., the mean and variance of a function of a single variable with mean \bar{y} and variance σ^2 are:

$$\bar{f} \approx f(\bar{y}) + \frac{1}{2} \sigma^2 f''(\bar{y}), \quad (22)$$

$$(f - \bar{f})^2 = \sigma^2 f'(\bar{y})^2. \quad (23)$$

Note that these equations will usually be applied by substituting some measured quantities, $\{\bar{y}\}$ say, for the true means, $\{\bar{y}\}$. If, as is often the case, $\bar{y}_n - \bar{y}_n$ is of order $\sqrt{V_{nn}}$, then there is no point in keeping the second order terms in Eq. (20) or (22) since the substitution itself introduces first order errors.

1. R. A. Fisher, *Statistical Methods for Research Workers* (Oliver and Boyd, Edinburgh and London, 1958).
2. M. Abramovitz and I. Stegun, eds., *Handbook of Mathematical Functions* (National Bureau of Standards, Applied Mathematics Series, Vol. 55, Washington, 1964).
3. W. T. Eadie, D. Drijard, F. E. James, M. Roos, and B. Sadoulet, *Statistical Methods in Experimental Physics* (North-Holland, Amsterdam and London, 1971).

Revised and expanded April 1974.

RELATIVISTIC KINEMATICS (Cont'd)

Application to the reaction $a + b \rightarrow c + d$:

1. General formulae

$$t \equiv t_{ac} = t_{bd} = t_{ac}^{\min} - 4p_a p_c \sin^2\left(\frac{\theta_{ac}}{2}\right),$$

with $t_{ac}^{\max} = (E_a - E_c)^2 - (p_a \mp p_c)^2$ (II-20)

which, after expansion in powers of $(1/s)$, gives

$$t_{ac}^{\min} = -\frac{(m_a^2 - m_c^2)(m_b^2 - m_d^2)}{s} - \frac{(m_a^2 + m_b^2 - m_c^2 - m_d^2)(m_a^2 m_b^2 - m_c^2 m_d^2)}{s^2} + O\left(\frac{1}{s^3}\right)$$
 (II-22)

In a similar way, defining $u_{ad} = (p_a - p_d)^2$, one finds

$$u \equiv t_{ad} = t_{bc} = u_{ad}^{\min} - 4p_a p_d \sin^2\left(\frac{\theta_{ad}}{2}\right).$$

$$u_{ad}^{\min} = -\frac{(m_a^2 - m_d^2)(m_b^2 - m_c^2)}{s} - \frac{(m_a^2 + m_b^2 - m_c^2 - m_d^2)(m_a^2 m_b^2 - m_c^2 m_d^2)}{s^2} + O\left(\frac{1}{s^3}\right)$$
 (II-23)

A general relation between the invariants is

$$s + t + u = m_a^2 + m_b^2 + m_c^2 + m_d^2.$$

IV. RELATIONS BETWEEN THE (j) PARTICLE AND (ij) PARTICLE REST FRAMES: LORENTZ TRANSFORMATION

The general Lorentz transformation has the matrix form

$$\begin{pmatrix} E \\ p_{\parallel} \\ p_{\perp} \end{pmatrix}^{(ij)} = \begin{pmatrix} \gamma & -\eta & 0 \\ -\eta & \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} E \\ p_{\parallel} \\ p_{\perp} \end{pmatrix}^{(j)}$$
 (IV-1)

If we define, in any frame, the quantity: $P_{\pm} = E \pm p_{\parallel}$,

$$\begin{pmatrix} P_{+} \\ P_{-} \\ p_{\perp} \end{pmatrix}^{(ij)} = \begin{pmatrix} (\gamma - \eta) & 0 & 0 \\ 0 & (\gamma + \eta) & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_{+} \\ P_{-} \\ p_{\perp} \end{pmatrix}^{(j)}$$
 (IV-3)

$$\gamma \mp \eta = \frac{s_{ij} + m_j^2 - m_i^2 \mp \sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{2m_j \sqrt{s_{ij}}} = e^{\mp \xi_{ij}}$$
 (IV-4)

$$\left\{ \begin{aligned} |\beta| &= \frac{|\eta|}{\gamma} = \tanh \xi_{ij} = \frac{\sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{s_{ij} + m_j^2 - m_i^2} = \frac{p_i^{(j)}}{E_i^{(j)} + m_j} \end{aligned} \right.$$
 (IV-5)

$$\left\{ \begin{aligned} \gamma &= \frac{1}{\sqrt{1 - \beta^2}} = \cosh \xi_{ij} = \frac{s_{ij} + m_j^2 - m_i^2}{2m_j \sqrt{s_{ij}}} = \frac{E_i^{(j)} + m_j}{\sqrt{s_{ij}}} \end{aligned} \right.$$
 (IV-6)

$$\left\{ \begin{aligned} |\eta| &= |\beta|\gamma = \sinh \xi_{ij} = \frac{\sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{2m_j \sqrt{s_{ij}}} = \frac{p_i^{(j)}}{\sqrt{s_{ij}}} \end{aligned} \right.$$
 (IV-7)

The spatial part of the matrix in Eq. (IV-1) above can be written with the classical vector form

$$\vec{p}^{(ij)} = \vec{p}^{(j)} + \vec{\eta} \left(\frac{\vec{\eta} \cdot \vec{p}}{\gamma + 1} - E^{(j)} \right)$$
 (IV-9)

A convenient way to write the latter expression is

$$\vec{p}^{(ij)} = \vec{p}^{(j)} - \vec{\eta} \left[\frac{E^{(ij)} + E^{(j)}}{\gamma + 1} \right]$$
 (IV-10)

2. Particular case

In the c.m.: $dt = 2p_a p_c d(\cos \theta_{ac})$.

i) $m_a = m_c, m_b \neq m_d$:

$$t_{ac}^{\min} \approx -\frac{m_a^2(m_b^2 - m_d^2)^2}{s^2}$$
 (II-24)

$$u_{ad}^{\min} \approx -\frac{(m_a^2 - m_d^2)(m_b^2 - m_c^2)}{s}$$
 (II-25)

ii) For elastic scattering ($a + b \rightarrow a' + b'$),

$$t_{aa'}^{\min} = 0,$$

and $t_{aa'} = -4p_a^2 \sin^2\left(\frac{\theta_{aa'}}{2}\right)$,

$$dt_{aa'} = 2p_a^2 d(\cos \theta_{aa'});$$
 (II-26)

also $u_{ab'}^{\min} = \frac{(m_a^2 - m_b^2)^2}{s}$ (II-27)

At high energy:

$$e^{\xi_{ij}} \approx \frac{\sqrt{s_{ij}}}{m_j} \quad \text{and} \quad \xi_{ij} \approx \ln \frac{\sqrt{s_{ij}}}{m_j}.$$

ξ_{ij} is called the *boost parameter* of the Lorentz transformation that connects the particle ij rest frame to the particle j rest frame.

Application

1. To reaction $a(\text{beam}) + b(\text{target})$

The transformation from the lab frame into the c.m. frame is given by:

$$\gamma = \frac{E_a^{\text{lab}} + m_b}{\sqrt{s}} \quad \text{and} \quad \eta = \frac{p_{\text{inc}}}{\sqrt{s}};$$

$$p_{\pm}^{\text{cm}} = e^{\mp \xi} p_{\pm}^{\text{lab}}$$

with

$$e^{\pm \xi} = \frac{s + m_b^2 - m_a^2 \mp \sqrt{\Delta(s, m_a^2, m_b^2)}}{2m_b \sqrt{s}}$$
 (IV-11)

2. The Lorentz transform of any vector

$\vec{p}_1 = (E_1, \vec{p}_1)$, given in a frame containing another vector $\vec{p}_2 = (E_2, \vec{p}_2)$, into the p_2 rest frame — where $p_2 = (m_2, 0)$, and $p_1 = (E_1, \vec{p}_1)$ — may be calculated using:

$$\gamma = \frac{E_2}{m_2} \quad \text{and} \quad \vec{\eta} = \frac{\vec{p}_2}{m_2};$$

$$\left. \begin{aligned} E_1' &= \gamma E_1 - \vec{\eta} \cdot \vec{p}_1 = \frac{E_1 E_2 - \vec{p}_1 \cdot \vec{p}_2}{m_2}, \\ \vec{p}_1' &= \vec{p}_1 - \frac{\vec{p}_2}{m_2} \left(\frac{E_1 + E_1'}{\gamma + 1} \right) = \vec{p}_1 - \vec{p}_2 \left(\frac{E_1 + E_1'}{E_2 + m_2} \right) \end{aligned} \right\}$$
 (IV-12)

RELATIVISTIC KINEMATICS (Cont'd)

V. RAPIDITY VARIABLE: DEFINITION AND KINEMATIC RELATIONS

Definition: in any system, for a particle of energy E and longitudinal momentum P_{\parallel} , the rapidity is defined by

$$y \equiv \frac{1}{2} \ln \left(\frac{E + P_{\parallel}}{E - P_{\parallel}} \right) = \tanh^{-1} \left(\frac{P_{\parallel}}{E} \right) \tag{V-1}$$

$$= \frac{1}{2} \ln \left(\frac{P_+}{P_-} \right) = \ln \left(\frac{P_+}{m_1} \right) = -\ln \left(\frac{P_-}{m_1} \right)$$

with the transverse mass $m_1 \equiv (m^2 + p_{\perp}^2)^{1/2}$.

From the last of the two equations above, one gets

$$P_+ = m_1 e^y \tag{V-2}$$

$$P_- = m_1 e^{-y} \tag{V-3}$$

so that
$$\begin{cases} E = m_1 \cosh y \tag{V-4} \\ P_{\parallel} = m_1 \sinh y \tag{V-5} \end{cases}$$

A. Relations between the rapidities of different frames:

$$y^{(ij)} = y^{(j)} - \xi_{ij} \tag{V-7}$$

B. Application to a reaction $a+b \rightarrow c+d$ (or $a+b \rightarrow c+X$):

In the lab frame,

$$\left. \begin{aligned} p_a^{\text{lab}} &= (m_a \cosh y_a^{\text{lab}}, m_a \sinh y_a^{\text{lab}}, 0) \\ p_b^{\text{lab}} &= (m_b, 0, 0) \\ p_c^{\text{lab}} &= (m_c \cosh y_c^{\text{lab}}, m_c \sinh y_c^{\text{lab}}, \vec{p}_{1c}) \end{aligned} \right\} \tag{V-8}$$

(and a similar formula for particle d or X).

In the c.m. frame,

$$\left. \begin{aligned} p_a^{\text{cm}} &= (m_a \cosh y_a^{\text{cm}}, m_a \sinh y_a^{\text{cm}}, 0) \\ p_b^{\text{cm}} &= (m_b \cosh y_b^{\text{cm}}, m_b \sinh y_b^{\text{cm}}, 0) \\ p_c^{\text{cm}} &= (m_c \cosh y_c^{\text{cm}}, m_c \sinh y_c^{\text{cm}}, \vec{p}_{1c}) \end{aligned} \right\} \tag{V-9}$$

(and a similar formula for particle d or X).

$$y_a^{\text{lab}} = \cosh^{-1} \left(\frac{s - m_a^2 - m_b^2}{2m_a m_b} \right) \xrightarrow{\text{at large } s} y_a^{\text{lab}} \approx \ln \frac{s}{m_a m_b} \tag{V-10}$$

$$y_a^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_a^2 - m_b^2}{2m_a \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_a^{\text{cm}} \approx \ln \frac{\sqrt{s}}{m_a} \tag{V-11}$$

$$y_b^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_b^2 - m_a^2}{2m_b \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_b^{\text{cm}} \approx -\ln \frac{\sqrt{s}}{m_b} \tag{V-12}$$

$$y_c^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_c^2 - m_d^2}{2m_{1c} \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_c^{\text{cm}} \approx \ln \frac{\sqrt{s}}{m_{1c}} \tag{V-13}$$

and m_d^2 (or m_X^2) $\ll s$

With $\xi \approx \ln \frac{\sqrt{s}}{m_b}$, one gets, at large s:

$$y_c^{\text{lab}} \approx \ln \frac{s}{m_b m_{1c}} \tag{V-14}$$

$$y^{\text{lab}} \approx \frac{1}{2} \ln \left(\frac{P^{\text{lab}} + P_{\parallel}^{\text{lab}}}{P^{\text{lab}} - P_{\parallel}^{\text{lab}}} \right) \approx -\ln \left(\tan \frac{\theta}{2} \right)$$

The maximum rapidity gap for the reaction $a+b \rightarrow c+d$ at large s is:

i) for the incoming particles (a,b):

$$y_{ab} = y_a^{\text{lab}} - y_b^{\text{lab}} = y_a^{\text{cm}} - y_b^{\text{cm}} \approx \ln \frac{s}{m_a m_b} \tag{V-15}$$

ii) for the outgoing particles (c,d):

$$y_{cd} = y_c^{\text{lab}} - y_d^{\text{lab}} = y_c^{\text{cm}} - y_d^{\text{cm}} \approx \ln \frac{s}{m_{1c} m_{1d}} \tag{V-16}$$

Application to reaction $a+b \rightarrow 1+2 \dots +N$, with the N outgoing particles being ordered by increasing rapidity.

1. The maximum rapidity gap between the outgoing particles has an approximate value

$$y_{1N} \approx \ln \frac{s}{m_{11} m_{1N}} \tag{V-17}$$

2. The rapidity gap $(y_i - y_j)$ between the rapidities y_i and y_j of particles i and j, respectively, is related to the invariant mass squared of particles i and j

$$s_{ij} = m_i^2 + m_j^2 + 2m_i \cdot m_j \cosh(y_i - y_j) - 2\vec{p}_{1i} \cdot \vec{p}_{1j} \tag{V-18}$$

$$y_i - y_j = \cosh^{-1} \frac{s_{ij} - m_i^2 - m_j^2 + 2\vec{p}_{1i} \cdot \vec{p}_{1j}}{2m_i m_j} \tag{V-19}$$

3. The rapidity gap between either of the extreme particles (1 or N) and its adjacent neighbor [2 or (N-1) respectively] is, on a statistical basis, and in first approximation

$$|y_1 - y_2| \approx \ln \frac{s}{M_1^2} \tag{V-20}$$

and

$$|y_{N-1} - y_N| \approx \ln \frac{s}{M_N^2} \tag{V-21}$$

M_1 (M_N) being the missing mass with respect to particle 1 (N) in the reaction $a + b \rightarrow 1 + M_1$ ($a + b \rightarrow N + M_N$).

4. The rapidity gap y_{a1} between the incoming particle a and the outgoing particle closest to a, 1, is related to the transverse momentum $t_{a1} = (p_a - p_1)^2$:

$$y_a - y_1 = \cosh^{-1} \left(\frac{t_{a1} - m_a^2 - m_1^2}{2m_a m_{11}} \right) \tag{V-23}$$

RELATIVISTIC KINEMATICS (Cont'd)

VI. INCLUSIVE REACTIONS AND SCALING VARIABLES: DEFINITIONS AND KINEMATIC RELATIONS

A. Scaling Variable: Definitions

• For the *inclusive reaction* $a + b \rightarrow c + X$

1) Feynman's definition:

$$x = p_{\parallel}^{cm} / p_{\parallel \max}^{cm} \quad (VI-3)$$

where
$$p_{\parallel \max}^{cm} = \frac{\sqrt{\Delta(s, m_c^2, m_{X, \min}^2)}}{2\sqrt{s}} \quad (VI-4)$$

is the center-of-mass momentum of particle c and the lightest possible particle(s) $m_{X, \min}$ which, consistent with conservation laws, could recoil against particle c.

A relationship between x and m_X^2 which becomes useful at high energy [see (VI-6)] derives from

$$m_X^2 = (p_a + p_b - p_c)^2 = s + m_c^2 - 2p_c \cdot (p_a + p_b)$$

2) An alternative definition:

$$x' = p_{\parallel}^{cm} / p_{\parallel \max}^{cm} \quad (VI-7)$$

with
$$p_{\parallel \max}^{cm} = \sqrt{(p_{\max}^{cm})^2 - p_{\perp}^2}$$

• For *electro- and photoproduction*:

1) Bjorken's definition:

$$\omega = \frac{2m_p v}{q^2}$$

where
$$\left\{ \begin{array}{l} v = e - e' = \text{the difference between the incident and the outgoing energies} \\ m_p = \text{proton mass} \\ q^2 = \text{momentum transfer squared} \end{array} \right.$$

(NOTE: ω^{-1} is sometimes also called x.)

2) Bloom-Gilman's definition:

$$\omega' = \omega + \frac{m_p^2 - q^2}{q^2}$$

B. For the inclusive reaction $a + b \rightarrow c + X$, the *invariant differential cross section* $E(d^3\sigma/dp_c^3)$ may be written in the following forms using various expressions of the phase space volume d^3p_c/E_c (and omitting the subscript c):

Longitudinal variable of particle c	Transverse variable of particle c	Corresponding invariant (differential) cross section $E(d^3\sigma/dp^3)$
p_{lab}	θ^{lab}	$\frac{E^{\text{lab}}}{(p^{\text{lab}})^2} \frac{d^3\sigma}{dp^{\text{lab}} d\Omega^{\text{lab}}} \quad (VI-10)$
p_{\parallel}	p_{\perp}	$\frac{E}{\pi} \frac{d^2\sigma}{dp_{\parallel} dp_{\perp}^2}$ (averaged over azimuth) $(VI-11)$
x	p_{\perp}	$\frac{E^{cm}}{p_{\max}^{cm}} \cdot \frac{1}{\pi} \frac{d^2\sigma}{dx dp_{\perp}^2} \approx \frac{2E^{cm}}{\pi\sqrt{s}} \frac{d^2\sigma}{dx dp_{\perp}^2}$ (averaged over azimuth) $(VI-12)$
y	p_{\perp} (or m_{\perp})	$\frac{1}{\pi} \frac{d^2\sigma}{dy dp_{\perp}^2} = \frac{1}{\pi} \frac{d^2\sigma}{dy dm_{\perp}^2}$ (averaged over azimuth) $(VI-13)$
$m_X^2 = (p_a + p_b - p_c)^2$	$t = (p_a - p_c)^2$	$\frac{\sqrt{\Delta(s, m_a^2, m_b^2)}}{\pi} \frac{d^2\sigma}{dm_X^2 dt}$ (averaged over azimuth) $(VI-14)$

• At high energy,

$$p_{\max}^{cm} \approx \frac{\sqrt{s}}{2} \quad \text{and} \quad x \approx \frac{2p_{\parallel}^{cm}}{\sqrt{s}} \quad (VI-5)$$

The relation between x and y is trivial:

$$x \approx 2 \frac{p_{\parallel}^{cm}}{\sqrt{s}} \approx \frac{2m_{\perp}}{\sqrt{s}} \sinh y^{cm} \quad (VI-8)$$

and, for y also large:

$$y^{cm} \approx \ln \frac{x\sqrt{s}}{m_{\perp}} \quad (VI-9)$$

$$x = 1 - \frac{m_X^2 - m_c^2}{s} \approx 1 - \frac{m_X^2}{s} \quad (VI-6)$$

(for $m_X^2 \gg m_c^2$).

For p_{\perp} fixed as $s \rightarrow \infty$,

$$x' \approx \frac{2p_{\parallel}^{cm}}{\sqrt{s}} \approx x$$

For q^2 and v large (Bjorken limit),

$$\omega' \approx \omega - 1$$

*See Denyse M. Chew, "Relativistic Kinematics", Particle Data Group Memo PDG-101, for details and references.

Revised and expanded April 1976.

PARTICLE DETECTORS, ABSORBERS, AND RANGES*

A. DETECTOR PARAMETERS

In this section we give various parameters for common detectors. The quoted numbers represent at best an order of magnitude useful only for preliminary design.

A.1 Scintillators: Photon yield $\approx 1\gamma/100$ eV in plastic scintillator¹ and $\approx 1\gamma/25$ eV in NaI.^{1,2}

A.2 Čerenkov³: Half angle θ_c of cone aperture in terms of velocity β and index of refraction n :

$$\theta_c = \arccos\left(\frac{1}{\beta n}\right) \sim \sqrt{2\left(1 - \frac{1}{\beta n}\right)}$$

Threshold velocity: $\beta_t = 1/n$; $\gamma_t = 1/\sqrt{1-\beta_t^2}$.

Therefore, $\beta_t \gamma_t = 1/\sqrt{2\delta + \delta^2}$, where $\delta = n-1$. Values of δ for various commonly used gases are given as a function of pressure and wavelength in Ref. 4; for values at atmospheric pressure, see the Table of Atomic and Nuclear Properties following.

Number of photons N per cm:

$$N = \frac{\alpha}{c} \int \left(1 - \frac{1}{\beta^2 n^2}\right) 2\pi \, d\nu = \frac{\alpha}{c} \beta_t^2 \int \left(\frac{1}{\beta_t^2 \gamma_t^2} - \frac{1}{\beta^2 \gamma^2}\right) 2\pi \, d\nu$$

$$\approx 500 \sin^2 \theta_c / \text{cm (visible spectrum)}$$

A.3 Photon Collection: In addition to the photon yield, one should take into account the light collection efficiency ($\approx 10\%$ for typical 1 cm thick scintillator), attenuation length (≈ 1 to 4 m for typical scintillators⁵), and quantum efficiency of the photomultiplier cathode ($\leq 25\%$).

A.4 Bubble, Streamer, Wire Chambers:

Chamber Type	Accuracy(rms)	Resolution Time	Dead Time
Bubble	$\pm 75\mu$	≈ 1 ms	$\approx 1/20$ s ^{a)}
Streamer	$\pm 300\mu$	≈ 2 μ s	≈ 100 ms
Optical Spark	$\pm 200\mu$ ^{b)}	≈ 2 μ s	≈ 10 ms
Magnetostrictive Spark	$\pm 500\mu$	≈ 2 μ s	≈ 10 ms
Proportional	$\approx \pm 300\mu$ ^{c, d)}	≈ 50 ns	≈ 200 ns
Drift	± 50 to 300μ	≈ 2 ns ^{e)}	≈ 100 ns

- a) Multiple pulsing time.
- b) 60μ for high pressure.
- c) 300μ is for 1 mm pitch.
- d) Delay line cathode readout can give $\pm 150\mu$ parallel to anode wire.
- e) For two chambers.

A.5 Shower Detectors: Typical energy resolution for incident electron in the 1 GeV range, E in GeV.

NaI (20 rad. lengths) ⁶ :	$\frac{2\%}{E^{1/4}}$	(FWHM)
Lead Glass (14 rad. lengths) ⁷ :	$\frac{8-12\%}{\sqrt{E}}$	(FWHM)
Lead Scintillator Sandwich (10 rad. lengths) ⁸ :	$\frac{22\%}{\sqrt{E}}$	(FWHM)

(10 lead plates, each 1 radiation length, and scintillators, each 1 cm in thickness).

A.6 Proportional Chamber Wire Instability: The limit on the voltage V for a wire tension T is given by⁹

$$V \leq \frac{sT^{1/2}}{lC}$$

where s , l , and C are the wire spacing, length, and capacitance per unit length. An approximation to C for a chamber half gap t and wire diameter d (good for $s \leq t$) gives¹⁰

$$V \leq 59T^{1/2} \left[\frac{t}{l} + \frac{s}{\pi l} \ln\left(\frac{s}{\pi d}\right) \right],$$

where V is in kV, and T is in grams.

B. COSMIC RAY FLUXES

The fluxes of particles of different types depend on the latitude, their energy, and on the conditions of measurement. Some typical sea-level values¹¹ are given below:

- I_V vertical flux
- J_1 total flux crossing a unit horizontal area
- J_2 total flux crossing a sphere of unit radius

	Total Intensity	Hard component	Soft component
I_V	1.14×10^{-2}	0.83×10^{-2}	0.31×10^{-2} cm ⁻² sec ⁻¹ sterad ⁻¹
J_1	1.79×10^{-2}	1.27×10^{-2}	0.52×10^{-2} cm ⁻² sec ⁻¹
J_2	2.41×10^{-2}	1.68×10^{-2}	0.73×10^{-2} cm ⁻² sec ⁻¹

Very approximately, about 75% of all particles at sea-level are penetrating, and are μ mesons. The absolute flux of protons at sea-level, in a momentum range 700-1100 MeV/c, is 1.5×10^{-5} cm⁻²sec⁻¹sterad⁻¹, or $\sim 0.1\%$ of all particles.

C. PASSAGE OF PARTICLES THROUGH MATTER

C.1 dE/dx ¹²: Ionization energy loss for particles heavier than an electron and having charge $z|e|$ (Bethe-Bloch equation):

$$-\frac{1}{\beta} \frac{dE}{dx} = z^2 D \frac{Z}{A \beta^2} \left[\ln\left(\frac{2m_e c^2 \beta^2 T'_{\max}}{(1-\beta^2) I^2}\right) - 2\beta^2 - K \right],$$

where $D = 2\pi N_e^2 m_e c^2 = 0.1535$ MeV cm²/g (see Physical and Numerical Constants Table), Z and A are charge and mass number of the material, $\beta = v/c$ for the incident particle, T'_{\max} is the maximum-energy δ ray considered to be ionization loss, I is the mean ionization potential ($I \approx 20Z$ eV for $Z = 1$, tending toward $I \approx 10Z$ for $Z \geq 20$), and K includes the shell correction (maximum $\approx 10\%$ around $\beta = 0.1$ in the heaviest elements), density effect (important at high velocities and high densities), and the z^3 correction¹³ (important at low velocity, a few percent at $\beta \approx 0.1$).

A point of minimum ionization occurs in the range $\beta \approx 0.95$ to 0.97, with the dependence on absorber confined to that range. Energy loss increases approximately as $\ln \gamma$ thereafter (relativistic rise), until modified by the density effect (which reduces the rate of increase). Values of dE/dx at minimum are listed in the Table of Atomic and Nuclear Properties of Materials (following).

Figures are given below for ionization energy loss and range in lead (with scaling indicated for copper, aluminum, and carbon), and in liquid hydrogen, with T'_{\max} taken to be T'_{\max} as defined in Sec. C.2 below. Energy losses due to bremsstrahlung, nuclear collisions, etc., are not included. For thin samples of absorber, fluctuations are very important.¹⁴ A review of the problem of the penetration of charged particles into matter may be found in Fano¹⁵.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mixtures or compounds may be treated in first (but fairly accurate) approximation using Bragg's additivity rule:

$$\left(\frac{1}{\rho} \frac{dE}{dx}\right)_{\text{mix}} = \sum_i \left(\frac{1}{\rho} \frac{dE}{dx}\right)_i f_i,$$

where f_i is the fractional weight ($= \rho_i / \rho_{\text{mix}}$) of element i .

C.2 δ rays¹⁶: Number N per g/cm^2 per MeV of δ rays of kinetic energy T produced by an incident particle of momentum p and velocity β (Rutherford formula):

$$\frac{dN}{dTdx} = D \frac{Z}{A} \frac{1}{\beta^2 T^2},$$

where D is defined in the dE/dx section. This expression is valid for any unit charge incident particle if $T \ll T_{\text{max}}$. T_{max} is the maximum δ -ray energy, given by

$$T_{\text{max}} = \frac{2 m_e c^2 (\text{pc})^2}{s} = \frac{2 \beta^2 \gamma^2 m_e c^2}{1 + 2\gamma \frac{m_e}{M} + \frac{m_e^2}{M^2}}$$

where m_e and M are the electron and incident particle masses and s is their center-of-mass energy squared.

C.3 Multiple Coulomb Scattering^{17, 18, 19}: Sketched below are the multiple scattering angle θ , displacement y , and sagitta s , of a particle passing through a scatterer of thickness L . If the incident particle has mass M , charge $z|e|$, velocity β , and energy E , and the scatterer has atomic mass m and radiation length L_R , the projection of the $1/e$ multiple scattering angle (i.e., the value of the scattering angle at which the distribution falls to $1/e$ of its peak value) is given by

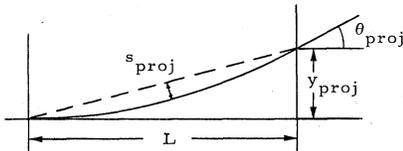
$$\theta_{\text{proj}}^{1/e} = z \frac{14 \text{MeV}/c}{\beta} \sqrt{\frac{L}{L_R}} \left[1 + \frac{1}{9} \log_{10} \left(\frac{L}{L_R} \right) \right] \left[1 + \frac{M^2}{E m_s} \right],$$

in the small angle approximation. The distribution is approximately Gaussian for $\theta < \theta^{1/e}$, but has long tails.²⁰ The accuracy of this expression is $\approx 5\%$ except for low z or β , where it is $\approx 10\text{-}20\%$.

For a given θ_{proj} ,

$$y_{\text{proj}} = \frac{L \theta_{\text{proj}}}{\sqrt{3}},$$

$$s_{\text{proj}} = \frac{L \theta_{\text{proj}}}{4 \sqrt{3}}.$$



C.4 Electron Range in Lead, Copper, Carbon and Hydrogen: See figure following.

C.5 Fractional Energy Loss for Electrons and Positrons in Lead: See figure following.

C.6 Contributions to Photon Cross Section in Lead: See figure following.

C.7 Photon Cross Section: See figure following.

D. ATOMIC AND NUCLEAR PROPERTIES OF MATTER

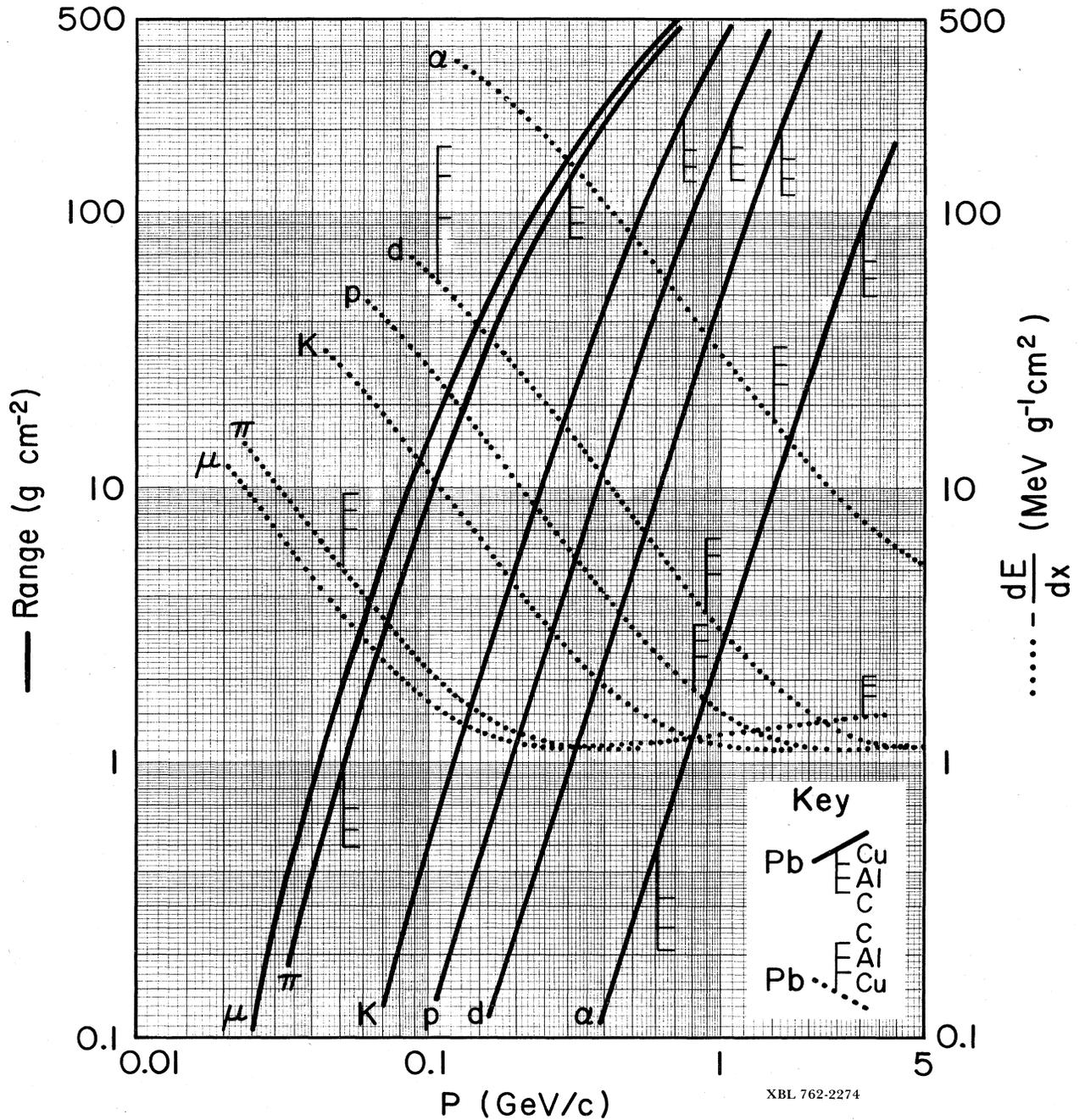
See table following.

*Prepared April 1974 by Sherwood Parker and Bernard Sadoulet. Revised April 1976.

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PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

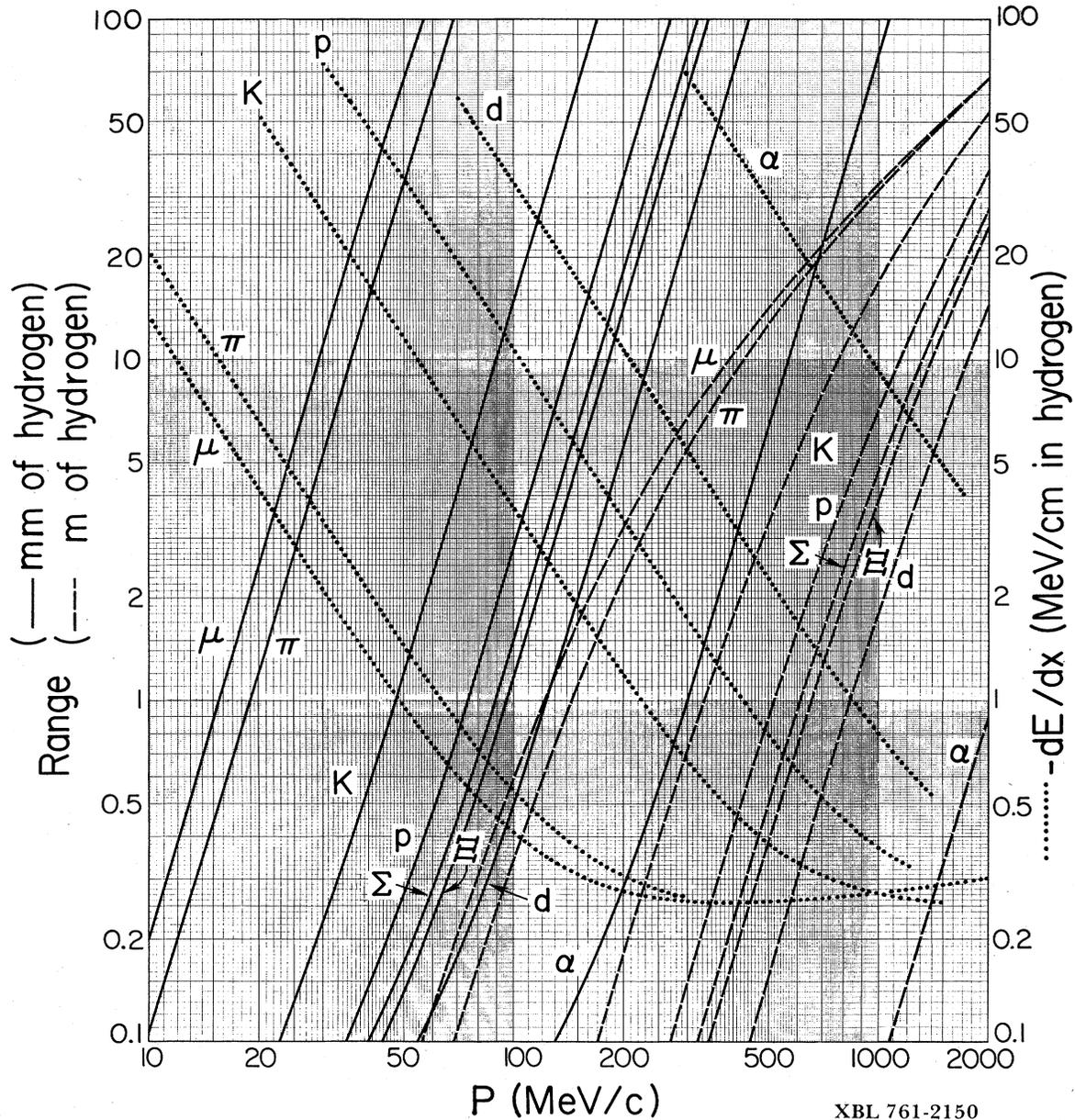
Mean Range and Energy Loss in Lead, Copper, Aluminum, and Carbon



Mean range and energy loss due to ionization for the indicated particles in Pb, with scaling to Cu, Al, and C indicated, using Bethe-Bloch equation (Section C.1 above) with corrections. Calculated using program of Hans Bichsel (UCRL-17538), with density correction added (Hans Bichsel, private communication). See also Joseph F. Janni [Air Force Weapons Laboratory Technical Report No. AFWL-TR-65-150 (1966)]. The average ionization potentials (I) assumed were: Pb (820 eV), Cu (320 eV), Al (166 eV), and C(77.5 eV). Figure indicates total path length; observed range may be smaller (by ~ 1% - 2% in heavy elements) due to multiple scattering, primarily from small energy-loss collisions with nuclei. The functional forms have not been experimentally verified to better than roughly ±1%. For higher energies refer to discussion by Cobb ["A Study of Some Electromagnetic Interactions of High Velocity Particles with Matter," University of Oxford Report HEP/T/55 (1973)] and by Turner ["Penetration of Charged Particles in Matter: A Symposium", National Academy of Sciences, Washington D. C. (1970), p. 48]. Scaling to other beam particles is, to a good approximation, described by the expression on the next page.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mean Range and Energy Loss in Liquid Hydrogen



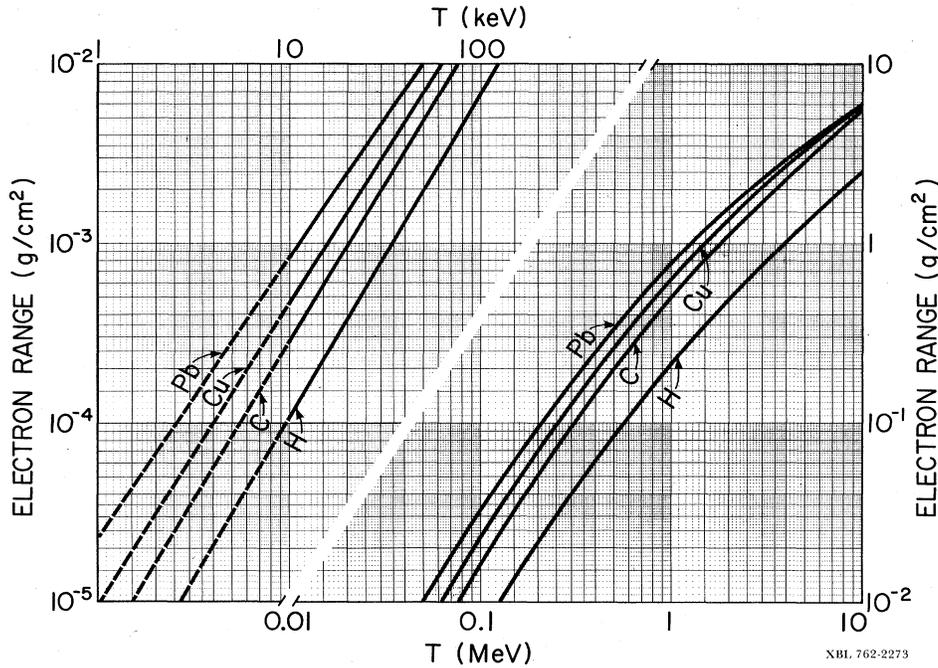
Range and energy loss in liquid hydrogen bubble chamber, based on Bethe-Bloch equation (Section C.1 above), using an average ionization potential for H_2 of $I = 20.0$ eV, which is an approximate average of the recent experimental result of Garbincius and Hyman [Phys. Rev. **A2**, 1834 (1970)] and the recent theoretical result of Ford and Browne [Phys. Rev. **A7**, 418 (1973)]: This is somewhat higher than used in the last edition. Bubble chamber conditions are chosen to be those of Garbincius and Hyman: parahydrogen of density = 0.0625 g/cm³ (note: range $\propto 1/\text{density}$), with vapor-pressure 60.8 lb/in² (absolute) and temperature 26.2 K. The functional dependence of the Bethe-Bloch equation is not experimentally verified to better than about $\pm 1\%$ over large momentum ranges. It should be noted that the number of bubbles per cm of a track in a bubble chamber is nearly proportional to $1/\beta^2$, not dE/dx . For the linear portions of the range curves, $R \propto p^{3.6}$.

Scaling law for particles of other mass or charge (except electrons): for a given medium, the range R_b of any beam particle with mass M_b , charge z_b , and momentum p_b is given in terms of the range R_a of any other particle with mass M_a , charge z_a , and momentum $p_a = p_b M_a / M_b$ (i. e., having the same velocity) by the expression

$$R_b(M_b, z_b, p_b) = \left[\frac{M_b/M_a}{z_b^2/z_a^2} \right] R_a(M_a, z_a, p_a = p_b M_a / M_b).$$

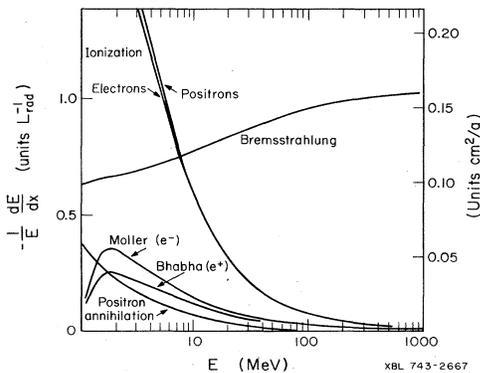
PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mean Electron Range in Lead, Copper, Carbon, and Liquid Hydrogen



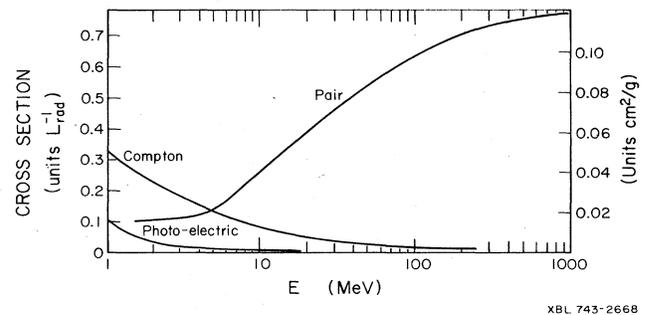
Mean range of electrons in the continuous-slowing-down approximation, taking into account energy loss by collisions with atomic electrons and by bremsstrahlung; strong fluctuations are to be expected for individual tracks. This range is the total path length; the practical range is shorter because of multiple Coulomb scattering, which becomes increasingly important as the electron slows down. E.g., for a fast electron the rms projected angle due to multiple Coulomb scattering reaches 1 radian by the time the electron has slowed to 0.4 MeV in hydrogen, 1.5 MeV in carbon, 9 MeV in copper, and 24 MeV (off scale) in lead. The "critical energy" (above which the energy loss due to bremsstrahlung exceeds that due to ionization and showering becomes important) is 400 MeV for hydrogen, 100 MeV for carbon, 25 MeV for copper, and 10 MeV for lead. The mean positron range may differ from the mean electron range by several percent. See Berger and Seltzer, NASA SP-3012 (1964) and SP-3036, and P. Trower, UCRL-2426, Vol. III Rev. (1966). 1-10 keV range was obtained by linear extrapolation.

Fractional Energy Loss for e⁺ and e⁻ in Lead



Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization when the energy loss per collision is below 0.255 MeV, and as Moller (Bhabha) scattering when it is above.

Contributions to Photon Cross Section in Lead

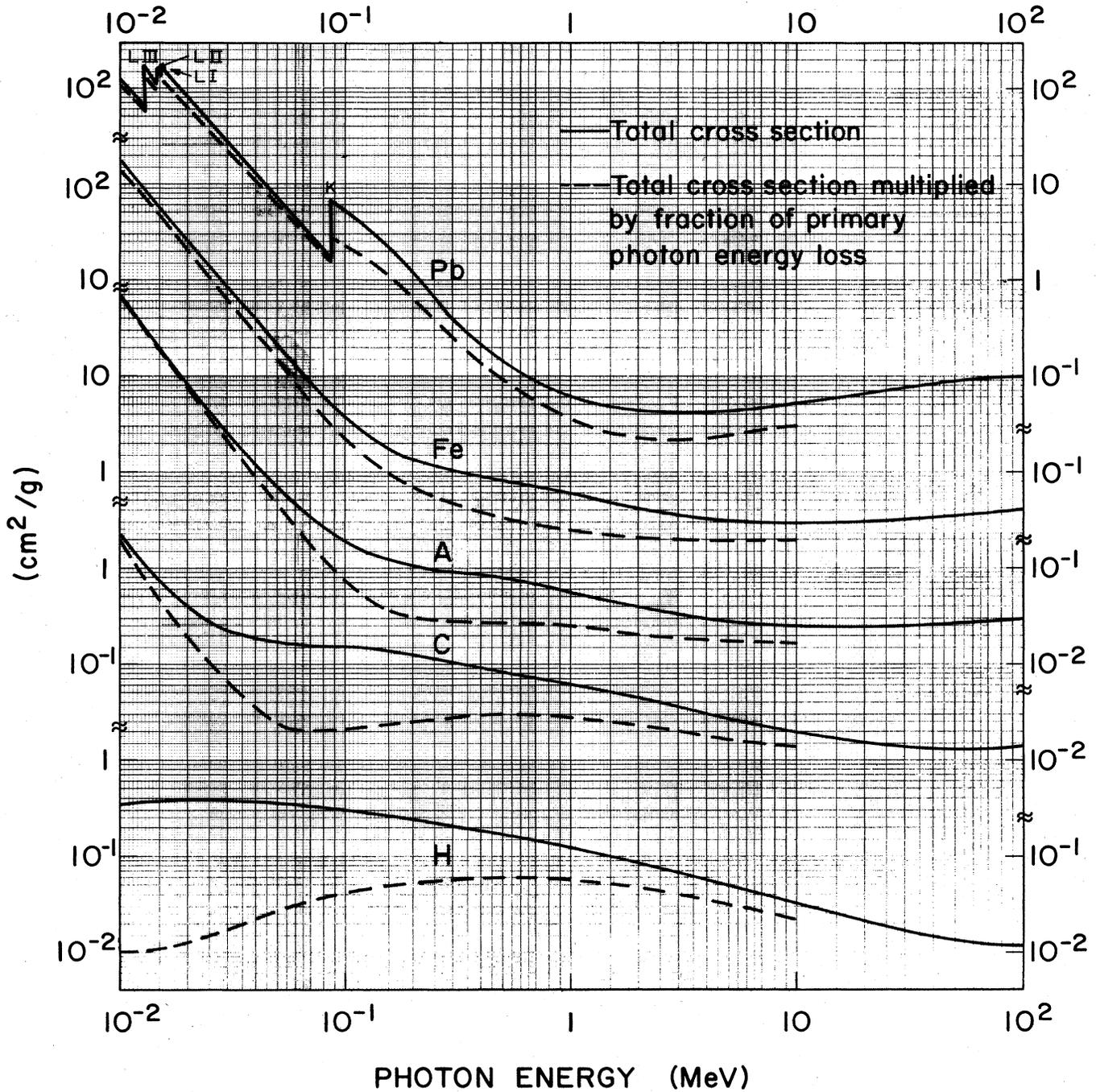


Photon cross section in lead in inverse radiation lengths as a function of photon energy. The intensity of photons can be expressed as $I = I_0 \exp(-\sigma x)$, where σ is read above and x is the path length in radiation lengths. See also figure following.

These figures are adapted from Fig. 3.2 and Fig. 3.3 from Messel and Crawford, *Electron-Photon Shower Distribution Function Tables for Lead, Copper and Air Absorbers*, Pergamon Press, 1970. Messel and Crawford use $L_r(\text{Pb}) = 5.82 \text{ g/cm}^2$, but we have modified the figures to reflect the value given in the Table of Atomic and Nuclear Properties of Materials (following), namely $L_r(\text{Pb}) = 6.4 \text{ g/cm}^2$. Note that the development of electron-photon cascades is approximately independent of absorber when the results are expressed in terms of inverse radiation lengths (i. e., scales on left of plots).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Photon Cross Sections



XBL 743-2662

The photon cross section as a function of photon energy for various substances. The solid curve for each substance gives the total cross section in (cm²/g). The dashed curve is the total cross section multiplied by the fraction of energy deposited by the photon (different from solid curve because of Compton scattering and bremsstrahlung of secondary electrons). See J. H. Hubbell, NSRDS-NBS 29 (1969).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Atomic and Nuclear Properties of Materials*

Material Z	A	Nuclear cross section σ^a [barns]	Nuclear collision length L_{coll}^b [cm]	Absorption length λ^b [cm]	dE/dx min ^c [$\frac{MeV}{g/cm^2}$] [$\frac{MeV}{cm}$]	Radiation length L_{rad}^d [cm]	Density [g/cm ³] [g/l]	Refractive index n_r^e () is for gas () is $(n-1) \times 10^6$ for gas				
H ₂	1	1.01	0.039	43.0	607	790	4.12	0.292	63.05	890	{ 0.0708 (0.090)	{ 1.112 (140)
D ₂	1	2.01	0.074	45.1	273	342	2.07	0.342	126.1	764	0.163	1.128
He	2	4.00	0.134	49.6	397	478	1.94	0.243	94.32	755	{ 0.125 (0.178)	{ 1.024 (35)
Li	3	6.94	0.215	53.6	100.4	120.6	1.65	0.881	82.76	155	0.534	-
Be	4	9.01	0.270	55.4	30.0	36.7	1.61	2.97	65.19	35.3	1.848	-
C	6	12.01	0.340	58.7	≈37.8	49.9	1.78	≈2.76	42.70	≈27.5	≈1.55 ^f	-
N ₂	7	14.01	0.390	59.7	73.8	99.4	1.82	1.47	37.99	47.0	{ 0.808 (1.25)	{ 1.205 (300)
Ne	10	20.18	0.520	64.4	53.7	74.9	1.73	2.08	28.94	24.0	{ 1.207 (0.90)	{ 1.092 (67)
Al	13	26.98	0.650	68.9	25.5	37.2	1.62	4.37	24.01	8.9	2.70	-
A	18	39.95	0.890	74.5	53.2	80.9	1.51	2.11	19.55	14.0	{ 1.40 (1.78)	{ 1.233 (283)
Fe	26	55.85	1.160	79.9	10.2	17.1	1.48	11.6	13.84	1.76	7.87	-
Cu	29	63.54	1.270	83.1	9.3	14.8	1.44	12.9	12.86	1.43	8.96	-
Sn	50	118.69	2.040	96.6	13.2	22.8	1.28	9.4	8.82	1.21	7.31	-
W	74	183.85	2.810	108.6	5.6	10.3	1.17	22.6	6.76	0.35	19.3	-
Pb	82	207.19	3.080	111.7	9.8	18.5	1.13	12.8	6.37	0.56	11.35	-
U	92	238.03	3.380	116.9	≈6.2	12.0	1.09	≈20.7	6.00	≈0.32	≈18.95	-
Air				60.2	50000 ^g	67500 ^g	1.82	0.0022 ^g	36.20	30050 ^g	{ 0.001205 ^g (1.29)	{ 1.000273 ^g (293)
H ₂ O				58.3	58.3	78.8	2.03	2.03	36.08	36.1	1.00	1.33
H ₂ (bubble chamber 26°K) ^h				43.0	≈683	887	4.12	≈0.26	63.05	≈1000	≈0.063 ^h	1.112
D ₂ (bubble chamber 31°K) ^h				45.1	≈322	403	2.07	≈0.29	126.1	≈900	≈0.140 ^h	1.110
H-Ne mixture (50 mole percent) ⁱ				62.9	154.5	215	1.84	0.75	29.70	73.0	0.407	1.092
Propane (C ₃ H ₈) ^j				55.0	134	176	2.28	0.98	45.38	111	{ 0.41 ^j (2.0)	{ 1.25 ^j (1005)
Freon 13B1 (CF ₃ Br) ^j				74.3	≈49.5	73.5	1.52	≈2.3	16.53	≈11	{ ≈1.50 ^j (8.71)	{ 1.238 ^j (750)
Iford emulsion				79.5	23.6	39.1	1.44	5.49	11.02	2.94	3.815	-
NaI				91.9	25.0	41.3	1.32	4.84	9.49	2.59	3.67	1.775
LiF				61.1	23.1	30.7	1.69	4.46	39.25	14.9	2.64	1.394
Polyethylene (CH ₂)				55.7	≈59.6	78.4	2.09	≈1.95	44.78	≈48	0.92-0.95	-
Mylar (C ₅ H ₄ O ₂)				58.5	42.1	56.1	1.91	2.65	39.95	28.7	1.39	-
Polystyrene, scintillator (CH) ^k				57.0	55.2	68.5	1.97	2.03	43.8	42.9	1.032	1.581
Lucite, Plexiglas (C ₅ H ₈ O ₂)				57.7	≈48.9	65.0	1.97	≈2.32	40.55	≈34.5	1.16-1.20	≈1.49
Spark or proportional chamber ^l				0.05%		0.03%	-	0.073		2.7%	0.046	-
Shielding concrete ^m				64.9	26.0	32.2	1.70	4.25	26.7	10.7	2.5	-
CO ₂ ⁿ				60.4	33800	46000	1.82	0.0033	36.2	20210	(1.79) ⁿ	(410) ⁿ
Freon 12 (CCl ₂ F ₂) ⁿ				68.1	13800	20200	1.64	0.0081	23.7	4810	(4.93) ⁿ	(1080) ⁿ
Freon 13 (CClF ₃) ⁿ				66.0	15000	21400	1.70	0.0072	27.15	6380	(4.26) ⁿ	(720) ⁿ
Silica Aerogel ^o				62.3	≈311	430	1.82	≈0.36	30	≈150	0.1-0.3	1.0+0.25 ρ

* Table revised January 1976 by J. Engler and F. Mönning. For details and references, see CERN NP Internal Report 74-1.

- a) σ of neutrons ($\approx \sigma$ of protons) at 20 GeV from Landolt-Börnstein, New Series I, Vol. 5. Energy dependence for all nuclei $\approx 1/2$ percent/GeV (from 5-25 GeV).
- b) $L_{coll} = A/(N \cdot \sigma)$. In the absorption length the elastic scattering is subtracted.
- c) For a minimum-ionizing, singly-charged particle in the material. From W.H. Barkas and M.J. Berger, Tables of Energy Losses and Ranges of Heavy Charged Particles, NASA-SP-3013 (1964).
- d) From Y.S. Tsai, Pair Production and Bremsstrahlung of Charged Leptons, SLAC-PUB-1365 (1974), Table III.6.
- e) Values for solids, or the liquid phase at boiling point, except where noted. Values in parentheses for gaseous phase STP (0°C, 1 atm.), except where noted.
- f) Density variable.
- g) Gas at 20°C.
- h) Density may vary about $\pm 3\%$, depending on operating conditions.
- i) Values for typical working condition with H₂ target: 50 mole percent, 29°K, 7 atm.
- j) Values for typical chamber working conditions: Propane $\sim 57^\circ\text{C}$, 8-10 atm. Freon 13B1 $\sim 28^\circ\text{C}$, 8-10 atm.
- k) Typical scintillator; e.g. PILOT B and NE 102A have an atomic ratio H/C = 1.10.
- l) Values for typical construction: 2 layers 50 μm Cu/Be wires, 8 mm gap, 60% argon, 40% isobutane or CO₂; 2 layers 50 μm Mylar/Aclar foils.
- m) Standard shielding blocks, typical composition O₂ 52%, Si 32.5%, Ca 6%, Na 1.5%, Fe 2%, Al 4% plus reinforcing iron bars. Attenuation length $l = 115 \pm 5$ g/cm², also valid for earth (typical $\rho = 2.15$) from CERN-LRL-RHEL Shielding exp. UCRL 17841 (1968).
- n) Used in Čerenkov counters, value at 26°C and 1 atm. Indices of refraction from E.R. Hayes, R.A. Schluter, and A. Tamosaitis, ANL-6916 (1964).
- o) $n(\text{SiO}_2) + 2n(\text{H}_2\text{O})$ used in Čerenkov counters, ρ = density in g/cm³. From M. Cantin et al., Nucl. Instr. Meth. 118, 177 (1974).

ELECTROMAGNETIC RELATIONS

Maxwell's Equations

Quantity	CGS (statcoul., statamp., sec cm ⁻¹)	MKSA (coul., amp., ohm)
Potentials:	$V = \frac{\sum \text{charges}}{r}$, $\vec{A} = \frac{1}{c} \text{curl} \sum \frac{\vec{I}}{r}$; c = speed of light in vacuum	$V = \frac{1}{4\pi\epsilon_0} \text{charges} \frac{q}{r}$, $\vec{A} = \frac{\mu_0}{4\pi} \text{curl} \sum \frac{\vec{I}}{r}$; $\epsilon_0 = \frac{1}{36\pi} 10^{-9}$ MKSA, $\mu_0 = 4\pi 10^{-7}$ MKSA
Fields:	$\vec{E} = -\vec{\nabla}V, \vec{B} = \vec{\nabla} \times \vec{A}$	$\vec{E} = -\vec{\nabla}V, \vec{B} = \vec{\nabla} \times \vec{A}$
Materials:	$\vec{D} = \epsilon \vec{E}, \vec{B} = \mu \vec{H}$	$\vec{D} = \epsilon \vec{E}, \vec{B} = \mu \vec{H}$
Force:	$\vec{F} = q(\vec{E} + \frac{\vec{v}}{c} \times \vec{B})$	$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$
Maxwell:	$\vec{\nabla} \cdot \vec{E} = 4\pi\rho$, $\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$, $\vec{\nabla} \cdot \vec{B} = 0$, $\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$	$\vec{\nabla} \cdot \vec{D} = \rho$, $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$, $\vec{\nabla} \cdot \vec{B} = 0$, $\vec{\nabla} \times \vec{H} = \vec{j} + \frac{\partial \vec{D}}{\partial t}$

Impedances: Alternating Currents (MKSA)

Ohm's law: $V = ZI, V = V_0 e^{i\omega t}$

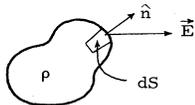
1. Impedance of self-inductance of inductance L : $Z = i\omega L$
2. Impedance of a capacitor of capacitance C : $Z = \frac{1}{i\omega C}$
3. Impedance of a flat conductor of width w at high frequency:
 $Z = \frac{(1+i)\rho}{w\delta}$;
 $\rho = \text{resistivity} \begin{cases} \sim 1.7 \times 10^{-8} \Omega \text{m for Cu} \\ \sim 2.8 \times 10^{-8} \Omega \text{m for Al} \end{cases}$
 $\delta = \text{effective skin depth}$
 $= \sqrt{\frac{\rho}{\pi\nu\mu}} \approx \frac{6.6 \text{ cm}}{\sqrt{\nu(\text{sec}^{-1})}}$ for Cu

Capacitance C and Inductance L per Unit Length (MKSA)

1. For flat plates of width w, separated by $d \ll w$:
 $C = \frac{\epsilon w}{d}$; $L = \mu \frac{d}{w}$
2. For coax cable of interior and exterior radii r_1 and r_2 :
 $C = \frac{2\pi\epsilon}{\ln(r_2/r_1)}$; $L = \frac{\mu}{2\pi} \ln(r_2/r_1)$;
 $\epsilon = \text{dielectric constant} \begin{cases} 2 \text{ to } 6 \text{ for plastics} \\ 4 \text{ to } 8 \text{ for porcelain, glasses} \end{cases}$
 $\mu = \text{magnetic susceptibility}$

Integral Forms (MKSA)

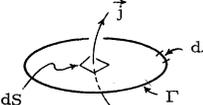
1. Gauss' theorem:



$$\int_{\text{surface}} \vec{E} \cdot \hat{n} \, dS = \int_{\text{volume}} \rho/\epsilon_0 \, dV$$

($\rho = \text{charge/volume}$)

2. Ampere's law:



$$\int_{\Gamma} \vec{H} \cdot d\vec{l} = \int_{\text{surface}} \vec{j} \cdot \hat{n} \, dS$$

($\vec{j} = \text{surface current density}$)

Transmission Lines (No Loss) (MKSA)

Velocity = $1/\sqrt{LC} = 1/\sqrt{\mu\epsilon}$

Impedance = $\sqrt{L/C}$

L, C are inductance and capacitance per unit length

Synchrotron Radiation (CGS)

Energy loss/revolution = $\frac{4\pi}{3} \frac{e^2}{\rho} \beta^3 \gamma^4$, $\rho = \text{orbit radius}$.

For electrons ($\beta \approx 1$), $\frac{\Delta E}{\text{rev.}} (\text{MeV}) = 0.0885 [E(\text{GeV})]^4 / \rho(\text{meter})$.

Critical frequency: $\omega_c = 3\gamma^3 c/\rho$

See J. D. Jackson, Classical Electrodynamics, 2nd Ed., John Wiley & Sons, New York, 1975, for more formulae and details. (Prepared April 1974; revised April 1976.)

RADIOACTIVITY AND RADIATION PROTECTION

Unit of activity = Curie:
1 Ci = 3.7×10^{10} disintegrations/sec
Unit of exposure dose for x and γ radiation = Roentgen:
1 R = $1 \text{ esu/cm}^2 = 87.8 \text{ erg/g} (5.49 \times 10^7 \text{ MeV/g})$ of air
Unit of absorbed dose = rad:
1 rad = $100 \text{ erg/g} (6.25 \times 10^7 \text{ MeV/g})$ in any material
Unit of dose equivalent (for protection) = rem:
rems (Roentgen equivalents for man) = rads \times QF,
where QF (quality factor) depends upon the type of radiation and other factors. For γ rays and HE protons, QF ≈ 1 ; for thermal neutrons, QF ≈ 3 ; for fast neutrons, QF ranges up to 10; and for α particles and heavy ions, QF ranges up to 20.
Maximum permissible occupational dose for the whole body:
5 rem/year (or ≈ 100 millirem/week)
Fluxes (per cm²) to liberate 1 rad in carbon:
 3.5×10^7 minimum ionizing singly charged particles
 1.0×10^9 photons of 1 MeV energy
(These fluxes are correct to within a factor of 2 for all materials.)
Natural background: 120 to 130 millirem/year
cosmic radiation (charged particles + neutrons) ~ 25
cosmic radiation (γ rays) ~ 25
radiation from rocks and air (γ rays) ~ 73 } mrem/yr
Cosmic ray background in counters: $\sim 1/\text{min/cm}^2/\text{ster}$

C.M. ENERGY AND MOMENTUM VS. BEAM MOMENTUM

$$E_{cm} dE_{cm} = m_p dT_{beam} = m_p v_{beam} dP_{beam} \approx m_p dP_{beam}$$

PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)				PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)				PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)			
	YP ep	TP	Kp	pp	YP ep	TP	Kp	pp		YP ep	TP	Kp	pp	YP ep	TP	Kp	pp		YP ep	TP	Kp	pp	YP ep	TP	Kp	pp
0	939	1078	1432	1877	0	0	0	0	1500	1922	1930	2022	2254	732	729	696	624	3.0	2.56	2.61	2.77	1.10	1.08	1.02		
20	958	1079	1432	1877	20	17	13	10	1520	1932	1940	2031	2261	738	735	702	631	3.2	2.63	2.68	2.83	1.14	1.12	1.06		
40	977	1083	1433	1877	38	35	26	20	1540	1942	1950	2039	2268	744	741	709	637	3.4	2.70	2.75	2.89	1.18	1.16	1.10		
60	996	1089	1434	1877	56	52	39	30	1560	1951	1959	2048	2275	750	747	715	643	3.6	2.77	2.82	2.96	1.22	1.20	1.14		
80	1015	1096	1436	1878	74	68	52	40	1580	1961	1969	2057	2282	756	753	721	650	3.8	2.83	2.88	3.02	1.26	1.24	1.18		
					T(P1) = PBEAM - 59 MEV																					
100	1033	1105	1439	1879	91	85	65	50	1600	1970	1978	2065	2289	762	759	727	656	4.0	2.90	2.95	3.08	1.29	1.27	1.22		
120	1051	1116	1441	1880	107	101	78	60	1620	1980	1988	2074	2296	768	765	733	662	4.2	2.96	3.01	3.14	1.33	1.31	1.26		
140	1069	1127	1445	1882	123	117	91	70	1640	1989	1997	2083	2304	773	770	739	668	4.4	3.03	3.07	3.19	1.36	1.34	1.29		
160	1087	1139	1449	1883	138	132	104	80	1660	1999	2006	2091	2311	779	776	745	674	4.6	3.09	3.13	3.25	1.40	1.38	1.33		
180	1104	1152	1453	1885	153	147	116	90	1680	2008	2016	2100	2318	785	782	751	680	4.8	3.15	3.19	3.31	1.43	1.41	1.36		
					T(P1) = PBEAM - 92 MEV																					
200	1121	1165	1457	1887	167	161	129	99	1700	2018	2025	2109	2325	791	788	756	686	5.0	3.21	3.25	3.36	1.46	1.44	1.40		
220	1137	1178	1462	1889	182	175	141	109	1720	2027	2034	2117	2332	796	793	762	692	5.2	3.26	3.30	3.40	1.49	1.47	1.43		
240	1154	1192	1468	1892	195	189	153	119	1740	2036	2043	2126	2339	802	799	768	698	5.4	3.31	3.35	3.45	1.51	1.49	1.45		
260	1170	1206	1474	1894	209	202	166	129	1760	2045	2053	2134	2346	807	805	774	704	5.6	3.36	3.40	3.50	1.54	1.52	1.48		
280	1186	1219	1480	1897	222	215	178	138	1780	2054	2062	2143	2353	813	810	779	710	5.8	3.41	3.45	3.57	1.57	1.55	1.51		
					T(P1) = PBEAM - 107 MEV																					
300	1201	1233	1486	1900	234	228	189	148	1800	2064	2071	2151	2360	818	816	785	716	6.0	3.46	3.50	3.61	1.60	1.58	1.54		
320	1217	1247	1493	1903	247	241	201	158	1820	2073	2080	2159	2367	824	821	791	721	6.2	3.51	3.55	3.67	1.63	1.61	1.57		
340	1232	1261	1500	1906	259	253	213	167	1840	2082	2089	2168	2374	829	827	796	727	6.4	3.56	3.60	3.72	1.66	1.64	1.60		
360	1247	1274	1507	1910	271	265	224	177	1860	2091	2098	2176	2381	835	832	802	733	6.6	3.61	3.65	3.77	1.69	1.67	1.63		
380	1262	1288	1514	1913	282	277	235	186	1880	2100	2107	2184	2388	840	837	808	739	6.8	3.66	3.70	3.82	1.72	1.70	1.66		
					T(P1) = PBEAM - 115 MEV																					
400	1277	1302	1522	1917	294	288	247	196	1900	2108	2115	2193	2395	845	843	813	744	7.0	3.71	3.75	3.87	1.75	1.73	1.69		
420	1292	1315	1530	1921	305	300	258	205	1920	2117	2124	2201	2402	851	848	818	750	7.2	3.76	3.80	3.92	1.78	1.76	1.72		
440	1306	1329	1538	1925	316	311	268	214	1940	2126	2133	2209	2409	856	853	824	756	7.4	3.81	3.85	3.97	1.81	1.79	1.75		
460	1320	1342	1546	1929	327	322	279	224	1960	2135	2142	2217	2416	861	859	829	761	7.6	3.86	3.90	4.02	1.84	1.82	1.78		
480	1335	1356	1554	1933	337	332	290	233	1980	2144	2150	2226	2423	867	864	835	767	7.8	3.91	3.95	4.07	1.87	1.85	1.81		
					T(P1) = PBEAM - 120 MEV																					
500	1349	1369	1563	1938	348	343	300	242	2000	2153	2159	2234	2430	872	869	840	772	8.0	3.96	4.00	4.11	1.90	1.88	1.84		
520	1362	1382	1572	1943	358	353	310	251	2020	2161	2168	2242	2437	877	874	845	778	8.2	4.01	4.05	4.16	1.93	1.91	1.87		
540	1376	1395	1580	1947	368	363	321	260	2040	2170	2176	2250	2444	882	879	851	783	8.4	4.06	4.10	4.21	1.96	1.94	1.90		
560	1390	1408	1589	1952	378	373	331	269	2060	2179	2185	2258	2451	887	885	856	789	8.6	4.11	4.15	4.26	1.99	1.97	1.93		
580	1403	1421	1598	1957	388	383	341	278	2080	2187	2194	2266	2458	892	890	861	794	8.8	4.16	4.20	4.31	2.02	2.00	1.96		
					T(P1) = PBEAM - 123 MEV																					
600	1416	1434	1607	1962	397	393	350	287	2100	2196	2202	2274	2465	897	895	866	799	9.0	4.21	4.25	4.36	2.05	2.03	1.99		
620	1430	1446	1616	1967	407	402	360	296	2120	2204	2211	2282	2472	902	900	872	805	9.2	4.26	4.30	4.41	2.08	2.06	2.02		
640	1443	1459	1625	1973	416	412	370	304	2140	2213	2219	2290	2479	907	905	877	810	9.4	4.31	4.35	4.46	2.11	2.09	2.05		
660	1456	1472	1634	1978	425	421	379	313	2160	2221	2227	2298	2486	912	910	882	815	9.6	4.36	4.40	4.51	2.14	2.12	2.08		
680	1468	1484	1644	1984	434	430	388	322	2180	2230	2236	2306	2493	917	915	887	821	9.8	4.41	4.45	4.56	2.17	2.15	2.11		
					T(P1) = PBEAM - 125 MEV																					
700	1481	1496	1653	1989	443	439	397	330	2200	2238	2244	2314	2500	922	920	892	826	10.0	4.46	4.50	4.61	2.20	2.18	2.14		
720	1496	1509	1662	1995	452	448	406	339	2220	2246	2253	2322	2507	927	925	897	831	10.2	4.51	4.55	4.66	2.23	2.21	2.17		
740	1506	1521	1671	2001	461	457	415	347	2240	2255	2261	2330	2514	932	930	902	836	10.4	4.56	4.60	4.71	2.26	2.24	2.20		
760	1519	1533	1681	2007	470	465	424	355	2260	2263	2269	2338	2520	937	934	907	841	10.6	4.61	4.65	4.76	2.29	2.27	2.23		
780	1531	1545	1690	2013	478	474	433	364	2280	2271	2277	2346	2527	942	939	912	846	10.8	4.66	4.70	4.81	2.32	2.30	2.26		
					T(P1) = PBEAM - 127 MEV																					
800	1543	1557	1699	2019	486	482	442	372	2300	2280	2286	2353	2534	947	944	917	852	11.0	4.71	4.75	4.86	2.35	2.33	2.29		
820	1555	1569	1709	2025	495	490	450	380	2320	2288	2294	2361	2541	951	949	922	857	11.2	4.76	4.80	4.91	2.38	2.36	2.32		
840	1567	1580	1718	2031	503	499	459	388	2340	2296	2302	2369	2548	956	954	927	862	11.4	4.81	4.85	4.96	2.41	2.39	2.35		
860	1579	1592	1728	2037	511																					

PERIODIC TABLE OF THE ELEMENTS

IA		IIA				1 H 1.008						IIIA		IVA		VA		VIA		VIIA		2 He 4.0026	
3 Li 6.94	4 Be 9.0122											5 B 10.81	6 C 12.011	7 N 14.0067	8 O 15.9994	9 F 18.9984	10 Ne 20.17						
11 Na 22.9898	12 Mg 24.305											13 Al 26.9815	14 Si 28.086	15 P 30.9738	16 S 32.06	17 Cl 35.453	18 Ar 39.948						
19 K 39.09	20 Ca 40.08	21 Sc 44.956	22 Ti 47.90	23 V 50.941	24 Cr 51.996	25 Mn 54.9380	26 Fe 55.847	27 Co 58.9332	28 Ni 58.71	29 Cu 63.55	30 Zn 65.38	31 Ga 69.72	32 Ge 72.59	33 As 74.9216	34 Se 78.96	35 Br 79.904	36 Kr 83.80						
37 Rb 85.47	38 Sr 87.62	39 Y 88.906	40 Zr 91.22	41 Nb 92.906	42 Mo 95.94	43 Tc 98.9062	44 Ru 101.07	45 Rh 102.906	46 Pd 106.4	47 Ag 107.868	48 Cd 112.40	49 In 114.82	50 Sn 118.69	51 Sb 121.75	52 Te 127.60	53 I 126.9045	54 Xe 131.30						
55 Cs 132.905	56 Ba 137.34	57-71 Rare Earths	72 Hf 178.49	73 Ta 180.947	74 W 183.85	75 Re 186.2	76 Os 190.2	77 Ir 192.22	78 Pt 195.09	79 Au 196.967	80 Hg 200.59	81 Tl 204.37	82 Pb 207.2	83 Bi 208.981	84 Po (209)	85 At (210)	86 Rn (222)						
87 Fr (223)	88 Ra 226.025	89- Acti- nides	104 (261)	105 (262)	106 (263)																		

57 La 138.91	58 Ce 140.12	59 Pr 140.908	60 Nd 144.24	61 Pm (145)	62 Sm 150.4	63 Eu 151.96	64 Gd 157.25	65 Tb 158.925	66 Dy 162.50	67 Ho 164.930	68 Er 167.26	69 Tm 168.934	70 Yb 173.04	71 Lu 174.97
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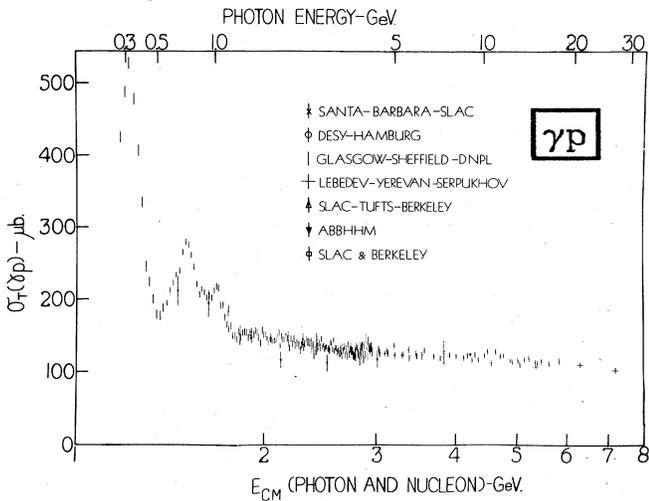
Rare earths (Lanthanide series).

89 Ac (227)	90 Th 232.038	91 Pa 231.036	92 U 238.03	93 Np 237.048	94 Pu (244)	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (254)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (260)
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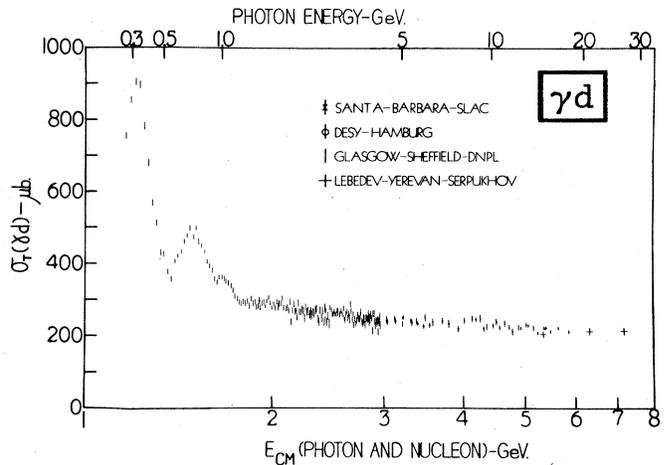
Actinide series

Numbers in parentheses are mass numbers of most stable isotope of that element. Adapted from the *Handbook of Chemistry and Physics, 1975-76.* (Particle Data Group update, April 1976.)

CROSS SECTION PLOTS

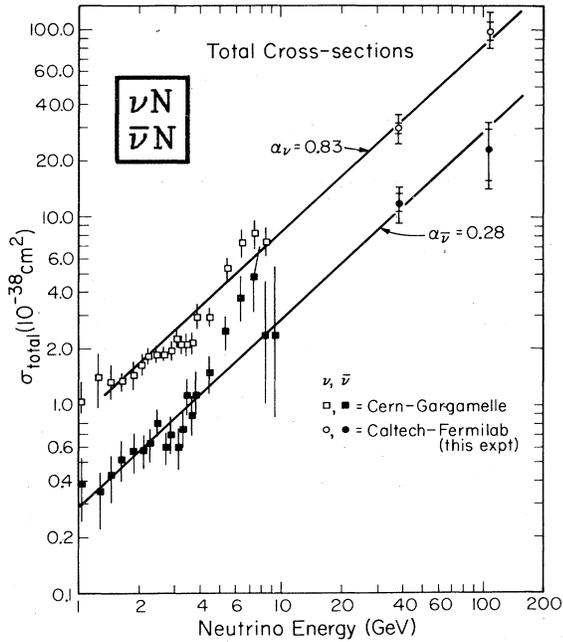


γp total cross section versus photon energy (top scale) and photon-plus-nucleon total center-of-mass energy (lower scale). Courtesy Gething M. Lewis, Glasgow.

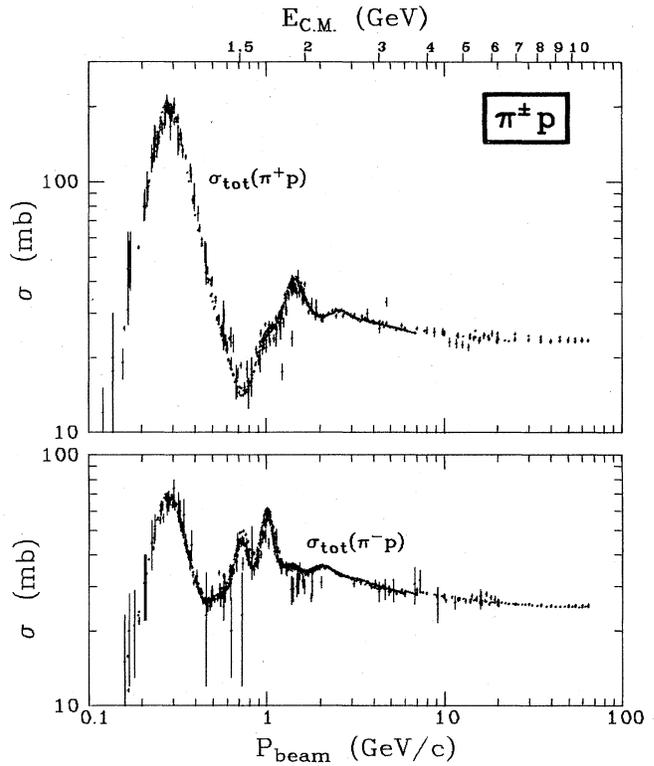


γd total cross section versus photon energy (top scale) and photon-plus-single-nucleon total center-of-mass energy (lower scale). Courtesy Gething M. Lewis, Glasgow.

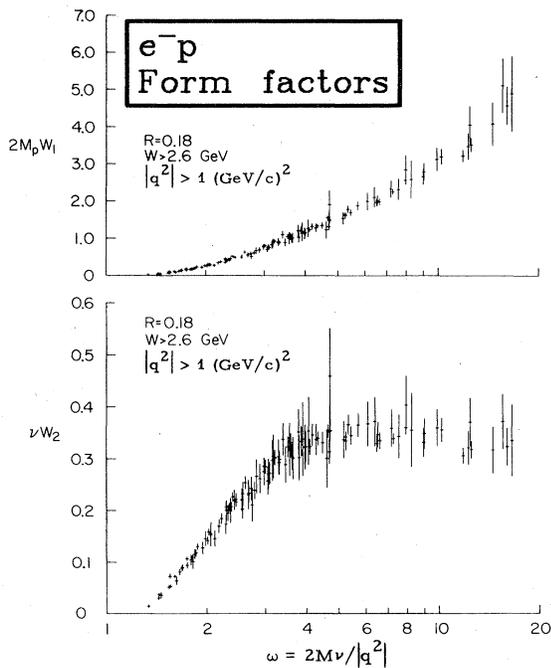
CROSS SECTION PLOTS (Cont'd)



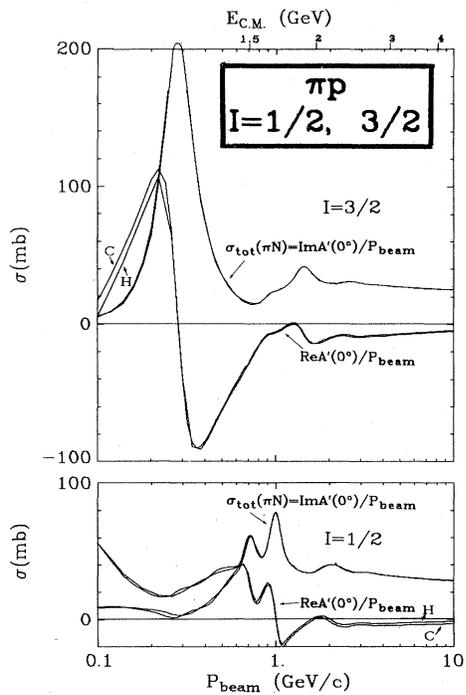
Neutrino and anti-neutrino total cross sections per nucleon [reproduced, with permission, from B.C. Barish et al., Phys. Rev. Lett. 35, 1316 (1975)]. The data are consistent with the form $\sigma_{\text{tot}} = \alpha E_{\nu, \bar{\nu}}$.



$\pi^\pm p$ total cross-section data from the Particle Data Group compilation "nN Two-Body Scattering Data," LBL-63 (1973).

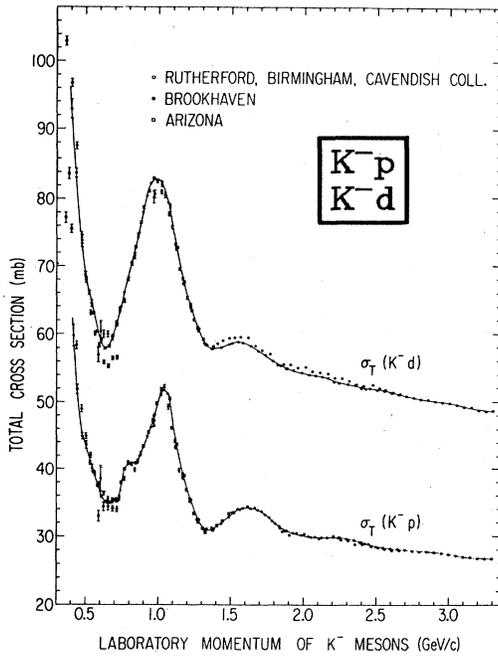


Inelastic electroproduction form factors for a virtual space-like photon scattering on a proton, from SLAC-MIT, Phys. Rev. D5, 528 (1972). $R = \sigma_L/\sigma_T$, the ratio of the longitudinal to transverse cross sections. See L. Hand, Phys. Rev. 129, 1834 (1963).

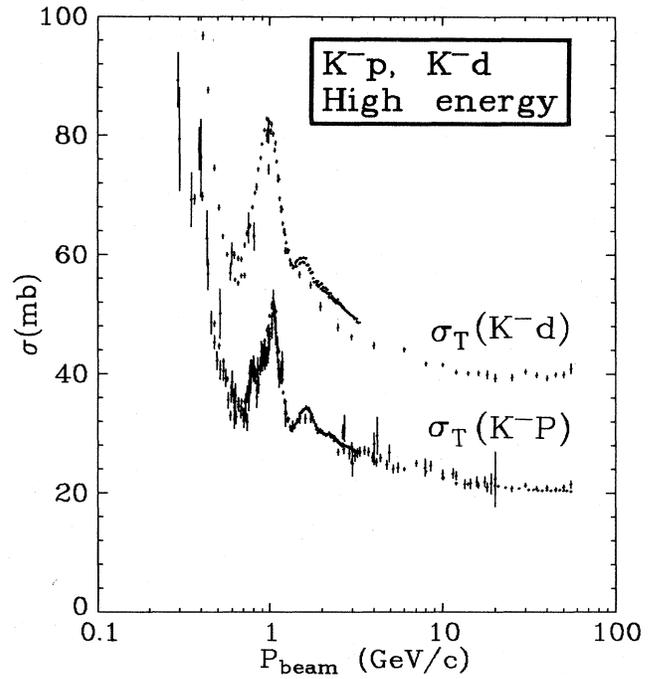


Interpolations of πN total cross sections for $I=3/2$ and $1/2$, and the corresponding real parts of the forward amplitudes as calculated from dispersion relations by A. A. Carter and J. R. Carter (RHEL ppt. RL-73-024, 1973; labeled C above) and by G. Hohler and H. P. Jakob (priv. comm., 1972; labeled H above). The normalization of the curves for each value of I is such that the sum of their squares divided by 19.6 gives $d\sigma/dt$ at 0° in $\text{mb}/(\text{GeV}/c)^2$.

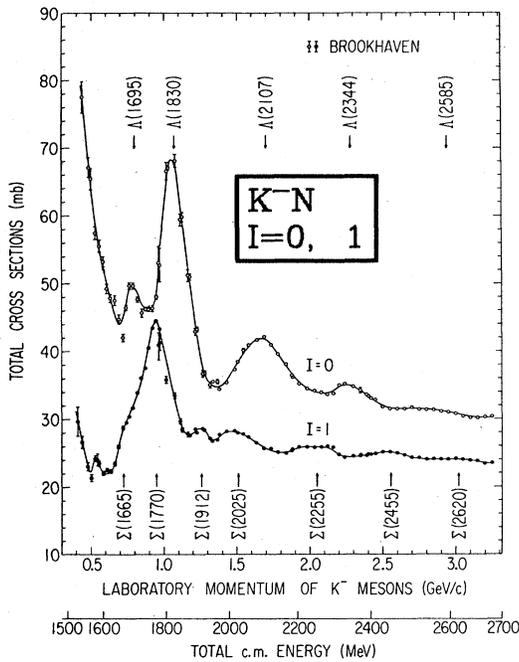
CROSS SECTION PLOTS (Cont'd)



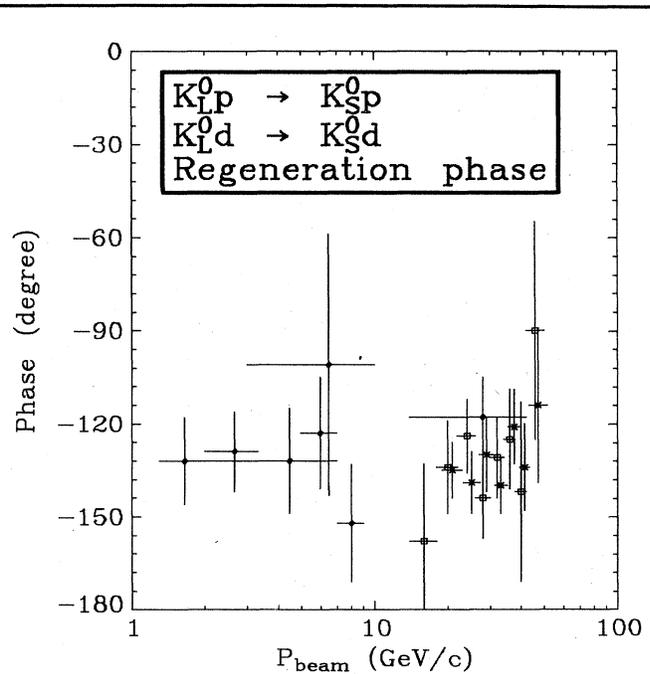
K^-p and K^-d total cross-section data compiled by Li et al., Proc. 1973 Purdue Conf. on Baryon Resonances. The solid curve passes through the Brookhaven data.



K^-p and K^-d total cross-section data. Compilation sources: E. Bracci et al., CERN/HERA 72-2, K^-p ; G.R. Lynch, K^-d (<3 GeV/c); Particle Data Group, K^-d (>3 GeV/c). The new BNL data below 1 GeV/c are not included.

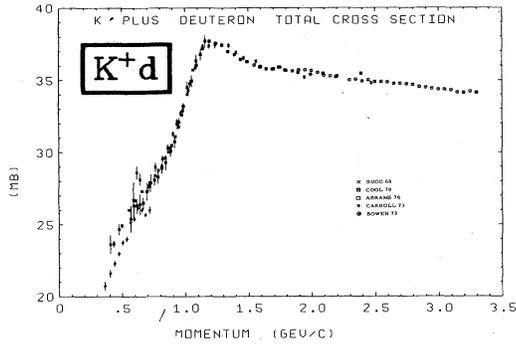
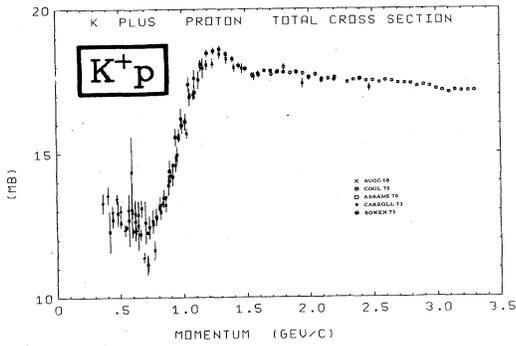


K^-N total cross sections for $I=0$ and $I=1$ below 3.3 GeV/c. Compiled and unfolded by Li et al., Proc. 1973 Purdue Conf. on Baryon Resonances.

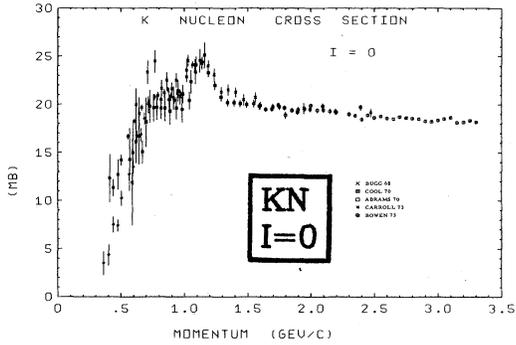


Phases of forward amplitudes for $K_{Lp}^0 \rightarrow K_{Sp}^0$ (\diamond, \square) and $K_{Ld}^0 \rightarrow K_{Sd}^0$ (*). The deuterium data are shifted 1 GeV/c higher to avoid overlap with the proton data. \diamond -p, compiled by Particle Data Group, LBL-55; \square -p, V. Birnlev et al; JINR-E1-6851 (1972); *-d, K.-F. Albrecht et al., Phys. Lett. 48B, 257 (1974).

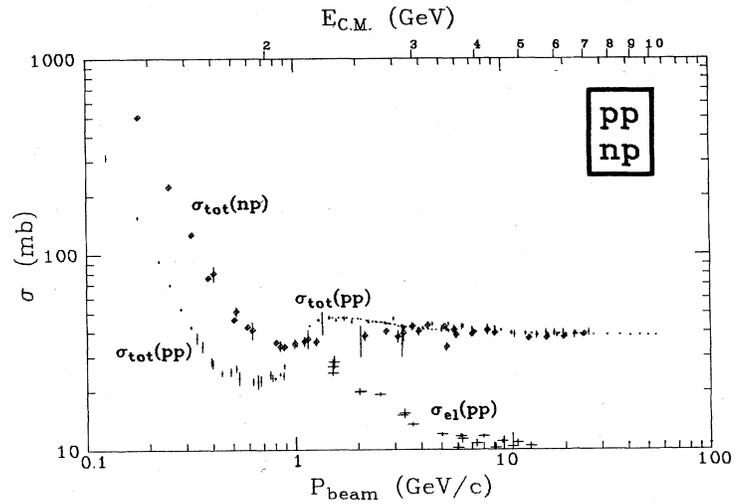
CROSS SECTION PLOTS (Cont'd)



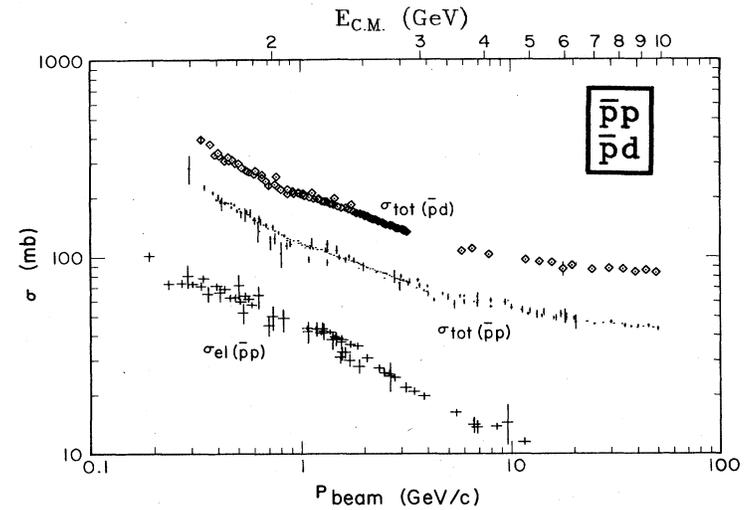
Compilation of recent K^+p and K^+d total cross-section measurements. References can be found in the Baryon Data Card Listings.



Total cross-section for isospin zero KN system. Unfolding of the BUGG 68 and BOWEN 70 and 73 data was done by G. R. Lynch (as in Proc. of 1970 Duke Conference). Tables of σ_0 were provided by the BNL authors. Lynch and BNL use the same method of unfolding; the BOWEN 73 unfolded distribution is obtained by a different method (see plot in Z^0 mini-review in the Baryon Data Card Listings).



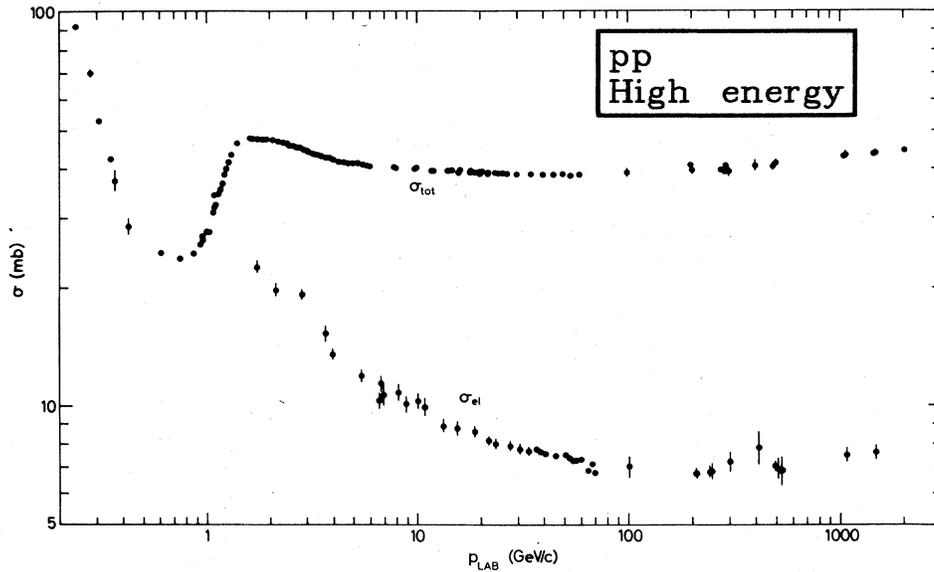
pp and np cross sections from Particle Data Group, "NN and ND Interactions -- A Compilation", UCRL-20 000 NN (August 1970); some points at higher energies added since original compilation.



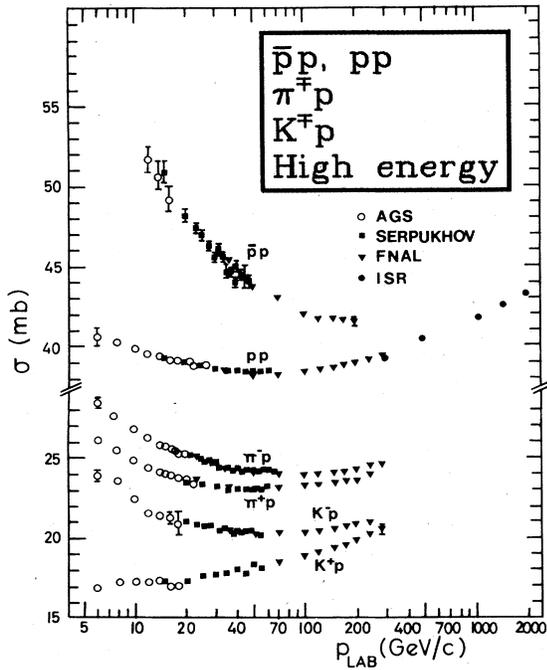
XBL 743-2661

$\bar{p}p$ and $\bar{p}d$ cross sections from Particle Data Group, "A Compilation of NN and ND Reactions", LBL-58 (1972); some points added since original compilation.

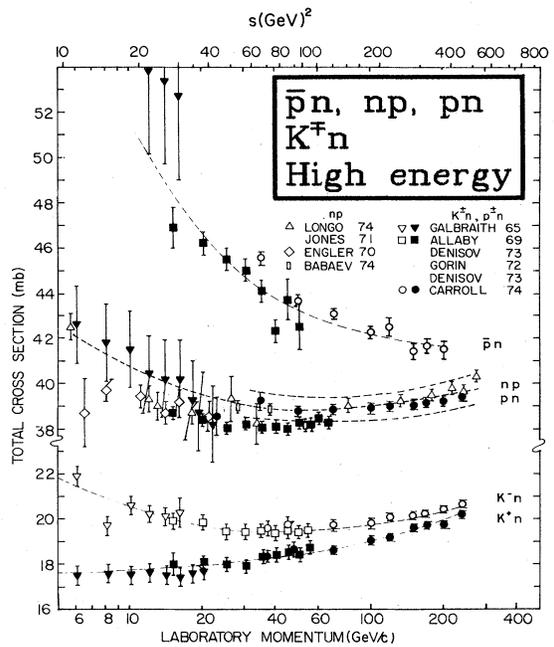
CROSS SECTION PLOTS (Cont'd)



Total and elastic pp cross-section data compiled by U. Amaldi, CERN.



$\bar{p}p$, pp , π^+p , π^-p , K^+p , and K^-p total cross sections versus $s(\approx 2m_p p_{lab})$, as compiled by U. Amaldi, CERN.



$\bar{p}n$, np , pn , K^+n , and K^-n total cross sections versus $s(\approx 2m_p p_{lab})$, as compiled by G. Giacomelli, CERN.

DATA CARD LISTINGS

Illustrative Key

Name of particle as it appears in table. **XX(1200)**

Arrow indicates this particle omitted from table. **OMITTED FROM TABLE**

Quantity tabulated below. **74 XX(1200) MASS (MEV)**

Code for quantity tabulated (M=mass, W=width, etc.)

Symbols used to key together data card and related comments.

Number of events above background.

Measured values (parentheses indicate value not used in average).

± Error in measured value (- field blank if error symmetric; parentheses on error only indicate data not used in average due to problems with error estimation).

Average value (and error) of quantity measured.

Vertical bar indicates average; width of horizontal bar on top is error (scaled) in average.

Value and error for each experiment.

Partial decay mode (labeled by P_i).

Branching ratio (labeled by R_j).

Value (and error) of quantity measured, as determined from constrained fit (using all measured branching ratios for this particle).

References listed by year, then author.

Abbreviated reference form used on data cards above.

Journal, report, preprint, etc. (see abbreviations on next page).

Particle name, and quantum numbers (if known). **XX MESON (1200, J^{PC}= -) I=1**

Particle code (for internal use only). **ORIGINALLY CALLED XXX**

General comments on particle. **LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION**

Abbreviated reference for this result; full reference given below.

Measurement technique (see abbreviations on next page.)

Charge(s) of particle detected.

Reaction producing particle, or comments.

Date this result punched (asterisk indicates result added or changed since previous edition).

Scale factor > 1 indicates inconsistent data.

Ideogram to display inconsistent data; curve is sum of Gaussians, one for each experiment (area of Gaussian = 1/error; width of Gaussian = ± error).

Contribution of experiment to χ^2 (if no entry present, experiment not used in calculating χ^2 or scale factor because of large error).

Representative masses of decay products (used for calculating last column of Particle Property Tables).

Author(s).

Quantum number determinations in this reference.

Institution(s) of author(s) (see abbreviations on next page).

REFERENCES FOR XX(1200)

A. MERRILL (TORINO+CERN)/TJP
 B. LYNCH (BNL)
 N. PIERCE (LRL)
 D. FENNER, B. BEANE (NYSE+AMEX)
 J. SMITH (SLAC)

Illustrative Key (cont'd)

Abbreviations

Journals

APAH	Acta Phys. Acad. Hungarica
ADVP	Advances in Physics
ANP	Annals of Physics
ARNNS	Annual Reviews of Nuclear Science
BAPS	Bulletin of the Amer. Phys. Soc.
JETP	English Transl. of Soviet Physics JETP
JETPL	Letters of Soviet Physics JETP
LNC	Letters to Nuovo Cimento
NC	Nuovo Cimento
NP	Nuclear Physics
PL	Physics Letters
PN	Particles and Nuclei
PPSL	Proc. of the Phys. Soc. of London
PR	Physical Review
PRL	Physical Review Letters
PRSL	Proc. of the Royal Soc. of London
RMI	Reviews of Modern Physics
SJNP	Soviet Journal of Nuclear Physics
ZPHV	Zeitschrift für Physik

Measurement techniques

ASPK	Automatic spark chambers	HYBR	Hybrid: BC + electronics
CC	Cloud chamber	IPWA	Energy-independent PWA
CNTR	Counters	MMS	Missing mass spectrometer
DASP	DESY double-arm spectrometer	MPWA	Model-dependent PWA
DBC	Deuterium bubble chamber	OSPK	Optical spark chamber
DPWA	Energy-dependent PWA	PBC	Propane bubble chamber
ELEC	Electronic combination	PLUT	DESY PLUTO detector
EMUL	Emulsions	PWA	Partial-wave analysis
PBC	Freon bubble chamber	RVUE	Review of previous data
FRAB	ADONE BB Group detector	SMAG	SPEAR magnetic detector
FRAG	ADONE YY Group detector	SPEC	Spectrometer
FRAM	ADONE MEA Group detector	SPRK	Spark chamber
HBC	Hydrogen bubble chamber	STRC	Streamer chamber
HEBC	Helium bubble chamber	WIRE	Wire chamber
HLBC	Heavy liquid bubble chamber	XEBC	Xenon bubble chamber

Conferences

Conferences are referred to by the location in which they were held (e.g., DUBNA, BOULDER, LUND, etc.).

Institutions

AACH	TECHNISCHE UNIV. AACHEN	AACHEN, GERMANY	LUND	UNIV. I LUND	LUND, SWEDEN
AARH	AARHUS UNIV.	AARHUS, DENMARK	MADR	JUNTA DE ENERGIA NUCLEAR	MADRID, SPAIN
AERE	ATOMIC ENERGY RES. ESTAB.	HARWELL, BERKS., ENGLAND	MADU	UNIV. AUTONOME DE MADRID	MADRID, SPAIN
ALBA	STATE UNIV. OF NEW YORK AT ALBANY	ALBANY, N. Y., USA	MANH	MANHATTAN COLLEGE	NEW YORK, N. Y., USA
AMST	UNIV. OF AMSTERDAM	AMSTERDAM, NETHERLANDS	MANI	UNIV. MANITOWA	MANITOWA, CANADA
ANKA	MIDDLE EAST TECHNICAL UNIV.	ANKARA, TURKEY	MANZ	UNIV. MAINZ	MAINZ, GERMANY
ANL	ARGONNE NAT. LAB.	ARGONNE, ILL., USA	MASA	UNIV. OF MASSACHUSETTS	AMHERST, MASS., USA
ARIZ	UNIV. OF ARIZONA	TUCSON, ARIZ., USA	MASB	BOSTON MASS.	BOSTON, MASS., USA
ATEA	NUCLEAR RES. CENTRE DEMOKRITOS	ATHENS, GREECE	MCGI	MCGILL UNIV.	MONTREAL, CANADA
ATHU	UNIV. OF ATHENS	ATHENS, GREECE	MCHS	UNIV. MANCHESTER	MANCHESTER, ENGLAND
BARC	UNIV. DE BARCELONA	BARCELONA, SPAIN	MCHC	MOUNT HOLYOKE COLL.	SOUTH HADLEY, MASS., USA
BARI	UNIV. DI BARI	BARI, ITALY	MICH	UNIV. OF MICHIGAN	ANN ARBOR, MICH., USA
BELG	INST. INTERUNIV. DES SCI. NUC.	BRUXELLES, BELGIUM	MILA	UNIV. DI MILANO	MILANO, ITALY
BERG	FYSISK INSTITUTT	BERGEN, NORWAY	MINN	UNIV. OF MINNESOTA	MINNEAPOLIS, MINN., USA
BERL	INST. HOCHENERGIEPHYS. CAW	ZELTEN/BERLIN, DDR	MIDA	MIAMI UNIV.	OXFORD, OHIO, USA
BERN	UNIV. BERNE	BERN, SWITZERLAND	MIEI	MASSACHUSETTS INST. OF TECHNOLOGY	CAMBRIDGE, MASS., USA
BGNA	UNIV. DI BOLOGNA	BOLGONA, ITALY	MODI	ISTITUTO DI FISICA DELLA UNIVERSITA	MONTREAL, CANADA
BING	STATE UNIV. OF NEW YORK AT BINGHAMTON	BINGHAMTON, N. Y., USA	MODE	UNIV. DE MONTREAL	MONTREAL, CANADA
BIRM	SIRKINGHAM UNIV.	BIRMINGHAM, ENGLAND	MONS	UNIV. DE MONTPELLIER	MONS, BELGIUM
BNL	BROOKHAVEN NATIONAL LAB.	UPTON, L.I., N. Y., USA	MONT	UNIV. OF MONTPELLIER	MONTPELLIER, FRANCE
BNHR	NIELS BOHR INSTITUTE	COPENHAGEN, DENMARK	MPH	MAX-PLANCK-INST. FUR PHYS.-ASTROPHYS.	HEIDELBERG, GERMANY
BONN	UNIV. BONN	BONN, GERMANY	MPHM	MAX-PLANCK-INST. FUR PHYS.-ASTROPHYS.	HEIDELBERG, GERMANY
BORD	UNIV. DE BORDEAUX	BORDEAUX, FRANCE	MSNA	INS. DI FISICA DELL'UNIV.	MESSINA, ITALY
BOST	BOSTON UNIV.	BOSTON, MASS., USA	MSU	MICHIGAN STATE UNIV.	EAST LANSING, MICH., USA
BRAN	BRANDEIS UNIV.	WALTHAM, MASS., USA	HUNG	UNIV. OF HUNGARY	MUNICH, GERMANY
BRIS	H. H. MILLS PHYS. LAB., U. OF BRISTOL	BRISTOL, ENGLAND	NAGC	NAGOYA UNIV.	NAGOYA, JAPAN
BROW	BROWN UNIV.	PROVIDENCE, R. I., USA	NAPL	UNIV. DI NAPOLI	NAPOLI, ITALY
BRUX	UNIV. LIBRE DE BRUXELLES	BRUXELLES, BELGIUM	NDAM	UNIV. OF NOTRE DAME	NOTRE DAME, IND., USA
BUCH	BUCHAREST STATE UNIV.	BUCHAREST, ROMANIA	NEUS	NORTH-EASTERN UNIV.	BOSTON, MASS., USA
BUDA	CENTRAL RESEARCH INSTITUTE OF PHYSICS	BUDAPEST, HUNGARY	NEUC	UNIV. DE NEUCHÂTEL	NEUCHÂTEL, SWITZERLAND
BUFF	STATE UNIV. OF NEW YORK AT BUFFALO	BUFFALO, N. Y., USA	RUFF	RUFFIN, N. Y., USA	WILKINGTON-IN-HUDSON, N.Y., USA
CAEN	LAB. DE PHYS. CORPUSCULAIRE	CAEN, FRANCE	NIJM	R. K. UNIV. NIJMEGEN	NIJMEGEN, NETHERLANDS
CANB	AUSTRALIAN NATIONAL UNIV.	CANBERRA, AUSTRALIA	NORD	NORDISK INS. FOR TEOR. ATOMFYS.	COPENHAGEN, DENMARK
CARL	CARLTON UNIV.	OTTAWA, CANADA	OTTA	OTTAWA UNIV.	OTTAWA, CANADA
CARN	CARNEGIE-MELLON UNIV.	PITTSBURGH, PA., USA	NRES	NAVAL RESEARCH LABORATORY	WASHINGTON, D.C., USA
CASE	CASE WESTERN RESERVE UNIV.	CLEVELAND, OHIO, USA	NRL	NORTHWESTERN UNIV.	EVANSTON, ILL., USA
CAVE	CAVENISH LAB., CAMBRIDGE UNIV.	CAMBRIDGE, ENGLAND	NEWY	NEW YORK UNIV.	NEW YORK, N. Y., USA
CCAC	COMMUNITY COLLEGE OF ALLEGHENY COUNTY	PITTSBURGH, PENN., USA	OHIO	OHIO UNIV.	ATHENS, OHIO, USA
CDEF	COLLEGE DE FRANCE	PARIS, FRANCE	OREG	UNIV. OF OREGON	EUGENE, ORE., USA
CEA	CENERGIE CENTRALE	CAMBRIEUX, MASS., USA	ORSA	GAK RIKU UNIV. LAB.	SAK RIDGE, TENN., USA
CERN	EUROPEAN ORG. FOR NUC. RES.	GENEVA, SWITZERLAND	ORSA	UNIV. DE PARIS, FAC. DES SCI.	ORSAY, FRANCE
CHIC	UNIV. OF CHICAGO	CHICAGO, ILL., USA	OSLO	OSLO UNIV.	OSLO, NORWAY
CINC	UNIV. OF CINCINNATI	CINCINNATI, OHIO, USA	OSU	OHIO STATE UNIV.	COLUMBUS, OHIO, USA
CIT	CALIF. INSTITUTE OF TECHNOLOGY	PASADENA, CALIF., USA	OXF	OXFORD UNIV.	OXFORD, ENGLAND
CNRC	CANADIAN NATIONAL RESEARCH COUNCIL	OTTAWA, CANADA	PAPO	UNIV. DI PADOVA	PADOVA, ITALY
COLD	COLUMBIA UNIV.	NEW YORK, N. Y., USA	PATR	UNIV. OF PATRAS	PATRAS, GREECE
COLU	CORNELL UNIV.	ITHACA, N. Y., USA	PENN	UNIV. OF PENNSYLVANIA	PHILADELPHIA, PA., USA
CRAC	INST. FOR NUCLEAR RESEARCH	CRACOW, POLAND	PISA	UNIV. DI PISA	PISA, ITALY
CUNY	CITY UNIV. OF NEW YORK	NEW YORK, N. Y., USA	PITT	UNIV. OF PITTSBURGH	PITTSBURGH, PA., USA
CURI	LABORATOIRE JOLIO-CURIE	PARIS, FRANCE	PPA	PRINCETON-PENN. PROTON ACCEL.	PRINCETON, N. J., USA
DARE	CARESBURY NUC. PHYS. LAB.	CARESBURY, ENGLAND	PRG	INSTITUTE OF PHYSICS, CSAV	PRAGUE, CZECHOSLOVAKIA
DART	DARTMOUTH COLLEGE	HANDOVER, N. H., USA	PRIN	PRINCETON UNIV.	PRINCETON, N. J., USA
DESY	DEUTSCHES ELEKTROEN-SYNCH.	HAMBURG, GERMANY	PURD	PURDUE UNIV.	LAFAYETTE, IND., USA
DUKE	DUKE UNIV.	DURHAM, N. C., USA	RHEL	WEIZMANN INST. OF SCI.	REHOVOT, ISRAEL
DURH	UNIV. OF DURHAM	DURHAM, ENGLAND	RISO	RUTHERFORD HIGH ENERGY LAB.	RUTHERFORD, BERKS., ENGLAND
DUIC	DUBLIN UNIVERSITY COLLEGE	DUBLIN, IRELAND	RISD	RESEARCH ESTAB. RISD	RISDALE, DENMARK
EDIN	UNIV. OF EDINBURGH	EDINBURGH, SCOTLAND	RMS	ROYAL MILITARY COLLEGE OF SCIENCE	SRIVENHAM, ENGLAND
EFTI	ENRICO FERMI INST. FOR NUCL. STUDIES	CHICAGO, ILL., USA	ROCH	UNIV. OF ROCHESTER	ROCHESTER, N. Y., USA
EPOL	ECOLE POLYTECHNIQUE	PARIS, FRANCE	RCMA	UNIV. DI ROMA	ROME, ITALY
ETHZ	SWISS FEDERAL INST. OF TECHNOLOGY	ZURICH, SWITZERLAND	RUTG	RUTGERS UNIV.	NEW BRUNSWICK, N. J., USA
FIRZ	UNIV. DI FIRENZE	FIRENZE, ITALY	CNRS	DI ETUDES NUC. SACLAY	SACLAY, FRANCE
FISK	FISK UNIV.	NASHVILLE, TENN., USA	SAGA	SAGA UNIV.	SAGA, JAPAN
FLOR	UNIV. OF FLORIDA	GAINESVILLE, FLA., USA	SEAT	SEATTLE PACIFIC COLLEGE	SEATTLE, WASH., USA
FNAL	FERMILAB NATIONAL ACCELERATOR LAB.	BATAVIA, ILL., USA	SEB	RESEARCH CENTER SEIBERSDORF	VIENNA, AUSTRIA
FRAS	LAB. NATIONALI DEL FINESTRONE	FRASCATI, ITALY	SERP	INST. OF HIGH EN. PHYS.	SERPUKOV, USSR
FSU	FLORIDA STATE UNIV.	TALLAHASSEE, FLA., USA	SETO	SETO UNIV.	SANTA CRUZ, N. J., USA
GENA	UNIV. DI GENOVA	GENOVA, ITALY	SFLA	UNIV. OF SOUTH FLORIDA	TAMPA, FLORIDA, USA
GESC	GENERAL ELECTRIC RES. AND DEV. CENTER	SCHENECTADY, N. Y., USA	SHEP	UNIV. OF SHEFFIELD	SHEFFIELD, ENGLAND
GEVO	UNIV. DE GENEVE	GENEVA, SWITZERLAND	SHMP	SHEPPARD UNIV. COLLEGE	SOUTHAMPTON, ENGLAND
GLAS	UNIV. OF GLASGOW	GLASGOW, SCOTLAND	SIN	SWISS INST. FOR NUCLEAR RESEARCH	VILLIGEN, SWITZERLAND
GRAZ	UNIV. GRAZ	GRAZ, AUSTRIA	SLAC	STANFORD LINEAR ACCEL. CENTER	STANFORD, CALIF., USA
GSCO	GEOLOGICAL SURVEY OF CANADA	OTTAWA, CANADA	SOFI	BULGARIAN ACAD. OF SCI.	SOFIYA, BULGARIA
HAIF	TECHNION ISRAEL INST. OF TECHNOLOGY	HAIFA, ISRAEL	STAN	STANFORD UNIV.	STANFORD, CALIF., USA
HAMB	UNIV. HAMBURG	HAMBURG, GERMANY	STEV	STEVENS INST. OF TECH.	HOBOKEN, N. J., USA
HARV	HARVARD UNIV.	CAMBRIDGE, MASS., USA	STLD	ST. LOUIS UNIV.	ST. LOUIS, MO., USA
HAWA	UNIV. OF HAWAII	HONOLULU, HAWAII, USA	STOH	STOCKHOLM UNIV.	STOCKHOLM, SWEDEN
HEID	UNIV. HEIDELBERG	HEIDELBERG, GERMANY	STON	STATE UNIV. OF NEW YORK AT STONY BROOK	STONY BROOK, L.I., N. Y., USA
HELS	HELSINKI UNIV. OF TECH.	HELSINKI, FINLAND	STRB	CENTRE DES RES. NUCLEAIRES	STRASBOURG, FRANCE
IDM	INTERNATIONAL BUSINESS MACHINES	PALO ALTO, CALIF., USA	SUSX	SUSSEX UNIV.	SUSSEX, ENGLAND
ILL	UNIV. OF ILLINOIS	URBANA, ILL., USA	SYRA	SYRACUSE UNIV.	SYRACUSE, N. Y., USA
ILLC	UNIV. OF ILLINOIS AT CHICAGO	CHICAGO, ILL., USA	TATA	TATA INST. OF FUNDAMENTAL RESEARCH	BOMBAY, INDIA
INDO	UNIV. OF INDIANA	BLOOMINGTON, IND., USA	TELA	UNIV. OF TEL-AVIV	TEL-AVIV, ISRAEL
INNS	PHYS. INST., UNIV. INNSBRUCK	INNSBRUCK, AUSTRIA	TEMP	TEMPLE UNIV.	PHILADELPHIA, PA., USA
IOWA	UNIV. OF IOWA	IOWA CITY, IOWA, USA	TENN	UNIV. OF TENNESSEE	KNOXVILLE, TENN., USA
IPNW	INST. DE PHYS. NUCLEAIRE	ORSAY, FRANCE	TORO	UNIV. OF TORONTO	TORONTO, CANADA
IPNC	INSTITUT DE PHYSIQUE NUCLEAIRE	PARIS, FRANCE	TOHO	TOHOKU UNIV.	SENDAI, JAPAN
IRAD	INSTITUTE OF RADIUM	CANADA	TOKY	UNIV. OF TOKYO	TOKYO, JAPAN
ISU	IOWA STATE UNIV.	AMES, IOWA, USA	TORI	UNIV. DI TORINO	TRINO, ITALY
ITEP	INST. FOR TECH. AND EXP. PHYS.	MOSCOW, USSR	TRST	UNIV. DI TRIESTE	TRIESTE, ITALY
INDIA	INDIAN UNIV. AT INDIANAPOLIS	INDIANAPOLIS, IN., USA	TUFT	TUFTS UNIV.	MEDFORD, MASS., USA
JACL	JACQUELIN UNIV.	CRACOW, POLAND	UCB	UNIV. OF BRITISH COLUMBIA	VANCOUVER, CANADA
JHU	JOHNS HOPKINS UNIV.	BALTIMORE, MD., USA	UCR	UNIV. OF CALIF. AT BERKELEY	BERKELEY, CALIF., USA
JINR	JOINT INST. FOR NUCL. RESEARCH	DUBNA, USSR	UCI	UNIV. OF CALIF. AT IRVINE	IRVINE, CALIF., USA
KANS	UNIV. OF KANSAS	LAWRENCE, KANSAS, USA	UCLA	UNIV. OF CALIF. AT LOS ANGELES	LOS ANGELES, CALIF., USA
KEK	TECHNICAL UNIVERSITY OF KARLSRUHE	KARLSRUHE, GERMANY	UCRB	UNION CARBIDE NUCLEAR DIVISION	YORKVILLE, TENN., USA
KIAC	NAT. LAB. FOR HIGH ENERGY PHYS., JAPAN	TSUKUBA-GUN., JAPAN	UCR	UNIV. OF CALIF. AT RIVERSIDE	RIVERSIDE, CALIF., USA
KIT	KIRCHHOFF INST. OF ATOMIC ENERGY	MOSCOW, USSR	UCSB	UNIV. OF CALIF. AT SANTA BARBARA	SANTA BARBARA, CALIF., USA
KNTY	UNIV. OF KENTING	LEIXINGTON, KY., USA	UCSD	UNIV. OF CALIF. AT SAN DIEGO	SANTA CRUZ, CALIF., USA
LALC	LINEAR ACCELERATOR LAB., ORSAY	ORSAY, FRANCE	UNC	UNIV. OF NORTH CAROLINA	CAMPDEN, N. C., USA
LANC	LANCASTER UNIV.	LANCASTER, ENGLAND	UNJN	UPSALA COLLEGE	EAST ORANGE, N. J., USA
LASL	U. C. LOS ALAMOS SCIENTIFIC LAB.	LOS ALAMOS, N. M., USA	UTAH	UNIV. OF UTAH	SALT LAKE CITY, UTAH, USA
LAUS	UNIV. OF LAUSANNE	LAUSANNE, SWITZERLAND	VAND	VANDERBILT UNIV.	NASHVILLE, TENN., USA
LEBD	LEBERDEY PHYSICS INST.	BERKELEY, CALIF., USA	VIEN	VIENNA UNIV.	VIENNA, AUSTRIA
LEHI	LEHIGH UNIV.	BETHLEHEM, PA., USA	VING	UNIV. OF VIRGINIA	CHARLOTTESVILLE, VA., USA
LEIT	INST. LORENTZ	LEIDEN, NETHERLANDS	WARS	UNIV. OF WARSAW	WARSAW, POLAND
LINZ	LINZ INSTITUT FUR PHYSIK, KEPLER HOCH.	LINZ, AUSTRIA	WASH	UNIV. OF WASHINGTON	SEATTLE, WASH., USA
LIVP	LIVERPOOL UNIV.	LIVERPOOL, ENGLAND	WIEN	UNIV. OF VIENNA	VIENNA, AUSTRIA
LIOC	IMPERIAL COL. OF SCI. AND TECH.	LONDON, ENGLAND	WILL	COLLEGE OF WILLIAM AND MARY	WILLIAMSBURG, VA., USA
LLOU	QUEEN MARY COLLEGE	LONDON, ENGLAND	WISC	UNIV. OF WISCONSIN	MADISON, WIS., USA
LOUC	UNIVERSITY COLLEGE	LONDON, ENGLAND	WOOD	WOODSTOCK COLLEGE	WOODSTOCK, MO., USA
LOWC	WESTFIELD COLLEGE	LONDON, ENGLAND	WASH	WASHINGTON UNIV.	ST. LOUIS, MO., USA
LPMF	LAB. DE PHYS. NUCL. ET HAUTES ENERGIES	PARIS, FRANCE	WYOM	UNIV. OF WYOMING	LARAMIE, WYOMING, USA
LPTP	LAB. DE PHYS. THEOR. ET HAUTES ENERGIES	PARIS, FRANCE	YALE	YALE UNIV.	NEW HAVEN, CONN., USA
LRI	U. C. LAWRENCE BERKELEY LAB.	BERKELEY, CALIF., USA	ZEEM	ZEEMAN LAB., UNIV. OF AMSTERDAM	AMSTERDAM, NETHERLANDS
LSU	LOUISIANA STATE UNIV.	BATON ROUGE, LA., USA			

Stable Particles

$\gamma, \nu_e, \nu_\mu, e, \mu$

Data Card Listings

For notation, see key at front of Listings.

γ

0 GAMMA(0,J=1)

0 GAMMA MASS (IN UNITS OF 10**--21 MEV)

M	P	(6.)	OR LESS	PATEL	65	SATELLITE DATA	10/69
M		6.	OR LESS	GINTSBURG	64	SATELLITE DATA	10/79
M		2.3	OR LESS	GOLDBABER	68	SATELLITE DATA	10/69
M	F	(0.06)	OR LESS	FRANKEN	71	LOW FREQ RES CIR	3/72
M		10.	OR LESS	WILLIAMS	71	TESTS GAUSS LAW	3/71
M		0.73	OR LESS	HOLLWEG	74	ALFVEN WAVES	7/74*
M	F		VALIDITY QUESTIONABLE ACCORDING TO AUTHORS AND KROLL 71.				3/72
M	P		SEE CRITICISM IN GOLDBABER 71				3/72

REFERENCES FOR GAMMA

GINTSBR 64 SOV. ASTR. AJ7 536 M. A. GINTSBRG (ACAD SCI, USSR)

PATEL 65 PL 14 105 V. L. PATEL (DURHAM)

GOLDBABER 68 PRL 21 567 A. GOLDBABER, M. NIETO (STONY BROOK)

FRANKEN 71 PRL 26 115 P A FRANKEN, G W AMPULSKI (MICH)

WILLIAMS 71 PRL 26 721 +FALLER, HILL (WESLEYAN)

HOLLWEG 74 PRL 32 961 J V HOLLWEG (NATL CENTER FOR ATMOS RESRCH)

PAPERS NOT REFERRED TO IN DATA CARDS

GOLDBABER 71 RMP 43 277 A S GOLDBABER, M M NIETO (STON+BOHR+UCSB)

KROLL 71 PRL 26 1395 N M KROLL (SLAC)

ν_e

1 E-NEUTRINO(0,J=1/2)

1 E-NEUTRINO MASS (KEV)

M		(0.25)	OR LESS	LANGER	52	CNTR ANTI-NEUT.(TRITIUM)	
M		(0.50)	OR LESS	HAMILTON	53	CNTR ANTI-NEUT.(TRITIUM)	11/73
M		(0.55)	(0.28)	FRIEDMAN	58	CNTR ANTI-NEUT.(TRITIUM)	
M		4.1	OR LESS	BECK	68	CNTR NEUTRINO SODIUM 22	11/73
M		0.5	OR LESS	DARIS	69	CNTR ANTI-NEUT.(TRITIUM)	11/73
M	D	0.32	OR LESS	SALGO	69	CNTR ANTI-NEUT.(TRITIUM)	11/73
M		0.06	OR LESS	BERGKVIS	72	CNT ANTI-NEUT.(TRITIUM)	11/73
M		(0.008)	OR LESS	COWSIK	72	THEOR. LIM. FROM COSMOLOGY	3/74
M		0.086	OR LESS	RODE	72	CNTR ANTI-NEUT.(TRITIUM)	11/73
M	D	DARIS 69 VALUE .075KEV(CL=.67)	DISAGREES WITH FIG.6. WE USE FIG.6.				11/73
M	C	450.	OR LESS	CLARK	74	ASPK KE3 DECAY	11/75*
M	C	LOWEST LIMIT FROM STRANGENESS	CHANGING DECAY.				11/75*

1 (E-NEUTRINO) - (E-ANTINEUTRINO) MASS DIFF. (KEV)

DM 450. OR LESS CL=.90 CLARK 74 ASPK KE3 DECAY 11/75*

REFERENCES FOR E-NEUTRINO

LANGER 52 PR 88 689 L M LANGER, P J D MOFFAT (INDIANA)

HAMILTON 53 PR 92 1521 D HAMILTON, W P ALFORD, L GROSS (PRINCETON)

FRIEDMAN 58 PR 109 2214 LEWIS FRIEDMAN, LINCOLN G SMITH (SNU)

BECK 68 ZPHY 216 229 F BECK, H DANIEL (MPIH)

DARIS 69 NP A138 545 N. DARIS, C ST-PIERRE (LAVAL-QUEBEC)

SALGO 69 NP A138 417 R C SALGO, H H STAUB (ZURICH)

BERGKVIS 72 NP 839 317 KARL-ERIK BERGKVIST (UNIV STOCKHOLM)

COWSIK 72 PRL 29 669 R COWSIK, J MC CLELLAND (UCB)

RODE 72 LNC 5 139 B RODE, H DANIEL (MUNICH+MPIH)

CLARK 74 PR D9 533 +ELIOFF, FRISCH, JOHNSON, KERTH, SHEN+ (LBL)

ν_μ

2 MU-NEUTRINO(0,J=1/2)

2 MU-NEUTRINO MASS (MEV)

M		3.5	OR LESS	BARKAS	56	EMUL	
M		4.0	OR LESS	DUZDIK	59	CNTR	
M		3.6	OR LESS	FEINBERG	63	RVUE	7/66
M		3.0	OR LESS	ALLCOCK	63	RVUE	7/66
M		2.5	OR LESS	BARDON	65	ASPK	
M		2.8	OR LESS	SHAHER	65	CNTR	5/71
M		1.6	OR LESS	BOOTH	67	CNTR	3/68
M		2.2	OR LESS	HYMAN	67	HEBC	0. K- HE 11/67
M	B M	(1.2)	OR LESS	BACKENSTOSS	71	CNTR	M**2=1.28+-1.24 10/71
M	S	1.15	OR LESS	SHRUM	71	CNTR	M**2=1.55+-1.14 12/71
M		(8 EV)	OR LESS	COWSIK	72	THEOR. LIM. FROM COSMOLOGY	3/74
M	B M	1.15	OR LESS	BACKENSTOSS	73	CNTR	M**2=0.29+-0.90 7/73
M		0.65	OR LESS	CLARK	74	ASPK	KMU3 DECAY 7/74*
M	D M	1.0	OR LESS	DAUM	76	SPEC	M**2=0.23+-0.54 1/76*
M	M		WE CALCULATE UPPER LIMIT AT CL=.90 FROM M**2.				1/76*
M	B		BACKENSTOSS 73 REPLACES BACKENSTOSS 71 AND USES THEIR NEW PI- MASS.				1/73
M	S		SHRUM 71 USES SHAFER 67 PI- MASS VALUE AND CRANE 71 MU MASS VALUE.				1/73
M	D		DAUM 76 USES OUR 1974 PI- AND MU MASSES.				1/76*

2 (MU-NEUTRINO) - (MU-ANTINEUTRINO) MASS DIFF. (MEV)

DM (0.45) OR LESS CL=.90 CLARK 74 ASPK KMU3 DECAY 11/75*

REFERENCES FOR MU-NEUTRINO

BARKAS 56 PR 101 778 M BARKAS, W BIRNBAUM, F M SMITH (LFL)

DUZDIK 59 PR 114 336 W F. DUZDIK, R SAGANE, J VEDDER (LRL)

FEINBERG 63 ARNS 13 431 G FEINBERG, L M LEDERMAN (COLUMBIA)

ALLCOCK 65 PPSL 85 875 G R ALLCOCK (LIVERPOOL)

BARDON 65 PRL 14 449 BARDON, NORTON, PEOPLES + (COLU+STONY BROOK)

SHAHER 65 PRL 14 923 R E SHAHER, C CROWE, JENKINS (LFL)

BOOTH 67 PL 268 39 BOOTH, JOHNSON, WILLIAMS, WORMALD (LIVERPOOL)

HYMAN 67 PL 258 376 HLOKEN, PENWIT, MCKENZIE+ (ANL+CARN+NWES)

BACKENSTOSS 71 PL 368 603 BACKENSTOSS, DANIEL+KOOCH+ (CERN+KARL+HEID)

SHRUM 71 PL 378 114 E V SHRUM, K O H ZIOCK (UNIV OF VIRGINIA)

COWSIK 72 PRL 29 669 R COWSIK, J MC CLELLAND (UCB)

BACKENSTOSS 73 PL 438 539 BACKENSTOSS, DANIEL+KOOCH+ (CERN+KARL+MUNICH)

CLARK 74 ERRORS 935 +ELIOFF, FRISCH, JOHNSON, KERTH, SHEN + (LBL)

DAUM 76 PUBL. IN PL B FEB +DUBAL, EATON, FROSCHE, MCCULLOUGH+ (STN, ETHZ)

e

3 ELECTRON(0,5,J=1/2)

3 ELECTRON MASS (MEV)

M		(.5110061(.00002)	COHEN	65	RVUE	
M		(.5110041(.0000016)	TAYLOR	69	RVUE	USING NEW E/H 7/70
M		.5110034 .0000014	COHEN	73	RVUE	3/74

3 ELECTRON MEAN LIFE (UNITS 10**21 YR)

T	S	2.0	OR MORE	MDE	65	CNTR	6/66
T	S	5.2	OR MORE	STEINBERG	75	CNTR	2/76*
T	S		STEINBERG 75 SENSITIVE TO ALL DECAY MODES IN WHICH DECAY PARTICLES ESCAPE FROM DETECTOR WITHOUT DEPOSITING ENERGY. TEST OF CHARGE CONSERVATION.				2/76*

3 ELECTRON MAGNETIC MOMENT(E/2ME)

MM SEE RICH 72 FOR A REVIEW OF THEORY AND EXPERIMENTS. 3/74

MM		(1.0011609)	+(24)*10**--7	SCHUPP	61	CNTR	
MM		(1.001159622)	+(27)*10**--9	WILKINSON	63	CNTR	8/66
MM		(1.001168)	+(22)*10**--6	RICH	66	CNTR +	POSITRON 8/66
MM	R	(1.001159555)	+(30)*10**--9	RICH	68	CNTR	6/68
MM		(1.0011596389)	+(31)*10**--10	TAYLOR	69	RVUE	2/71
MM		(1.001159644)	+(7)*10**--9	WESLEY	70	CNTR	6/70
MM		(1.0011596577)	+(35)*10**--10	WESLEY	71	CNTR	2/72
MM		(1.0011603)	+(12)*10**--7	GILLENLAND	72	CNTR +	2/72
MM	R	1.0011596567	+(25)*10**--10	COHEN	73	CNTR	2/74

MM R RICH 68 IS REEVALUATION OF WILKINSON 63.

REFERENCES FOR ELECTRON

SCHUPP 61 PR 121 1 A A SCHUPP, R W PIDD, H R CRANE (MICH)

WILKINSON 63 PR 130 852 D T WILKINSON, H R CRANE (MICH)

COHEN 65 RMP 37 537 COHEN, JUMOND (N.A. AVIATION SCI. CENTER+CIT)

MDE 65 PR 140 B 992 M K MDE, F REINES (CASE INST TECHNOLOGY)

RICH 66 PRL 17 271 A RICH, H R CRANE (MICH)

RICH 68 PRL 20 967 A RICH (MICH)

TAYLOR 69 RMP 41 375 +PARKER, LANGENBERG (PRIN+UCI+PENN)

WESLEY 70 PRL 24 1320 J C WESLEY, A RICH (MICH)

WESLEY 71 PR 44 1341 J C WESLEY, A RICH (MICH)

GILLENLAND 72 PR 45 38 J GILLENLAND, A RICH (MICH)

RICH 72 RMP 44 250 A RICH, J C WESLEY (MICH)

COHEN 73 J. PHYS. CHEM. REF. DATA 2, P. 663, E. R. COHEN, B. N. TAYLOR (UMD)

STEINBERG 75 PR D12 2582 STEINBERG, KWIAKOWSKI, MAENHAUT, WALL (UMD)

μ

4 MUON(106,J=1/2)

4 MUON MASS (MEV)

M		(105.659)	(0.002)	FEINBERG	63	RVUE	
M		(105.65999)	(0.0014)	TAYLOR	69	RVUE	USING NEW E/H 7/70
M	C	(105.6597)	(0.0005)	CRANE	72	CNTR	INCLUDED IN COHEN73 1/73
M	D	(105.6594)	(0.0006)	CRANE	72	CNTR	INCLUDED IN COHEN73 2/72
M		105.65948	0.00035	COHEN	73	RVUE	3/74
M	C	CRANE 71 GIVES MU/ME=206.76878(85). WE USE ME=5110041(16)MEV.					1/73
M	D	CRANE 72 GIVES MU/ME=206.7682(5) AND USES ME=5110041(16)MEV.					1/73
M	FIT	105.65948	0.00035	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)			2/76*

4 MUON MEAN LIFE (UNITS 10**--6)

T		2.198	0.001	0.001	FARLEY	62	CNTR	
T		2.203	0.004		LUNDY	62	CNTR	CONLEV=.98 11/67
T		2.202	0.003	0.003	ECKHAUSE	63	CNTR	
T		2.197	0.005	0.002	MEYER	63	CNTR +	
T		2.198	0.002	0.002	MEYER	63	CNTR -	7/66
T	W	(2.20206)	(0.00081)		WILLIAMS	72	CNTR +	2/76*
T		2.1973	0.0003		DUCLOS	73	CNTR +	1/76*
T		2.19711	0.00008		BALANDIN	74	CNTR +	1/76*
T	W	WILLIAMS 72 MEAN LIFE MEASUREMENT WAS NOT THE PRIMARY PURPOSE OF THEIR EXPERIMENT AND DISAGREES STRONGLY WITH LATER EXPTS. NOT AVGD.						1/76*
T	AVG	2.19713	0.000077	0.000077	AVERAGE ERROR INCL. SCALE FACTOR OF 1.01			
T	STUDENT	2.197133	0.000084	0.000084	AVG. USING STUDENT10(H/1.11) -- SEE TEXT			

4 MU+MU- MEAN LIFE RATIO

DT		1.000	0.001		MEYER	63	CNTR	MEAN LIFE MU+/MU- 7/66
----	--	-------	-------	--	-------	----	------	------------------------

4 MUON ANOMALOUS MAGN. MOMENT (10**--6*(2*MU MASS))

MM SEE RICH 72 FOR A REVIEW OF THEORY AND EXPERIMENTS. 3/74

MM		(1162.0)	(5.0)	CHARPAK	62	CNTR +	
MM	B	(1165.75)	(0.71)	BAILEY	68	CNTR +	STOR. RINGS 5/69
MM	B	(1166.25)	(0.24)	BAILEY	68	CNTR +	STOR. RINGS 5/69
MM	B	ERRORS COMBINED TO GIVE MU- VALUE BELOW					5/69
MM		1166.16	0.31	BAILEY	68	CNTR +	STOR. RINGS 5/69
MM		1165.895	0.027	BAILEY	75	CNTR +	STORAGE RING 11/75*
MM	AVG	1165.897	0.027	AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01			
MM	STUDENT	1165.897	0.029	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

Data Card Listings

Stable Particles

For notation, see key at front of Listings.

μ

4 MUON TO PROTON MAGNETIC MOMENT RATIO

Table listing muon to proton magnetic moment ratio data. Includes columns for experiment name (e.g., MMR, MMR D), description, and numerical values.

4 MUON PARTIAL DECAY MODES

Table listing muon partial decay modes. Columns include mode name (P1-P5), description, and branching ratios.

4 MUON BRANCHING RATIOS

Table listing muon branching ratios. Columns include mode name (R1-R4), description, and numerical values.

4 MUON DECAY PARAMETERS

RELATED TEXT SECTION VI A

Table listing muon decay parameters and related text. Columns include parameter name (RHO, RHO P, RHO C, etc.), description, and values.

Table listing Delta parameter and other decay constants. Columns include parameter name (DEL, GS, GA, etc.), description, and values.

REFERENCES FOR MUON

Table listing references for muon. Columns include author names (e.g., COFFIN, LUNDY, BARDOON, etc.) and publication details.

Stable Particles

π^\pm, π^0

π^\pm

8 CHARGED PION(140,JPG=0--) I=1

Table with 4 columns: M, S, B, S. Rows include mass measurements from various experiments like CROWE, BARKAS, SHAFER, BACKENSTO, etc.

8 (PI+) - (MU+) MASS DIFFERENCE (MEV)

Table with 4 columns: D, D, D, D. Rows include mass difference measurements from BARKAS, HYMAN, BOOTH, etc.

8 ((PI+) - (PI-))/AVG., MASS DIFFERENCE (PERCENT)

Table with 4 columns: DM, D, D, D. Rows include mass difference percentage measurements from AYRES, etc.

8 CHARGED PION MEAN LIFE (UNITS 10**--9)

Table with 4 columns: T, T, T, T. Rows include mean life measurements from CROWE, ANDERSON, ASHKIN, etc.

8 ((PI+) - (PI-))/AVG., MEAN LIFE DIFF. (PERCENT)

Table with 4 columns: DT, DT, DT, DT. Rows include mean life difference percentage measurements from LOBKOWICZ, BARDON, etc.

8 CHARGED PION PARTIAL DECAY MODES

Table with 4 columns: P1, P2, P3, P4. Rows include partial decay modes like PION INTO MU, PION INTO E, etc.

8 CHARGED PION BRANCHING RATIOS

Table with 4 columns: R1, R2, R3, R4. Rows include branching ratios for various decay channels like PION INTO MU NEU GAMMA, etc.

Data Card Listings
For notation, see key at front of Listings.

REFERENCES FOR CHARGED PION

Table listing references for charged pion, including authors like CROWE, ANDERSON, BARKAS, SHAFER, etc.

PAPERS NOT REFERRED TO IN DATA CARDS
MERRISON 62 ADVP 11 1
SHAPIRO 62 PR 121 1022
CZIRR 63 PR 130 341

π^0

9 NEUTRAL PION(135,JPG=0--) I=1

Table with 4 columns: D, D, D, D. Rows include mass difference measurements for neutral pion from PANOFKY, CHINOWSKY, etc.

9 NEUTRAL PION MEAN LIFE (UNITS 10**--16)

Table with 4 columns: T, T, T, T. Rows include mean life measurements for neutral pion from GLASSER, TITGE, etc.

9 NEUTRAL PION PARTIAL DECAY MODES

Table with 4 columns: P1, P2, P3, P4. Rows include partial decay modes for neutral pion like PION INTO 2 GAMMA, etc.

9 NEUTRAL PION BRANCHING RATIOS

Table with 4 columns: R1, R1, R1, R1. Rows include branching ratios for neutral pion like PION INTO (GAMMA E+ E-), etc.

Data Card Listings
For notation, see key at front of Listings.

Stable Particles

pi0, K±

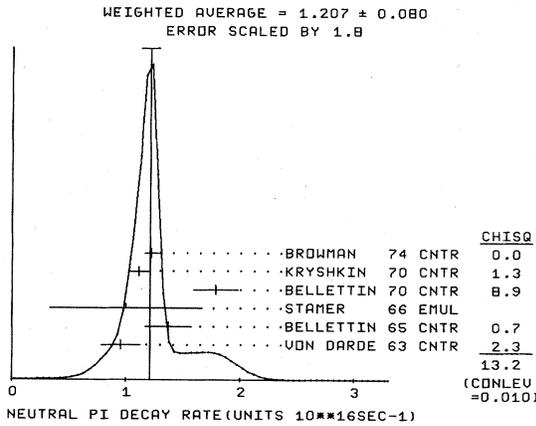


Table with columns for name, count, and CHISQ. Includes entries for BROWMAN, KRYSHKIN, BELLETTIN, STAMER, BELLETTIN, and UDN DARDE.

Table of experimental data for neutral pion decays, including parameters like P10 INTO (3 GAMMA)/(2 GAMMA) and various counts and error values.

REFERENCES FOR NEUTRAL PION

Bibliographic references for neutral pion decays, listing authors and publication details.



10 CHARGED K (494, JP=0-) I=1/2

10 CHARGED K MASS (MEV)

Table listing charged K meson masses and decay parameters, including names like COHEN, BARKAS, GREINER, and KUNSELMAN.

10 (K+) - (K-) MASS DIFFERENCE (MEV)

Table showing mass differences between K+ and K- mesons, including parameters like DM F and error values.

10 CHARGED K MEAN LIFE (UNITS 10**--8)

Table of charged K mean life measurements, listing experimenters like ILOFF, EISENBERG, BURROWS, and others.

WEIGHTED AVERAGE = 0.8084 ± 0.0021, ERROR SCALED BY 2.4

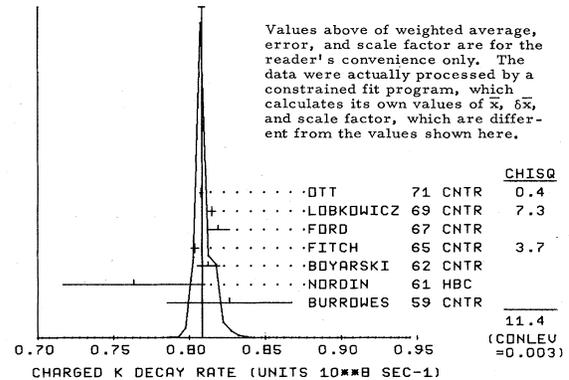


Table with columns for name, count, and CHISQ. Includes entries for DTT, LOBKOWICZ, FORD, FITCH, BOYARSKI, NORDIN, and BURROWS.

10 ((K+) - (K-))/AVG., MEAN LIFE DIFFERENCE (PERCENT)

Table showing the mean life difference between K+ and K- mesons, including parameters like DT N and error values.

10 CHARGED K PARTIAL DECAY MODES

Table listing various partial decay modes for charged K mesons, such as K INTO MU NEU, K INTO PI P10, etc.

CHARGED K CONSTRAINED FIT

OVERALL FIT OF MEAN LIFE, WIDTHS AND BRANCHING RATIOS USES 57 DATA POINTS TO DETERMINE SIX QUANTITIES. OVERALL FIT HAS CHISQ=78-2. MAIN CONTRIBUTION (13.3) COMES FROM R19 OF HAIDT 71 (WE SEE NO REASON TO REJECT THIS EXPERIMENT AT THIS TIME)

Stable Particles

K±

Data Card Listings

For notation, see key at front of Listings.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_i δP_i), while the off-diagonal elements are the normalized correlation coefficients (δP_i δP_j)/(δP_i δP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 6 columns (P 1 to P 6) and 6 rows of branching fraction data and error values.

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., G_i ≡ Γ_i = Γ_total P_i, in appropriate units. In analogy to the matrix above, the diagonal elements are G_i ± δG_i, where δG_i = √(δG_i δG_i), while the off-diagonal elements are the normalized correlation coefficients (δG_i δG_j)/(δG_i δG_j). Note that, because of the error in Γ_total, the errors and correlations here are not directly derivable from those above.

Table with 6 columns (G 1 to G 6) and 6 rows of rate data and error values.

10 CHARGED K DECAY RATES

Table listing decay rates for charged K particles, including parameters like MU NEU, TAU, and (MU P10 NEU) + (E P10 NEU).

10 ((K+) - (K-))/AVG., DECAY RATE DIFFERENCE (PERCENT)

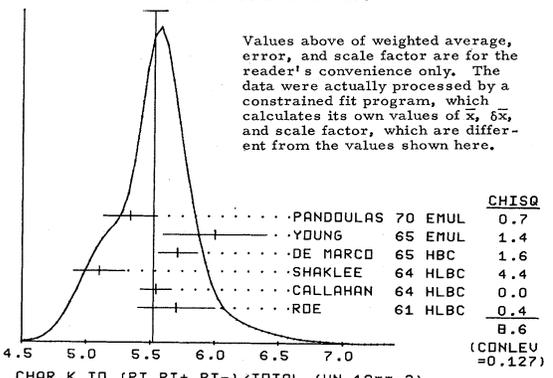
Table showing decay rate differences in percent for various K decay modes, including tau prime rates and pi2 rates.

10 CHARGED K BRANCHING RATIOS

Table listing branching ratios for charged K particles, including data from older experiments and fits.

Table listing experimental data for K decay rates, including names of researchers (e.g., BIRGE, ALEXANDER, CALLAHAN) and their respective measurements.

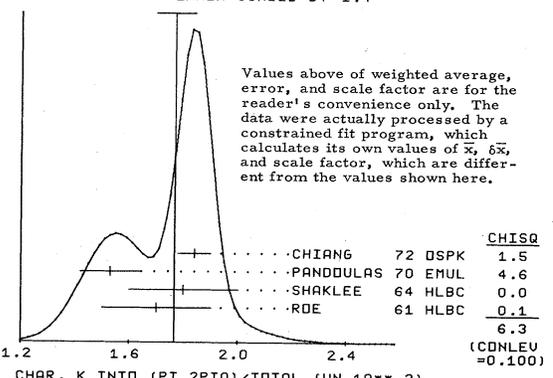
WEIGHTED AVERAGE = 5.521 ± 0.098 ERROR SCALED BY 1.3



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x-bar, delta x-bar, and scale factor, which are different from the values shown here.

Table listing experimental data for K decay rates, including names of researchers (e.g., BIRGE, ALEXANDER, TAYLOR) and their respective measurements.

WEIGHTED AVERAGE = 1.767 ± 0.071 ERROR SCALED BY 1.4



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x-bar, delta x-bar, and scale factor, which are different from the values shown here.

Data Card Listings

Stable Particles

For notation, see key at front of Listings.

K±

R5 CHAR. K INTO (MU P10 NEU)/TOTAL (UNITS 10**--2) (P5)
R5 O (2.8) (1.0) BIRGE 56 EMUL +
R5 O (5.9) (1.3) ALEXANDER 57 EMUL +
R5 O (2.8) (0.4) TAYLOR 59 EMUL +
R5 O EARLIER EXPERIMENTS NOT AVERAGED
R5 2345 3.33 0.16 CHIANG 72 OSPK + 1.84 GEV/C K+ 9/72
R5 FIT 3.197 0.087 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

R6 CHAR. K INTO (E P10 NEU)/TOTAL (UNITS 10**--2) (P6)
R6 O (3.2) (1.3) BIRGE 56 EMUL +
R6 O (5.1) (1.3) ALEXANDER 57 EMUL +
R6 O EARLIER EXPERIMENTS NOT AVERAGED
R6 429 4.7 0.3 SHAKLEE 61 HLBC + 11/67
R6 3516 4.86 0.10 CHIANG 72 OSPK + 1.84 GEV/C K+ 9/72
R6 AVG 4.849 0.093 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R6 STUDENT 4.85 0.10 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R6 FIT 4.823 0.052 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R7 CHAR. K INTO (P12 + MU3)/TOTAL (UNITS 10**--2) (P2+P5) 11/67
R7 WE COMBINE THESE TWO MODES FOR EXPTS MEASURING THEM IN XENON BC
R7 BECAUSE OF DIFFICULTIES OF SEPARATING THEM THERE
R7 23.4 1.1 ROE 61 HLBC + 11/67
R7 886 25.4 0.9 SHAKLEE 64 HLBC + 11/67
R7 AVG 24.60 0.98 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
R7 STUDENT 24.61 0.83 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R7 FIT 24.25 0.15 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R8 K+ INTO (P1+ PI+ E- NEU)/TOTAL (UNITS 10**--7) (P8)
R8 20. OR LESS CL=.95 BIRGE 65 FBC + 8/66
R8 0 6.9 OR LESS CL=.95 ELY 69 HLBC + 10/69
R8 0 9.0 OR LESS CL=.95 SCHWEINBE 71 HLBC + 9/71

R9 K+ INTO (P1+ PI- MU+ NEU)/TOTAL (UNITS 10**--5) (P9)
R9 1 0.77 0.54 0.50 CLINE 65 FBC + 8/66

R10 K+ INTO (P1+ PI+ MU- NEU)/TOTAL (UNITS 10**--6) (P10)
R10 0 3.0 OR LESS CL=.95 BIRGE 65 FBC + 8/66

R11 CHAR. K INTO (E NEU)/TOTAL (UNITS 10**--5) (P11) 11/67
R11 160.0 OR LESS CL=.95 BORREANI 64 HBC + 11/67
R11 4 (2.4) (1.5) (1.1) LJUNG 73 HLBC + P1+ KE 55-90 MEV 9/73
R11 BOWEN RESULT SHOULD BE CORRECTED TO 1.9(+1.7,-1.2) BECAUSE OF
R11 K+ TO E+ NEU GAMMA DECAYS BEFORE COMPARING WITH BOTTERILL 67 R28

R12 CHAR. K INTO (PI GAMMA)/TOTAL (UNITS 10**--4) (P17)
R12 ALL VALUES GIVEN HERE ASSUME A PHASE SPACE PION ENERGY SPECTRUM
R12 -0.1 0.6 CHEN 68 OSPK + T1P1 60-90 MEV 8/73
R12 0 0.5 OR LESS CL=.90 KLEMS 71 OSPK + T1P1GT 117 MEV 9/71
R12 0 0.35 OR LESS CL=.90 LJUNG 73 HLBC + 6-102,114-127MEV 9/73

R13 CHAR. K INTO (PI P10 GAMMA)/TOTAL (UNITS 10**--4) (P13)
R13 18 2.2 0.7 CLINE 64 FBC + P1+ KE 55-80 MEV 8/66
R13 0 1.9 OR LESS CL=.90 EMERSON 69 OSPK P1+ KE 55-80 MEV 10/69
R13 M 0 1.0 OR LESS CL=.90 MALTSEV 70 HLBC + P1+ KE LT 55 MEV 12/75*
R13 A2100 2.71 0.19 ABRAMS 72 ASPK -- P1+ KE 55-90 MEV 1/73
R13 24 2.4 0.8 EDWARDS 72 OSPK P1+ KE 58-90 MEV 8/72
R13 L (1.5) (1.1) (0.6) LJUNG 73 HLBC + P1+ KE 55-80 MEV 9/73
R13 L (2.4) (1.5) (1.1) LJUNG 73 HLBC + P1+ KE 55-90 MEV 9/73
R13 L 17 6.8 3.7 2.1 LJUNG 73 HLBC + P1+ KE 55-102MEV 9/73
R13 M MALTSEV TO SELECTS LOW P1+ ENERGY TO ENHANCE DIRECT EMISSION CONTR. 1/76*
R13 L THE LJUNG 73 VALUES ARE NOT INDEPENDENT. 9/73
R13 A ABRAMS 72 OBSERVES DIRECT EMISSION BR. RATIO OF (1.56+-0.35)*10**--5 1/73
R13 A +-0.5*10**--5 ADDNL. SYST. ERROR AND INNER BREMSSTRAHLUNG BR. RATIO 1/73
R13 A OF (2.55+-0.18)*10**--4. WE QUOTE THE SUM OF THESE BR. RATIOS. 1/73
R13 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

R14 CHAR. K INTO (PI P1+ PI- GAMMA)/TOTAL (UNITS 10**--4) (P14)
R14 1.0 0.4 STAMER 65 EMUL + EGAM GT 11MEV 8/66

R15 CHAR. K INTO (PI E-)/TOTAL (UNITS 10**--6) (P15)
R15 1 2.45 OR LESS CL=.90 CAMERINI 64 FBC + 8/66
R15 4.4 OR LESS CL=.90 BISI 67 DBC + 11/67
R15 C (0.4) OR LESS CLINE1 67 FBC + 11/67
R15 C 0.88 OR LESS CL=.90 CLINE2 67 FBC + 2/74
R15 32.0 OR LESS CL=.90 BETIER 72 OSPK -- 9/72
R15 1.7 OR LESS CL=.90 CENCE 74 ASPK + THREE TRACK EVTS 10/74*
R15 0.27 OR LESS CL=.90 GENCE 74 ASPK + TWO TRACK EVENTS 10/74*
R15 C CLINE2 REPLACES CLINE1. CLINE1 IS NOT FOR CL=.90 2/74

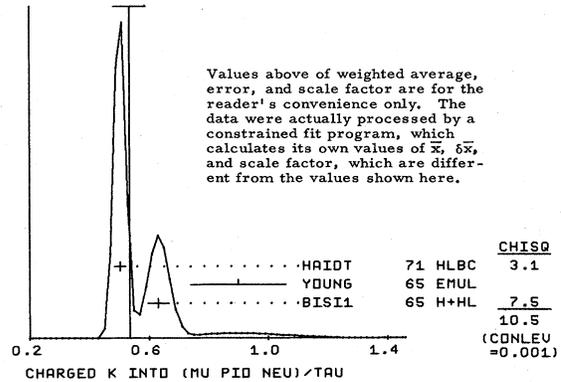
R16 CHAR. K INTO (PI MU+ MU-)/TOTAL (UNITS 10**--6) (P16)
R16 3.0 OR LESS CL=.90 CAMERINI 65 FBC + 8/66
R16 2.4 OR LESS CL=.90 BISI 67 DBC + 11/67

R17 CHAR. K INTO (PI P10)/TAU (P21)/(P3)
R17 134 3.24 0.34 YOUNG 65 EMUL + 8/66
R17 1045 3.96 0.15 CALLAHAN 66 FBC + 9/66
R17 AVG 3.84 0.27 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)
R17 STUDENT 3.86 0.16 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R17 FIT 3.768 0.033 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R18 CHAR. K INTO (PI ZP10)/TAU (P41)/(P3)
R18 2027 0.303 0.009 BISI 65 H+HL + 8/66
R18 1.7 0.393 0.099 YOUNG 65 EMUL + 8/66
R18 AVG 0.3037 0.0090 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R18 STUDENT 0.3037 0.0097 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R18 FIT 0.3098 0.0079 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

R19 CHAR. K INTO (MU P10 NEU)/TAU (P51)/(P3)
R19 2845 0.632 0.035 BISI 1 65 H+HL + 8/66
R19 38 0.90 0.16 YOUNG 65 EMUL + 8/66
R19 H 1505 (0.510) (0.017) EICHTEN 68 HLBC + 11/68
R19 H1505 0.503 0.019 HAIDT 71 HLBC + 12/70
R19 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
R19 AVG 0.536 0.054 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.2)
R19 STUDENT 0.527 0.025 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R19 FIT 0.572 0.016 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.536 ± 0.054
ERROR SCALED BY 3.2



R20 CHAR. K INTO (E P10 NEU)/TAU (P61)/(P3)
R20 230 0.90 0.06 BORREANI 64 HBC + 8/66
R20 37 0.90 0.16 YOUNG 65 EMUL + 8/66
R20 854 0.94 0.09 BELLOTT2 67 HLBC 11/67
R20 H 4385 (0.846) (0.021) EICHTEN 68 HLBC + 11/68
R20 H4385 0.850 0.019 HAIDT 71 HLBC + 12/70
R20 2827 0.856 0.040 BRAUN 75 HLBC + 12/75*
R20 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
R20 AVG 0.858 0.016 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R20 STUDENT 0.858 0.018 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R20 FIT 0.8632 0.0098 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R21 K+ INTO (P1+ PI- E+ NEU)/TAU (UNITS 10**--4) (P71)/(P3)
R21 69 6.7 1.5 BIRGE 65 FBC + 8/66
R21 269 5.83 0.63 ELY 69 HLBC + 11/68
R21 500 7.36 0.68 BURQUIN 71 ASPK 12/71
R21 106 7.0 0.9 SCHWEINBE 71 HLBC + 9/71
R21 AVG 6.64 0.40 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R21 STUDENT 6.66 0.48 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

R22 K+ INTO (P1+ PI- MU+ NEU)/TAU (UNITS 10**--4) (P91)/(P3)
R22 1 (2.5) APPROX GREINER 64 EMUL + 8/66
R22 7 2.57 1.55 BISI 67 DBC + 11/67

R23 CHAR. K INTO (E P10 NEU)/(MU2+P12) (UNITS 10**--2) (P61)/(P1+P2)
R23 1679 5.89 0.21 CESTER 66 OSPK + 8/67
R23 5110 6.16 0.22 ESCHSTRUT 68 OSPK + 3/68
R23 AVG 6.02 0.15 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R23 STUDENT 6.02 0.17 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R23 FIT 5.997 0.067 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R24 CHAR. K INTO (PI P10)/(MU NEU) (P21)/(P1)
R24 A4917 0.3277 0.0065 AUERBACH 67 OSPK + 1/74
R24 1600 0.305 0.018 ZELLER 69 ASPK + 10/69
R24 25K 0.328 0.005 WEISSENBE 74 STRC + 7/74*
R24 A AUERBACH 67 CHANGED FROM .3253+-0.0065. SEE COMMENT WITH RATIO R26. 1/74
R24 AVG 0.3268 0.0039 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R24 STUDENT 0.3269 0.0042 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R24 FIT 0.3309 0.0029 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R25 CHAR. K INTO (E P10 NEU)/(MU NEU) (P61)/(P1)
R25 A 295 0.0791 0.0054 AUERBACH 67 OSPK + 1/74
R25 960 0.0775 0.0033 BOTTERILL 68 ASPK + 5/68
R25 561 0.069 0.006 GARLAND 68 OSPK + 4/68
R25 350 0.069 0.006 ZELLER 69 ASPK + 10/69
R25 A AUERBACH 67 CHANGED FROM .0797+-0.0054. SEE COMMENT WITH RATIO R26. 1/74
R25 A THE VALUE .0785+-0.0025 GIVEN IN AUERBACH 67 IS AN AVERAGE OF
R25 A AUERBACH 67 R25 AND CESTER 66 R23. 3/74
R25 AVG 0.0752 0.0024 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R25 STUDENT 0.0753 0.0027 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R25 FIT 0.07582 0.00091 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R26 CHAR. K INTO (MU P10 NEU)/(MU NEU) (P51)/(P1)
R26 A 307 0.0486 0.0040 AUERBACH 67 OSPK + 1/74
R26 G 424 0.0480 0.0037 GARLAND 68 OSPK 1/74
R26 240 0.054 0.009 ZELLER 69 ASPK + 10/69
R26 A AUERBACH 67 CHANGED FROM .0602+-0.0046 BY ERRATUM WHICH BRINGS THE 1/76*
R26 A MU-SPECTRUM CALCULATION INTO AGREEMENT WITH GAILLARD TO APPENDIX B. 1/74
R26 G GARLAND 68 CHANGED FROM .055+-0.004 IN AGREEMENT WITH MU-SPECTRUM 1/74
R26 G CALCULATION OF GAILLARD 70 APPENDIX B. L.G.PONDROM, PRIV.COMM.(73) 1/74
R26 AVG 0.0488 0.0026 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R26 STUDENT 0.0487 0.0028 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R26 FIT 0.0503 0.0014 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

R27 CHAR. K INTO (MU NEU)/TAU (P11)/(P3) 9/66
R27 R 427 (10.38) (0.82) YOUNG 65 EMUL +
R27 R DELETED FROM OVERALL FIT BECAUSE YOUNG 65 CONSTRAINS HIS RESULTS.
R27 R TO ADD UP TO 1. ONLY YOUNG MEASURED MU2 DIRECTLY.
R27 FIT 11.385 0.072 FROM FIT

R28 CHAR. K INTO (E NEU)/(MU NEU) (UNITS 10**--5) (P111)/(P1)
R28 10 1.9 0.7 0.5 BOTTERILL 67 ASPK + 11/67
R28 8 1.8 0.8 0.6 MACEK 69 ASPK + 4/69
R28 112 2.42 0.42 CLARK 72 OSPK + 1/73
R28 534 2.37 0.17 HEARD2 75 SPEC + 11/75*
R28 404 2.51 0.15 HEINTZE 76 SPEC + 2/76*
R28 AVG 2.42 0.11 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R28 STUDENT 2.42 0.12 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

Stable Particles

K±

Data Card Listings

For notation, see key at front of Listings.

R29	CHAR. K INTO (MU P10 NEU)/(E P10 NEU)	(P51)/(P6)	
R29	C1509 0.703 0.056	CALLAHAN 66 HLCB	6/68
R29	5601 0.667 0.017	BOTTERI 68 ASPK +	6/68
R29	H 1398 (0.634) (0.022)	EICHTEH 68 HLCB	10/68
R29	H (0.596) (0.025)	HAIDT 71 HLCB +	12/70
R29	D3480 0.698 0.025	CHIANG 72 OSPK + 1.84 GEV/C K+	9/72
R29	L 554 0.705 0.063	LUCAS 73 HBC - DALITZ PRS ONLY	11/73
R29	B 1585 (0.608) (0.014)	BRAUN 75 HLCB +	1/76*
R29	COMMENTS		
R29	C FROM CALLAHAN 66 WE USE ONLY THE MU3/E3 RATIO AND DO NOT		
R29	C INCLUDE IN THE FIT THE RATIOS MU3/TAU AND E3/TAU, SINCE THEY SHOW		
R29	C LARGE DISAGREEMENTS WITH THE REST OF THE DATA.		
R29	H HAIDT 71 IS A REANALYSIS OF EICHTEH 68.		
R29	H ONLY INDIVIDUAL RATIOS INCLUDED IN FIT (SEE R19 AND R20).		
R29	D CHIANG 72 R29 IS STATISTICALLY INDEPENDENT OF CHIANG 72 R5 AND R6.		
R29	L LUCAS 73 GIVES N(MU3)=554+-7.6PCT, N(E3)=786+-3.1PCT. WE DIVIDE.		
R29	B BRAUN 75 VALUE IS FROM FORM FACTOR FIT. ASSUMES MU-E UNIVERSALITY.		
R29	AVG 0.680 0.013 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R29	STUDENT 0.680 0.015 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT		
R29	FIT 0.663 0.018 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)		
R30	CHAR. K INTO (P10 E NEU GAMMA)/(P10 E NEU) (UNITS 10**-2)		
R30	R13 1.2 0.8 BELLOTTI 67 HLCB + EGAM GT 30MEV (P181)/(P6) 11/67		
R30	R (0.76) (0.28) ROMANO 71 HLCB EGAM GT 10MEV 10/71		
R30	R (0.53) (0.22) ROMANO 71 HLCB + EGAM GT 30 MEV 9/73		
R30	L 16 0.48 0.20 LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R30	L (0.22) (0.19) (0.10) LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R30	L FIRST LJUNG VALUE IS FOR COS(ELECT-GAMMA)=0.9, SECOND VALUE IS		
R30	L FOR COS(ELECT-GAMMA) BETW 0.6 AND 0.9 FOR COMPARISON WITH ROMANO.		
R30	R BOTH ROMANO VALUES ARE FOR COS(ELECT-GAMMA) BETW 0.6 AND 0.9.		
R30	R SECOND VALUE IS FOR COMPARISON WITH SECOND LJUNG VALUE.		
R30	R WE USE LOWEST EGAM CUT FOR TABLE VALUE. SEE ROMANO FOR EGAM DEPEND.		
R30	R30 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)		
R31	K- INTO (PI+ E-)/TOTAL (UNITS 10**-5) (P19)		
R31	TEST OF LEPTON NUMBER CONSERVATION.		
R31	1.5 OR LESS CHANG 68 HBC - 3/68		
R32	CHAR. K INTO (PI NEU NEU)/TOTAL (UNITS 10**-6) (P20)		
R32	C (1.4) OR LESS CL=.90 KLEMS 71 OSPK + T(P1) 117-127MEV 3/74		
R32	C (0.94) OR LESS CL=.90 CABLE 73 CNTR + T(P1) 60-105 MEV 2/74		
R32	C 0.56 OR LESS CL=.90 CABLE 73 CNTR + T(P1) 60-127 MEV 2/74		
R32	L 0 57.0 OR LESS CL=.90 LJUNG 73 HLCB + 9/73		
R32	C KLEMS 71 AND CABLE 73 ASSUME PI SPECTRUM SAME AS KE3 DECAY.		
R32	C SECOND CABLE LIMIT COMBINES CABLE AND KLEMS DATA FOR VECTOR INT.		
R32	L LJUNG 73 ASSUMES VECTOR INTERACTION. 9/73		
R33	CHAR. K INTO (E NEU GAMMA)/TOTAL (UNITS 10**-5) (P21)		
R33	M 7.1 OR LESS MACEK 70 OSPK + P(E) 234 TO 247 12/70		
R33	M ABOVE IS MEASUREMENT OF STRUCTURE-DEPENDENT DECAY ONLY.		
R34	CHAR. K INTO (PI GAMMA)/TOTAL (UNITS 10**-6) (P22)		
R34	4.0 OR LESS CL=.90 KLEMS 71 OSPK + 8/71		
R35	CHAR. K INTO (TAU)/(TAU PRIME)		
R35	USED FOR DELTA I=1/2 TEST. (P3/P4)		
R35	FIT 3.227 0.083 FROM FIT		
R36	CHAR. K INTO (PI 3GAMMA)/TOTAL (UNITS 10**-4) (P23)		
R36	3.0 OR LESS CL=.90 KLEMS 71 OSPK + T(P1) GT 117MEV 8/71		
R37	K+ INTO (PI+ PI+ E- NEU)/(PI+ PI- E+ NEU) (P8)/(P7)		
R37	0 0.013 OR LESS CL=.95 BOURQUIN 71 ASPK 12/71		
R38	CHAR. K INTO (P10 P10 E NEU)/KE3 (UNITS 10**-6) (P24)/(P6)		
R38	0 37.0 OR LESS CL=.90 ROMANO 71 HLCB 12/71		
R38	2 3.8 5.0 1.2 LJUNG 73 HLCB + 9/73		
R39	K+ INTO (PI- E+ MU-)/TOTAL (UNITS 10**-8) (P25)		
R39	K- INTO (PI- E- MU-)/TOTAL IS ALSO INCLUDED HERE		
R39	2.8 OR LESS CL=.90 BEIER 72 OSPK +- 9/72		
R40	K+ INTO (PI+ E+ MU-)/TOTAL (UNITS 10**-8) (P26)		
R40	K- INTO (PI- E- MU-)/TOTAL IS ALSO INCLUDED HERE		
R40	1.4 OR LESS CL=.90 BEIER 72 OSPK +- 9/72		
R41	CHAR. K INTO (MU 3NEU)/TOTAL (UNITS 10**-6) (P27)		
R41	P 0 6.0 OR LESS CL=.90 PANG 73 CNTR + 11/73		
R41	P PANG 73 ASSUMES MU SPECTRUM FROM NEU-NEU INTERACTION OF BARDIN 70. 3/74		
R42	CHAR. K INTO (P10 MU NEU GAM)/TOTAL (UNITS 10**-5) (P28)		
R42	0 6.1 OR LESS CL=.90 LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R43	CHAR. K INTO (E P10 NEU)/(PI P10) (P6)/(P2)		
R43	L 786 0.221 0.012 LUCAS 73 HBC - DALITZ PRS ONLY 11/73		
R43	L LUCAS 73 GIVES N(E3)=786+-3.1PCT, N(P12)=3564+-3.1PCT. WE DIVIDE. 11/73		
R43	FIT 0.2291 0.0031 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R44	CHAR. K INTO (PI 2P10)/(PI P10) (P4)/(P2)		
R44	L 574 0.081 0.005 LUCAS 73 HBC - DALITZ PRS ONLY 11/73		
R44	L LUCAS 73 GIVES N(P1 2P10)=574+-5.9 PCT, N(P12)=3564+-3.1 PCT. 11/73		
R44	L WE QUOTE 0.5*(PI 2P10)/(PI 2P12) WHERE 0.5 IS BECAUSE ONLY DALITZ 11/73		
R44	L PAIR P10'S WERE USED. 11/73		
R44	FIT 0.0822 0.0022 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R45	CHAR. K INTO (MU NEU GAMMA)/TOTAL (UNITS 10**-3) (P12)		
R45	12 5.8 3.5 WEISSENBE 74 STRC + E-GAMMA GT 9 MEV 7/74*		
R46	CHAR. K INTO (PI+ E-)/(PI+ PI- E NEU) (UNITS 10**-3) (P15)/(P7)		
R46	B 41 7.0 1.3 BLOCH 75 SPEC + 11/75*		
R46	B BLOCH 75 QUOTES THIS RESULT MULTIPLIED BY OUR 1974 KE4 BR.FRAC. 11/75*		
R47	CHAR. K INTO (E NEU GAM)/(E NEU) (P21)/(P11)		
R47	H 56 1.05 0.25 0.30 HEARD 75 SPEC + P(E) 236 TO 247 11/75*		
R47	H ABOVE IS SENSITIVE ONLY TO STRUCTURE DECAY TERM. 11/75*		

Note on Slope Parameter for K → 3π Decays

As was discussed in Section VI B.1 of the text, for the 3π decays of the K mesons we list the slope parameter "g" which is defined, as in that section, by

$$|M|^2 \propto 1 + g \frac{(s_3 - s_0)}{m_{\pi^+}^2} + h \left(\frac{s_3 - s_0}{m_{\pi^+}^2} \right)^2 + j \frac{(s_2 - s_1)}{m_{\pi^+}^2} + k \left(\frac{s_2 - s_1}{m_{\pi^+}^2} \right)^2 + \dots \quad (1)$$

where

$$s_i = (p_K - p_i)^2 = (m_K - m_i)^2 - 2m_K T_i \quad (2)$$

$$s_0 = \frac{1}{3} \sum s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2) \quad (3)$$

p_K, p_i are the four-vectors for the K and the i^{th} pion, and the index 3 refers to the odd pion, i.e., the third pion in the decays listed below.

We refer to the three possible charged decays as $\tau, \tau',$ and τ^0 :

$$\begin{aligned} \tau^\pm & K^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp \\ \tau^\pm & K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm \\ \tau^0 & K_L^0 \rightarrow \pi^+ \pi^- \pi^0 \end{aligned}$$

The measurements of g vary considerably beyond the authors' quoted errors as can be seen in the ideograms associated with the GT+, GT-, and GTP subsections of the K± Data Card Listings and the GTO subsection of the K_L⁰ Listings. Appendix I discusses tests of the ΔI = 1/2 rule utilizing these slopes.

There is no indication of a CP-violating asymmetry in K_L⁰ decay as measured by the asymmetry parameter σ_\pm (equivalent to j) given in subsection A of the K_L⁰ Listings.

There is conflicting evidence regarding the presence or absence of the second order terms in Eq. (1), even when only the high-statistics experiments are considered. In K_L⁰ → π⁺π⁻π⁰, MESSNER 74 find a significant h and k, while others find them smaller or consistent with zero. The results on h for the high-statistics experiments are as follows:

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

Experiment	Number of events	Value of h
MESSNER 74	509 K	+0.079 ± 0.007
BUCHANAN 75	56 K	+0.041 ± 0.024
BUCHANAN 70	36 K	consistent with 0
ALBROW 70	29 K	-0.009 ± 0.016
BISI 74	20 K	0.000 ± 0.009

All values other than MESSNER 74 have been converted by us from the authors' parametrizations into the parametrization of Eq. (1) above.

FORD 72 (1.5M events) have studied $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ and find that the χ^2/DF goes from 1.38 to 1.20 for $DF \approx 150$ when the second order and the CP violation terms are added. However, the authors state that since their Coulomb correction is larger than the experimental errors and is not well known, it is difficult to interpret these results. LUCASI 73 (81K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ events) note a quadratic behavior in their data, but find that it can be explained by normal measurement error and subsequent fitting, which together result in a depletion of events in the center of the Dalitz plot and a bunching of events near the kinematic limit. Because of this they state that it is not appropriate for them to quote a value of the quadratic coefficient. MAST 69 (51K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ events), HOFFMASTER 72 (40K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ events) and SMITH 75 (27K $K^{\pm} \rightarrow \pi^0 \pi^0 \pi^{\pm}$ events) also find no evidence for quadratic behavior.

HEUSSE 70 have studied the $K_L^0 \rightarrow \pi^0 \pi^0 \pi^0$ decay where only a second order term could explain deviation from uniformity of the Dalitz plot. They also get results consistent with a zero coefficient.

We list the value of g obtained when the quadratic terms were dropped from the fit unless otherwise noted in the data card footnotes.

In the literature other definitions of slope parameters have appeared. We have converted to the definition of g in Eq. (1) whatever experimental quantity has been reported. We give the conversion to the definition (1) for the most widely used parametrizations and tabulate the conversion factors for the reader's convenience.

a) For analysis of charged K's the expression often used is:

$$|M|^2 = 1 + a_y y$$

with

$$y = \frac{3T_3 - Q}{Q}, \quad Q = m_K - \sum m_i$$

The relevant formulae are:

$$y = -\frac{3}{2} \frac{s_3 - s_0}{m_K Q} + \Delta$$

with

$$\Delta = \frac{m_1 - m_3}{Q} \left(2 - \frac{m_3 + m_1}{m_K} \right)$$

and

$$g = \frac{-c_y a_y}{1 + a_y \Delta}, \quad \text{with } c_y = \frac{3}{2} \frac{m_{\pi^+}^2}{m_K Q}$$

b) For the analysis of K^0 decay the expression often used is:

$$|M|^2 = 1 + 2a_t \frac{m_K}{m_{\pi^+}^2} (2T_3 - T_{3\max})$$

with

$$T_{3\max} = \frac{m_K^2 + m_3^2 - 4m_{12}^2}{2m_K} - m_3$$

The relevant transformations are

$$T_3 = -\frac{s_3 - s_0}{2m_K} + \frac{Q}{3} (1 + \Delta)$$

and

$$g = \frac{-2a_t}{1 + a_t c_t}$$

with

$$c_t = \frac{2m_K}{m_{\pi^+}^2} \left[\frac{2}{3} Q(1 + \Delta) - T_{3\max} \right]$$

c) Other K^0 authors use the same form of matrix element as given in b) above, but define

$$T_{3\max} = \frac{2}{3} Q$$

Stable Particles

K±

Data Card Listings

For notation, see key at front of Listings.

The relevant transformation is then

$$g = \frac{-2a_u}{1 + a_u c_u}, \text{ with } c_u = \frac{4m_K}{3m_{\pi^+}} Q\Delta.$$

d) Older K⁰ analyses were done using

$$|M|^2 = 1 + a_v \frac{T_3}{m_K}.$$

The relevant transformation is then

$$g = \frac{-c_v a_v}{1 + d_v a_v}, \text{ with } c_v = \frac{m_{\pi^+}^2}{2m_K^2}$$

and

$$d_v = \frac{Q}{3m_K} (1 + \Delta).$$

For the reader's convenience we give a table of numerical values for Q, T_{3max}, Δ, c_y, and c_t, obtained using the masses from the current edition.

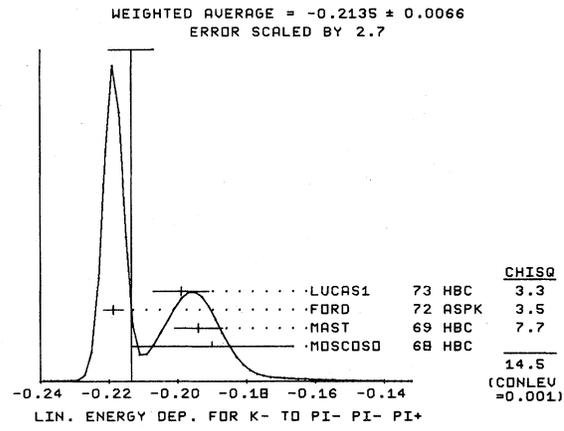
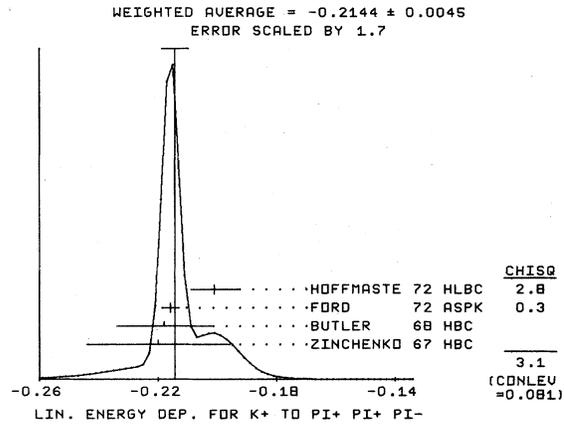
	τ [±]	τ [±]	τ ⁰
Q	75.00	84.21	83.60
T _{3max}	48.10	53.22	53.91
Δ	0.0	-0.0789	0.0798
c _y	0.7891	0.7028	0.7023
c _t	0.0963	-0.0768	0.3204
c _u	0.0	-0.2247	0.2272
c _v	0.0400	0.0400	0.0393
d _v	0.0506	0.0524	0.0605

10 CHARGED K ENERGY DEPENDENCE OF DALITZ PLOT

RELATED TEXT SECTION VI B.1, APPENDIX I, AND MINI-REVIEW ABOVE
MATRIX ELEMENT SQUARED = 1 + G (S3-S0)/(MPI**2)

GT+ LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS K+ INTO PI+ PI+ PI-
 GT+ THESE EXPTS FIT **2-1+AY. WE LIST G IN THE MAIN LISTING AND
 GT+ GIVE AY AT RIGHT. G=-1.5*AY*(MPI**2)/(MK*Q). SEE NOTE ABOVE.
 GT+Z 5428 -0.22 0.024 ZINCHENKO 67 HBC + AY=0.28+-03 10/69
 GT+ 9994 -0.218 0.016 BUTLER 68 HBC + AY=0.277+-020 10/69
 GT+ 617898 (-0.196) (0.012) GRAUMAN 70 HLCB + AY=0.228+-030 8/70
 GT+Q 750K -0.2158 0.0028 FORD 72 ASPK + AY=0.2734+-0035 4/72
 GT+ 39819 -0.201 0.008 HOFFMASTE 72 HLCB + INCLUDES GRAUMAN 1/71
 GT+ Q THIS VALUE OF AY IS FROM A QUADRATIC FIT WITH Y**2 COEF=.030+-010. 4/72
 GT+ Q A LINEAR FIT IS QUOTED ONLY FOR THEIR COMBINED K+ AND K- SAMPLE. 4/72
 GT+ Q IT GIVES AY=0.2737+-0032. THE QUADRATIC FIT TO THE COMBINED 4/72
 GT+ Q SAMPLE GIVES AY=0.2752+-0033 AND Y**2 COEFF=0.025+-010. 4/72
 GT+ Q (CHISO/DF)=1.38 FOR LINEAR FIT AND 1.20 FOR QUADRATIC FIT. 1/73
 GT+ G EMULS. DATA ADDED - ALL EVENTS INCLUDED BY HOFFMASTE 72 1/71
 GT+ Z ALSO INCLUDES DBC EVENTS.
 GT+ AVG -0.2144 0.0045 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)
 GT+ STUDENT -0.2146 0.0029 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
 (SEE IDEOGRAM BELOW)

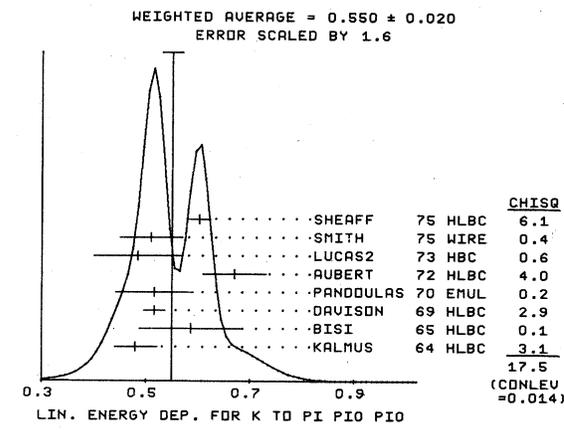
GT- LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS K- INTO PI- PI- PI+
 GT- FOR DEFINITION OF AY SEE NOTE IN ABOVE SECTION GT+.
 GT- F 1347 (-0.220) (0.035) FERRO-LUZ 61 HBC - AY=0.28+-045 10/69
 GT- M 5778 -0.190 0.023 MOSCOSO 68 HBC - AY=0.242+-029 10/69
 GT- 50919 -0.194 0.007 MAST 69 HBC - AY=0.247+-009 10/69
 GT-Q 750K -0.2187 0.0028 FORD 72 ASPK - AY=0.2770+-0035 4/72
 GT- 81K -0.199 0.008 LUCAS1 73 HBC - AY=0.252+-011 10/72
 GT- Q THIS VALUE OF AY IS FROM A QUADRATIC FIT WITH Y**2 COEF=-0.20+-010. 4/72
 GT- Q SEE ALSO THE NOTE Q IN THE GT+ SECTION ABOVE.
 GT- F NO RADIATIVE CORRECTIONS INCLUDED.
 GT- M ALSO INCLUDES DBC EVENTS.
 GT- AVG -0.2135 0.0066 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.7)
 GT- STUDENT -0.2142 0.0039 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
 (SEE IDEOGRAM BELOW)



DG ((GT+)-(GT-))/((GT+)+(GT-)) IN PERCENT
 DG A NON-ZERO VALUE FOR THIS QUANTITY INDICATES CP VIOLATION
 DG 3.2M -0.70 0.53 FORD 70 ASPK 11/70

GTP LINEAR ENERGY DEPENDENCE (G) FOR TAU PRIME DECAY CHA-K INTO PI PI0
 GTP 1792 0.48 0.04 KALMUS 64 HLCB + 10/69
 GTP 1874 0.586 0.098 BISI 65 HLCB + ALSO HBC 10/69
 GTP 4948 0.516 0.020 DAVIDSON 69 HLCB + ALSO EMUL 10/69
 GTP 198 0.516 0.074 PANDOLAS 70 EMUL + 10/70
 GTP A1365 0.67 0.06 AUBERT 72 HLCB 1/73
 GTP 574 0.484 0.084 LUCAS2 73 HBC - DALITZ PRS ONLY 9/73
 GTP S 27K 0.510 0.060 SMITH 75 WIRE + 12/75*
 GTP 5635 0.602 0.021 SHEAFF 75 HLCB + 2/76*
 GTP A WE GIVE LINEAR TERM OF HIGHER ORDER FIT. EQ.1 OF APP.II,AUBERT 72. 1/73
 GTP S SMITH 75 MEASURES QUADRATIC COEFFICIENT H=0.009+-0.040. 12/75*

GTP AVG 0.550 0.020 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
 GTP STUDENT 0.542 0.019 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
 (SEE IDEOGRAM BELOW)



Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

Note on K_{l3}[±] and K_{l3}⁰ Form Factors

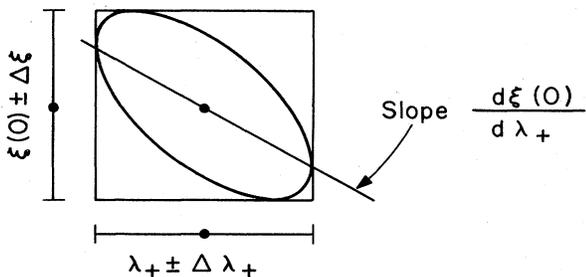
Definitions of the parameters λ₊, ξ(0), λ₀, |f_S/f₊| and |f_T/f₊| and a general discussion of the methods of analysis are given in Section VI B.2 of the text.

This note describes the contents of the Data Card Listings for the two K_{μ3} parametrizations, (λ₊, ξ(0)) and (λ₊, λ₀), which were discussed in the text. Problems related to our data entries for individual experiments are discussed and a comparison of results is given.

K_{μ3} Experiments

The matrix element for K_{μ3} decay, assuming a pure vector current, is given by Eq. (2) in Section VI B.2 of the text. Most experiments appear to be compatible with the assumption that f₊ depends linearly on t and that f₋ is constant. Only DALLY 72 (K_{μ3}⁰) appears to require λ₋ ≠ 0 (by about three standard deviations). A single data bin at low q² seems to be responsible. The effect is not observed in the high-statistics experiment of DONALDSON 74 (also K_{μ3}⁰).

λ₊, ξ(0) Parametrization: λ₊ data from K_{μ3} decay are entered into the K[±] and K_L⁰ sections of the Data Card Listings in subsection L+M. The corresponding ξ(0) values are entered in subsection XIA, XIB, or XIC, depending on whether Method A, B, or C, discussed below and in the text, was used. The data cards contain the values, one-standard-deviation errors Δλ₊ and Δξ(0), as well as the correlation dξ(0)/dλ₊, all indicated on the e^{-1/2} likelihood contour below. The correlations are given on the right side of the ξ(0) data cards.



XBL 743-2682

λ₊, λ₀ Parametrization: This parametrization is used in recent K_{μ3} analyses. To facilitate comparison between experiments, we convert earlier experiments from the (λ₊, ξ(0)) parametrization to (λ₊, λ₀) whenever possible (i.e., when λ₊ and ξ(0) values, errors, and correlations are given). The transformation between these parametrizations is:

$$\lambda_0 = \lambda_+ + a\xi(0) ,$$

$$\Delta\lambda_0^2 = (1 + 2a \frac{d\xi(0)}{d\lambda_+}) \Delta\lambda_+^2 + a^2 \Delta\xi^2 ,$$

$$\frac{d\lambda_0}{d\lambda_+} = 1 + a \frac{d\xi(0)}{d\lambda_+} ,$$

where $a = m_\pi^2 / (m_K^2 - m_\pi^2)$. The λ₀ value, the one-standard-deviation error Δλ₀, and the correlation dλ₀/dλ₊ are given in subsection L0 of the data cards.

We also convert (λ₊, λ₀) results into the (λ₊, ξ(0)) parametrization whenever possible so that subsection L0 is essentially equivalent to the three subsections XIA, XIB, and XIC.

Individual analyses have used a variety of parametrizations. Problems arise when trying to express their results in terms of the parametrizations used here. The discussion of these problems is divided into three sections corresponding to the three methods of analyses discussed in the text.

Method A: Dalitz plot analyses and pion spectrum analyses usually determine λ₊ and ξ(0) (or λ₀) values, errors, and correlation. Such measurements are entered in the L+M, XIA, and L0 subsections. They give rise to the error ellipses shown in Figs. 1 and 2.

Some analyses of this type fix λ₊ and determine ξ(0), e.g., CARPENTER 66 and PEACH 73 (both K_{μ3}⁰). We enter ξ(0) and dξ(0)/dλ₊ in the XIA section and give the fixed λ₊ value in the data card footnote. The ξ(0) error is parenthesized because it does not include the uncertainty in the value of λ₊. These results, transformed to λ₀ measurements, give rise to bands in Fig. 2.

In some cases, we alter an error from its published value in order to obtain an error ellipse with a width which matches the error in ξ(0) for

Stable Particles

K^\pm

Data Card Listings

For notation, see key at front of Listings.

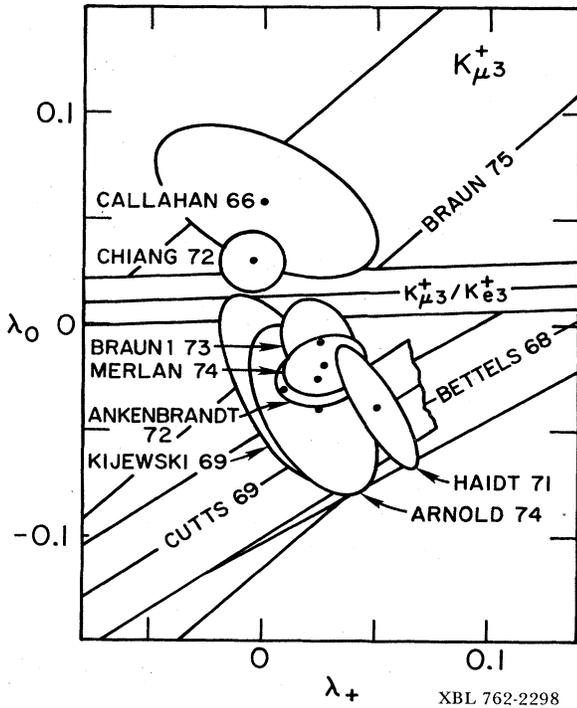


Fig. 1. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$.

fixed λ_+ . These adjustments are noted in the $\xi(0)$ data card footnotes, e.g., for CALLAHAN 66 and HAITD 71 (K^+ subsection XIA), where the published errors and correlation violate the constraint $|C_{\lambda\xi}| < 1$ on the normalized correlation coefficient $C_{\lambda\xi}$ given by

$$C_{\lambda\xi} = \frac{\Delta\lambda_+}{\Delta\xi} \frac{d\xi(0)}{d\lambda_+}$$

In some cases, e.g., BRAUN 73, the parametrization used is $\lambda_+, \xi(0), \xi(t^*)$, where t^* is the weighted average of t with weighting according to the sensitivity to ξ . In this case we do not use $\xi(0)$. It is a badly determined parameter comparable to λ_- or the slope of $\xi(t)$. Instead, we use

$$\xi(0) = \xi(t^*) (1 + \lambda_+ t^*)$$

$$\frac{d\xi(0)}{d\lambda_+} = \frac{d\xi(t^*)}{d\lambda_+} (1 + \lambda_+ t^*) + \xi(t^*) t^*$$

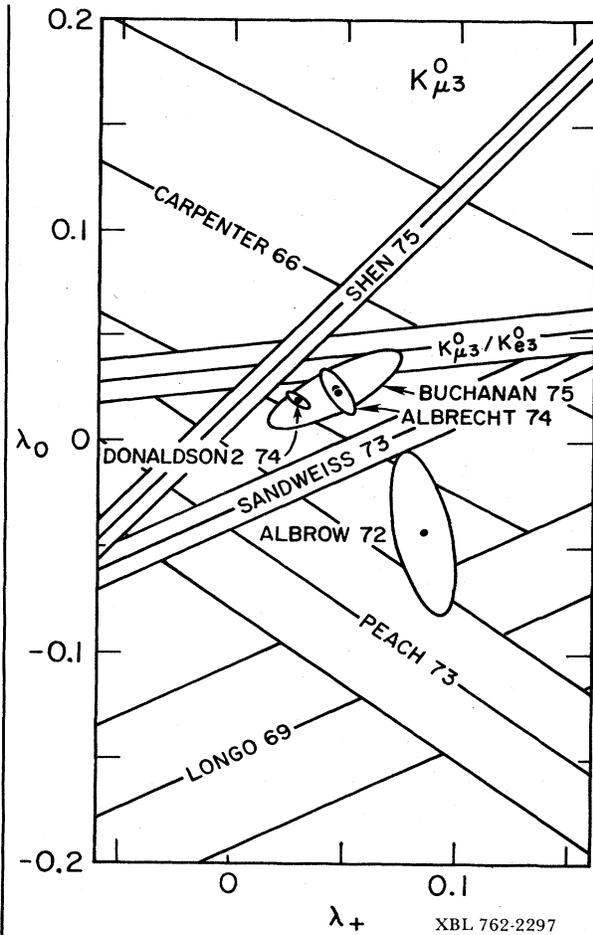


Fig. 2. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^0$.

With the BRAUN 73 values, $\lambda_+ = 0.027$, $\xi(6.6) = -0.34 \pm 0.20$, and $d\xi(6.6)/d\lambda_+ = -14$, we obtain

$$\xi(0) = (-0.40 \pm 0.24) - 19(\lambda_+ - 0.027);$$

or for their fitted $\lambda_+ = 0.025 \pm 0.017$, we get $\xi(0) = -0.36 \pm 0.40$.

Method B: Branching ratio experiments cannot determine λ_+ and $\xi(0)$ simultaneously, but simply fix a relationship between them, given in Section VI B.2 of the text. Results are usually quoted as values of $\xi(0)$ at fixed λ_+ . We list these results in subsection XIB, but we do not average them because the λ_+ values differ. Instead, we compute a combined result by using the relations in the

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

text and our fitted values of $\Gamma(K_{\mu 3}^+)/\Gamma(K_{e 3}^+)$ and $\Gamma(K_{\mu 3}^0)/\Gamma(K_{e 3}^0)$, which include the branching ratios from these experiments. The branching ratios from our current edition and the results for $\xi(0)$ and λ_0 evaluated at $\lambda_+ = 0.030$ are

	K [±]	K _L ⁰
$\Gamma(K_{\mu 3})/\Gamma(K_{e 3})$	0.663 ± .018 (S=1.7)	0.696 ± .017
$\xi(0)$	-0.20 ± .15 (S=1.7)	+0.09 ± .13
$d\xi(0)/d\lambda_+$	-11.9	-10.3
λ_0	+0.014 ± .012 (S=1.7)	+0.038 ± .011
$d\lambda_0/d\lambda_+$	+0.04	+0.12

The scale factor S is the amount by which the error has been multiplied in order to compensate for discrepancies in the branching ratios. These λ_0 results give rise to the $K_{\mu 3}/K_{e 3}$ bands in Figs. 1 and 2.

Method C: Polarization experiments measure $\langle \xi(t) \rangle$, the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$. Such measurements are entered in subsection XIC.

To reinterpret these results in the $(\lambda_+, \xi(0))$ parametrization, we recognize that $\lambda_+ = 0$ corresponds to $\xi(t)$ constant (always assuming $\lambda_- = 0$) so that

$$\xi(0) \Big|_{\lambda_+=0} \equiv \langle \xi(t) \rangle .$$

The correlation with λ_+ is given by the following relations (valid for small λ_+):

$$\xi(0) \approx \langle \xi(t) \rangle \left(1 + \lambda_+ \left\langle \frac{t}{m_\pi^2} \right\rangle \right) ,$$

$$\frac{d\xi(0)}{d\lambda_+} \approx \langle \xi(t) \rangle \left\langle \frac{t}{m_\pi^2} \right\rangle ,$$

where $\langle t/m_\pi^2 \rangle$ is the average value of t weighted by the sensitivity to $\xi(t)$. These results, transformed to λ_0 and $d\lambda_0/d\lambda_+$ values, are entered in subsection I0 and give rise to bands in Figs. 1 and 2.

In Figs. 1 and 2, we disregard those polarization measurements for which $d\xi(0)/d\lambda_+$ is not obtainable. Also we disregard MERLAN 73 because

the signs of $\xi(0)$ and $d\xi(0)/d\lambda_+$ are opposite, whereas the above equation requires them to be the same (since $t > 0$).

Comparison of $K_{\mu 3}$ Experiments: Figures 1 and 2 show the likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$ and $K_{\mu 3}^0$ respectively.

The $K_{\mu 3}^+$ Dalitz plot results (ellipses) shown are fairly consistent and appear to cluster between the $K_{\mu 3}/K_{e 3}$ result and the polarization results of BETTELS 68 and CUTTS 69. The $K_{\mu 3}^0$ results are much less consistent with a small cluster appearing in the neighborhood of the DONALDSON2 74 result.

χ^2 fits to the results shown in Fig. 1 and Fig. 2 yield the following values for λ_+ and λ_0 . The corresponding values of $\xi(0)$ are also given.

	K _{μ3} ⁺	K _{μ3} ⁰
λ_+	+0.027 ± .008*	+0.034 ± .006*
λ_0	-0.009 ± .007*	+0.021 ± .006*
$d\lambda_0/d\lambda_+$	-0.15	-0.24
χ^2/DF	34/18	72/12
S	1.4	2.5
.....		
$\xi(0)$	-0.45 ± .14*	-0.17 ± .10*
$d\xi(0)/d\lambda_+$	-14.	-15.

*All errors have been increased by the scale factor $S = (\chi^2/DF)^{1/2}$ to take into account the discrepancies between measurements.

In view of the large χ^2/DF of these fits, especially $K_{\mu 3}^0$, the fit results should be taken with a grain of salt. The largest contributors to χ^2 in the $K_{\mu 3}^+$ case are CHIANG 72 with 10.4, and the polarization results, BETTELS 68 with 5.4 and CUTTS 69 with 4.2. In the $K_{\mu 3}^0$ case the largest contributors are the polarization results of SANDWEISS 73 with 17, LONGO 69 with 14, and SHEN 75 with 12, and the Dalitz plot results of ALBROW 72 with 11, ALBRECHT 74 with 7.3, and PEACH 73 with 5.3. All other χ^2 values were less than 4.

The DONALDSON2 74 result

$$\lambda_+ = .030 \pm .003$$

$$\lambda_0 = .019 \pm .004$$

Stable Particles

K±

clearly dominates the statistics in the K0_u3 case. The lambda_+ value is consistent with the K_e3 value of lambda_+, and with the pole approximation

f_+(t) = f_+(0) * (m_K^2 / (m_K^2 - t))

Their f_0(t) extrapolates linearly to the Callan-Treiman point. It is less than two standard deviations from the K_u3/K_e3 result.

K_e3 Experiments

The f_- term of the matrix element [Eq. (2) text Section VI.B.2] can be neglected for K_e3 because it is proportional to the lepton mass. The f_+ term is usually assumed to be linear in t=q^2 = (P_K - P_pi)^2, the square of the four-momentum transfer, i.e., the effective mass of the lepton pair. We quote the linear coefficient lambda_+^e (L+E on the data cards).

There has been some suggestion of departure from linearity [CHIEN 71 (K_e3^0) and Chounet, Gaillard, and Gaillard^1 - Review] but no compelling evidence. The lambda_+ results are fairly consistent and the average values are

K_e3^+ : lambda_+ = 0.0285 +/- 0.0043
K_e3^0 : lambda_+ = 0.0288 +/- 0.0028 (S=1.4)

where the K_e3^0 error has been multiplied by the scale factor 1.4 to compensate for inconsistencies (see ideogram in K_L^0 section L+E).

Many of the changes in this section were stimulated by the comments and criticism of L. M. Chounet, and by the excellent reviews of Gaillard and Chounet^1 and Chounet, Gaillard, and Gaillard.^2

References

- 1. M. K. Gaillard and L. M. Chounet, K_L3 Form Factors, CERN 70-14 (May 1970), and Phys. Letters 32B, 505 (1970).
2. L. M. Chounet, J. M. Gaillard, and M. K. Gaillard, Physics Reports 4C, 199 (1972).

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle ID, form factor values, and experimental details. Includes sections for CHARGED K FORM FACTORS, RELATED TEXT SECTION VI B.2 AND MINI-REVIEW ABOVE, and IMAGINARY PART OF XI (TEST OF T REVERSAL).

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K±

LAMBDA 0 (LINEAR ENERGY DEPENDENCE OF FO IN KMU3 DECAY) WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(O) INTO VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KE3 DECAY) FOR RAD. COR. OF KE3 DP SEE GINSBURG 67 AND BECHERRAWY 70.

FS/F+ RATIO OF SCALAR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE) FT/FT+ RATIO OF TENSOR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE)

FT/FT+ RATIO OF TENSOR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE) FT/FT+ RATIO OF TENSOR TO F+ COUPLINGS FOR KE3 DECAY

KE4 KE4 DECAY FORM FACTORS ARE GIVEN IN THE FOLLOWING PAPERS

REFERENCES FOR CHARGED K

BIRGE 56 NC 4 834 BIRGE, PERKINS, PETERSON, STORK, WHITEHEA (LRL)
LDFD 56 PR 102 927 ILOFF, GOLDHABER, LAMNUTTI, GILBERT (LRL)
ALEXANDER 57 NC 6 478 ALEXANDER, JOHNSTON, OCEALLAIGH (DUBLIN INST)
COHEN 57 FUND. CONS. PHYS. +CROWE, DUMOND (ATOMICS INTER. +LRL+CIT)
EISENBERG 58 NC 8 663 EISENBERG, KOCH, LOHRMANN, NIKOLIC + (BERN)
BURROUGHS 59 PRL 2 117 BURROUGHS, CALDWELL, FRISCH, HILL + (MIT)
TAYLOR 59 PR 114 359 S TAYLOR, HARRIS, OREAR, LEE, BAUMEL (COLUMBIA)
FREDEN 60 PR 118 564 S C FREDEN, F C GILBERT, R S WHITE (LRL)
BARKAS 61 PR 121 1209 BARKAS, DYER, HASON, NORRIS, NICKOLS, SMIT (LRL)
BHOWMICK 61 NC 20 857 B BHOWMICK, G JAIN, P C MATHUR (DELHI UNIV)
FERRO-LU 61 NC 22 1087 FERRO-LUZZI, MILLER, MURRAY, ROSENFELD+ (LRL)
NORDIN 61 PR 123 2166 PAUL NORDIN JR (LRL)
ROE 61 PRL 7 346 ROE, SINCLAIR, BROWN, GLASER + (MICH+LRL)
BOYARSKI 62 PR 128 2398 BOYARSKI, LIDH, NITEMELA, RITSON (MIT)
BROWN 62 PRL 8 450 BROWN, KADYK, TRILLING, ROE+ (LRL, MICH)
BARKAS 63 PRL 11 26 W H BARKAS, J N DYER, H H HECKMAN (LFL)
BORREANI 64 PL 12 123 G BORREANI, G RINAUDO, A WERBROUCK (TURIN)
CALLAHAN 64 PR 136 B 1463 A CALLAHAN, R MARCH, R STARK (WISCONSIN)
CAMERINI 64 PRL 13 318 CAMERINI, CLINE, FRY, POWELL (WISCONSIN+LRL)
CLINE 64 PRL 13 101 D CLINE, W F FRY (LRL)
GIACOMELLI 64 NC 34 1134 GIACOMELLI, MONTI, QUARENI + (BOLOGNA, MUNICH)
GREINER 64 PR 13 288 D GREINER, L W OSBORNE, R BARKAS (LRL)
JENSEN 64 PR 136 B1431 JENSEN, SHAKLEE, ROE, SINCLAIR (MICH)
KALMUS 64 PRL 13 99 +KERNAN, J U, POWELL, DOWD (LRL, WISC)
SHAKLEE 64 PR 136 B 1423 SHAKLEE, JENSEN, ROE, SINCLAIR (MICH)
BIRGE 65 PR 139 B 1600 BIRGE, ELY, GIDAL, CAMERINI, CLINE + (LRL+WISC)
BIST 65 NC 35 768 BISI, BORREANI, CESTER, FERRARO + (TORINO)
BIST I 65 PR 139 B 1068 BORREANI, MARZARI, CHIESA, R RINAUDO+ (TORINO)
BORREANI 65 PR 139 B 1066 BORREANI, MARZARI, CHIESA, R RINAUDO+ (TORINO)
CALLAHAN 65 PRL 15 129 A CALLAHAN, D CLINE (WISCONSIN)
CAMERINI 65 NC 37 1795 +CLINE, GIDAL, KALMUS, KERNAN (WISC+LRL)
CLINE 65 PL 15 293 A CLINE, W F FRY (WISCONSIN)
CUTTS 65 PR 138 B969 CUTTS, ELIOFF, STIENING (LRL)
DE MARCO 65 PR 140 B 1430 DE MARCO, GROSSO, RINAUDO (TORINO+CERN)
FITCH 65 PR 140 B 1088 FITCH, QUARLES, WILKINS (PRINCETON+MT HOLYOKE)
GREINER 65 ARNS 15 67 QUOTED BY BARKAS (LRL)
STAMER 65 PR 138 B 440 STAMER, HUETTER, KOLLER, TAYLOR, GRAUMAN (STEVE)
TRILLING 65 UCRL 16473 GEORGE H TRILLING (LFL)
UPDATED FROM 1965 ARGONNE CONF., PAGE 5.
YOUNG 65 UCRL 16362 PDH-SHIEN YOUNG (THEISIS, BERKELEY) (LRL)
ALSO 67 PR 156 1464 P-S YOUNG, W. Z. OSBORNE, W. H. BARKAS (LFL)

CALLAHA1 66 PR 150 1153 CALLAHAN 66 NC 44A 90 CALLAHAN, CAMERINI + (WISC, LRL, RIVERSIDE, BARI)
CESTER 66 PL 21 343 CESTER, ESCHSTRUTH, ONEILL + (PRINCETON-PENN)
ALSO 67 AUERBACH, FOOTNOTE 1.

AUERBACH 67 PR 155 1505 +DOBBS, MANN, MCFARLANE, WHITE+ (PENN, PRIN)
ERRATUM
BELLOTTI 67 HEIDELBERG CONF BELLOTTI, PULLIA (MILAN)
BELLOTTI 67 NC 52A 1287 BELLOTTI, FIORINI, PULLIA (MILAN)
ALSO 66 PL 20 690 BELLOTTI, FIORINI, PULLIA+ (MILAN)
BISI 67 PL 25B 572 BISI, CESTER, CHIESA, VIGONE (TORINO)
BOTTERILL 67 PRL 19 982 BOTTERILL, BROWN, CORBETT, CULLIGAN + (OXFORD)
ALSO 68 BOTTERILL
BOWEN 67 PR 154 1314 BOWEN, MANN, MCFARLANE, HUGHES+ (PENN-PRINCETON)

CLINE1 67 HEIDELBERG CONF CLINE, HAGGERTY, SINGLETON, FRY+ (WISCONSIN)
CLINE2 67 HERCEG NOVI TBL.4 D. CLINE, PROC. INT'L SCH. ON ELEM. PART. PHYSICS
FLETCHER 67 PRL 19 98 FLETCHER, BEIER, EDWARDS + (ILLINOIS)
FORD 67 PRL 18 1214 +LEMONICK, NAUENBERG, PIROUE (PRINCETON)
IMLAY 67 PR 160 1203 IMLAY, ESCHSTRUTH, FRANKLIN+ (PRINCETON)
KALMUS 67 PR 159 1187 KALMUS, KERNAN (LRL)
ZINCHENK 67 RUTGERS (THEISIS) ZINCHENKO (RUTGERS)

BETTELS 68 NC 56A 1106 AACHEN-BARI-IBERGEN-CERN-EP-NIJMEGEN-ORSAY+
ALSO 71 HAIDT
BOTTERILL 68 PR 171 1402 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL 68 PR 174 1661 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL 68 PRL 21 766 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BUTLER 68 UCRL-18420 +BLAND, GOLDHABER, GOLDHABER, HIRATA+ (LRL)
CHANG 68 PRL 20 510 CHANG, YODH, EHRLICH, PLANO+ (MARYLAND, RUTGERS)

CHEN 68 PRL 20 73 CHEN, CUTTS, KIJEWSKI, STIENING + (LRL, MIT)
EICHEN 68 PL 27B 586 AACHEN-BARI-CERN-EP-ORSAY-PADOVA-VALENCIA
EISLER 68 PR 169 1090 EISLER, FUNG, MARATECK, MEYER, PLANO (RUTGERS)
ESCHSTRUTH 68 PR 165 1487 ESCHSTRUTH, FRANKLIN, HUGHES+ (PRINCETON, PENN)
GARLAND 68 PR 167 1225 +MSTOW, SOLODOVNIKOVA, FADEEV + (LJUNG)
MOSCOSO 68 THESIS M L MOSCOSO (UNIV PARIS ORSAY)
CUTTS 69 PR 184 1380 +STIENING, WIEGAND, DEUTSCH (LRL, MIT)
ALSO 68 PRL 20 955 CUTTS, STIENING, WIEGAND, DEUTSCH (LRL, MIT)
DAVISON 69 PR 180 1333 +BACASTOW, BARKAS, EVANS, FUNG, PORTER+ (UCR)
ELY 69 PR 180 1319 ELY, GIDAL, HAGOPIAN, KALMUS+ (LOUCHWISC+LRL)
EMMERSON 69 PRL 23 393 EMMERSON, QUIRK (OXFORD)

HERZO 69 PR 186 1403 +BANNER, BEIER, BERTRAM, EDWARDS + (ILL)
KIJEWSKI 69 UCRL-18433 THESIS P K KIJEWSKI (LRL)
LOBKOWICZ 69 PR 185 1676 +MELISSINOS, NAGASHIMA, TEWSBURY+ (ROCH, BNL)
ALSO 66 PRL 17 548 LOBKOWICZ, MELISSINOS, NAGASHIMA+ (ROCH+BNL)
MACER 69 PRL 22 32 MCKER, MANN, MCFARLANE, ROBERTS+ (PENN, TEMPLE)
MAESTRO 69 PR 183 1200 +GERSHWIN, ALSTON-GARJOST, BANGERTER+ (LRL)
ZELLER 69 PR 182 1420 ZELLER, HADDOCK, HELLAND, PAHL+ (UCLA, LRL)

ALSO 67 AUERBACH, FOOTNOTE 1.
AUERBACH 67 PR 155 1505 +DOBBS, MANN, MCFARLANE, WHITE+ (PENN, PRIN)
ERRATUM
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BELLOTTI 67 NC 52A 1287 BELLOTTI, FIORINI, PULLIA (MILAN)
ALSO 66 PL 20 690 BELLOTTI, FIORINI, PULLIA+ (MILAN)
BISI 67 PL 25B 572 BISI, CESTER, CHIESA, VIGONE (TORINO)
BOTTERILL 67 PRL 19 982 BOTTERILL, BROWN, CORBETT, CULLIGAN + (OXFORD)
ALSO 68 BOTTERILL
BOWEN 67 PR 154 1314 BOWEN, MANN, MCFARLANE, HUGHES+ (PENN-PRINCETON)

CLINE1 67 HEIDELBERG CONF CLINE, HAGGERTY, SINGLETON, FRY+ (WISCONSIN)
CLINE2 67 HERCEG NOVI TBL.4 D. CLINE, PROC. INT'L SCH. ON ELEM. PART. PHYSICS
FLETCHER 67 PRL 19 98 FLETCHER, BEIER, EDWARDS + (ILLINOIS)
FORD 67 PRL 18 1214 +LEMONICK, NAUENBERG, PIROUE (PRINCETON)
IMLAY 67 PR 160 1203 IMLAY, ESCHSTRUTH, FRANKLIN+ (PRINCETON)
KALMUS 67 PR 159 1187 KALMUS, KERNAN (LRL)
ZINCHENK 67 RUTGERS (THEISIS) ZINCHENKO (RUTGERS)

BETTELS 68 NC 56A 1106 AACHEN-BARI-IBERGEN-CERN-EP-NIJMEGEN-ORSAY+
ALSO 71 HAIDT
BOTTERILL 68 PR 171 1402 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL 68 PR 174 1661 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL 68 PRL 21 766 BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BUTLER 68 UCRL-18420 +BLAND, GOLDHABER, GOLDHABER, HIRATA+ (LRL)
CHANG 68 PRL 20 510 CHANG, YODH, EHRLICH, PLANO+ (MARYLAND, RUTGERS)

CHEN 68 PRL 20 73 CHEN, CUTTS, KIJEWSKI, STIENING + (LRL, MIT)
EICHEN 68 PL 27B 586 AACHEN-BARI-CERN-EP-ORSAY-PADOVA-VALENCIA
EISLER 68 PR 169 1090 EISLER, FUNG, MARATECK, MEYER, PLANO (RUTGERS)
ESCHSTRUTH 68 PR 165 1487 ESCHSTRUTH, FRANKLIN, HUGHES+ (PRINCETON, PENN)
GARLAND 68 PR 167 1225 +MSTOW, SOLODOVNIKOVA, FADEEV + (LJUNG)
MOSCOSO 68 THESIS M L MOSCOSO (UNIV PARIS ORSAY)
CUTTS 69 PR 184 1380 +STIENING, WIEGAND, DEUTSCH (LRL, MIT)
ALSO 68 PRL 20 955 CUTTS, STIENING, WIEGAND, DEUTSCH (LRL, MIT)
DAVISON 69 PR 180 1333 +BACASTOW, BARKAS, EVANS, FUNG, PORTER+ (UCR)
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LOBKOWICZ 69 PR 185 1676 +MELISSINOS, NAGASHIMA, TEWSBURY+ (ROCH, BNL)
ALSO 66 PRL 17 548 LOBKOWICZ, MELISSINOS, NAGASHIMA+ (ROCH+BNL)
MACER 69 PRL 22 32 MCKER, MANN, MCFARLANE, ROBERTS+ (PENN, TEMPLE)
MAESTRO 69 PR 183 1200 +GERSHWIN, ALSTON-GARJOST, BANGERTER+ (LRL)
ZELLER 69 PR 182 1420 ZELLER, HADDOCK, HELLAND, PAHL+ (UCLA, LRL)

BARON 71 PL 36B 619 +BREHIN, DIAMANT-BERGER, KUNZ+ (SACL+GEVA)
BOURQUIN 71 PL 36B 615 +BOYMOND, EXTERMANN, MARASO+ (GEVA, SACL)
HAIDT 71 PR D3 10 AACHEN+BARI+CERN+EP+NIJMEGEN+ORSAY+PADOVA+
ALSO 69 PL 29B 691 +(AACH, BARI, CERN, EPOL, NIJMEGEN, ORSAY, PADO, TORI)
KLEMS 71 PR D4 66 +HILDEBRAND, STIENING (CHIC, LRL)
ALSO 70 PRL 24 1086 KLEMS, HILDEBRAND, STIENING (LRL, CHIC)
ALSO 70 PRL 25 473 KLEMS, HILDEBRAND, STIENING (LRL, CHIC)

ALSO 67 AUERBACH, FOOTNOTE 1.
AUERBACH 67 PR 155 1505 +DOBBS, MANN, MCFARLANE, WHITE+ (PENN, PRIN)
ERRATUM
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BETTELS 68 NC 56A 1106 AACHEN-BARI-IBERGEN-CERN-EP-NIJMEGEN-ORSAY+
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EICHEN 68 PL 27B 586 AACHEN-BARI-CERN-EP-ORSAY-PADOVA-VALENCIA
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BARON 71 PL 36B 619 +BREHIN, DIAMANT-BERGER, KUNZ+ (SACL+GEVA)
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ALSO 69 PL 29B 691 +(AACH, BARI, CERN, EPOL, NIJMEGEN, ORSAY, PADO, TORI)
KLEMS 71 PR D4 66 +HILDEBRAND, STIENING (CHIC, LRL)
ALSO 70 PRL 24 1086 KLEMS, HILDEBRAND, STIENING (LRL, CHIC)
ALSO 70 PRL 25 473 KLEMS, HILDEBRAND, STIENING (LRL, CHIC)

Stable Particles

K±, K0, KS0

Data Card Listings

For notation, see key at front of Listings.

Table listing experimental data for K particles, including experiment names like CERN, BNL, and various researchers.

***** K0 *****

Table 11: NEUTRAL K MASS (MEV). Includes data for various experiments and a weighted average of 497.87 ± 0.32.

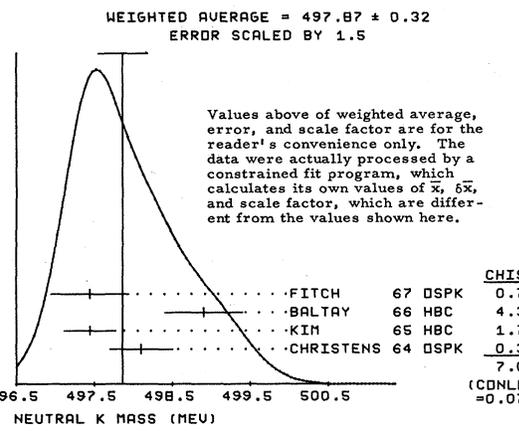


Table 11 (K0) - (K±) MASS DIFFERENCE (MEV). Lists mass differences for various experiments like CRAWFORD, ROSENFELD, etc.

REFERENCES FOR NEUTRAL K. Lists references for various experiments and authors.

***** KS0 *****

Table 12: SHORT-LIVED NEUTRAL K (498, JP=0-) I=1/2. Lists mean life data for various experiments.

Note on the KS0 Mean Life

From 1968 until 1972 our average value for the KS0 mean life was (0.862 ± 0.006) × 10^-10 second. Since then three high-precision experiments, SKJEGGESTAD 72 (Oslo, CERN, Saclay), GEWENIGER 74 (CERN, Heidelberg), and CARITHERS 75 (Columbia, NYU)

have obtained results compatible with each other which average (0.8930 ± 0.0023) × 10^-10 second. This is about five standard deviations above the previous average. The origin of this discrepancy is not known.

The corrections for systematic biases in SKJEGGESTAD 72 and HILL 68 (including the correction for the new value of η+-) amount to +1% and 0.7%, respectively. Similar corrections, if applied to the older bubble chamber results, would probably increase the old average by only about one standard deviation and would not account for the discrepancy.

The two experiments which contribute most to the χ^2 (see ideogram below) are KIRSCH 66 and DONALD 68. The combined result of all other experiments is (0.8914 ± 0.0021) × 10^-10 second with a χ^2 of 11.5 for 10 degrees of freedom. This value does not differ significantly from the average of the newer (post-1971) experiments.

Since the newer experiments are in principle superior - that is, they have higher statistics, better acceptance, and easier trigger conditions - we have chosen to average them separately from the older experiments as is seen in the Data Card Listings below. In the Stable Particle Table, we quote the new value, but give the old value in a warning footnote.

Table 12: K0S MEAN LIFE (UNITS 10^-10 SEC). Lists mean life data for various experiments like CRAWFORD, BOWEN, etc.

Table 12: K0S MEAN LIFE (POST-1971 EXPERIMENTS). Lists mean life data for experiments like SKJEGGESTAD, FACKLER, etc.

COMMENTS. Lists comments on the data, such as corrections and systematic errors.

Data Card Listings

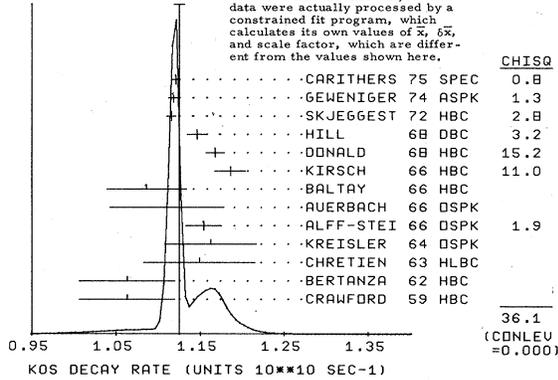
For notation, see key at front of Listings.

Stable Particles

K⁰

WEIGHTED AVERAGE = 1.1257 ± 0.0065
ERROR SCALED BY 2.5

Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of X, δX, and scale factor, which are different from the values shown here.



12 KOS PARTIAL DECAY MODES
Table with columns: Mode, Decay Masses. Rows include KOS INTO PI+ PI-, KOS INTO MU+ MU-, KOS INTO E+ E-, KOS INTO PI+ PI- GAMMA, KOS INTO GAMMA GAMMA, KOS INTO 3PI0.

12 KOS BRANCHING RATIOS
Table with columns: Mode, Ratio, Author. Rows include KOS INTO (PI+ PI-)/TOTAL, KOS INTO (PI0 PI0)/TOTAL, KOS INTO 3PI0, CP VIOLATING/(KOL INTO PI+ PI- PI0).

WEIGHTED AVERAGE = 0.316 ± 0.014
ERROR SCALED BY 1.3

Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of X, δX, and scale factor, which are different from the values shown here.

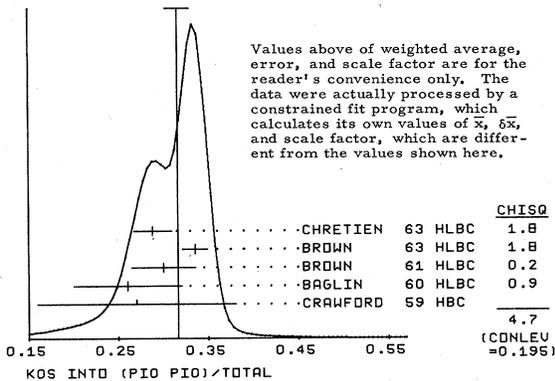


Table of experimental data for K0. Columns include mode (e.g., KOS INTO (PI+ PI-)/(PI0 PI0)), author, value, error, and scale factor. Rows include data from BOZOKI, GBBRI, MORFIN, MOFFETT, BALTAY, ALITTI, MORSE, NAGY, HILL, CONWELL, NAGY 72 IS A FINAL RESULT WHICH INCLUDES BOZOKI 69, MOFFETT 70 IS A FINAL RESULT WHICH INCLUDES GBBRI 69, etc.

REFERENCES FOR KOS
List of references for K0 decay modes. Includes authors like BOLDT, CRAWFORD, BAGLIN, BOWEN, COLUMBIA, BROWN, ANDERSON, CHRETIEN, ALFF-STEI, AUERBACH, BALTAY, BEHR, BELLOTTI, BOTT-BOD, KIRSCH, BOTT-BOD, DONALD, HILL, BANNER, BOZOKI, DOYLE, etc.

Stable Particles

K_S⁰, K_L⁰

Data Card Listings

For notation, see key at front of Listings.

GOBBI 69 PRL 22 682
HYAMS 69 PL 298 521
MORFIN 69 PRL 23 660
STUTZKE 69 PR 177 2009
HOFFETT 70 BAPS 15 512
WEBBER 70 PR D1 1967
ALSO 69 UCRL 19226 THESIS
BALTAY 71 PRL 27 1678
ALSO 71 NEVIS-187 THESIS
CHO 71 PR D3 1557
JAMES 71 PL 35B 265
MEISNER 71 PR D3 59
REPELLIN 71 PL 36B 603
ALITTI 72 PL 39B 568
BANNER 72 PR 29 237
JAMES 72 NP 849 1
JONES 72 NC 9A 151
METCALF 72 PL 40B 703
HORSE 72 PRL 28 388
NAGY 72 NP 847 94
ALSO 69 PL 30B 498
SKJEGGES 72 NP 848 343
BARMINI 73 PL 46B 465
BARMINI 73 PL 47B 463
BURGUN 73 PL 46B 481
FACKLER 73 PRL 31 847
GJESDAL 73 PL 48B 217
HILL 73 PR D8 1290
MALLARY 73 PR D7 1953
BOBISUT 74 LNC 11 646
COMWELL 74 PR D8A 2083
GENEIGE 74 PL 48B 487
GJESDAL 74 PL 52B 119
BALDOCEO 75 NC 25A 688
CARITHERS 75 PRL 34 1244

BIRGE 60 ROCH CONF 601
MULLER 60 PRL 4 418
FITCH 61 NC 22 1160
GOOD 61 PR 124 1223
CRAWFORD 62 CERN CONF 827
AUERBACH 65 PRL 14 192
TRILLING 65 PRL 14 192
UPDATED FROM 1965 ARGONNE CONF., PAGE 115.



13 LONG-LIVED NEUTRAL K (498, JP=0) I=1/2

WE GIVE (KOL-KOS MASS DIFFERENCE / HBAR) IN UNITS OF 10**10 SEC-1
D TX (2.20) (0.35) FITCH 61 CNTR
D X (0.84) (0.29) (0.22)GOOD 61 HLBC
D TXC (1.02) (0.23) CAMERINI 62 HLBC 8/67
D TX (0.55) (0.24) AUBERT 65 HLBC 6/66
D X (0.26) (0.36) (0.26)BALDO-CEO 65 HLBC ASSUMES CP CONS.
D TXA (0.64) (0.12) CHRISTENS 65 OSPK 6/66
D TX (0.70) OR LESS FITCH 65 OSPK CF. MEISNER 66 7/66
D V 130 (0.89) (0.15) VISHNEVSKY 65 OSPK CU AND AL REGEN 8/67
D X (0.54) (0.099) ALF-STEI 65 OSPK 8/67
D X 84 (0.42) (0.24) (0.36) BALDO-CEO 66 HLBC K0+N INTO HYPER. 8/67
D B (0.531) (0.027) BOTT-BODE 66 OSPK C REGEN 9/66
D TX 77 (0.58) (0.17) CAMERINI 66 HBC, DBC K0+N INTO HYPER 8/67
D N 72 (0.44) (0.18) CANTER 66 DBC KO SCATTER IN D2 11/66
D X 95 (0.62) (0.10) (0.16) CHANG 66 HBC K0+P INTO HYPER. 9/66
D X (0.81) (0.17) FUJII 66 OSPK IRON REGENERATOR 9/66
D X 59 (0.74) (0.34) MEISNER1 66 HBC 9/66
D X + SIGMA FAVORED MEISNER2 66 HBC 9/66
D X (0.38) (0.16) JYANDOVIC 66 OSPK CHURANUM REGEN. 11/66
D TX 136 (+0.64) (0.19) CANTER 47 DBC K0+D INTO HYPER. 11/67
D X (0.65) (0.11) MISCHKE 67 OSPK 11/67
D X 590 (0.59) (0.13) BALATZ 68 OSPK AL REGENERATOR 3/68
D X (0.520) (0.044) CARNEGIE 68 HBC GAP METHOD 3/68
D TX (+0.487) (0.046) MELHOP 68 OSPK ST-STEEL REGEN 6/68
D BX (0.547) (0.024) BOTT-BODE 69 OSPK C REGEN 1/71
D FX (0.555) (0.020) FAISSNER 69 ASPK REGEN IN CU 10/69
D D 0.542 0.006 CULLEN 70 CNTR 1/71
D R (0.542) (0.006) ARNOLD 70 ASPK GAP METHOD 1/71
D X (0.481) (0.052) (0.075)BALATS 71 OSPK 9/71
D R (0.534) (0.007) CARNEGIE 71 ASPK GAP METHOD 8/71
D TH 119 (+0.67) (0.14) HILL 71 DBC 10/71
D S 1757 (0.557) (0.038) FACKLER 73 OSPK 11/73
D 0.5340 0.0030 GEMENIG 74 SPEC GAP METHOD 11/75*
D 0.5334 0.0040 GJESDAL 74 SPEC CHG ASYMMETRY 11/75*

D AVG 0.5349 0.0022 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
D STUDENT 0.5348 0.0025 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT COMMENTS
D X NO ATTEMPT HAS BEEN MADE TO CORRECT OLDER EXPERIMENTS WITH LARGE D X ERRORS FOR THE SUBSEQUENT CHANGES IN THE KOS MEAN LIFE OR IN ETAA--
D T A KOS MEAN LIFE OF 0.862 10**10 SEC WAS USED IN CONVERTING THE D T MASS DIFFERENCE FROM UNITS OF INVERSE KOS MEAN LIVES TO ABSOLUTE D T UNITS, VALUES NOT BEARING THIS FOOTNOTE EITHER WERE GIVEN IN D T ABSOLUTE UNITS OR WERE CONVERTED USING THE AUTHORS' VALUE OF THE D T KOS MEAN LIFE.
D C CAMERINI 62 VALUE CHANGED FROM 1.7 (SEE TABLE 1 OF CAMERINI 66)
D A CHRISTENSON 65 CORRECTED FOR INTERFERENCE BY FITCH 65 FOOTNOTE.
D V VISHNEVSKY 65 NOT CORRECTED FOR INTERFERENCE EFFECTS.
D N CANTER 66 ERROR IGNORES UNCERTAINTY OF PHASE SHIFTS. THESE EVENTS D N ARE USED IN HILL 71.
D B BOTT-BODENHAUSEN 69 IS A REEVALUATION OF BOTT-BODENHAUSEN 66.
D F FAISSNER 69 HAS ADDNL. SYSTEMATIC ERROR LESS THAN TWO PERCENT.
D R ARNOLD 70 AND CARNEGIE 71 USE KOS MEAN LIFE=0.862+-0.06 E-10 SEC.
D R WE HAVE NOT ATTEMPTED TO ADJUST THESE VALUES FOR THE SUBSEQUENT D R CHANGE IN THE KOS MEAN LIFE OR IN ETAA--
D H HILL 71 PRIMARY RESULT IS THAT DM IS POSITIVE.
D H THE MAGNITUDE MAY HAVE AN ADDITIONAL SYSTEMATIC ERROR OF ABOUT 0.12
D S NOT AVERAGED BECAUSE ERROR IS LARGE AND SYSTEMATICS NOT DISCUSSED.

13 KOL MEAN LIFE (UNITS 10**8 SEC)
T 34 8.1 3.2 2.4 BARDON 58 CNTR
T ASSUMED DS=DO AND DELTA I=1/2 CRAWFORD 59 HBC
T 15 5.1 2.4 1.3 DARMON 62 FBC
T 5.3 0.6 FUJII 66 OSPK
T 1700 6.1 1.5 1.2 ASTBURY3 65 CNTR
T 5.15 0.14 DEVLIN 67 CNTR
T L (5.0) (0.5) LOWMY 67 HLBC
T .4M 5.154 0.044 VOSBURGH 72 CNTR 2/71
T L SUM OF PARTIAL DECAY RATES.
T AVG 5.158 0.042 0.042 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.0)
T STUDENT 5.158 0.046 0.045 AVG. USING STUDENT10(H/1.11) -- SEE TEXT
T FIT 5.181 0.040 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

13 KOL PARTIAL DECAY MODES
P1 KOL INTO 3PI0 TAU 0 PRIME 134+ 134+ 134
P2 KOL INTO PI+ PI- P10 TAU 0 139+ 139+ 134
P3 KOL INTO PI MU NEUTRINO KL MU3 139+ 105+ 0
P4 KOL INTO PI E NEUTRINO KL E3 139+ 5+ 0
P5 KOL INTO PI+ PI- KL PI+ PI- 139+ 139
P6 KOL INTO MU+ MU- KL 2MU 105+ 105
P7 KOL INTO E+ E- KL 2E 5+ 5
P8 KOL INTO E MU KL EMU 5+ 105
P9 KOL INTO TWO GAMMAS KL 2GAMMA 0+ 0
P10 KOL INTO PI+ PI- GAMMA KL PI+ G 139+ 139+ 0
P11 KOL INTO P10 P10 KL P10 134+ 134 0
P12 KOL INTO PI E NEU GAMMA KL E3GAM 139+ 5+ 0+ 0
P13 KOL INTO P10 TWO GAMMAS KL P10GAMMA 134+ 0+ 0
P14 KOL INTO E+ E- GAMMA KL E+ G 5+ 5+ 0
P15 KOL INTO MU+ MU- GAMMA KL 2MUGAM 105+ 105+ 0
P16 KOL INTO MU+ MU- P10 KL 2MUP10 105+ 105+ 134
P17 KOL INTO PI+ PI- E+ E- KL 2PI2E 139+ 139+ 5+ 5
P18 KOL INTO P10 PI+ E+ NEU KL 2PIENEU 134+ 139+ 5+ 0

NEUTRAL K CONSTRAINED FIT
OVERALL FIT OF MEAN LIFE, WIDTHS AND BRANCHING RATIOS USES 62 DATA POINTS TO DETERMINE SIX QUANTITIES. OVERALL FIT HAS CHI-SQUARED=67.7 2/76*

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i + delta P_i, where delta P_i = sqrt(delta P_i delta P_i), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.
P 1 P 2 P 3 P 4 P 5 P 11
P 1 .2140+-0.0067
P 2 -.4051 .1225+-0.0019
P 3 -.5365 .0677 .2710+-0.0047
P 4 -.6408 -.0893 -.2416 .3894+-0.0053
P 5 -.2195 .5350 .0342 .0453 .0020+-0.0001
P 11 .1476 -.0732 -.0938 -.1119 -.0396 .0009+-0.0002

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., G_i = Gamma_i = Gamma_i total P_i, in appropriate units. In analogy to the matrix above, the diagonal elements are G_i + delta G_i, where delta G_i = sqrt(delta G_i delta G_i), while the off-diagonal elements are the normalized correlation coefficients (delta G_i delta G_j) / (delta G_i delta G_j). Note that, because of the error in Gamma_i total, the errors and correlations here are not directly derivable from those above.
G 1 G 2 G 3 G 4 G 5 G 11
G 1 .0413+-0.0014
G 2 -.2195 .0236+-0.0004
G 3 -.3386 .1977 .0523+-0.0010
G 4 -.3854 .2367 -.0345 .0752+-0.0011
G 5 -.1255 .5630 .1098 .1315 .0004+-0.0000
G 11 .1569 -.0445 -.0644 -.0733 -.0254 .0002+-0.0000

13 KOL DECAY RATES

W1 KOL INTO P10 P10 P10 (UNITS 10**6 SEC-1) (G1)
W1 54 5.22 1.03 0.84 BEHR 66 HLBC ASSUMES CP 8/66
W1 FIT 4.13 0.14 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
W2 KOL INTO PI+ PI- P0 (UNITS 10**6 SEC-1) (G2)
W2 18 3.26 0.77 ANDERSON 65 HBC 8/66
W2 14 1.4 0.4 FRANZINI 65 HBC 8/66
W2 136 2.62 0.28 0.27 BEHR 66 HLBC ASSUMES CP 8/66
W2 53 2.20 0.35 WEBBER 70 HBC ASSUMES CP 10/71
W2 99 2.71 0.28 CHO 71 DBC ASSUMES CP 4/71
W2 J 98 (2.5) (0.3) JAMES 71 HBC ASSUMES CP 6/71
W2 50 2.12 0.33 MEISNER 71 HBC ASSUMES CP 10/71
W2 J 180 2.35 0.20 JAMES 72 HBC ASSUMES CP 1/73
W2 192 2.32 0.13 0.15 BALDOCEO 75 HLBC ASSUMES CP 1/76*
W2 IN THE OVERALL FIT THIS RATE IS WELL DETERMINED BY THE MEAN LIFE AN W2 THE BRANCHING RATIO R2. FOR THIS REASON THE DISCREPANCY BETWEEN THE W2 MEASUREMENTS DOES NOT AFFECT THE SCALE FACTOR OF THE OVERALL FIT
W2 JAMES 72 IS A FINAL MEASUREMENT AND INCLUDES JAMES 71. 11/73
W2 10/71
W2 AVG 2.34 0.11 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)
W2 STUDENT 2.35 0.10 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
W2 FIT 2.364 0.039 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1) (SEE IDEOGRAM BELOW)

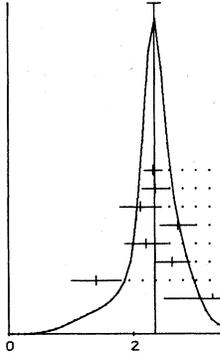
Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K⁰

WEIGHTED AVERAGE = 2.34 ± 0.11
ERROR SCALED BY 1.2



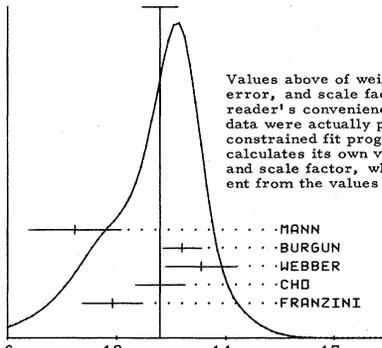
Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

Table listing particle names and their corresponding values: BALDOCELDL 75 HLBC 0.0, JAMES 72 HBC 0.0, MEISNER 71 HBC 0.4, CHD 71 DBC 1.7, WEBBER 70 HBC 0.2, BEHR 66 HLBC 1.0, FRANZINI 65 HBC 5.5, ANDERSON 65 HBC.

CHI SQ 0.0, B.9 (CONLEU = 0.177)

Table with 5 columns: W3, W4, W5, W6, W7. Rows include KOL INTO PI E NEUTRINO (UNITS 10**6 SEC-1) (G4), KOL INTO CHARGED (3-BODY) (UNITS 10**6 SEC-1) (G2+G3+G4), KOL INTO LEPTONIC (KMU3+KE3) (UNITS 10**6 SEC-1) (G3+G4), and KOL INTO PI MU NEUTRINO UNITS 10**6 SEC-1 (G3).

WEIGHTED AVERAGE = 11.60 ± 0.65
ERROR SCALED BY 1.5



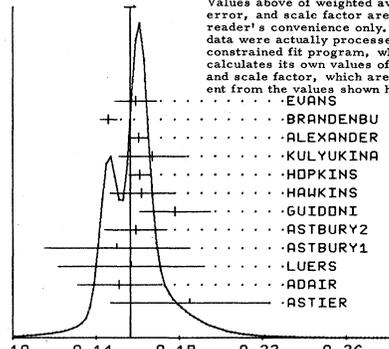
Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

Table listing particle names and their corresponding values: MANN 72 HBC 3.4, BURGUN 72 HBC 1.3, WEBBER 71 HBC 1.3, CHD 70 DBC 0.0, FRANZINI 65 HBC 2.6, B.6 (CONLEU = 0.072)

Table with 5 columns: W6, W7, W8, W9, W10. Rows include KOL INTO PI MU NEUTRINO UNITS 10**6 SEC-1 (G3), KOL BRANCHING RATIOS (P1)/(P2+P3+P4), and KOL INTO (PI+ PI- PI0)/CHARGED (P1)/(P2+P3+P4).

Table with 5 columns: R2, R3, R4, R5, R6. Rows include KOL INTO (PI+ PI- PI0)/CHARGED (P2)/(P2+P3+P4) for various particles like ASTIER, ADAIR, LUERS, ASTBURY1, ASTBURY2, GUIDONI, HOPKINS, HAWKINS, HOPKINS, KULYUKINA, ALEXANDER, BRANDENBU, EVANS, and AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2).

WEIGHTED AVERAGE = 0.1564 ± 0.0026
ERROR SCALED BY 1.2



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

Table listing particle names and their corresponding values: EVANS 73 HLBC 0.1, BRANDENBU 73 HBC 6.7, ALEXANDER 73 HBC 1.2, KULYUKINA 68 CC 0.4, HOPKINS 67 HBC 0.8, HAWKINS 66 HBC 0.1, GUIDONI 65 HBC 1.6, ASTBURY2 65 CC 0.0, ASTBURY1 65 CC, LUERS 64 HBC, ADAIR 64 HBC 0.1, ASTIER 61 CC.

CHI SQ 0.1, 11.1 (CONLEU = 0.194)

Table with 5 columns: R3, R4, R5, R6, R7, R8, R9. Rows include KOL INTO (PI MU NEUTRINO)/CHARGED (P3)/(P2+P3+P4), KOL INTO (PI E NEUTRINO)/CHARGED (P4)/(P2+P3+P4), KOL INTO (PI E NEU)/(PI E NEU+PI MU NEU) (P4)/(P3+P4), KOL INTO (PI+ PI- PI0)/TOTAL (P2), KOL INTO (LEPTON PI NEUTRINO)/TOTAL (P3+P4), KOL INTO (2 GAMMA)/TOTAL (UN. 10**4) (P9), THIS VALUE USES (E00/E+)**2=1.05**0.14, IN GENERAL, S13R = (4.32+-0.55)*(10**4)**(E00/E+)**2, OLD EXPERIMENTS EXCLUDED FROM FIT, AVERAGE ETA+ OF THESE EXPERIMENTS AND FOR NOTE ON DISCREPANCY, FROM SAME DATA AS R27 MESSNER 73, BUT WITH DIFFERENT NORMALIZATION, and AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT.

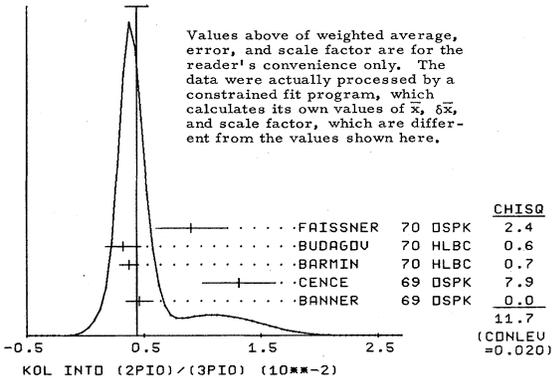
Stable Particles

K_L^0

Data Card Listings

For notation, see key at front of Listings.

R10	KOL INTO (PI MU NEU)/(PI E NEU)	(P3)/(P4)			
R10	0.81	0.19	ADAIR	64 HBC	
R10	0.82	0.10	DEBOUARD	67 OSPK	
R10	273	0.7	HANKINS	67 HBC	
R10	0.81	0.08	HOPKINS	67 HBC	
R10	770	0.71	BUDAGOV	68 HLBC	
R10 K	(0.67)	(0.13)	KULYUKINA	68 CC	
R10 B	569	(0.71)	(0.04)	BEILLIERE 69 HLBC	
R10	1309	(0.648)	(0.030)	EVANS 69 HLBC REPL. BY EVANS 73	
R10	3548	0.68	0.08	BASILE 70 OSPK	
R10	5700	0.741	0.024	BRANDENBU 73 HBC	
R10	1309	0.662	0.030	EVANS 73 HLBC	
R10	10K	0.662	0.037	WILLIAMS 74 ASPK	
R10 K	KULYUKINA 68 R10 IS NOT MEASURED INDEPENDENTLY FROM R2 AND R4.				
R10 B	BEILLIERE 69 IS A SCANNING EXPT USING SAME EXPOSURE AS BUDAGOV 68				
R10	AVG	0.695	0.019	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)	
R10	STUDENT	0.695	0.021	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R10	FIT	0.696	0.017	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R11	KOL INTO (MU+MU-)/CHARGED (UNITS 10**=-6)	(P6)/(P2+P3+P4)			
R11	100.0	OR LESS	ANIKINA	65 CC	
R11	250.0	OR LESS	CL=.90	ALFF-STEI 66 OSPK	
R11	5	OR LESS	CL=.90	BOTT-BODE 67 OSPK	
R11	35.0	OR LESS	CL=.90	FITCH 67 OSPK	
R12	KOL INTO (PI+ PI- GAMMA)/TOTAL (UNITS 10**=-3)	(P10)			
R12	15.0	OR LESS	ANIKINA	65 CC	
R12	0	OR LESS	BELLOTTI	66 HLBC	
R12	1	3.0	OR LESS	NEFKENS 66 OSPK	
R12	0.4	OR LESS	CL=.90	THATCHER 68 OSPK	
R12	3.2	OR LESS	CL=.90	BOBISUT 74 HLBC	
R12 D	24	0.062	0.021	DONALD S1 74 SPEC	
R12	0.46	OR LESS	CL=.90	WOD 74 SPEC	
R12 D	USES KOL TO PI+PI-PI0/ALL KOL DECAYS = 0.126				
R13	KOL INTO (E+ E-)/CHARGED (UNITS 10**=-6)	(P7)/(P2+P3+P4)			
R13	1000.0	OR LESS	ANIKINA	65 CC	
R13	200.0	OR LESS	CL=.90	ALFF-STEI 66 OSPK	
R13	23.0	OR LESS	CL=.90	BOTT-BODE 67 OSPK	
R14	KOL INTO (E MU)/CHARGED (UNITS 10**=-4)	(P8)/(P2+P3+P4)			
R14	10.0	OR LESS	ANIKINA	65 CC	
R14	1.0	OR LESS	CL=.90	CARPENTER 66 OSPK	
R14	0.1	OR LESS	CL=.90	BOTT-BODE 67 OSPK	
R14	0.08	OR LESS	CL=.90	FITCH 67 OSPK	
R15	KOL INTO (E+ PI- NEU)/(E- PI+ NEU)				
R15 O	97	(0.90)	(0.18)	NEAGU 61 CC	
R15 O	(1.01)	(0.16)		LUERS 64 HBC	
R15 O	894	(0.92)	(0.023)	KULYUKINA 66 CC	
R15 O	1539	(1.06)	(0.05)	VERHEY 66 OSPK	
R15 O	LOW PRECISION EXPTS NOT AVERAGED. FOR MORE PRECISE VALUE, SEE S13A2 (BENNETT 70, MARX 70)				
R16	KOL INTO (MU+ PI- NEU)/(MU- PI+ NEU)				
R16	1M	1.0081	0.0027	DORFAN 67 OSPK	
R16	SEE ALSO S13A2 AND S13AL IN THE CP VIOLATION SECTION				
R17	KOL INTO (PI0 PI0)/TOTAL (UNITS 10**=-3)	(P11)			
R17 C	7	(1.2)	(1.5)	(1.2) CRIEGEE 66 OSPK	
R17 C	CRIEGEE EXPT NOT DESIGNED TO MEASURE 2 PI0 DECAY MODE				
R17 G	189	(2.5)	(0.8)	E00=3.6+-0.6	
R17 G	LATEST RESULT OF THIS EXPERIMENT GIVEN BY FAISSNER TO R19				
R17	AVG	0.94	0.19	FROM FIT	
R18	KOL INTO (3PI0)/(PI+PI-PI0)	(P11)/(P2)			
R18	188	2.0	0.6	ALEKSANYA 64 FBC	
R18	1010	1.80	0.13	BUDAGOV 68 HLBC	
R18	883	(1.65)	(0.07)	BARMIN 72 HLBC	
R18	ERROR STAT. ONLY				
R18	AVG	1.81	0.13	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R18	STUDENT	1.81	0.14	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R18	FIT	1.747	0.070	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	
R19	KOL INTO (2PI0)/(3PI0)	(P11)/(P1)			
R19 C	109	(1.89)	(0.31)	CRONIN 1 67 OSPK	
R19 C	(1.36)	(0.18)		CRONIN 2 67 OSPK	
R19 C	CRONIN 2 IS FURTHER ANALYSIS OF CRONIN 1, NOW BOTH WITHDRAWN				
R19	NO EVENTS SEEN				
R19	57	0.46	0.11	BARTLETT 68 OSPK	
R19	133	1.31	0.31	BANNER 69 OSPK	
R19	29	0.37	0.08	CENCE 69 OSPK	
R19	30	0.32	0.15	BARMIN 70 HLBC	
R19	F	172	0.90	0.30	BUDAGOV 70 HLBC
R19 F	FAISSNER TO CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69 R17				
R19	AVG	0.439	0.098	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)	
R19	STUDENT	0.425	0.065	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R19	FIT	0.439	0.088	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.5)	
				(SEE IDEOGRAM BELOW)	
				WEIGHTED AVERAGE = 0.439 ± 0.098	
				ERROR SCALED BY 1.7	



R20	KOL INTO (PI+ PI-)/(KE3 + KMU3) (UNITS 10**=-3)	(P5)/(P3+P4)		
R20 O	309	(2.51)	(0.23)	DEBOUARD 67 OSPK
R20 O	525	(2.35)	(0.19)	FITCH 67 OSPK
R20 O	OLD EXPERIMENTS EXCLUDED FROM FIT. SEE SUBSECTION E+ BELOW FOR			
R20 O	AVERAGE ETA+- OF THESE EXPERIMENTS AND FOR NOTE ON DISCREPANCY.			
R20	3.041	0.090	FROM FIT	
R21	KOL INTO (2GAMMA)/(3 PI0) (UNITS 10**=-3)	(P9)/(P1)		
R21	16	2.5	0.7	ARNOLD 68 HLBC
R21	BANNER 69 IS NEW EXPT. NOT TO BE CONF WITH R8 OF CRONIN 67			
R21	115	2.24	0.28	BANNER 69 OSPK
R21	28	2.13	0.43	BARMIN 71 HLBC
R21	AVG	2.24	0.22	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R21	STUDENT	2.24	0.24	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

Note on the $K_L^0 \rightarrow \mu^+ \mu^-$ Controversy

The $K_L^0 \rightarrow \mu^+ \mu^-$ branching ratio upper limit of CLARK 71 does not agree with the results of subsequent experiments.

A total of twelve examples of this decay have been observed in three experiments: CARITHERS1 73 (Columbia-CERN-NYU), CARITHERS2 73 (Columbia-BNL-CERN), and FUKUSHIMA 76 (Princeton-Univ. of Mass.-Amherst). Background from all sources (primarily from $K_L^0 \rightarrow \pi \mu \nu$ with the pion either penetrating the muon counter or decaying to $\mu \nu$ in flight) has been estimated to be negligible for these experiments. The rates from these three experiments and the upper limits of FOETH 69 and DARRIULAT 70 are all quite consistent (see subsection R22 below). The combined result of these five experiments is

$$R = \frac{\Gamma(K_L^0 \rightarrow \mu^+ \mu^-)}{\Gamma(K_L^0 \rightarrow \pi^+ \pi^-)} = (5.0^{+1.6}_{-1.4}) \times 10^{-6}$$

where we have divided the observed number ($12^{+3.8}_{-3.3}$) by the sum of the effective denominators (which include the $\pi \pi / \mu \mu$ relative acceptances) given in the data card comments field.

CLARK 71 found no events and obtained the limit $R < 1.53 \times 10^{-6}$ at 90% confidence level. This result is from the FIELD 74 reanalysis of the experiment. The reanalysis uncovered no major flaws, but did find a few problems which increased the upper limit from 1.2×10^{-6} to the above. The combined result of the other five experiments leads us to expect 7.5 events for the sensitivity of the CLARK 71 experiment (9.5 events before reanalysis).

The theoretical lower limit for $\Gamma(K_L^0 \rightarrow \mu^+ \mu^-) / \Gamma(K_L^0 \rightarrow \pi^+ \pi^-)$ based on unitarity considerations is 3×10^{-6} , in agreement with the combined result, but above the CLARK 71 upper limit.

Data Card Listings
For notation, see key at front of Listings.

Stable Particles
K10

In the Stable Particle Table we quote the above combined value, but warn in a footnote that the value does not include CLARK 71.

Reference

- 1. C. Quigg and J. D. Jackson, UCRL-18487 (1968).

Table of data cards (R22, R23, R24, R25, R26, R27, R28, R29, R30, R31, R32, R33) containing particle properties and experimental details.

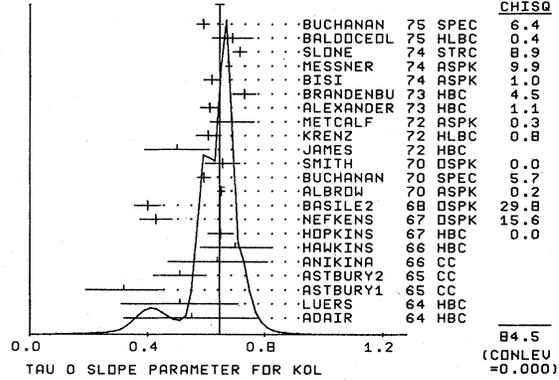
13 KOL ENERGY DEPENDENCE OF DALITZ PLOT

RELATED TEXT SECTION VI B.1, APPENDIX 1, AND MINI-REVIEW ON SLOPE PARAMETERS IN THE CHARGED K SECTION OF THE DATA CARD LISTINGS ABOVE

MATRIX ELEMENT SQUARED = 1 + G (S3-S0)/(MPI**2)

Table of linear energy dependence (G) for tau decays, listing parameters like KLONG, INTO, PI+, PI-, P10 and various fit results.

WEIGHTED AVERAGE = 0.646 ± 0.014
ERROR SCALED BY 2.5



13 KOL FORM FACTORS

RELATED TEXT SECTION VI B.2 AND MINI-REVIEW ON FORM FACTORS IN THE CHARGED K SECTION OF THE DATA CARD LISTINGS ABOVE.

IN THE FORM FACTOR COMMENTS, THE FOLLOWING ABBREVIATIONS ARE USED.
F+ AND F- ARE FORM FACTORS FOR THE VECTOR MATRIX ELEMENT.
FS AND FT REFER TO THE SCALAR AND TENSOR TERM.
FO = (F+) + (F-)*T/(MK**2-MPI**2)
L+, L- AND LO ARE THE LINEAR EXPANSION COEFFS. OF F+, F- AND FO.
L+ REFERS TO THE KMU3 VALUE EXCEPT IN THE KE3 SECTIONS.
DXI/DL IS THE CORRELATION BETWEEN XI(0) AND L+ IN KMUS.
DLO/DL+ IS THE CORRELATION BETWEEN LO AND L+ IN KMUS.
T = MOMENTUM TRANSFER TO THE PI IN UNITS OF MPI**2.
DP = DALITZ PLOT ANALYSIS
PI = PI SPECTRUM ANALYSIS
MU = MU SPECTRUM ANALYSIS
POL = MU POLARIZATION ANALYSIS
BR = KMU3/KE3 BRANCHING RATIO ANALYSIS
E = POSITRON OR ELECTRON SPECTRUM ANALYSIS
RC = RADIATIVE CORRECTIONS

Table of KOL form factors (XIA, XIB, XIC, etc.) listing parameters like F+, F-, F0, L+, L-, LO, DXI/DL, DLO/DL+, and various fit results.

Stable Particles

K_L⁰

Data Card Listings

For notation, see key at front of Listings.

XIC XIC = F/F+ (DETERMINED FROM MU POLARIZATION IN KMU3)
XIC THE MU POLARIZATION IS A MEASURE OF XI(T). NO ASSUMPTIONS ON L+
XIC NECESSARY, T (WEIGHTED BY SENSITIVITY TO XI) SHOULD BE SPECIFIED.

IXI IMAGINARY PART OF XI (TEST OF T REVERSAL)
IXI -0.02 0.08 ABRAMS 68 OSPK POLARIZATION 10/69
IXI 2.2M -0.060 0.045 SANDWEISS 73 CNTR POL, T=3.3 1/74

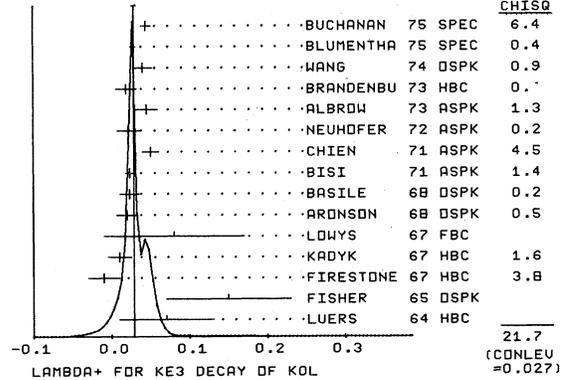
L+M LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KMU3 DECAY)
L+M SEE ALSO THE CORRESPONDING ENTRIES AND NOTES IN SECTION XIA AND LO.
L+M FOR RAD. CORR. OF KMU3 DP SEE GINSBURG 70 AND BECHERRAWY 70.

LO LAMBDA 0 (LINEAR ENERGY DEPENDENCE OF F0 IN KMU3 DECAY)
LO WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(I) INTO
LO VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

L+E LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KO E3 DECAY)
L+E FOR RAD. CORR. OF KE3 DP SEE GINSBURG 67 AND BECHERRAWY 70.
L+E 153 +0.07 .08 LUERS 64 HBC DP, NO RC 8/67

FS FS/F+ RATIO OF SCALAR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE)
FS 5600 0.19 OR LESS CL=.68 KULYUKINA 67 CC 10/69
FS 25K 0.04 OR LESS CL=.68 BLUMENHA 75 SPEC 7/75*

WEIGHTED AVERAGE = 0.0288 ± 0.0028
ERROR SCALED BY 1.4



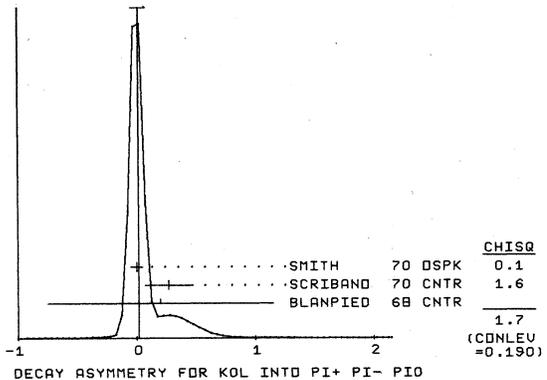
13 CP VIOLATION PARAMETERS IN KOL DECAYS
RELATED TEXT SECTION VI B.3 AND MINI-REVIEW BELOW

13 CHARGE ASYMMETRY IN TAU DECAYS
TEXT SECTION VI B.3 B

SEE SCRIBANO 70 FOR DEFINITION (HIS SIGMA=-1, A=1 FOR MAX ASYMMETRY
(M)*2 = 1+ SIG- (2/SQRT(3)) * ((T+)-(T-))/TMAX) AS SCRIBANO 70

Table with 4 columns: A, DE, parameter, value. Rows include DECA, BLANPIED, SCRIBANO, SMITH, AVG, and STUDENT.

WEIGHTED AVERAGE = 0.016 ± 0.063
ERROR SCALED BY 1.3



13 CHARGE ASYMMETRY IN LEPTONIC DECAYS (PERCENT)
TEXT SECTION VI B.3 C

SUCH ASYMMETRY VIOLATES CP. IT IS RELATED TO REAL(EPSILON).

Table with 4 columns: A, D, parameter, value. Rows include KOL INTO, DORFAN, PACIOTTI, PICCIONI, MCCARTHY, GEMENIGI, PACIOTTI, AVG, STUDENT, KOL INTO, BENNETT, SAAL, MARX, ASHFORD, FITCH, GEMENIGI, AVG, and STUDENT.

Data Card Listings

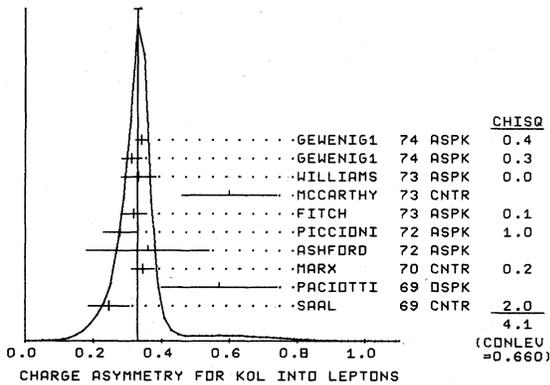
For notation, see key at front of Listings.

Stable Particles

K_L^0

AL	KOL INTO ((L+)-(L-))/((L+)+(L-)) (COMBINED A1 AND A2) (PERCENT)			
AL B	10M	0.246	0.059	SAAL 69 CNTR KE3 2/71
AL D	1M	0.57	0.17	PACIOTTI 69 DSPK KMU3 1/73
AL	10M	0.346	0.033	MARX 70 CNTR KE3 2/71
AL	600K	0.36	0.18	ASHFORD 72 ASPK KE3 2/72
AL	7.7M	0.278	0.051	PICCIONI 72 ASPK KMU3 1/73
AL	40M	0.318	0.038	FITCH 73 ASPK KE3 12/73
AL	4.1M	0.60	0.14	MCCARTHY 73 CNTR KMU3 6/73
AL	33M	0.333	0.050	WILLIAMS 73 ASPK KMU3*KE3 12/73
AL	15M	0.313	0.029	GEWENIGER 74 ASPK KMU3 7/74*
AL	34M	0.341	0.018	GEWENIGER 74 ASPK KE3 7/74*
AL	SEE FOOTNOTES IN SECTIONS A1 AND A2 ABOVE.			
AL	AVG	0.330	0.012	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AL	STUDENT	0.330	0.013	AVERAGE USING STUDENT(10/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.330 ± 0.012
ERRDR SCALED BY 1.0



-----13 PARAMETERS FOR KOL INTO 2PI DECAY-----
TEXT SECTION VI B.3 D
ETA+- = A(KL TO PI+PI-)/A(KS TO PI+PI-)
ETA0 = A(KL TO PIP0)/A(KS TO PIP0)

Note on $|\eta_{+-}|$

There is a very large discrepancy between old and new results for $|\eta_{+-}|$, the ratio of amplitudes $A(K_L^0 \rightarrow \pi^+\pi^-)/A(K_S^0 \rightarrow \pi^+\pi^-)$. The average of the older $|\eta_{+-}|$ results of CHRISTENSON 64, GALBRAITH 65, BASILE 66, BOTT-BODENHAUSEN 66, DE BOUARD 67, and FITCH 67 (see subsections R9, R20, and E+- of the K_L^0 Data Card Listings) is

$$|\eta_{+-}| = (1.95 \pm 0.03) \times 10^{-3}$$

with very good consistency.

In 1973 a CERN-Heidelberg K_L^0, K_S^0 interference experiment and a Colorado-SLAC-Santa Cruz branching ratio measurement gave the following results:

$$|\eta_{+-}| = (2.30 \pm 0.035) \times 10^{-3} \quad \text{GEWENIGER 74,}$$

$$|\eta_{+-}| = (2.23 \pm 0.05) \times 10^{-3} \quad \text{MESSNER 73,}$$

which are around 11 standard deviations above the previous average but are in good agreement with each

other. GEWENIGER 74 also made a check measurement in order to confirm their interference result by measuring the branching ratio $K_L^0 \rightarrow \pi^+\pi^-/K_L^0 \rightarrow \pi e \nu$. They obtained the result $|\eta_{+-}| = (2.30 \pm 0.06) \times 10^{-3}$, which, although systematically less reliable than the interference result, is in good agreement with it. A Columbia experiment¹ and two Chicago experiments (Cronin group² and Telegdi group³) have also yielded unpublished measurements which corroborate the higher values of $|\eta_{+-}|$.

The origin of the discrepancy between old and new results is not known. The effect of the change in the K_S^0 mean life value (see note in K_S^0 section of the Stable Particle Data Card Listings) is insufficient to explain it, raising $|\eta_{+-}|$ by only about 2%. The MESSNER 73 result above was evaluated for the old K_S^0 mean life and branching fractions for $K_S^0 \rightarrow \pi^+\pi^-$ and $K_L^0 \rightarrow \pi^+\pi^0$. The new K_S^0 mean life and the current branching fractions increase the MESSNER 73 $|\eta_{+-}|$ value slightly to $(2.246 \pm 0.032) \times 10^{-3}$.

We are troubled by this large unexplained discrepancy. We feel that our normal procedure of averaging and increasing the error by a scale factor S to account for the discrepancy is not adequate for this case. The two new results, when combined with the average of the earlier results by that procedure, give 2.15 ± 0.11 ($S = 6.0$). While this value-and-error makes some sense in that it nearly spans both incompatible sets of data, we choose not to quote it. Instead, since the newer experiments are in principle superior (higher statistics, better acceptance, easier trigger conditions), we have chosen to average them separately from the earlier experiments as is seen in the E+- subsection of the Data Card Listings below.

The entry referenced as GKL/GKS 71 is the average of all seven experiments before 1971. This average is determined from the branching ratios quoted for these experiments in subsections R9 and R20 above. It was quoted by us as the $|\eta_{+-}|$ value in our 1971 edition. The average and fit values at the end of the E+- subsection below do not include

Stable Particles

K_L^0

the pre-1971 $|\eta_{+-}|$ results. The reader may thus utilize the pre-1971 and post-1971 results as desired.

The fitted values in the EOS, E+-, and ER subsections below are the result of a fit to the unparenthesized values of $|\eta_{00}|$, $|\eta_{+-}|$, and $|\eta_{00}/\eta_{+-}|$ in these subsections. We quote the fitted values of $|\eta_{+-}|$ and $|\eta_{00}|$ in the Addendum to the Stable Particle Table but warn in a footnote that they exclude pre-1971 $|\eta_{+-}|$ results.

See also D. Nygren's "Review of K^0 Decays"⁴ for additional discussion of this and other K^0 controversies.

References

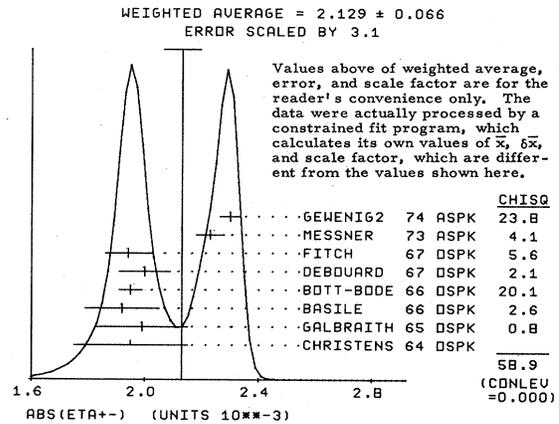
1. D. R. Nygren, private communication.
2. R. DeVoe, private communication.
3. S. Aronson, private communication to D. Nygren reported in reference 4.
4. D.R. Nygren, Review of K^0 Decays, presented at the Berkeley A.P.S. Meeting of the Division of Particles and Fields, Berkeley, California, August 13-17, 1973; LBL-2407.

THE FITTED VALUES OF η_{+-} AND η_{00} GIVEN BELOW ARE THE RESULTS OF A FIT TO η_{+-} , η_{00} AND η_{00}/η_{+-} RESULTS. THE VALUES LISTED BELOW WHICH ARE NOT PARENTHESIZED ENTER THE FIT AS SHOWN. THE VALUES WHICH ARE PARENTHESIZED AND BEAR THE FOOTNOTE X DO NOT ENTER THE FIT AS SHOWN. THESE EXPERIMENTS GIVE BRANCHING RATIOS AND ENTER THE FIT VIA THE QUANTITY ACTUALLY MEASURED -- BRANCHING RATIOS R₉, R₂₀ AND R₂₇ (η_{+-}) AND R₁₇ AND R₁₉ (η_{00}). THESE BRANCHING RATIOS ARE COMBINED WITH CURRENT NORMALIZATIONS AND CURRENT K_L AND K_S MEAN LIVES TO OBTAIN $\pi^+\pi^-$ RATES. THE η_{+-} AND η_{00} VALUES OBTAINED FROM THESE RATES ARE ENTERED BELOW WITH THE NAME 'GKL/GKS'.

EOS	(ETA00)**2 = (A(KL TO 2PI0)/A(KS TO 2PI0))**2 (UNITS 10**--6) ---	
EOS X	0 (-2.) (7.0)	BARTLETT 68 DSPK 10/69
EOS X	57 (4.9) (1.2)	BANNER 69 DSPK 2/72
EOS X	133 (14.1) (3.4)	CENCE 69 DSPK 10/69
EOS XF	180 (13.1) (4.)	GAILLARD 69 DSPK 10/69
EOS X	29 (4.08) (0.9)	BARMIN 70 HLBC 12/70
EOS X	30 (3.61) (1.9)	BUDAGOV 70 HLBC 10/70
EOS C	8.7 3.7	CHOLLET 70 DSPK CU REG.,4 GAMMAS 2/72
EOS XF	172 (9.9) (3.4)	FAISSNER 70 DSPK 12/70
EOS C	56 2.0	WOLFF 71 DSPK CU REG.,4 GAMMAS 12/71
EOS X	5.2 1.1	GKL/GKS 76 RVUE BR SCALE FACTOR=1.5 2/76*
EOS X	SEE NOTE ABOVE REGARDING FITTED VALUES OF η_{+-} AND η_{00} .	
EOS C	CHOLLET 70 GIVES $\eta_{00}=(1.23+-0.24)*(\text{REGEN AMPL, 2GEV/C CUI}/10000\text{MB})$	2/72
EOS C	WOLFF 71 GIVES $\eta_{00}=(1.13+-0.12)*(\text{REGEN AMPL, 2GEV/C CUI}/10000\text{MB})$	2/72
EOS C	WE COMPUTE BOTH η_{00} VALUES FOR (REGEN AMPL, 2GEV/C CUI)=24+-2MB.	2/72
EOS C	THIS REGEN AMPL RESULTS FROM AVERAGING OVER FAISSNER 69.	2/72
EOS C	EXTRAPOLATED USING OPTICAL MODEL CALCULATIONS OF BOHM ET AL.	2/72
EOS C	PL 278 594 (1968) AND THE DATA OF BALATS 71. (FROM H. FAISSNER,	2/72
EOS C	PRIVATE COMMUNICATION)	2/72
EOS F	FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69	
EOS	
EOS AVG	5.90 0.93	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
EOS STUDENT	5.9 1.1	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
EOS FIT	5.41 0.22	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
EOO	THIS FIT VALUE CORRESPONDS TO $\eta_{00}=2.325+-0.093$	
E+-	$\eta_{+-} = A(KL TO \pi^+\pi^-)/A(KS TO \pi^+\pi^-)$ UNITS 10**--3 -----	
E+- X	45 (1.95) (0.20)	CHRISTENS 64 DSPK 2/76*
E+- X	54 (1.99) (0.16)	GALBRAITH 65 DSPK 2/76*
E+- X	(1.92) (0.13)	BASILE 66 DSPK 2/76*
E+- X	(1.95) (0.04)	BOTT-BODE 66 DSPK 2/76*
E+- X	(2.00) (0.09)	DEBOUARD 67 DSPK 2/76*
E+- X	(1.94) (0.08)	FITCH 67 DSPK 2/76*
E+- AX	(1.95) (0.03)	GKL/GKS 71 RVUE EXPTS. BEFORE 71 2/76*
E+- A	AVERAGE OF ABOVE EXPTS. EXCLUDED FROM FIT. SEE TYPED NOTE ABOVE.	
E+- X	4200 (2.23) (0.05)	MESSNER 73 ASPK 11/75*
E+-	2.30 0.035	GENIEG2 74 ASPK 3/74
E+- B X	2.246 0.032	GKL/GKS 76 RVUE BR EXP. AFTER 71 2/76*
E+- B	CURRENTLY INCLUDES ONLY MESSNER 73 + NEW KL TAU AND KS $\pi^+\pi^-$ RATES	
E+- X	SEE NOTE ABOVE REGARDING FITTED VALUES OF η_{+-} AND η_{00} .	
E+-	
E+- AVG	2.271 0.027	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
E+- STUDENT	2.270 0.027	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
E+- FIT	2.272 0.023	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
	(SEE IDEOGRAM BELOW)	
ER	RATIO OF η_{00} OVER η_{+-}	
ER	124 1.03 0.07	BANNER1 72 DSPK 8/72
ER	167 1.00 0.06	HOLDER 72 ASPK 8/72
ER	
ER AVG	1.013 0.046	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
ER STUDENT	1.013 0.049	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
ER FIT	1.023 0.040	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

Data Card Listings

For notation, see key at front of Listings.



Note on $K_L^0 \rightarrow 2\pi$ and K_S Regeneration

Some experiments obtain ϕ_{+-} (the phase of η_{+-}) using K_S , $K_L \rightarrow \pi^+\pi^-$ interference behind a regenerator. In these interference experiments the measured quantity is the difference of ϕ_{+-} and the regeneration phase ϕ_R , as shown in the expression below. After the regenerator, the intensity of the $\pi^+\pi^-$ decays in the forward direction is

$$I(t, p) = S(p) \left[|R(p)|^2 e^{-\Gamma_S t} + |\eta_{+-}|^2 e^{-\Gamma_L t} + 2|R(p)| |\eta_{+-}| \times e^{-(\Gamma_S + \Gamma_L)t/2} \times \cos(\Delta m t + \phi_R(p) - \phi_{+-}) \right], \quad (1)$$

where:

t is the decay time in the K^0 rest frame,
 $\Delta m = m_L - m_S$, and $m_L, \Gamma_L, m_S, \Gamma_S$ are the masses and decay rates of the long- and short-lived K^0 ,

$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}}$ is the ratio of decay amplitudes $A(K_L \rightarrow \pi^+\pi^-)/A(K_S \rightarrow \pi^+\pi^-)$,

$S(p)$ is proportional to the K_L momentum spectrum, and

$R(p) = |R(p)| e^{i\phi_R(p)}$ is the transmission-regenerated K_S amplitude (relative to the K_L):

$$R(p) = \pi N i \frac{[f_0(p) - \bar{f}_0(p)]}{p} \left\{ \frac{-\frac{1}{2}\Gamma_S \ell(p) [1 - 2i\Delta m/\Gamma_S]}{1 - e^{-\frac{1}{2}\Gamma_S \ell(p)}} \right\}, \quad (2)$$

where

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

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$\lambda(p)$ is the thickness of regenerator measured in units of the mean decay length of K_S ,

N is the number of nuclei per cubic centimeter,

Λ is the K_S mean decay length, and

$f_0(p)$, $\bar{f}_0(p)$ are the forward scattering amplitude of K^0 and \bar{K}^0 .

From (1) above it is clear that the value of ϕ_{+-} is highly correlated with the value of Δm and ϕ_R . In addition, a less obvious but significant correlation exists between ϕ_{+-} and Γ_S . Usually Δm is a parameter of the fit and ϕ_R is determined by some other means (optical model calculations, time dependence of the charge asymmetry in K_{e3} decay, etc.).

We list ϕ_{+-} and give in comment cards both the value of ϕ_R used by the authors and the Δm dependence of ϕ_{+-} .

F+-	PHASE OF ETA +- (DEGREES)			
F+-	THE DEPENDENCE OF THE PHASE ON THE KOL-KOS MASS DIFFERENCE IS GIVEN FOR EACH EXPERIMENT IN THE COMMENTS BELOW, WHERE DM IS (MASS DIFF./HBAR) IN UNITS 10**10 SEC-1. WE HAVE EVALUATED THESE MASS DEPENDENCES USING OUR APRIL 1976 VALUE, DM=0.5349+-0.0022 TO OBTAIN THE VALUES AND AVERAGE QUOTED BELOW.			
F+- O	(45.0) (50.0)	FITCH	65 OSPK	BE REGEN 11/67
F+- O	(30.0) (45.0)	FIRESTONE	66 HBC	11/67
F+- O	(70.0) (21.0)	BOTT-BODE	67 OSPK	C REGEN 11/67
F+- O	(25.0) (35.0)	MISCHKE	67 OSPK	CU REGEN 7/68
F+- O	OLD EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE. 2/76*			
F+- N	(51.0) (11.0)	BENNETT2	68 CNTR	CU REG. USES 8/68
F+- C	34.2 10.0	BENNETT	69 CNTR	CU REGEN 2/71
F+- B	45.3 12.0	BOHM	69 OSPK	VACUUM REGEN 2/71
F+- F	45.2 7.4	FAISSNER	69 ASPK	CU REGEN 2/71
F+- J	40.6 4.2	JENSEN	70 ASPK	VACUUM REGEN 2/71
F+- D	37.2 12.0	BALATS	71 OSPK	CU REGEN 9/71
F+- P	36.2 6.1	CARNEGIE	72 ASPK	CU REGEN 1/73
F+- G	46.5 1.6	GEWENIG2	74 ASPK	VACUUM REGEN 3/74
F+- H	45.5 2.8	CARITHERS	75 SPEC	C REGEN 7/75*
F+-	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 11/69			
F+- AVG	45.0 1.3	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
F+- STUDENT	45.1 1.4	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/76*
F+- FIT	45.0 1.2	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/76*
COMMENTS				
F+- N	BENNETT 69 IS A REEVALUATION OF BENNETT2 68. 11/69			
F+- C	BENNETT 69 USES MEASUREMENT OF (F+-)-(PHI) OF ALFF-STEINBERGER 66. 2/71			
F+- B	BOHM 69 F+-=(34.9+-10.0)+69*(DM-.545) DEG. FR=-49.9+-5.4 DEG. 2/71			
F+- F	FAISSNER 69 F+-=(41.0+-12.0)+479*(DM-.526) DEG. 2/71			
F+- J	JENSEN 70 ERROR ENLARGED TO INCLUDE ERROR IN REGENERATOR PHASE. 11/69			
F+- D	BALATS 71 F+-=(42.4+-4.0)+576*(DM-.538) DEG. FR=-42.7+-5.0 DEG. 2/71			
F+- P	CARNEGIE 72 F+- IS INSENSITIVE TO DM. FR=-43.0+-4.0 DEG. 9/71			
F+- G	GEWENIG2 74 F+-=(49.4+-1.0)+565*(DM-.540) DEG. FR=-56.2+-5.2 DEG. 1/73			
F+- H	CARITHERS 75 F+-=(45.5+-2.8)+224*(DM-.5348) DEG. FR=-40.9+-2.6 DEG. 11/75*			
F00	PHASE OF ETA 00 (DEGREES)			
F00	FIRST QUADRANT PREFERRED GORBI 69 OSPK CU REG.,4 GAMMAS 10/70			
F00 C	51.0 30.0	CHOLLET	70 OSPK	CU REG.,4 GAMMAS 10/70
F00 W	56 38.0 25.0	WOLFF	71 OSPK	CU REG.,4 GAMMAS 12/71
F00 C	CHOLLET 70 USES REGENERATOR PHASE FR=-46.5+-4.4 DEG. 1/73			
F00 W	WOLFF 71 USES REGENERATOR PHASE FR=-48.2+-3.5 DEG. 1/73			
F00	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*			
F00 AVG	43.3 19.2	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
F00 STUDENT	43.3 20.7	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/76*
F00 FIT	48.3 13.2	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/76*
DF	PHASE DIFFERENCE F00 - F+- (DEGREES)			
DF B	7.6 18.0	BARBIELLI 73 ASPK		7/73
DF B	INDEPENDENT OF REGENERATOR MECHANISM, DM, AND LIFETIMES. 7/73			
DF	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*			
DF FIT	3.3 13.1	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		2/76*

Superweak Model Predictions for ϕ_{+-} and $Re\epsilon$

The superweak model of Wolfenstein, Phys. Letters 13, 562 (1964) predicts that

$$\phi_{+-} = \phi_{00} = \tan^{-1}\left(\frac{2\Delta m\tau_S}{h}\right)$$

and

$$Re\epsilon = |\eta_{+-}| \left[1 + \left(\frac{2\Delta m\tau_S}{h}\right)^2 \right]^{-1/2}$$

These expressions and the values of the $K_L^0 - K_S^0$ mass difference $\Delta M = (0.5349 \pm 0.0022) \times 10^{10} h \text{ sec}^{-1}$, the K_S^0 mean life $\tau_S = (0.8930 \pm 0.0023) \times 10^{-10} \text{ sec}$, and the magnitude of the $K_L^0 \rightarrow \pi^+\pi^-/K_S^0 \rightarrow \pi^+\pi^-$ amplitude ratio $|\eta_{+-}| = (2.272 \pm 0.023) \times 10^{-3}$, all from the current edition, result in the predictions that

$$\phi_{+-} = \phi_{00} = (43.69 \pm 0.13)^\circ$$

and

$$Re\epsilon = (1.643 \pm 0.017) \times 10^{-3}$$

These can be compared with the experimental values

$$\phi_{+-} = (45.0 \pm 1.2)^\circ$$

$$\phi_{00} = (48.3 \pm 13.2)^\circ$$

$$Re\epsilon = (1.624 \pm 0.088) \times 10^{-3}$$

where $Re\epsilon$ has been computed using the relation

$$Re\epsilon = \frac{\delta}{2} \left(\frac{|1-x|^2}{|1-|x||^2} \right)$$

and our current values of the charge asymmetry parameter for leptonic K_L^0 decay

$\delta = (0.330 \pm 0.012)\%$ and the $\Delta S = -\Delta Q$ amplitude $(Re\epsilon, Im\epsilon) = (0.008 \pm 0.020, -0.003 \pm 0.027)$.

The superweak predictions are in good agreement with the data.

The values of τ_S and $|\eta_{+-}|$ have undergone major revision in the past few years (see notes on these parameters in the K_S^0 and K_L^0 sections of the Data Card Listings). We have utilized only the post-1971 experimental values of these parameters and have made necessary changes in the related parameters ΔM and ϕ_{+-} (see footnotes in the related data card subsections D and F+- above).

Stable Particles

K_L⁰

Data Card Listings

For notation, see key at front of Listings.

13 X = (DS-DQ AMPLITUDE)/(DS+DQ AMPLITUDE)
RELATED TEXT SECTION VI B.4

Table with columns: REAL PART OF X, RESEARCHER, PARTICLE TYPE, and DATE. Includes entries for BALDO-CE, AUBERT, MALLARY, etc.

WEIGHTED AVERAGE = 0.008 ± 0.020
ERROR SCALED BY 1.4

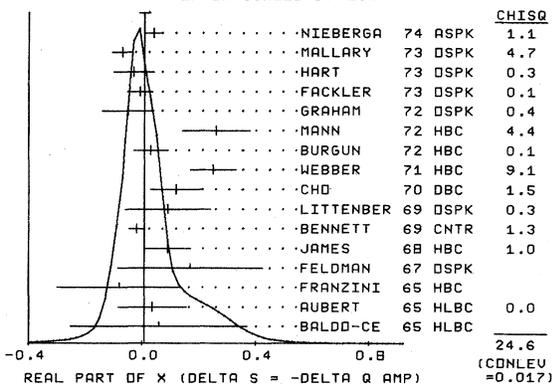
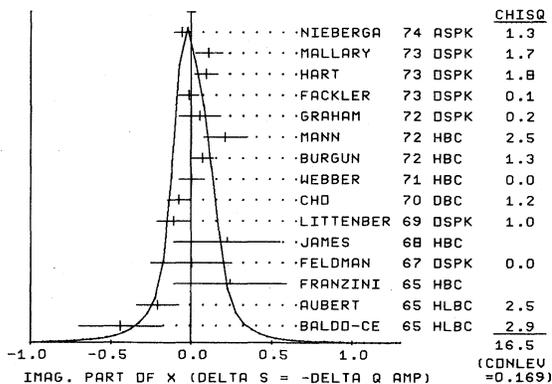


Table with columns: IMX, IMAGINARY PART OF X (ASSUMES MK(L)-MK(S) POSITIVE), RESEARCHER, PARTICLE TYPE, and DATE. Includes entries for BALDO-CE, AUBERT, MALLARY, etc.

WEIGHTED AVERAGE = -0.003 ± 0.027
ERROR SCALED BY 1.2

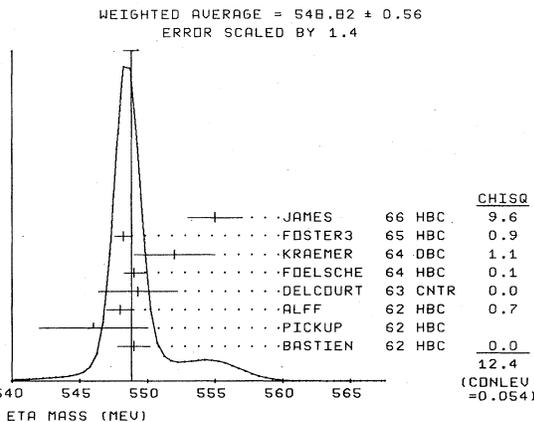


REFERENCES FOR K_L

List of references for K_L particles, including authors like BARON, CRAWFORD, FITCH, GOOD, NEAGU, CAMERINI, ADAIR, ALEKSANYAN, ANIKINA, ASTBURY, FUJII, LUERS, ANIKINA, ANDERSON, etc.

Stable Particles

η



14 ETA WIDTH

Mode	Method	Value	Source
91	FRM MASS SPECTRUM	110.0	ALFF 62 HBC
148	FRM MASS SPECTRUM	OR LESS	FDELSCHKE 64 HBC
31	FRM MASS SPECTRUM	OR LESS	JAMES 66 HBC
4.0	FRM MASS SPECTRUM	OR LESS	BALTAY 66 DBC
1.9	FRM MASS SPECTRUM	OR LESS	JONES 66 CNTR
ETA WIDTH DETERMINED FROM DECAY RATE (UNITS KEV)			
THIS IS THE PARTIAL DECAY RATE (W1) FOR THE MODE (ETA INTO 2GAMMA)			
DIVIDED BY THE FITTED BRANCHING FRACTION (P1) FOR THAT MODE.			
FIT		0.85	0.12

14 ETA PARTIAL DECAY MODES

Mode	Description	Decay Masses
P1	ETA INTO 2GAMMA	0+ 0
P2	ETA INTO 5PI0	134+ 134+ 134
P3	ETA INTO PI+ PI- PI0	139+ 139+ 134
P4	ETA INTO PI+ PI- GAMMA	139+ 139+ 0
P5	ETA INTO E+ E- PI0 (VIOLATES C IN E.M.I.)	134+ .5+ .5
P6	ETA INTO E+ E- PI+ PI-	139+ 139+ .5+ .5
P7	ETA INTO PI0 2GAMMA	154+ 0+ 0
P8	ETA INTO E+ E- GAMMA	.5+ .5+ 0
P9	ETA INTO 2PI0 GAMMA (VIOLATES C)	134+ 134+ 0
P10	ETA INTO PI+ PI- PI0 GAMMA	139+ 139+ 134+ 0
P11	ETA INTO PI+ PI- 2GAMMA	139+ 139+ 0+ 0
P12	ETA INTO MU+ MU- GAMMA	105+ 105
P13	ETA INTO MU+ MU- GAMMA	105+ 105+ 0
P14	ETA INTO MU+ MU- PI0	105+ 105+ 134
P15	ETA INTO PI+ PI-	139+ 139

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \cdot \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 7	P 8
P 1	.3799+-0.0098					
P 2	-.2691	.2990+-0.106				
P 3	-.3224	-.2353	.2358+-0.0056			
P 4	-.2866	-.2095	.8201	.0489+-0.0013		
P 7	-.4271	-.5781	-.0939	-.0801	.0314+-0.0109	
P 8	-.0434	-.0326	-.0494	-.0501	-.0036	.0050+-0.0012

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., $G_i = \Gamma_i = \Gamma_{total} P_i$, in appropriate units. In analogy to the matrix above, the diagonal elements are $G_i \pm \delta G_i$, where $\delta G_i = \sqrt{(\delta G_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta G_i \delta G_j) / (\delta G_i \cdot \delta G_j)$. Note that, because of the error in Γ_{total} , the errors and correlations here are not directly derivable from those above.

	G 1	G 2	G 3	G 4	G 7	G 8
G 1	.3240+-0.0460					
G 2	-.9386	.2550+-0.0385				
G 3	-.9689	.9513	.2011+-0.0294			
G 4	-.9646	.9471	1.0097	.0417+-0.0061		
G 7	-.3661	-.3028	.4240	.4226	.0268+-0.0104	
G 8	-.5178	-.5082	.5233	.5199	.2218	.0043+-0.0012

Data Card Listings

For notation, see key at front of Listings.

14 ETA DECAY RATES

Note on Conflicting Results for $\eta \rightarrow \gamma\gamma$ Decay Rate

The BROWMAN 74 result for $\Gamma_{\gamma\gamma}$ is three standard deviations below the earlier BEMPORAD 67 result. This is because the BEMPORAD 67 analysis finds a negligible strong-production amplitude and attributes all of their observed η photoproduction cross section to Coulomb production. BROWMAN 74, on the other hand, find that the strong-production contribution to the cross section is significant, especially at lower energies, and that the Coulomb contribution is smaller, resulting in a smaller value of $\Gamma_{\gamma\gamma}$.

The η photoproduction data of BROWMAN 74 were taken at incident γ -ray energies of 5.8, 9.0, and 11.45 GeV on beryllium, aluminum, copper, silver, and uranium targets, while the BEMPORAD 67 data were taken at 4.0 and 5.5 GeV on lead, silver, and zinc. The higher energies, higher statistics, better angular resolution, and larger range of atomic weights of the BROWMAN 74 experiment result in a more reliable separation of the Coulomb amplitude from the strong-production amplitude. In addition, the bad approximation mentioned in BROWMAN 74 footnote 4 was present in the theoretical work utilized by BEMPORAD 67 and this may account for some of the discrepancy.

Browman et al. state that the BROWMAN 74 result is compatible with the BEMPORAD 67 data, although the agreement is not as good as that given by the BEMPORAD 67 fit to these data. On the other hand, the BEMPORAD 67 result appears to be incompatible with the 9 GeV data (especially uranium) shown in Fig. 2 of BROWMAN 74.

We quote the BROWMAN 74 result.

Mode	Rate	Source	Notes
W1	ETA INTO 2GAMMA (UNITS KEV)	(G1)	
W1 B	(1.00) (0.22)	BEMPORAD 67 CNTR	PRIMAKOFF EFFECT 11/75*
W1	0.324 0.046	BROWMAN 74 CNTR	PRIMAKOFF EFFECT 7/74*
W1 B	BEMPORAD 67 GIVES W1=1.21+-26 KEV ASSUMING THAT W1/TOTAL=0.314.		11/75*
W1 B	BEMPORAD PRIVATE COMMUNICATION GIVES MORE GENERAL RESULT AS		11/75*
W1 B	W1*W1/TOTAL=.380+-0.083. WE EVALUATE THIS USING W1/TOTAL=.38+-0.1.		11/75*
W1 B	NOT INCLUDED IN AVERAGE BECAUSE THE UNCERTAINTY RESULTING FROM THE		2/76*
W1 B	SEPARATION OF THE COULOMB AND NUCLEAR AMPLITUDES HAS APPARENTLY		2/76*
W1 B	BEEN UNDERESTIMATED. SEE NOTE ON DISCREPANCY ABOVE.		2/76*
W1	0.324 0.046	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

η

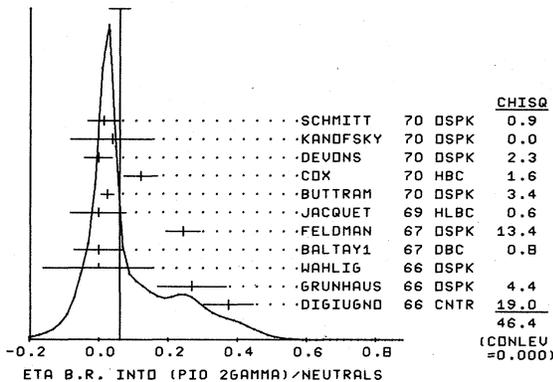
14 ETA BRANCHING RATIOS

R1	ETA INTO NEUTRALS/CHARGED				(P1+P2+P7)/(P3+P4+P8)
R1 N	10 (2.5)	11.0	PICKUP	62 HBC	
R1 N	53 (3.20)	(1.26)	BASTIEN	62 HBC	
R1 N	(2.7)	(0.8)	SHAFFER	62 HBC	
R1	2.6	.9	BUSCHBECK	63 HBC	7/66
R1 N	280 (4.5)	(1.0)	JAMES	66 HBC	6/66
R1 N THESE EXPERIMENTS HAVE NOT BEEN USED IN COMPUTING THE AVERAGES					
R1 N AS THEY WERE UNABLE TO SEPARATE CLEARLY PARTIAL MODES (3) AND (4)					
R1 N FROM EACH OTHER. THE REPORTED VALUES THUS PROBABLY CONTAIN					
R1 N SOME (UNKNOWN) FRACTION OF MODE (4).					
R1	2.64	0.23	BALTAY2	67 DBC	11/67
R1				
R1 AVG	2.64	0.22	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R1 STUDENT	2.64	0.24	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R1 FIT	2.452	0.081	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R2	ETA INTO 2 GAMMA/CHARGED				(P1)/(P3+P4+P8)
R2	0.99	0.48	CRAWFORD	63 HBC	
R2 FIT	1.311	0.053	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		

Note on $\eta \rightarrow \pi^0 \gamma \gamma$

The discrepancies between various measurements of branching ratios involving $\eta \rightarrow \pi^0 \gamma \gamma$ are displayed in the ideogram below, in which all relevant experiments have been converted to a common ratio, $\pi^0 \gamma \gamma$ /neutrals. Our branching ratio fit does not include DIGIUGNO 66, FELDMAN 67, or the upper limit measurements. See page 43 of "Review of Particle Properties", Physics Letters 39B, No. 1 (1972) for more discussion.

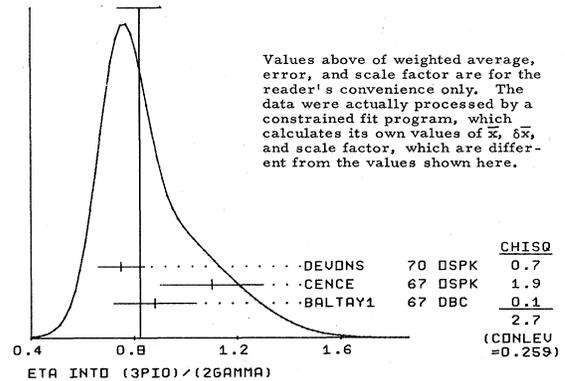
WEIGHTED AVERAGE = 0.061 ± 0.031
ERROR SCALED BY 2.3



R3	ETA INTO (PIO 26AMMA)/NEUTRALS				(P7)/(P1+P2+P7)	
R3 S	(0.375)	(0.072)	DIGIUGNO	66 CNTR	ERROR DOUBLED	
R3 THE ERRORS OF DIGIUGNO+66 HAVE BEEN INCREASED BY A FACTOR						
R3 OF TWO, TO TAKE INTO ACCOUNT POSSIBLE SYSTEMATIC ERRORS, AS						
R3 SUGGESTED BY THE AUTHORS.						
R3	.27	.10	GRUNHAUS	66 DSPK	8/67	
R3 R	(.028)	(.044)	BUNIA TOV	67 DSPK	11/67	
R3 S	(.244)	(.051)	FELDMAN	67 DSPK	8/67	
R3 S SEE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.						
R3	.026	.019	BUTTRAM	70 DSPK	12/70	
R3	.122	.052	COX	70 HBC	6/70	
R3	(.07)	OR LESS	CL=.90	DEVONS	70 DSPK	
R3 R	.016	.047	SCHMITT	70 DSPK	12/70	
R3 R	SCHMITT 70 IS A REANALYSIS BUNIA TOV 67					
R3 E	(0.11)	(0.03)	STRUGALSK	71 HLBC	5/71	
R3 E THIS MEASUREMENT HAS BEEN EXCLUDED BECAUSE THE ERROR APPEARS						
R3 E TO BE SERIOUSLY UNDERESTIMATED.						
R3					
R3 AVG	0.042	0.023	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)			
R3 STUDENT	0.039	0.019	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R3 FIT	0.044	0.015	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)			
R4	ETA INTO (PI+ PI- GAMMA)/(PI+ PI- PIO)				(P4)/(P3)	
R4	0.14	0.08	FOELSCHKE	64 HBC		
R4 M	(.75)	(0.25)	PAULI	64 DBC		
R4 M THIS EXPERIMENT HAS NOT BEEN INCLUDED IN THE AVERAGES SINCE IT IS						
R4 M NOT CLEAR THAT THEIR CLASS B EVENTS ARE ACTUALLY FROM ETAS.						
R4	0.30	0.06	CRAWFORD	66 HBC	6/66	
R4	.10	.10	KRAEMER	64 DBC	7/66	
R4	.196	.041	FOSTER3	65 HBC	7/66	
R4	.25	.035	LITCHFIELD	67 DBC	8/67	
R4	0.28	0.04	BALTAY2	67 DBC	11/67	
R4	.201	.006	CORNLEY	70 ASPK	6/70	
R4	18K	0.209	0.004	THALER	73 ASPK	6/73
R4					
R4 AVG	0.2074	0.0037	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)			
R4 STUDENT	0.2074	0.0037	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R4 FIT	0.2075	0.0033	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)			

R5	ETA INTO (3PIO) + 2/3(PIO 26AMMA)/ PI+PI-PIO				(P2+2/3P7)/P3
R5	0.83	0.32	CRAWFORD	63 HBC	7/66
R5	2.0	1.0	FOELSCHKE	64 HBC	7/66
R5	0.90	0.24	FOSTER1	65 HBC	7/66
R5				
R5 AVG	0.91	0.19	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R5 STUDENT	0.91	0.20	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R5 FIT	1.357	0.057	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)		
R6	ETA INTO 3PIO/26AMMA				(P2)/(P1)
R6	(.90)	OR MORE	CHRETIEN	62 HBC	
R6	0.88	0.16	BALTAY1	67 DBC	11/67
R6	1.1	0.2	CENCE	67 DSPK	1/68
R6	0.75	0.09	DEVONS	70 DSPK	12/70
R6				
R6 AVG	0.824	0.085	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)		
R6 STUDENT	0.821	0.085	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R6 FIT	0.787	0.039	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
(SEE IDEOGRAM BELOW)					

WEIGHTED AVERAGE = 0.824 ± 0.085
ERROR SCALED BY 1.2



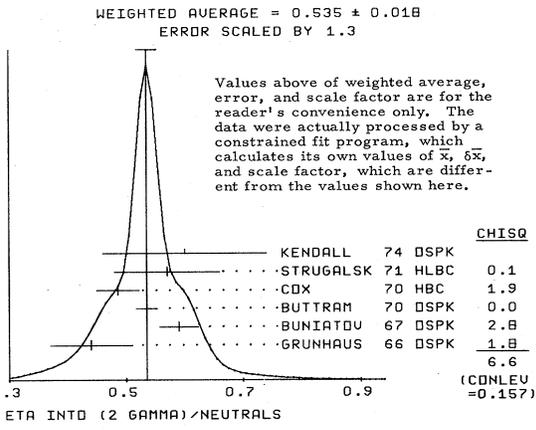
R7	ETA INTO 26AMMA/(PI+ PI- PIO)				(P1)/(P3)
R7	401	1.72	.25	FOSTER1	65 HBC
R7				BAGLIN	69 HLBC
R7				
R7 AVG	1.69	0.21	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R7 STUDENT	1.57	0.23	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R7 FIT	1.611	0.065	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R8	ETA INTO NEUTRAL/(PI+ PI- PIO)				(P1+P2+P7)/(P3)
R8	50	3.6	0.8	KRAEMER	64 DBC
R8				PAULI	64 DBC
R8		2.89	0.56	ALFF-STEI	66 HBC
R8	244	3.6	0.6	FLATTE2	67 HBC
R8	29	3.4	1.1	AGULLAR-B	72 HBC
R8 B	70	2.83	0.80	BLOODWORTH	72 HBC
R8 B	74	2.54	1.89	KENDALL	74 DSPK
R8 B	ERROR INCREASED FROM PUBLISHED VALUE 0.5 BY BLOODWORTH, PRIV. COMM.				
R8				
R8 AVG	3.26	0.30	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R8 STUDENT	3.26	0.33	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R8 FIT	3.01	0.10	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R9	ETA INTO (E+E-PIO)/(PI+PI-PIO) (UNITS 10**4)				(P5)/(P3)
R9	110.	OR LESS	PRICE	65 HBC	
R9	0	77.	OR LESS	FOSTER2	65 HBC
R9	42.	OR LESS	CL=.90	BAGLINI	67 HLBC
R9	0	16.	OR LESS	CL=.90	BILLING
R9	1.9	OR LESS	CL=.90	JANE1	75 DSPK
R10	ETA INTO (E+E-PI+PI-)/TOTAL (UNITS 10**2)				(P6)
R10	(0.7)	OR LESS	RITTENBER	65 HBC	
R11	ETA INTO (E+E-PI+PI-)/(PI+PI-GAMMA)				(P6)/(P4)
R11	1	0.026	0.026	GROSSMAN	66 HBC
R12	ETA INTO 2 GAMMA/NEUTRALS				(P1)/(P1+P2+P7)
R12 S	(0.416)	(0.044)	DIGIUGNO	66 CNTR	ERROR DOUBLED
R12	.44	.07	GRUNHAUS	66 DSPK	8/67
R12 S	(.579)	(.052)	FELDMAN	67 DSPK	8/67
R12 S SEE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.					
R12 T	(0.39)	(0.06)	JONES	66 CNTR	8/67
R12 T THIS RESULT FROM COMBINING CROSS SECTIONS FROM TWO DIFFERENT EXPTS.					
R12	.59	.033	BUNIA TOV	67 DSPK	11/67
R12	.535	.018	BUTTRAM	70 DSPK	12/70
R12	.486	.036	COX	70 HBC	6/70
R12	.32	0.09	STRUGALSK	71 HLBC	5/71
R12	113	0.60	0.14	KENDALL	74 DSPK
R12				
R12 AVG	0.535	0.018	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R12 STUDENT	0.535	0.015	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R12 FIT	0.535	0.015	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)		
(SEE IDEOGRAM BELOW)					
R13	ETA INTO 3PIO/NEUTRALS				(P2)/(P1+P2+P7)
R13 S	(0.209)	(0.054)	DIGIUGNO	66 CNTR	ERROR DOUBLED
R13 R	(.29)	(.10)	GRUNHAUS	66 DSPK	8/67
R13 S	(.177)	(.035)	FELDMAN	67 DSPK	8/67
R13 S SEE THE NOTE ON ETA DECAY INTO NEUTRALS ABOVE.					
R13 R	REDUNDANT INFORMATION FROM THIS EXPERIMENT.				
R13 R	(.439)	(.024)	BUTTRAM	70 DSPK	12/70
R13	.392	.042	COX	70 HBC	6/70
R13	0.32	0.09	STRUGALSK	71 HLBC	5/71
R13	75	0.40	0.14	KENDALL	74 DSPK
R13				
R13 AVG	0.397	0.025	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R13 STUDENT	0.397	0.027	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		
R13 FIT	0.421	0.014	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		

Stable Particles

η

Data Card Listings

For notation, see key at front of Listings.



Label	Parameter	Value	Source	Notes
R14	ETA INTO P10 (2 GAMMA)/2GAMMA	(.51) OR LESS CL=.90	WAHLIG 66 SPRK	(P71)(P1) 7/66
R14		0.0 0.14	BALTAY1 67 DBC	11/67
R14 P		(0.05) (0.04)	BONAMY 67 SPRK	PRELIMINARY RESULT 11/67
R14		0.083 0.030		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
R14 FIT				
R15	ETA INTO (E+e-P10)/TOTAL (UNITS 10**=-2)	(0.71) OR LESS CL=.90	BAZIN 68 DBC	(P5) 6/66
R15		(0.084)OR LESS CL=.90	BAZIN 68 DBC	6/68
R16	ETA INTO 2GAMMA/(3P10 + P10 2GAMMA)	0.80 .25	BACCI 63 CNTR	(P11)/(P2+P7) 7/66
R16		1.150 0.060		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
R16 FIT				
R17	ETA INTO (PI+PI-P10 GAMMA)/(PI+PI-P10) (UNITS 10**=-2)	(7.0) OR LESS CL=.95	PRICE 67 HBC	(P101)/(P3) 8/67
R17		(0.9) OR LESS CL=.95	PRICE 67 HBC	8/67
R17		(1.6) OR LESS CL=.95	BALTAY2 67 DBC	11/67
R17		(1.7) OR LESS CL=.90	ARNOLD 68 HLBC	9/68
R17 T		0.24 OR LESS CL=.90	THALER 73 ASPK	6/73
R17 T		0 0 0.15		LIMIT ABOVE RESTATED FOR AVERAGING 6/73
R18	ETA INTO (PI+PI-2GAMMA)/(PI+PI-P10)	(.0091)OR LESS CL=.95	BALTAY2 67 DBC	(P111)/(P3) 8/67
R18		(.0161)OR LESS CL=.95	BALTAY2 67 DBC	11/67
R19	ETA INTO 3P10/(PI+ PI- P10)	1.47 0.20 0.17	BAGLIN2 67 HLBC	(P21)/(P3) 8/67
R19		1.50 .15 .29	BAGLIN 69 HLBC	9/68
R19		1.46 0.13		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) 7/69
R19 AVC		1.46 0.14		AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R19 STUDENT		1.268 0.060		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R19 FIT				
R20	ETA INTO 2GAMMA/((3P10)+2/3(P10 2GAMMA))	1.10 0.5	MULLER 63 DBC	(P11)/(P2+3P7) 7/66
R20		1.187 0.058		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
R20 FIT				
R21	ETA INTO NEUTRALS/TOTAL	.79 .08	BUNIATOV 67 DSPK	(P1+P2+P7) 11/67
R21		.705 .008	BASILE 71 CNTR MM SPECTROMETER	8/71
R21		0.7058 0.0080		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R21 STUDENT		0.7058 0.0086		AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R21 FIT		0.7103 0.0068		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R22	ETA INTO (P10 2GAMMA)/TOTAL	.12 OR LESS CL=.95	JACQUET 69 HLBC	(P7) 6/70
R22		0.031 .011		FROM FIT
R22 FIT				
R23	ETA INTO MU+MU-/TOTAL (UNITS 10**=-5)	0 2. OR LESS CL=.95	WEHMANN 68 OSPK	(P12) 4/68
R24	ETA INTO MU+MU-P10/TOTAL (UNITS 10**=-4)	5. OR LESS	WEHMANN 68 OSPK	(P14) 4/68
R25	ETA INTO MU+MU-/2GAMMA (UNITS 10**=-5)	5.9 2.2	HYAMS 69 OSPK	(P12)/(P11) 7/69
R26	ETA INTO (P10 2GAMMA)/(3P10 + P10 2GAMMA)	0.1 0.3	KANDFSKY 70 OSPK	(P71)/(P2+P7) 2/71
R26 N		0.1 0.3		WE HAVE CHANGED THE ERROR ON THIS EXPERIMENT FROM +0.3, -0.1 TO THE ABOVE +0.3, -0.3 SINCE IT IS CLEAR FROM FIGURE 7 IN THE ARTICLE THAT A CENTRAL VALUE OF 0.0 IS ABOUT AS PROBABLE AS THE QUOTED VALUE OF 0.1
R26		0.095 0.032		FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
R26 FIT				
R27	ETA INTO (PI+ PI-)/TOTAL (UNITS 10**=-2)	(0.15) OR LESS	THALER 73 ASPK CON. LEV. NOT GIVEN	(P15) 6/73
R28	ETA INTO (E+e-GAMMA)/(PI+PI-P10) (UNITS 10**=-2)	2.11 0.50	JANZ 75 OSPK	(P81)/(P3) 2/76*
R28				VALUE CHANGED BY ERRATUM. 2/76*
R28 FIT				

14 ETA C-NONCONSERVING DECAY PARAMETERS

RELATED TEXT SECTION VI C.1

A1	LEFT-RIGHT ASYMMETRY PARAMETER FOR PI+ PI- P10 (UNITS 10**=-2)	8/66
A1	1351 7.2 2.8	BALTAY 66 DBC
A1	1300 5.8 3.4	CLPHY 66 HBC
A1	10665 (0.3) (1.0)	CNDPS 66 OSPK REPL BY MULLER 69
A1	705 -6.1 4.0	LARRIBE 66 HBC
A1	G36800 (1.5) (1.5)	GORMLEY3 68 ASPK
A1	10709 1.3 1.1	MULLER 69 OSPK
A1	1138 -1.4 3.	CARPENTR 70 HBC
A1	349 3.2 5.4	DANBURG 70 DBC
A1	220K -0.05 0.22	LAYER 72 ASPK
A1	165K 0.23 0.26	JANE1 74 OSPK
A1 G	GORMLEY3 68 ASYMMETRY PROBABLY DUE TO UNMEASURED (E X B) SPK. CH.	3/74
A1 G	EFFECTS. NEW EXPTS. WITH (E X B) CONTROLS DONT OBSERVE ASYMMETRY.	3/74
A1	AVG .12 0.17	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
A1	STUDENT 0.11 0.19	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
A2	LEFT-RIGHT ASYMMETRY PARAMETER FOR PI+ PI- GAMMA (UNITS 10**=-2)	11/66
A2	33 -2. 17.	CRAWFORD 66 HBC
A2	-4. 8.	LITCHFIELD 67 DBC
A2 N	1620 1.5 2.5	MULLER 69 OSPK
A2	7257 1.22 1.56	GORMLEY 70 ASPK
A2	36K 0.5 0.6	THALER 72 ASPK
A2	35K 1.2 0.6	JANE2 74 OSPK
A2 N	MULLER 69 IS SENSITIVE ONLY TO UPPER .4 OF GAMMA-RAY SPECTRUM.	3/74
A2	AVG 0.88 0.40	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
A2	STUDENT 0.88 0.45	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
AS	SEXTANT ASYMMETRY PARAMETER FOR PI+ PI- P10 (UNITS 10**=-2)	12/75*
AS	1300 6.8 3.3	CLPHY 66 HBC
AS	705 -2.4 4.0	LARRIBE 66 HBC
AS	37K 0.5 0.5	GORMLEY 68 WIRE
AS	220K 0.10 0.22	LAYER 72 ASPK
AS	165K 0.20 0.25	JANE1 74 OSPK
AS	AVG 0.19 0.16	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AS	STUDENT 0.19 0.17	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
AQ	QUADRANT ASYMMETRY PARAMETER FOR PI+ PI- P10 (UNITS 10**=-2)	12/75*
AQ	220K -0.07 0.22	LAYER 72 ASPK
AQ	165K -0.30 0.25	JANE1 74 OSPK
AQ	AVG -0.17 0.17	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AQ	STUDENT -0.17 0.18	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
BET	BETA FOR ETA TO PI+ PI- GAMMA. SENSITIVE TO D-WAVE CONTRIBUTION.	12/75*
BET	DN/DCOS THETA = SIN**2 THETA * (1 + BETA * COS**2 THETA)	12/75*
BET	7250 -0.060 0.065	GORMLEY 70 WIRE
BET L	0.12 0.06	THALER 72 ASPK
BET	35K 0.11 0.11	JANE1 74 OSPK
BET L	AUTHORS DONT BELIEVE THIS TO INDICATE D-WAVE BECAUSE DEPENDENCE OF BETA ON GAMMA ENERGY INCONSISTENT WITH THEOR. PREDICTION.	12/75*
BET L	COS**2 DEPENDENCE MAY ALSO COME FROM P AND F-WAVE INTERFERENCE.	12/75*
BET	AVG 0.047 0.062	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
BET	STUDENT 0.053 0.053	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

14 ENERGY DEPENDENCE OF ETA DALITZ PLOT

RELATED TEXT SECTION VI C.2

THE FOLLOWING EXPTS FIT TO ONE OR MORE OF THE COEFFICIENTS

DP	A, B, C, D, OR E FOR ETA TO PI+ PI- P10	
DP	MATRIX ELEMENT**2+1 + A**2 + B**2 + C**2 + D**2 + E**2 + EX**2 + EY**2	
DP	1300 SEE TEXT SEC VI C.2	CLPHY 66 HBC
DP	705 SEE TEXT SEC VI C.2	LARRIBE 66 HBC
DP	7170 SEE TEXT SEC VI C.2	CNDPS 68 OSPK
DP	37K SEE TEXT SEC VI C.2	GORMLEY3 68 WIRE
DP	826 SEE TEXT SEC VI C.2	BAGLIN 69 HLBC
DP	1138 SEE TEXT SEC VI C.2	CARPENTR 70 HBC
DP	349 SEE TEXT SEC VI C.2	DANBURG 70 DBC
DP	7250 SEE TEXT SEC VI C.2	GORMLEY 70 WIRE
DP	220K SEE TEXT SEC VI C.2	LAYER 72 ASPK
DP	81K SEE TEXT SEC VI C.2	LAYER 73 ASPK
AO	ALPHA PARAMETER FOR ETA TO 3 P10	
AO	MATRIX ELEMENT**2 = 1 + 2*ALPHA*HZ (SEE TEXT SEC VI C.3)	
AO	192 -0.32 0.37	BAGLIN 70 HLBC

REFERENCES FOR ETA

PEVSNER 61 PRL 7 421	PEVSNER, KRAEMER, NUSSBAUM, RICHARDSON + (JHU)
ALFF 62 PRL 9 322	ALFF, BERLEY, COLLEY, BRUGGER + (COLU+RUTGERS)
BASTIEN 62 PRL 8 114	BASTIEN, BERGE, DAHL, FERRO-LUZZI + (LRL)
CHRETIEN 62 PRL 9 127	CHRETIEN+ (BRAN+BROWN+HARVARD+MIT+PADOVA)
PICKUP 62 PRL 8 329	E PICKUP, ROBINSON, SALANT (CNRC+BNL)
SHAFFER 62 CERN CONF 307	J SHAFFER, FERRO-LUZZI, MURRAY + (UCB+BNL)
BACCI 63 PRL 11 37	BACCI, PENSO, SALVINI + (ROMA+FRAS)
BUSCHBECK 63 SIENA CONF 1 166	BUSCHBECK-CZAPP, COOPER + (VIENNA, CERN, AMST)
CRAWFORD 63 PRL 10 546	F S CRAWFORD, LLOYD, FOWLER (LRL+DUKE)
DEL COURT 63 PL 7 215	F S CRAWFORD, L LLOYD, E FOWLER (LRL+DUKE)
MULLER 63 SIENA CONF 99	DEL COURT, LEFRANCOIS, PEREZ Y JORBA + (ORSAY)
FOELSCH 64 PR 134 B 1138	MULLER, PAULI + (SACL+ROMA)
KRAEMER 64 PR 136 B 496	H W FOELSCH, H L KRAYBILL (LYALE)
PAULI 64 PL 13 351	KRAEMER, MADANSKY, FIELDS + (JHU+NWES+WOOD)
FOSTER 65 PR 138 B 652	E PAULI, A MULLER (ISACLAY)
FOSTER 65 THESIS	FOSTER, PETERS, MEER, LOEFFLER + (WISC+PURDUE)
FOSTER 65 PRL 15 123	FOSTER, GOOD, MEER (WISCONSIN)
RITTENBERG 65 PRL 15 556	M.C. FOSTER (WISCONSIN)
ALFF-STE 66 PR 145 1072	L.R. PRICE, F.S. CRAWFORD (LRL)
BALTAY 66 PR 16 1224	RITTENBERG, KALBFLEISCH (LRL+BNL)
CLPHY 66 PR 149 1044	ALFF-STEINBERGER, BERLEY + (COLUMBIA+RUTGERS)
CNDPS 66 PL 22 546	+FRANZINI, KIM, KIRSCH+ (COLUMBIA+STONY BROOK)
CRAWFORD 66 PR 16 333	COLUMBIA, LRL, PURDUE, WISCONSIN, YALE
DIGIUGNO 66 PRL 16 767	CNDPS, FINOCCHIARIS, LASSALLE, (CERN, ETHZ, SACL)
GROSSMAN 66 PR 146 993	F.S. CRAWFORD, L.R. PRICE (LRL)
GRUNHAUS 66 THESIS	DIGIUGNO, GIORGI, SILVESTRI + (NAPL, TRST, FRAS)
JAMES 66 PR 142 896	R GROSSMAN, L PRICE, F CRAWFORD (LRL)
JONES 66 PL 23 597	J. GRUNHAUS (COLUMBIA)
LARRIBE 66 PL 23 600	F E JAMES, H L KRAYBILL (COLUMBIA)
WAHLIG 66 PRL 17 221	JONES, BINNIE, DUANE, HORSEY, MASON, (LOIC, RHEL)
	LARRIBE, LEVEQUE, MULLER, PAULI, + (SACL+RHEL)
	WAHLIG, SHIBATA, MANNELLI (MIT+PSIA)

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

n, p, n

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include BAZIN, COHEN, HARRISON, etc.

Table listing particle data for protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include COHEN, TAYLOR, etc.

Table listing particle data for antiprotons and other related particles. Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include BUTTON, ROBERTS, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include HARRISON, SANDARS, WRIGHT, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include TAYLOR, COHEN, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include MATTAUCH, COHEN, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include COHEN, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include BAIRD, DRESS, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include SOSNOVSKI, CHRISTENSEN, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include CONFORTO, KRUPF, etc.

Table listing particle data for neutrons (n) and protons (p). Includes columns for name, mass, magnetic moment, and electric dipole moment. Entries include COHEN, HARRISON, etc.

Stable Particles

Data Card Listings

A

For notation, see key at front of Listings.



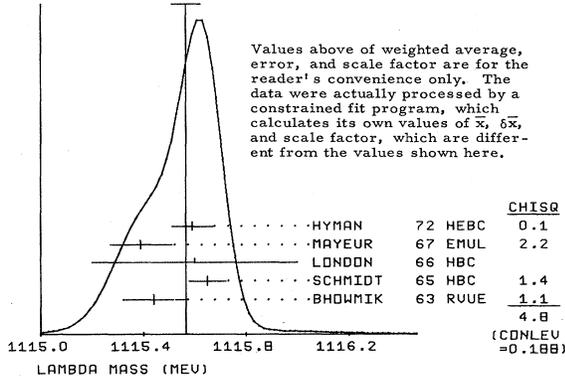
18 LAMBDA(1115,JP=1/2+) I=0

18 LAMBDA MASS (MEV)

M N SINCE OUR FINAL VALUES FOR THE SIGMA AND LAMBDA MASSES COME FROM... WE HAVE USED THE UNCORRELATED MEASUREMENTS FROM SCHMIDT 65 RATHER THAN THE ONES COMING FROM THE OVERALL FIT REPORTED IN THAT PAPER.

Table of LAMBDA MASS measurements with columns for mass, error, and source. Includes entries like BHOWMIK 63 RVUE, BALTAY 65 HBC, etc.

WEIGHTED AVERAGE = 1115.566 ± 0.056
ERROR SCALED BY 1.3



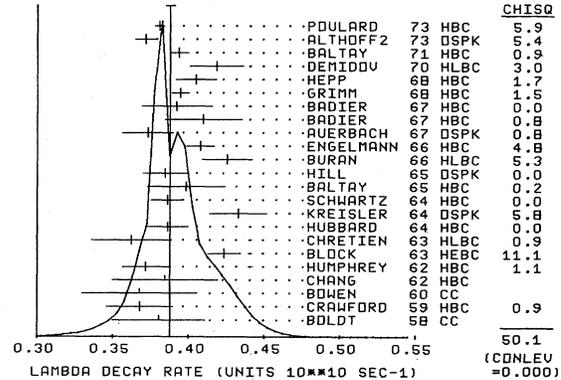
Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program...

Table for 18 LAMDA - ANTI LAMBDA MASS DIFFERENCE (MEV) with columns for difference, error, and source.

18 LAMBDA MEAN LIFE (UNITS 10**-10)

Table of LAMBDA MEAN LIFE measurements with columns for mean life, error, and source. Includes entries like BOLDT 58 CC, CRAWFORD 59 HBC, etc.

WEIGHTED AVERAGE = 0.3879 ± 0.0031
ERROR SCALED BY 1.6



18 (LAMBDA - ANTI LAMBDA)/AVG., MEAN LIFE DIFFERENCE

Table for 18 (LAMBDA - ANTI LAMBDA)/AVG., MEAN LIFE DIFFERENCE with columns for difference, error, and source.

18 LAMBDA MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

Table of LAMBDA MAGNETIC MOMENT measurements with columns for moment, error, and source.

18 LAMBDA ELECTRIC DIPOLE MOMENT (UNITS 10**-14 E CM)
NONZERO VALUE IMPLIES VIOLATION OF T AND P

Table of LAMBDA ELECTRIC DIPOLE MOMENT measurements with columns for moment, error, and source.

18 LAMBDA PARTIAL DECAY MODES

Table of LAMBDA PARTIAL DECAY MODES with columns for decay mode and branching ratio.

18 LAMBDA BRANCHING RATIOS

Table of LAMBDA BRANCHING RATIOS with columns for ratio, error, and source. Includes entries like R1 LAMBDA INTO (P PI-)/((P PI-)+(N P10)), etc.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Λ , Σ^+

R4	LAMBDA INTO (P MU- NEU)/TOTAL (UNITS 10**4)	(P31)/(P1+P2)	
R4	1 (0.2) OR MORE	GOOD	62 HBC
R4	1 (1.0) OR LESS	ALSTON	63 HBC
R4	2 (1.0) OR LESS	KERNAN	64 FBC
R4	BETWEEN 1.3 AND 6.0	LIND	64 HBC
R4	3 1.3	LIND	64 RVUE
R4	2 1.5	RONNE	64 FBC
R4	9 2.4	CANTERL	71 HBC STOPPED K-P
R4	14 1.4	BAGGETT2	72 HBC STOP K-
R4	AVG	1.57	0.35 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R4	STUDENT	1.56	0.38 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

18 LAMBDA DECAY PARAMETERS

RELATED TEXT SECTION VI D AND APPENDIX III

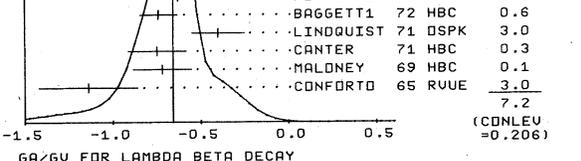
A-	ALPHA LAMBDA- (LAMBDA INTO PI- PROTON)		
A-	1156	0.62	0.07 CRONIN 63 CNTR LAMBDA FROM PI-P
A-	(0.663)	(0.022)	BERGE 66 RVUE INCLUDES ABOVE
A-	10130	0.645	0.017 OVERSETH 67 OSPK LAMBDA FROM PI-P
A-	M 2529	(0.747)	(0.086) MERRILL 68 HBC REPL BY DAUBER 68
A-	3520	0.67	0.06 DAUBER 69 HBC FROM XI DECAY
A-	10325	0.649	0.023 CLELAND 72 OSPK LAMBDA FROM PI-P
A-	AVG	0.647	0.013 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
A-	STUDENT	0.647	0.014 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

A0	ALPHA/A - FOR LAMBDA (L INTO P0 N/L INTO PI- P)		
A0	4760	1.000	0.068 CORK 60 CNTR LAMBDA
A0	AVG	1.006	0.066 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
A0	STUDENT	1.006	0.071 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
A0	0	DONE BY COMPARING PROTON DISTR. WITH N DISTR. FROM LAMBDA DECAY.	

F-	PHI ANGLE (SIN(PHI)/COS(PHI)=BETA/GAMMA) (DEGREES)		
F-	1156	13.0	17.0 CRONIN 63 OSPK LAMBDA FROM PI-P
F-	10130	-8.0	6.0 OVERSETH 67 OSPK LAMBDA FROM PI-P
F-	7377	(-9.2)	(5.2) CLELAND 67 OSPK REPL BY CLELAND 72
F-	10325	-7.0	4.5 CLELAND 72 OSPK LAMBDA FROM PI-P
F-	AVG	-6.5	3.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
F-	STUDENT	-6.6	3.8 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

AV	GA/GV FOR LAMBDA BETA DECAY (SEE TEXT SEC. VI D.1 FOR SIGN CONV.)		
AV	C 22	(-1.03)	LIND 64 HBC
AV	C 102	(0.6)	OR MORE BAGLIN 65 HBC NO SIGN GIVEN
AV	C	BETW 0. AND -1.1	BARLOW 65 OSPK
AV	C 102	(0.7)	OR MORE CL=.95 ELY 65 HBC ABS. VALUE
AV	C	EXPERIMENTS INCLUDED IN CONFORTO 65, RVUE	
AV	148	-0.72	0.14 0.19 MALONEY 69 HBC
AV	M 1078	(-0.62)	(0.08) (0.09) ALTHOFF2 71 OSPK POLARIZED LAMBDA
AV	M 141	-0.75	0.15 0.18 CANTER 71 HBC
AV	L 173	(-0.32)	(0.13) (0.17) LINDQUIST 71 OSPK UP-DOWN ASYMMETRY
AV	M 173	(-0.68)	(0.27) (0.54) LINDQUIST 71 OSPK E-NEU CORRELATION
AV	L 173	-0.40	0.13 0.17 LINDQUIST 71 OSPK E-NEU AND UP-DOWN
AV	M 352	-0.74	0.09 0.12 BAGGETT1 72 HBC STOP K-
AV	A 817	-0.65	0.06 ALTHOFF1 71 OSPK POLARIZED LAMBDA
AV	A	ALTHOFF1 73 INCLUDES DATA OF ALTHOFF2 71. USES PROT SPECTRUM AND	
AV	A	THREE SPIN ASYMMETRIES.	
AV	M	EXPT MEASURES ONLY THE ABSOLUTE VALUE OF A/V	
AV	L	LINDQUIST 71 GETS THREE VALUES. WE AVERAGE THE ONE THAT USES	
AV	L	ALL DATA.	
AV	AVG	-0.658	0.054 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)
AV	STUDENT	-0.660	0.051 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

WEIGHTED AVERAGE = -0.658 ± 0.054
ERROR SCALED BY 1.2



REFERENCES FOR LAMBDA

EISLER 57 NC 5 1700	EISLER, PLANO, SAMIOS, SCHWARTZ + (COLU+BNL)
BOLDT 58 PRL 1 148	E BOLDT, D O CALDWELL, Y PAL (MIT)
CRAWFORD 59 PRL 2 266	CRAWFORD, CRESTI, DOUGLASS, GOOD + (LRL)
BAGLIN 60 NC 18 1043	BAGLIN, BLOCH, BRITSON, HENNESSY + (EPOL)
BOWEN 60 PR 119 2030	BOWEN, HARDY, REYNOLDS, SUN + (PRINCETON)
CORK 60 PR 120 1000	CORK, KERTH, WENZEL, CRCNIN+ (LRL+PRIN+BNL)
COLUMBIA 60 ROCH CONF 726	M SCHWARTZ + (COLUMBIA)
HUMPHREY 61 PRL 6 478	HUMPHREY, KIRZ, ROSENFELD, RHEE + (LRL+SYRA)
ANDERSON 62 CERN CONF-832	ANDERSON, CRAWFORD, GOLDEN, LLOYD + (LRL)
AUBERT 62 NC 25 479	AUBERT, BRISSON, HENNESSY, SIX + (EPOL)
CHANG 62 THESIS DUKE	CHUEN CHUEN CHANG (DUKE)
COOL 62 PR 127 2223	COOL, HILL, MARSHALL + (BNL+MIT+NYU+ANL)
GOOD 62 PRL 9 518	M L GOOD, V G LIND (WISCONSIN)
HUMPHREY 62 PR 127 1305	W E HUMPHREY, R ROSS (LRL)
ALSTON 63 UCRL 10926	ALSTON, KIRZ, NEUFELD, SOLMITZ, WOHLMUT (LRL)
BHOWMIK 63 NC 28 1494	B BHOWMIK, D P GOYAL (DELHI)
BLOCK 63 PR 130 766	BLOCK, GESSAROLI, RATTI + (NWS+BGNA+SYRA+ORNL)
BROWN 63 PR 130 769	BROWN, KADYK, TRILLING, ROE + (LRL+MICH)
CHRISTEN 63 PR 131 2208	CHRISTEN, CROUCH + (BRAN+BROWN+HARVARD+MIT)
CRONIN 63 PR 129 1795	J W CRONIN, D E OVERSETH (PRINCETON)
ELY 63 PR 131 868	ELY, GIDAL, KALMUS, OSWALD, POWELL + (LRL)
KERNAN 63 PR 129 870	KERNAN, NOVEY, MARSHAW, WATTENBERG (ANL+LTL)
ANDERSON 64 PRL 13 167	J A ANDERSON, F S CRAWFORD (LRL)
BAGLIN 64 NC 35 977	BAGLIN, BINGHAM+ (EPOL+CERN+LOUC+RHEL+BERG)
HUBBARD 64 PR 135 B 183	HUBBARD, BERGE, KALBFLEISCH, SHAFER + (LRL)
KERNAN 64 PR 135 B 1271	KERNAN, POWELL, SANDLER + (LRL+LOUC)
KREISLER 64 PR 136 B 1074	M N KREISLER, D OVERSETH, J CRONIN + (LRL+LOUC)
LIND 64 PR 135 B 1483	LIND, S INFORD, GOOD, STERN (WISCONSIN)
RONNE 64 PL 11 357	RONNE+ (CERN+EPOL+LOUC+UNIV. BERGEN)
SCHWARTZ 64 UCRL 1360 THESIS	JOSEPH ADAM SCHWARTZ (LRL)
BAGLIN 65 NC 35 977	BAGLIN + (EPOL, CERN, LOUC, RHEL, BERGEN)
BALTAY 65 PR 140 B 1027	BALTAY, SANDWEISS, CULWICK, KOPP + (YALE+BNL)
BARLOW 65 PL 18 64	J BARLOW, BLAIR, CONFORTO+ (CERN+RHEL+PENN)
CHARRIERE 65 PL 15 86	CHARRIERE, GIBSON+ (EPOL+BRIS+CERN+MPIM)
ALSO 66 NC 46A 205	CHARRIERE, GIBSON+ (EPOL+BRIS+CERN+MPIM)
CONFORTO 65 EC INT HERZEGNOVI	G CONFORTO (CERN)
ELY 65 PR 137 B1302	ELY, GIDAL, KALMUS, POWELL + (LPL, LOUC)
HILL 65 PRL 15 85	HILL, L I, JENKINS, KYCIA, RUDERMAN (MIT, BNL)
SCHMIDT 65 PR 140 B 1328	P SCHMIDT (COLUMBIA)
BERGE 66 BERKELEY 46	BERGE, CABIBBO (RVUE) LRL, CERN)
BURAN 66 PL 20 318	BURAN, ETIVINDSON, SKJEGGESTAD, TOFFE + (OSLO)
CITEN 66 PR 152 B 1171	+LACH, SANDWEISS, TAFT, YEN, OREN + (YALE+BNL)
ENGELMAN 66 NC 45A 1038	ENGELMANN, FILTHUTH, ALEXANDER+ (HEID, REHD)
GIBSON 66 NC 45A 882	W M GIBSON, K GREEN (BRIS)
LONDON 66 PR 143 1034	LONDON, RAU, GOLDBERG, LICHTMAN+ (BNL, SYRA)
AUERBACH 67 NC 47A 19	AUERBACH, BRIAN, DOBBS, LANDE, MANN+ (PENN)
BADIER 67 PL 25B 152	+BONNET, BIANDET, SADOULET (EPOL)
CLELAND 67 PL 26B 45	CLELAND, BIENLEIN, CONFORTO+ (CERN+GEVA+LUND)
MAYEUR 67 ULLBR. BRUX. BUL32	G. MAYEUR, E. TOMPA, J. WICKENS (BELG, LOUC)
OVERSETH 67 PRL 19 391	O E OVERSETH, R F ROTH (MICH+PRIN)
GRIMM 68 NC 54A 187	H.-J. GRIMM (HEIDELBERG)
HEPP 68 ZPHYS 214 71	V. HEPP, H. SCHLEICH (HEIDELBERG)
MERRILL 68 PR 167 1202	MERRILL, SHAFER (LRL)
DAUBER 69 PR 179 1262	+BERGE, HUBBARD, MERRILL, MILLER (LRL)
DOYLE 69 UCRL 18139-THESIS	J. C. DOYLE (LRL)
MALONEY 69 PRL 23 425	MALONEY, SECHI-ZORN (UNIV MARYLAND)
BOHM 70 NC 70A 384	+ KRECKER + (BERL+BRUX+DUUC+LOUC+LOWC+HARS)
DEMIDOV 70 S JNP 10 481	+KRILL, V. UGRYUMOV, PONOSOV, PROTASOV+ (ITEP)
OLSEN 70 PRL 24 843	+PCNDROM, HANDLER, LIMON, SMITH + (MISC, MICH)
ALTHOFF1 71 PL 378 531	+BROWN, FREYTAG, HEARD, HEINTZE + (CERN, HEID)
ALTHOFF2 71 PL 378 535	+BROWN, FREYTAG, HEARD, HEINTZE + (CERN, HEID)
BALTAY 71 PR 04 670	+BRIDGEWATER, COPPER, HABIBI + (COLUMBING)
BARKOV 71 JETPL 14 60	+GUREVICH, MAKARINA, MARTEMYANOV+ (ITEP)
BARONI 71 LNC 2 1256	G BARONI, S PETREIRA, G ROMANO (ROMA)
CANTER 71 PRL 26 868	+COLE, LEE-FRANZINI, LOVELESS + (STON+COLL)
CANTER1 71 PRL 27 59	+COLE, LEE-FRANZINI, LOVELESS+ (STON+COLL)
DAHLJENS 71 NC 3A 1	+CERN+ANKA+LAUS+MPI+ROMA)
HILL 71 PR 04 1979	+L I, JENKINS, KYCIA, RUDERMAN (MIT, BNL)
ALSO 65 PRL 15 85	HILL, L I, JENKINS, KYCIA, RUDERMAN (MIT, BNL)
LINDQUIST 71 PRL 27 612	LINDQUIST, SUMNER+ (EFI, WUSL, OSU, ANL)
BAGGETT1 72 ZPHY 249 279	+BAGGETT, EISELE, FILTHUTH, FREHSE+ (HEID)
BAGGETT2 72 ZPHY 252 362	+BAGGETT, EISELE, FILTHUTH, FREHSE+ (HEID)
BAGGETT3 72 PL 428 579	+BAGGETT, EISELE, FILTHUTH, FREHSE, HEPP+HEID)
CLELAND 72 NP 840 221	+CONFORTO, EATON, GEPBER+ (CERN+GEVA+LUND)
HYMAN 72 PR 05 1063	+BUNNELL, DERRICK, FIELDS, KATZ+ (ANL+CERN)
ALTHOFF1 73 PL 438 237	+BROWN, FREYTAG, HEARD, HEINTZE+ (CERN+HEID)
ALTHOFF2 73 NP 866 29	+BROWN, FREYTAG, HEARD, HEINTZE+ (CERN+HEID)
POULARD 73 PL 468 135	+GIVERNAUD, BORG (SACL)

PAPERS NOT REFERRED TO IN DATA CARDS

ARMENTEROS 62 CERN CONF 236	ARMENTEROS+ (CERN+EPOL+LOUC+BRN+CEN-SACLAY)
BALTAY 62 CERN CONF 233	BALTAY, FOWLER, SANDWEISS, CULWICK+ (YALE+BNL)
BERGE 63 THESIS (BERKELEY)	J PETER BERGE (LRL)

Σ^+

19 SIGMA+(1189, JP=1/2+ I=1)

19 SIGMA+ MASS (MEV)

M N SEE NOTE PRECEDING LAMBDA MASS LISTINGS

M	144 1189.38	0.15	BARKAS 63 EMUL + SEE NOTE 5 BELOW
M	58 1189.48	0.22	BHOWMIK 64 EMUL + SEE NOTE 5 BELOW
M	S	ABOVE SIGMA+ MASSES HAVE BEEN RAISED 30 KEV TO ACCOUNT FOR 46 KEV	
M	S	INCREASE IN PROTON MASS AND 21 KEV DECREASE IN PION MASS	
M	4205 1189.61	0.08	SCHMIDT 65 HBC SEE NOTE N
M	1189.16	0.12	HYMAN 67 HBC
M	R 607 1189.33	0.04	89HM 72 EMUL
M	B	BOHM 72 UPDATED WITH PDG APR. 73 K-, PI- AND P0 MASSES.	
M	AVG	1189.371	0.060 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)
M	STUDENT	1189.354	0.041 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
M	FIT	1189.366	0.057 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)

(SEE IDEOGRAM BELOW)

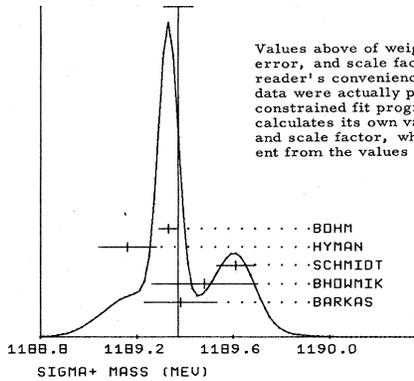
Stable Particles

Σ^+

Data Card Listings

For notation, see key at front of Listings.

WEIGHTED AVERAGE = 1189.371 ± 0.060
ERROR SCALED BY 1.8



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{x} , $\bar{\sigma}$, and scale factor, which are different from the values shown here.

Table with 3 columns: Name, Particle Type, CHISO value. Includes entries for BDHM, HYMAN, SCHMIDT, BHDWMK, BARKAS.

19 SIGMA+ MEAN LIFE (UNITS 10**+10)

Table listing mean life measurements for Sigma+ particles from various experiments (GLASER, PUSCHEL, EVANS, etc.) with associated error bars and CHISO values.

19 SIGMA+ MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

Table listing magnetic moment measurements for Sigma+ particles (COOK, KOTELCHUC, SULLIVAN, etc.) with associated error bars and CHISO values.

19 SIGMA+ PARTIAL DECAY MODES

Table listing partial decay modes for Sigma+ particles (SIGMA+ INTO PROTON P, SIGMA+ INTO NEUTRON N, etc.) with associated branching ratios and CHISO values.

19 SIGMA+ BRANCHING RATIOS

Table listing branching ratios for Sigma+ particles (SIGMA+ INTO NEUTRON P+/(NUCLEON P), SIGMA+ INTO NEUTRON P+ GAMMA, etc.) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (CARRARA, BAZIN, QUARENI, etc.) with associated error bars and CHISO values.

WEIGHTED AVERAGE = 0.240 ± 0.035
ERROR SCALED BY 1.4

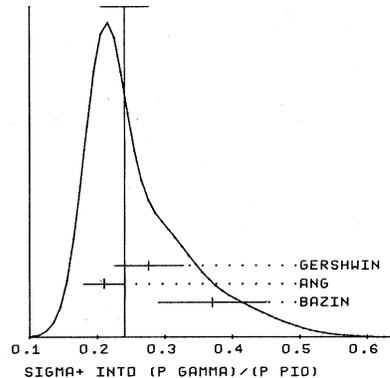


Table with 3 columns: Name, Particle Type, CHISO value. Includes entries for GERSHWIN, ANG, BAZIN.

Table listing Sigma+ data for various experiments (R5 ED, R5 U, R5 UA, etc.) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (R6, R6 E, R6 U) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (R7, R7 E) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (R8, R8 A) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (R9, R9 E) with associated error bars and CHISO values.

Table listing Sigma+ data for various experiments (R10, R10 E) with associated error bars and CHISO values.

19 SIGMA+ DECAY PARAMETERS

Table listing decay parameters for Sigma+ particles (ALPHA+ALPHA FOR SIGMA+ TO PI+ N, SIGMA+ TO PI+ P, etc.) with associated error bars and CHISO values.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Σ⁺, Σ⁻

ALPHA SIGMA0 (SIG+ INTO P10 PROTON) BEALL 62 CNTR REPLAC. BY BANGE
AO 0.80 0.16
AO (-0.901) (0.25) TRIPP 62 HBC
AO D 5200 (-0.986) (0.072) BANGERTER 66 HBC K-P TO SIG+ PI- 7/66
AO 32000 -0.999 0.022 BANGERTER 69 HBC 10/69
AO H 1335 -0.98 0.05 0.02 HARRIS 70 OSPK 5/70
AO 16K -0.940 0.045 BELLAMY 72 ASPK PI+P TO SIG+ K+ 11/72
AO L 1259 -0.945 0.055 0.042 LIPMAN 73 OSPK PI+P TO SIG+ 7/73
AO L DECAY PROTONS SCATTERED OFF ALUMINUM.
AO H DECAY PROTONS SCATTERED OFF CARBON.
AO
AO AVG -0.979 0.016 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AO STUDENT -0.979 0.018 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
F+ PHI+ ANGLE (SIG+ INTO N PI) SIN(PHI)/COS(PHI)=BETA/GAMMA (DEGREE)
F+ O 370 (180.) (30.) BERLEY 66 HBC + NEUTRON RESCATT. 9/66
F+ 560 143. 29. BANGERTI 69 HBC 10/69
F+ C105% 18%. 24. BERLEY 70 HBC K-P AT 400 MEV/C 11/69
F+ C CHANGED FROM 176 TO 184 TO AGREE WITH SIGN CONVENTION.
F+
F+ AVG 167.3 20.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
F+ STUDENT 167.5 21.2 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
AG ALPHA SIGMA0 (SIG+ INTO PROTON GAMMA)
AG 61 -1.03 0.52 0.42 GERSHWIN 69 HBC K-P TO SIG PI 11/69
FO PH10 ANGLE (SIG+ INTO P10 PROTON) SIN(PHI)/COS(PHI)=BETA/GAMMA (DEG)
FO H 22.0 90.0 HARRIS 70 OSPK PI+P TO SIG+ K+ 5/70
FO L 1259 38.1 35.7 37.1 LIPMAN 73 OSPK PI+P TO SIG+ K+ 7/73
FO L DECAY PROTON SCATTERED OFF ALUMINUM.
FO H DECAY PROTONS SCATTERED OFF CARBON.
FO
FO AVG 15.8 33.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
FO STUDENT 35.8 36.3 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

REFERENCES FOR SIGMA+

CORK 60 PR 120 1000
EVANS 60 NC 15 873
FREDEN 60 NC 16 611
KAPLON 60 ANP 9 139
PUSCHEL 60 NP 20 254
BARKAS 61 PR 124 1209
BERTHELO 61 NC 21 693
CHIESA 61 NC 19 1171
BEALL 62 PRL 8 75
GRARD 62 PR 127 407
GALTIERI 62 PRL 9 26
HUMPHREY 62 PR 127 1305
TRIPP 62 PRL 9 66
BARKAS 63 PRL 11 26
ALSO 61 UCRL 9450
BHOWMIK 64 NP 53 22
CARRARA 64 PL 12 72
COURANT 64 PR 136 B 1791
MURPHY 64 PR 134 B 188
NAUENBERG 64 PRL 12 679
WILLIS 64 PRL 13 291
BALTAY 65 PR 140 B 1027
BAZIN 65 PRL 14 154
BAZIN 65 PRL 14 154
CARAYAN 65 PR 138 B 433
QUARENI 65 NC 40 A 928
SCHMIDT 65 PR 140 B 1328
BANGERTER 66 PRL 17 495
BERLEY 66 PRL 17 1071
CHANG 66 PR 151 1081
ALSO 65 NEVIS 145 THESIS
CHEN 66 PR 152 1171
COOK 66 PRL 17 223
BAGGETT 67 PRL 19 1458
ALSO 68 VIENNA ABS. 374
ALSO 68 PRIVATE COMM.
BARASH 67 PRL 19 181
EISELE 67 ZPHYS 205 409
HYMAN 67 PL 25 B 376
KOTELCHU 67 PRL 18 1166
SULLIVAN 67 PRL 18 1163
ALSO 64 PRL 13 246
BIERMAN 68 PRL 20 1459
COMBE 68 NC 57A 54
MAST 68 PRL 20 1312
ANG 69 ZPHYS 228 151
BAGGETT 69 MDDP-TR-973
BALTAY 69 PRL 22 615
BANGERTER 69 UCRL-19244
BANGERTI 69 PR 187 821
BARLOUTA 69 NP B14 153
EISELE 69 ZPHYS 221 1
EISELE 69 ZPHYS 221 401
GERSHWIN 69 PR 188 2077
ALSO UCRL 19246 THESIS
NORTON 69 NEVIS 175 (THESIS)
BERLEY 70 PR D1 2015
EISELE 70 ZPHY 238 372
HARRIS 70 PRL 24 165
ALLEY 71 PR 03 75
BAKKE 71 LNC 1 37
COLE 71 PR D4 631
TOVEE 71 NP 833 493
BELLAMY 72 PL 398 299
BOHM 72 NP 848 1
ALSO 73 TME-73-2 NOV
EBENHOH 73 ZPHY 264 413
LIPMAN 73 PL 43B 89
SAHA 73 PR D7 3295
SCHIZOR 73 PR D8 12
EBENHOH 74 ZPHY 266 367
REUCROFT 76 PREPRINT
CORK, KERTH, WENZEL, CRONIN, COOL (LRL+PRIN+BNL)
BRIST+BRUSS+IAS-UJ, COL-DUBL IN-LON+MILAN-PAD
S FREDEN+H KORNBUM, R WHITE (LRL)
M KAPLON, A MELISSINOS, YAMANOUCHI (ROCH)
W PUSCHEL (MAX PLANCK INST)
BARKAS, DYER, HASON, NICHOLS, SMITH (LRL)
BERTHELOT, DAUDIN, GOUSSU (SACLAY+ORSAY)
CHIESA, QUASSIATI, RINAUDO (INFN-TURIN)
BEALL, CORK, KEEFE, MURPHY, WENZEL (LRL)
F GRARD, G SMITH (LRL)
GALTIERI, BARKAS, HECKMAN, PATRICK, SMITH (LRL)
W E HUMPHREY, R R ROSS (LRL)
R D TRIPP, M B WATSON, M FERRO-LUZZI (LRL)
W H BARKAS, J N DYER, H H HECKMANN (LRL)
JOHN DYER (THEISIS, BERKELEY) (LRL)
B BHOWMIK, P JAIN, P MATHUR, LAKSHMI (DELHI)
CARRARA, CRESTI, GRIGOLETTO, PERUZZO (PADOVA)
COURANT, FILTHUTH+ (CERN+HEID+UMD+NRL+BNL)
C THORNTON MURPHY (WISCONSIN)
NAUENBERG, MARATECK+ (COLU+RUTG+PRIN)
WILLIS+ COURANT, ENGELMAN+ (BNL, CERN, HEID, UMD)
BALTAY, SANDHEISS, CULWICK, KOPP + (YALE+BNL)
BAZIN, BLUMENFELD, NAUENBERG + (PRIN+COLU)
BAZIN, PLANO, SCHMIDT+ (PRIN, RUTG, COLU)
CARAYAN+POPOULOS, TAUFEST, WILLMAN (PURDUE)
QUARENI, CARTACCI + (BGN, FIRZ, GENO, PARMA)
P SCHMIDT (COLUMBIA)
BANGERTER, GALTIERI, BERGE, MURRAY+ (LRL)
+HERZBACH, KOFLER, YAMAMOTO + (BNL+MAS+YALE)
CHUNG YUN CHANG (COLUMBIA)
CHUNG YUN CHANG (COLUMBIA)
+LACH, SANDHEISS, TAFT, YEH, OREN + (YALE+BNL)
V COOK, EWART, MASEK, ORR, PLATNER (WASHINGTON)
BAGGETT, DA V, GLASSER, KEHOE, KNOP+ (MARYLAND)
BAGGETT, KEHOE (MARYLAND)
N. BAGGETT (MARYLAND)
BARASH, DAY, GLASSER, KEHOE, KNOP + (MARYLAND)
+ENGELMANN, FILTHUTH, FOHLSCH, HEPP+ (HEID)
+LOKEN, PEWITT, MCKENZIE, + (ANL+CARM+MWS)
KOTELCHUCK, GOZA, SULLIVAN, ROSS (VANDERBILT)
SULLIVAN, MCINTURFF, KOTELCHUCH (VANDERBILT)
A D MCINTURFF, C E ROOS (VANDERBILT)
BIERMAN, KOUNOSU, NAUENBERG + (PRINCETON)
CERN-BRISTOL-LAUSANNE-MUNICH-ROME-COLLABOR
MAST, GERSHWIN, ALSTON-GARNJUST + (HEID)
+EBENHOH, EISELE, ENGELMANN, FILTHUTH+ (UMD)
N V BAGGETT (THEISIS) (UMD)
BALTAY, FRANZINI, NEWMAN, NORTON+ (COLU, STON)
ROGER ODELL BANGERTER (THEISIS) (LRL)
BANGERTER, GARNJUST, GALTIERI, GERSHWIN+ (LRL)
BARLOUTA, BELLEFON, GRANET+ (SACL+GERN+HEID)
+ENGELMANN, FILTHUTH, FOHLSCH, HEPP+ (HEID)
+ENGELMANN, FILTHUTH, FOHLSCH, HEPP+ (HEID)
+ALSTON-GARNJUST, BANGERTER + (LRL)
LAWRENCE, K GERSHWIN (LRL)
HERBERT NORTON (COLUMBIA)
+YAMIN, HERTZBACH, KOFLER + (BNL, MAS, YALE)
+FILTHUTH, HERR, PRESER, ZIEHL (HEID, WISC)
+OVERSETH, PONDROM, DETTMANN (MICH, WISC)
+BENBROOK, COOK, GLASS, GREEN, HAGUE + (WASH)
+SABRE COLLAB. (ZEEM+SACL+BGNA+REHO+EPDL)
+LEE-FRANZINI, LOVELESS, BALTAY+ (STON, COLU)
LOUC, BELGRADE, BERL, BRUX, DOUBLIN, WARS COLLAB
+ANDERSON, CRAWFORD, OSMON+ (RHEL+MEL+SUSS)
BERLIN+BELGRADE+BRUX+DOUBLIN+LUDC+WARSAW
BRUSSELS BULLETIN, SAME COLLABORATION
+EISELE, FILTHUTH, HEPP, LEITNER, THOU+ (HEID)
+UTO, WALKER, MONTGOMERY+ (RHEL+SUSS+LOWC)
+FEYKOWICH, HEINTZELMAN, MELTZER + (GARM)
B. SECHI-ZORN, G. SNOW (UMD)
+EISELE, ENGELMANN, FILTHUTH, HEPP + (HEID)
+ROOS, WATERS, WEBSTER, HANSL+ (VAND, MPIH)

PAPERS NOT REFERRED TO IN DATA CARDS
GLASER 58 CERN CONF 270 GLASER, GOOD, MORRISON (MICH+LFL)
QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS
TRIPP 62 PRL 8 175 R TRIPP, M WATSON, M FERRO-LUZZI (LRL) P
ALFF 63 SIENA CONF 1 205 ALFF, NAUENBERG, KIRSCH, + (COLU+RUTG+BNL)
ALSO 65 PR 137 B 1105 ALFF, GELFAND, BRUGGER, BERLEY+ (COLU+RUTG+BNL)
COURANT 63 SIENA CONF 1 73 COURANT, FILTHUTH, BURNSTEIN, DAY+ (CERN+UMD)

Σ- 20 SIGMA- (1198, JP=1/2+) I=1

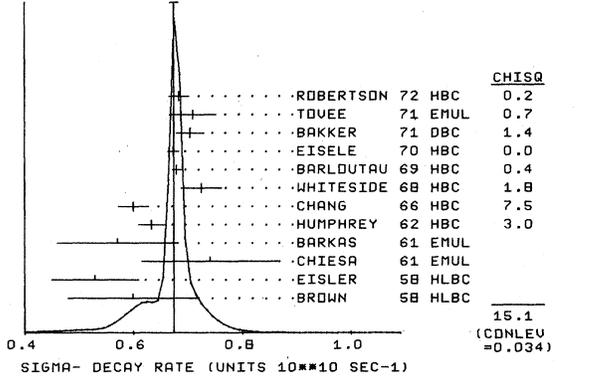
20 SIGMA- MASS (MEV)
M N SEE NOTE PRECEDING LAMBDA MASS LISTINGS
M 3000 1197.43 0.08 SCHMIDT 65 HBC SEE NOTE N 3/74
M FIT 1197.35 0.06 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*

20 (SIGMA-) - (SIGMA+) MASS DIFFERENCE (MEV)
D 87 8.25 0.40 BARKAS 63 EMUL -
D 2500 8.25 0.25 DOSCH 65 HBC
D 86 7.91 0.23 BOHM 72 EMUL 1/73
D AVG 8.09 0.16 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
D STUDENT 8.10 0.18 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
D FIT 7.98 0.38 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2) 2/76*

20 (SIGMA-) - (LAMBDA) MASS DIFFERENCE (MEV)
DL N SEE NOTE PRECEDING LAMBDA MASS LISTINGS.
DL 85 81.70 0.19 BURNSTEIN 64 HBC 9/66
DL 2279 81.80 0.13 SCHMIDT 65 HBC SEE NOTE N 3/74
DL 81.64 0.09 HEPP 68 HBC 8/68
DL AVG 81.693 0.069 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
DL STUDENT 81.692 0.077 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
DL FIT 81.750 0.054 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*

20 SIGMA- MEAN LIFE (UNITS 10**-10)
T 1.67 0.40 0.28 BROWN 58 HLBC
T 1.89 0.33 0.25 EISLER 58 HLBC
T 45 1.35 0.32 0.17 CHIESA 61 EMUL
T 41 1.75 0.39 0.30 BARKAS 61 EMUL
T 1208 1.58 0.06 0.06 HUMPHREY 62 HBC STOP. K- 6/66
T C 3267 1.666 0.075 CHANG 66 HBC STOP. K- 9/67
T S 61 (2.08) (0.22) CHEN 66 HBC + 6.9 PBAR P, ANTI 9/67
T S 64 (1.46) (0.31) CHEN 66 HBC + 6.9 PBAR P, ANTI 9/67
T 506 1.38 0.07 WHITESIDE 68 HBC STOP. K- 6/68
T 10253 1.472 0.016 BARLOUTA 69 HBC K-P 1.4-1.2 GEV/C 11/69
T 1M 1.485 0.022 EISELE 70 HBC K-P AT REST 2/71
T 1383 1.42 0.05 BAKKER 71 DBC - K-N TO SIG- 2PI 10/71
T 1.41 0.09 TOVEE 71 EMUL 12/71
T 2400 1.463 0.039 ROBERTSON 72 HBC K-P +25 GEV/C 3/74
T C CHANG ERROR 0.018 RAISED BY US. SEE 1970 EDITION, RMP 42, 1231 (1970) 1/73
T S ERROR PURELY STATISTICAL.
T AVG 1.482 0.017 0.017 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.5)
T STUDENT 1.479 0.013 0.013 AVG. USING STUDENT10(H/1.11) -- SEE TEXT
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.6746 ± 0.0076
ERRR SCALED BY 1.5



20 SIGMA- MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)
MM R, BTWN -1.6 AND +0.8 FIX 73 CNTR SIG-ATOM FINE ST 3/74
MM R -1.48 0.37 ROBERTS 74 CNTR SIG-ATOM FINE ST 12/75*
MM R ROBERTS 74 INCLUDES DATA FROM FOX 73. 12/75*

Stable Particles

Data Card Listings

For notation, see key at front of Listings.

Σ^-, Σ^0

20 SIGMA- PARTIAL DECAY MODES

Table with columns P1, P2, P3, P4, P5 and decay modes like SIGMA- INTO NEUTRON PI- GAMMA, SIGMA- INTO NEUTRON MU- NEUTRINO, SIGMA- INTO LAMBDA E- NEUTRINO.

20 SIGMA- BRANCHING RATIOS

Table with columns R1, R2, R3, R4 and branching ratios for various decay channels, including average error factors.

20 SIGMA- DECAY PARAMETERS

RELATED TEXT SECTION VI D AND APPENDIX III

Table with columns A- and decay parameters like ALPHA SIGMA-, PHI ANGLE, and average error factors.

Table with columns AV and decay parameters for SIGMA TO LAMBDA BETA DECAY, including average error factors.

Table with columns AV1 and decay parameters for SIGMA TO NEUTRON BETA DECAY, including average error factors.

REFERENCES FOR SIGMA- BROWN 58 CERN CONF 270, EITSLER 58 NC SER10 10 150, BROWN, GLASER, GRAVES, PERL, CRONIN + (MICH) EITSLER, BASSI, CONVERSI + (COLU, BNL, BGNA, PISA)

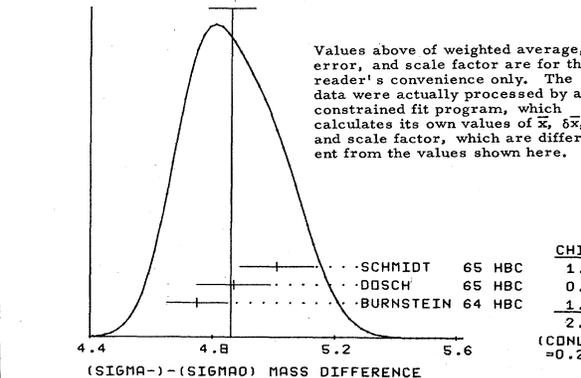
Table listing particle properties for SIGMA- and SIGMA0, including authors, dates, and experimental details.

***** SIGMA0(1193, JP=1/2+) I=1 *****

21 SIGMA0(1193, JP=1/2+) I=1

Table with columns D1, N and mass difference data for SIGMA- and SIGMA0.

WEIGHTED AVERAGE = 4.860 ± 0.076 ERROR SCALED BY 1.2



Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Ξ^0, Ω^-

23 XI0 DECAY PARAMETER

RELATED TEXT SECTION VI D AND APPENDIX III

Table with columns for particle name, alpha value, error, and reference. Includes entries for ALPHA XI 0, PIERROU, BERGE, LONDON, MERRILL, DAUBER, BRIDGEWATER, MAYEUR, BALTAY, etc.

WEIGHTED AVERAGE = -0.441 ± 0.078
ERROR SCALED BY 1.3

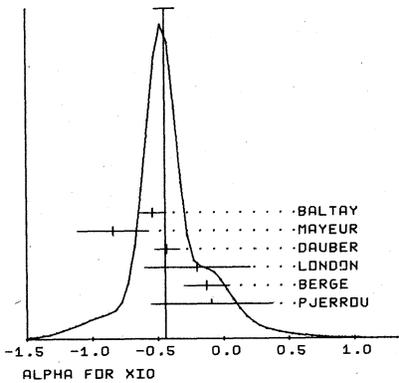


Table with columns for name, value, and CHISQ. Includes entries for BALTAY, MAYEUR, DAUBER, LONDON, BERGE, PIERROU, and overall average values.

Table with columns for name, phi angle, and other parameters. Includes entries for PH I ANGLE, BERGE, MERRILL, DAUBER, BRIDGEWATER, MAYEUR, BALTAY, etc.

REFERENCES FOR XI0

List of references for XI0 decay parameter, including authors like ALVAREZ, JAUNEAU, TICH0, CARMONY, HUBBARD, PIERROU, etc.



24 OMEGA-(1675, JP=3/2+) I=0
QUANTUM NUMBERS ASSIGNED FROM SLB

Table with columns for name, mass (MEV), and other parameters. Includes entries for OMEGA- MASS (MEV), EISENBERG, FRY1, FRY2, ABRAMS, PALMER, SCHULTZ, SCOTTER, SPETH, ABCLV, etc.

24 ANTI-OMEGA+ MASS (MEV)

Table with columns for name, mass (MEV), and other parameters. Includes entry for FIRESTONE 71 HBC.

24 OMEGA- MEAN LIFE (UNITS 10**-10 SEC)

Table with columns for name, mean life, and other parameters. Includes entries for ABRAMS, BARNES, COLLEY, RICHARDSON, ABCLV, SCHULTZ, SCOTTER, etc.

24 OMEGA- PARTIAL DECAY MODES

Table with columns for name, decay mode, and other parameters. Includes entries for OMEGA- INTO LAMBDA K-, OMEGA- INTO XI0 PI-, OMEGA- INTO XI- PI0.

24 OMEGA- BRANCHING RATIOS

Table with columns for name, branching ratio, and other parameters. Includes entries for OMEGA- INTO LAMBDA K-, OMEGA- INTO XI0 PI-, OMEGA- INTO XI- PI0.

24 OMEGA- DECAY PARAMETERS

Table with columns for name, decay parameter, and other parameters. Includes entry for ALPHA FOR OMEGA- TO K- LAMBDA.

REFERENCES FOR OMEGA-

List of references for OMEGA- decay parameters, including authors like EISENBERG, FRY1, FRY2, ABRAMS, BARNES, COLLEY, RICHARDSON, SAMIOS, etc.

Stable Particles

 Ω^- , PARTICLE SEARCHES

ABCLV CO 68 NUC PHYS 84 326	AACHEN+BERLIN+CERN+LONDON IMP. COLL.+VIENNA
ALLISON 68 PRIV. COMM.	JOHN ALLISON (LANCASTER)
PALMER 68 PL 268 323	PALMER, RADDJICIC, RAU, RICHARDSON+ (BNL, SYRA)
SCHULTZ 68 PR 168 1509	SCHULTZ+ (ILL, ARGONNE, NORTHWESTERN, WISC)
SCOTTER 68 PL 268 474	SCOTTER+ (BIRM, GLASGOW, LOIC, MUNICH, OXF)
SPETH 69 PL 298 252	SPETH+ (AACHEN, BERLIN, CERN, LOIC, VIEN)
FRESTON 71 PRL 26 410	*GOLDHABER, LISSAUER, SHELTON, TRILLING (LBL)
ABCLV 73 NP 861 102	AACHEN+BERLIN+CERN+LONDON+VIENNA COLLABOR.
ALVAREZ 73 PR D8 702	LUIS W. ALVAREZ (LBL)
KOCHER 74 PL 518 193	KOCHER, WERNHARD (INNS+VIEN)

PARTICLE SEARCHES

Heavy Leptons

Two types of searches for heavy leptons have been made. The first type uses the fact that we can calculate the transition rate for the decay of a heavy lepton to states involving other leptons, with or without pions. Production mechanisms are also considered to be well known, proceeding via the electromagnetic interaction. BERNARDINI 73 and BARISH 73 are of this type. The second method (BUSHNIN 73) is to look for charged, weakly interacting, long-lived particles.

Several experiments have observed events which could arise from the decay of a heavy lepton. PERL 75 in a SLAC-LBL experiment have observed 64 events of the form

$$e^+ + e^- \rightarrow e^\pm + \mu^\mp + \geq 2 \text{ undetected particles}$$

at or above 4 GeV c.m. energy. Such events have no conventional explanation, but could arise from the production and decay of heavy lepton pairs or charmed boson pairs.

Dimuon events such as those of the Harvard, Pennsylvania, Wisconsin, Fermilab experiment (see BENVENUTI1-5 75 entries in the charmed hadron searches section) could involve heavy lepton production and decay. However, the characteristics of their events do not favor the heavy lepton interpretation.

The Kolar Gold Mines cosmic ray experiment (KRISHNASWAMY 75) reported 5 events which typically have three observed tracks emerging from a common origin about 70 cm from the rock walls of the mine. The authors suggest that these may be massive, long-lived particles, possibly heavy leptons produced by neutrino interactions in the rock walls. Another hypothesis (see DE RUJULA 75) which was put forth to account for their apparent copious production but slow decay is that the observed particle is a

Data Card Listings

For notation, see key at front of Listings.

very long-lived neutral heavy lepton which is a decay product of a charged heavy lepton, itself pair-produced by cosmic rays in the upper atmosphere.

An FNAL neutrino experiment (BENVENUTI 75) roughly simulated the Kolar Gold Mines experiment for the most likely case that the new particle is neutral. Their result appears to contradict (at least fails to confirm) the KRISHNASWAMY 75 result.

Intermediate Bosons (W^{\pm} and ϕ^0)

Experimentally, W bosons could be produced by hadrons, muons, or neutrinos. They can decay into leptons or into hadrons. At present, there is no positive evidence for the existence of W bosons, although there have been several searches. We list upper limits on the mass and on production cross section.

All gauge theories postulate, in addition to the vector boson W, at least one scalar boson (Higgs' scalar), ϕ^0 . The couplings of the Higgs' scalar are model-dependent, and limits on its mass depend on M_W . We list these limits and add in comment cards the assumptions made.

Quarks

Three main techniques have been used in the experimental search for quarks: accelerators, cosmic rays, and searches in stable matter. The most recent limits using each of these techniques are given in the Data Card Listings. Accelerator experiments generally measure quark production cross sections (we quote these in Section C) and differential cross sections (D). Cosmic ray experiments measure quark flux (F), and searches in stable matter measure quark concentration (RHO). Some of the accelerator and cosmic ray experiments have looked for fractionally charged particles, and some have looked for high-mass, low-velocity particles. We give the fractional charge, mass ranges, velocity ranges, and other information in the notes on the right-hand side of the card and below the data cards.

Magnetic Monopoles

Searches have been made for magnetic monopoles using accelerators, cosmic rays, and searches in matter.

In accelerator searches strong magnetic fields

Data Card Listings

For notation, see key at front of Listings.

are used to extract magnetic monopoles trapped in beam dumps. Cross-section limits are calculated using the reaction $pp \rightarrow pp MM$.

Searches in matter have been made on sediments at the bottom of the Atlantic Ocean (FLEISCHER 69) and at the bottom of the Pacific Ocean (KOLM 71). Both experiments used high magnetic fields in attempting to extract monopoles from the matter. ROSS 73 searched a lunar rock sample using an electromagnetic detector which measures a current change induced in a superconducting circuit when traversed by a magnetically charged object.

PRICE 75, a cosmic ray experiment, reported the observation of a magnetic monopole track in a stack of Lexan sheets and emulsion. They report its mass as $m > 200 m_p$ and strength $g = 137e$. However, a more conventional interpretation of this event exists. According to ALVAREZ 75, the data can be well fit by the assumption of a fragmenting heavy nucleus. See also EBERHARD 75 for a discussion of the conflict with other experiments.

Charmed Hadrons

The interpretation that the J/ψ and ψ mesons (see the New Heavy Mesons section of the Meson Data Card Listings) are charm-anticharm bound states has encouraged experimenters to search directly for charmed hadrons. Several experiments have obtained results which strongly support the charm hypothesis.

BENVENUTI 75 (see Charmed Particle references) in a Harvard, Pennsylvania, Wisconsin, Fermilab experiment have observed 70 events of the form

$$\nu(\text{or } \bar{\nu}) + N \rightarrow \mu^- + \mu^+ + \text{anything}.$$

The authors state that these events involve the production and decay of one or more new particles. The observed properties of these events indicate that the new particle could be a charmed hadron but it is not likely to be a heavy lepton or intermediate vector boson.

The 64 events of the form

$$e^+ + e^- \rightarrow e^\pm + \mu^\mp + \geq 2 \text{ undetected particles}$$

reported by PERL 75 (see note on Heavy Leptons above) could arise from the pair production of charmed bosons or heavy leptons.

Two heavy liquid bubble chamber neutrino experiments have found events of the form

$$\nu + N \rightarrow \mu^- + e^+ + \nu^0 + \text{anything}$$

where the ν^0 is a K^0 or Λ . The appearance of a strange particle with an e^+ is an expected signature of a charmed hadron.¹ The CERN Gargamelle experiment (DEDEN 75 and BLIETSCHAU 75) sees two events in which the ν^0 could be K^0 or Λ . An additional four $\mu^- e^+$ events without a ν^0 were observed. The FNAL 15' chamber experiment (VON KROGH 76) sees four $\mu^- e^+ K^0$ events and no $\mu^- e^+$ events without a K^0 . This is a preliminary result. Their full exposure may yield as many as 50 such events.

A BNL neutrino experiment (CAZZOLI 75) observes a three-constraint event

$$\nu + p \rightarrow \mu^- + \Lambda + 3\pi^+ + \pi^-$$

which they interpret either as a large violation of the $\Delta S = \Delta Q$ rule or as a charmed baryon decaying to $\Lambda + 3\pi^+ + \pi^-$.

See the review of Gaillard et al.¹ for a discussion of charmed particle phenomenology with an eye toward searches for these states.

Other Stable Particle Searches

Searches for stable particles (i.e., particles immune to decay via the strong interaction) which do not fit clearly into the above categories are collected here. They are described in the data cards.

Reference

1. M. K. Gaillard, B. J. Lee, and J. L. Rosner, Rev. Mod. Phys. **47**, 277 (1975).

HEAVY LEPTON SEARCHES

M	HEAVY LEPTON MASS LIMITS (GEV)						
M	0	0.6	OR MORE	BACCI	73 ELEC E+E- TO E+E-GAM	1/76*	
M	0	2.2	OR MORE	BACCI	73 ELEC E+E- TO E+E-GAM	1/76*	
M	0	2.0	OR MORE CL=.90	BARISH	73 ASPK + 50,145GEV NEU,NAL	2/74	
M	0	1.4	OR MORE CL=.95	BERNARDIN	73 ASPK E+E-, FRASCATI	2/74	
M	0	1.0	OR MORE CL=.95	BERNARDIN	73 ASPK E+E-, FRASCATI	2/74	
M	N	NONE BETWEEN 0.55 AND 4.5		BUSHNIN	73 CNTR TOGEV P,SERPUKOV	2/74	
M	E	0	2.4	OR MORE CL=.90	EICHTEN	73 HLBC + 1-10 GEV NEU	3/74
M	A	1.8	OR MORE CL=.90	ASRATYAN	74 HLBC -- CERN NEU EPT	11/75*	
M	C	8.4	OR MORE CL=.90	BARISH	74 SPEC + 50,135GEV NEU	7/74*	
M	0	1.15	OR MORE CL=.95	ORITO	74 ASPK -- E+E- FRASCATI	11/75*	
M	P	64 BTM	1.6 AND 2.0	PERL	75 SPEC	1/76*	
M	B	BACCI	73 LOOKS FOR E GAMMA DECAY MODE OF HEAVY ELECTRON. MASS LIMIT			1/76*	
M	B	B	DEPENDS ON COUPLING CONSTANT LAMBDA FOR THIS DECAY. FIRST VALUE			1/76*	
M	B	B	ABOVE IS FOR LAMBDA**2 LT 9*10**5, 2ND IS FOR LAMBDA**2 LT 10**3.			1/76*	
M	C	BARISH	LOOKED FOR DECAY TO (W+ 2NEUTRINS). ASSUMES BR.FRAC.=.3			11/75*	
M	D	BERNARDINI	73 FIRST VALUE ASSUMES UNIVERSAL COUPLING TO ORDINARY			2/74	
M	D	LEPTONS. SECOND VALUE ALSO ASSUMES COUPLING TO HADRONS.				2/74	
M	N	BUSHNIN	73 MASSES ASSUME MEAN LIFE ABOVE 7E-10 AND 3E-8 RESPECTIVLY.			2/74	
M	N	N	CALCULATED FROM CROSS SEC(C BELOW) AND 30 GEV MUON PAIR PROD. DATA.			2/74	
M	E	EICHTEN	73 LIMIT IS FOR GEORGI-GLASHOW TYPE HVY.LEP. M+ PRODUCED IN			2/76*	
M	E	E	NEU NUCL --> M+ HADRONS ASSUMING 15 PERCENT DECAY TO E+ NEU.			2/76*	
M	A	ASRATYAN	74 LIMIT IS FOR HEAVY LEPTON L+ PRODUCED IN THE REACTION			2/76*	
M	A	ANTINEU NUCL --> L+ HADRONS.				2/76*	
M	O	ORITO	74 LOOKED FOR M+- PAIRS GIVING MU-E PAIRS. MASS LIMIT REFERS			3/74	
M	O	O	TO HEAVY ELECTRON, HEAVY MUON OR HEAVY LEPTON WITH OWN L NUMBER.			3/74	
M	O	O	COUPLING TO HADRONS ASSUMED FROM THEORETICAL MODELS.			3/74	
M	P	PERL	75 EVENTS ARE EK E- TO E+ MU-+ AND TWO OR MORE MISSING			1/76*	
M	P	P	PARTICLES DONE AT SLAC (SPEAR). AUTHORS CLAIM NO CONVENTIONAL			1/76*	
M	P	P	EXPLANATION FOR THESE EVENTS. TOTAL CM ENERGY 3-7.5 GEV.			1/76*	

Stable Particles
PARTICLE SEARCHES

Data Card Listings

For notation, see key at front of Listings.

HEAVY LEPTON PRODUCTION CROSS SECTION (E+ E-) (UNITS 10**=-35 CM**2)
PERL 75 SPEC M=1.6-2.0 GEV
MORE MISSING PARTICLES. CS IS FOR EVENTS ARE E-GEV AND THETA=50-130DEG.
CROSS SECTION RISES FROM 5X10**=-36 AT E=4 TO ABOVE MAX. THEN DROPS
TO 6E-36 AT E=7.5. AUTHORS SAY THESE EVENTS HAVE NO CONVENTIONAL
EXPLANATION. SUGGEST HEAVY LEPTON OR CHARMED HADRON, M=1.6-2.0GEV.

QUARK SEARCHES

QUARK PRODUCTION CROSS SECT. FROM ACCELERATOR EXPTS (CM**2)
Y 0 3.2E-39 OR LESS CL=90 ALLBAY 69 CNTR Q=-1/3 M=2GEV
Y 0 1.0E-35 OR LESS CL=90 ALLBAY 69 CNTR Q=+1/3 M=2GEV
Y 0 1.0E-35 OR LESS CL=90 ALLBAY 69 CNTR Q=+2/3 M=2GEV
Y 0 4. E-37 OR LESS CL=90 ANTIPOV1 69 CNTR Q=-2/3 M=0-5GEV
Y 0 3. E-37 OR LESS CL=90 ANTIPOV2 69 CNTR Q=-1/3 M=4-5GEV
Y 0 1. E-34 OR LESS CL=90 ANTIPOV 71 CNTR Q=-4/3 M=4+GEV
Y 0 3. E-36 OR LESS BOTT-BODE 72 CNTR Q=+1/3 M=0-22GEV
Y 0 6. E-34 OR LESS BOTT-BODE 72 CNTR Q=+2/3 M=0-13GEV
Y 0 1. E-32 OR LESS CL=90 ALPER 73 SPEC Q= 4/3 M=4-24 GEV
Y 0 1. E-35 OR LESS LEIPUNER 73 CNTR Q= 1/3 M=0-12GEV
Y 0 1. E-35 OR LESS LEIPUNER 73 CNTR Q= 2/3 M=0-12GEV
Y 0 5. E-31 OR LESS LEIPUNER 73 CNTR Q= 4/3 M=0-12GEV
Y ALLBAY 69 IS A CERN 27 GEV P+BE EXPT. STUDYS MASSES 0-2.7GEV
Y ASSUMING NN=NNQQ. CROSS SECTIONS ASSUME ISOTROPIC PROD. IN CM.
Y CROSS SECTIONS AT 2GEV ARE GIVEN HERE. SEE FIG-9 FOR MASS DEPEN.
ANTIPOV1 69 IS A SERPUKHOV 70 GEV P EXPT. MASS LIMIT FROM NN=NNQQ.
ANTIPOV1 69 AND ANTIPOV2 69 ARE SERPUKHOV 70GEV P EXPTS. ANTIPOV2
GIVES RESULTS FOR M=2-5GEV ASSUMING NN--NNQQ, HADRONIC OR LEPTONIC
QUARKS. WE QUOTE TYPICAL VALUES.
ANTIPOV 71 IS A SERPUKHOV 70 GEV P+AL EXPT. STUDIES DIQUARK MASSES
1.9-4.4GEV. WE SHOW 4GEV VALUES. SEE THEIR FIG-2 FOR MASS DEPEN.
BOTT-BODENHAUSEN 72 IS A CERN ISR 26+26 GEV P+P EXPERIMENT.
ALPER 73 IS CERN ISR 26+26 GEV P+P EXPT. ASSUMES ISOTROPIC C.M.
P PRODUCTION. SENSITIVE TO ANY Q2/3.
LEIPUNER 73 IS AN AL 300 GEV P EXPERIMENT.

REFERENCES FOR HEAVY LEPTON SEARCHES
+GRACHEV,KHODYREV,KUBAROVSKY+ (SERP)
+PARTI,PEVEDI,SALVINI,STELLA+ (ROMANFRAS)
+BARTLETT,BUCHHOLZ,HUMPHREY+ (CIT+FNAL)
+BOLLINI,BRUNINI+ (CERN+BGNA+FRAS)
ALLES-BORELLI,BERNARDINI,BOLLINI+ (CERN)
DUNAYTZEV,GOLOVKIN,KUBAROVSKY+ (SERP)
GOLOVKIN,GRACHEV,SHODYREV+ (SERP)
+ODEDAN+IAACH+BELG+CERN+EPOL+MILA+LALO+LOUCI
+GERSHTEIN,KAPTANOV,KUBANTZEV,LAPINA+ (SERP)
+BARTLETT,BUCHHOLZ,HERRITT+ (CIT+FNAL)
+FRISCH,SIOCHET,BOYNDON,MERMOD+ (EP1+PRIN)
+VISENTIN,CERARDINI,CONVERSI+ (FRAS+ROMA)
BENVENUTI,CLINE,FORD+ (HARV+PENN+WISC)
BINT INGER,CURRY+ (EP1+HARV+PENN+WISC)
KRISHNASWAMY,MENDH+ (SRM+HARV+SARAK)
DE RUJULA,GEORGI,GLASHOW
+ABRAMS,BOYARSKI,BREIDENBACH+ (LBL+SLAC)

QUARK FLUX FROM COSMIC RAY EXPERIMENTS (NUMBER/CM**2-SR-SEC)
D 0 3. E-10 OR LESS CHUPKA 68 CNTR M=5GEV OR MORE
D 0 2.4E-8 OR LESS KASHA 68 CNTR V=5-9C M=2-15GEV UP
D 0 5. E-11 OR LESS CL=90 FUKUSHIMA 69 CNTR Q=1/3 SEA LEVEL
D 0 7.5E-10 OR LESS CL=90 FUKUSHIMA 69 CNTR Q=2/3 SEA LEVEL
D 1 EVENT CLAIMED MCCUSKER 69 CC Q=2/3
D B 0 1. E-10 OR LESS CL=90 BOHM 72 CNTR Q=1/3
D B 0 1. E-10 OR LESS CL=90 BOHM 72 CNTR Q=2/3
D H 0 1.7E-8 OR LESS CL=90 HICKS 73 CNTR Q= 1/3
D H 0 1.7E-8 OR LESS CL=90 HICKS 73 CNTR Q= 2/3
D C BJORNBOE 68 LOOKED FOR DELAYED PARTICLES AFTER AIR SHOWERS.
D M MCCUSKER 69 CLAIMS 1 CANDIDATE. LATER SIMILAR EXPTS. SEE NONE.
D B BOHM 72 IS FLUX IN 10**14 TO 10**15 EV AIR SHOWERS
D H HICKS 73 LOCKED AT LARGE ZENITH ANGLES, THUS USING THE ATMOSPHERE
AS AN EXTENDED FILTER FOR HADRONIC QUARKS. THEIR SEARCH PUTS AN
UPPER LIMIT ON LEPTONIC QUARK FLUX IN COSMIC RAYS.

INTERMEDIATE BOSON SEARCHES

W BOSON MASS LIMITS (GEV)
M B 1.7 OR MORE CL=99 BERNARDINI 65 HYBR + NEU N, CERN
M B 0 2.0 OR MORE CL=90 BURNS 65 OSPK + NEU N, BNL
M C 0 3.8 OR MORE CL=90 BARISH 73 ASPK + W TO LEP+NEU=.2
M C 0 4.5 OR MORE CL=90 BARISH 73 ASPK + W TO LEP+NEU=.5
M E 0 4.0 OR MORE CL=95 BERGESON 73 ELEC
M B LOOKED FOR (NEU N) TO (W+ MU- N), W+ TO (MU+ NEU, E+ NEU, OR HOPNS)
M C BARISH 73 LOOKED FOR (NEU N) TO (W+ MU- N), W+ TO (MU+ NEU) AT NAL.
M C RESULT GIVEN FOR THREE ASSUMED BR.FRACS. W+ TO (LEPTON NEU)/ALL.
M E BERGESON 73 LOOKED AT ENERGY DISTR OF NEU-INDUCED MUON FLUX UNDER-
GROUND. SCALE INVARIANCE OF THE INELASTIC STRUCT FN ASSUMED.

QUARK CONCENTRATION IN MATTER (QUARKS PER NUCLEON)
RHO 0 3. E-29 OR LESS CHUPKA 66 CNTR SEAWATER
RHO 0 1. E-16 OR LESS GALLINARO 66 CNTR GRAPHITE
RHO 0 4. E-19 OR LESS STOVER 67 CNTR IRON
RHO 0 1. E-20 OR LESS RANK 68 CNTR OIL
RHO 0 1. E-24 OR LESS COOK 69 CNTR SEAWATER
RHO 0 1. E-23 OR LESS COOK 69 CNTR ROCK SAMPLES

REFERENCES FOR INTERMEDIATE BOSON SEARCHES
+BIENLEIN,BOHM,DARDEL,FAISSNER+ (CERN)
+GOUL LANGS,HYMAN,LEDERMAN,LEE+ (COLU+BNL)
ANKENBRANDT,LARSEN,LEIPUNER+ (BNL+YALE)
+BARTLETT,BUCHHOLZ,HUMPHREY+ (CIT+FNAL)
+CASSIDAY,HENDRICKS (UTAH)
+D'ANGELO,GATTO,PAULUZI (RCMA)

REFERENCES FOR QUARK SEARCHES
+EADES,LEDERMAN,LEE,TING (COLU)
+SCHIFFER,STEVENS (ANL)
GALLINARO,MORPURGO (GENO)
+MORAN,TRISCHKA (SVRA)
+DAMGARD,HANSEN,CHATTERJEE+ (BOHR+BERN)
FRANZINI,SULMAN (COLU)
+STEFANSKI (BNL+YALE)
D.W.RANK (MICH)
+BIANCHINI,DIDENDS,DORINSON,HARTUNG+ (CERN)
+KARPOV,KHROMOV,LANDSBERG,LAPSHIN+ (SERP)
+BOLOTOV,DEVISHEV,DEVISHEVA,ISAKOV+ (SERP)
+DEPASQUALI,FRAUENFELDER,PEACOCK+ (ILL)
FUKUSHIMA,KIFUNE,KONDO,KOSHIBA+ (TOKY)
MCCUSKER,CAIRNS (SYDNEY)
+KACHANDU,KUTJIN,LANDSBERG,LEBDEEV+ (SERP)
+DIEMONT,FAISSNER,FASSOLD,KRISOR+ (AACH)
+CALDWELL,BARN,GRUBB,PEAKS+ (CERN+MPHM)
BRITISH-SCANDINAVIAN COLLABORATION
+PLINT,STANDIL (MANI)
+LARSEN,SESSOMS,SMITH,WILLIAMS+ (BNL+YALE)
JOVANOVIICH+ (MANI+IACH+CERN+GENO+HARV+TORTI)

Data Card Listings
For notation, see key at front of Listings.

Stable Particles
PARTICLE SEARCHES

MAGNETIC MONOPOLE SEARCHES

CHARMED HADRON SEARCHES

C MONOPOLE PROD. CROSS SECTION - ACCELERATOR EXP. (CM**2)/NUCLEON
A 0 1 E=40 OR LESS CL=.95 AMALDI 63 EMUL M=0 TO 5 GEV 12/75*

CE CHARMED HADRON PRODUCTION CROSS SECTION (E+ E-) (UNITS 10**-35 CM**2)
CE B 0 18. OR LESS CL=.90 BOYARSKI 75 SPEC K+ PI+, K+ PI- 2/76*

REFERENCES FOR MAGNETIC MONOPOLE SEARCHES
+BARONI, MANFREDINI, BRADNER+ (ROMAUCS+D+ERN)

REFERENCES FOR CHARMED HADRON SEARCHES
AUBERT 75 PRL 35 416 +BECKER, BIGGS, BURGER, CHEN+ (MIT+BNL)

Stable Particles
PARTICLE SEARCHES

Data Card Listings

For notation, see key at front of Listings.

OTHER STABLE PARTICLE SEARCHES

SEARCHES FOR STABLE PARTICLES (I.E. PARTICLES IMMUNE TO DECAY VIA THE STRONG INTERACTION) NOT AMONG ABOVE SEARCHES ARE LISTED HERE.

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D HEAVY PARTICLE PRODUCTION DIFFERENTIAL CROSS SECTION (CM**2/SR-GEV)
D A 0 1. E-31 OR LESS CL=.90 APPEL 74 CNTR +- M=3.2-7.2 GEV 2/76*
D J 0 8. E-35 OR LESS CL=.90 JOVANDVIC 75 CNTR +- M=15-26 GEV 2/76*
D J 0 1.5E-34 OR LESS CL=.90 JOVANDVIC 75 CNTR Q=+-2, M=3-10 GEV 2/76*
D J 0 6. E-35 OR LESS CL=.90 JOVANDVIC 75 CNTR Q=+-2, M=10-26 GEV 2/76*
D A APPEL 74 IS NAL 300 GEV P+W EXPERIMENT. STUDIES FORWARD PRODUCTION 2/76*
D A OF HEAVY (UP TO 24 GEV) CHARGED PARTICLES WITH MOMENTA 24-200GEV(-) 2/76*
D A AND 40-150GEV (+CHG). ABOVE TYPICAL VALUE IS FOR 75 GEV. 2/76*
D J JOVANDVIC 75 IS A CERN ISR 26+26 AND 15+15 GEV P+P EXPERIMENT. 2/76*
D J FIG.4 COVERS RANGES Q=1/3 TO 2 AND M=3 TO 26 GEV. 2/76*

F HEAVY PARTICLE FLUX IN COSMIC RAYS (NUMBER/CM**2-SEC-SR) 1/76*
F Y 5 6 E-9 OR MORE YOCK 74 CNTR M GT 6 GEV 1/76*
F Y YOCK 74 EVENTS COULD BE TRITONS. 1/76*

C LIGHT (BETWEEN MU AND E MASSES) PARTICLE MASS (UNITS-ELECTRON MASSES)
C 0 NONE BETWEEN 2 AND 13 BLAGOV 75 CNTR SPINOR,TAU=2E-10SEC 2/76*
C 0 NONE BETWEEN 2 AND 10.6 BLAGOV 75 CNTR SCALAR,TAU=2E-10SEC 2/76*

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REFERENCES FOR OTHER STABLE PARTICLE SEARCHES

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APPEL 74 PRL 32 428 +BOURQUIN,GAINES,LEDERMAN,PAAR+ (COLU+NAL)
YOCK 74 NP 876 175 P.C.M.YOCK (UNIV OF AUCKLAND)
BLAGOV 75 YAD.FIZ. 21,300 +KOMAR,MURASHOVA,SYREISHCHIKOVA+ (LEBD)
JOVANDVI 75 PL 56B 105 JOVANDVIC+ (MANI+AACH+CERN+GEND+HARV+TORI)

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Mesons

$\rho(770)$

Data Card Listings

For notation, see key at front of Listings.

WEIGHTED AVERAGE = 150.3 ± 2.7
ERROR SCALED BY 1.4

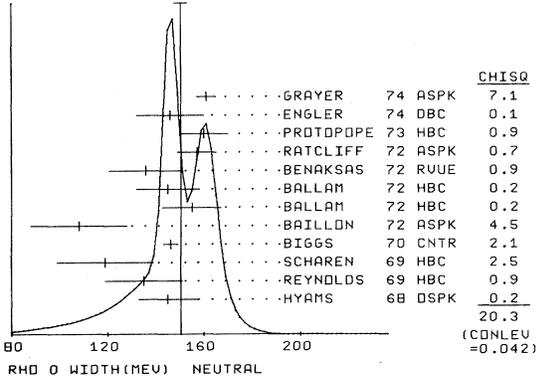


Table with 2 columns: RHO PARTIAL DECAY MODES and DECAY MASSES. It lists various decay modes like RHO INTO 2PI, RHO INTO 4PI, etc., and their corresponding decay masses.

Table with 2 columns: RHO BRANCHING RATIOS and (P21)/(P1). It lists branching ratios for various decay modes and the ratio of partial widths.

Note on $\rho^0 \rightarrow e^+e^-$

Extraction of a ratio for $\rho^0 \rightarrow e^+e^-$ is complicated by interference with ω decay. In photoproduction, $\gamma A \rightarrow e^+e^-A$, there is substantial interference between the allowed $(\rho^0, \omega) \rightarrow e^+e^-$ decays. The interference in the colliding-beam reaction $e^+e^- \rightarrow \pi^+\pi^-$ is due to G-parity-violating mixing of the overlapping ρ^0 and ω resonances; it alters the results for the rate $\Gamma(\rho^0 \rightarrow e^+e^-)$ only by a small amount. Therefore at present we average only the values from the $e^+e^- \rightarrow \pi^+\pi^-$ experiments.

Table with 2 columns: RHO INTO e^+e^- / (PI+PI-) (UNITS 10**+4) and (P4)/(P1). It lists experimental results for the ratio of partial widths.

Table with 2 columns: RHO INTO (PI+PI-)/PI (UNITS 10**+4) and (P61)/(P1). It lists experimental results for the ratio of partial widths.

REFERENCES FOR RHO

A large table listing references for the rho meson, including authors, journals, and years.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho(770)$, $\omega(783)$

Table listing particle data for various mesons including names like ALVENSLE, BATON, BIGGS, BINGHAM, GALLOWAY, ABRAMS, BLOODWORTH, DEERY, BALLON, BALLAM, BASDEVAN, BENAKSAS, DRIVER, EISENBERG, GRAY, GRAY, JACOBS, RATCLIFF, TAKAHASHI, BYERLY, CHARLES, GLADDING, HYAMS, PROTOPOP, CARROLL, ENGLER, ESTABROO, GOBBI, GRAYER, HABER, NORDBERG, SPITAL, ROOS.

$\omega(783)$

Table for $\omega(783)$ listing OMEGA MASS (MEV) and OMEGA FULL WIDTH (MEV) with columns for mass, width, and various decay channels like ARMENTERO, ALFF, BUSHBECK, KRAEMER, MILLER D, JAMES, BALTAY, BARASH, KEY, BIZZARRI, ABRAMOVIC, ATHERTON, BIGGS, etc.

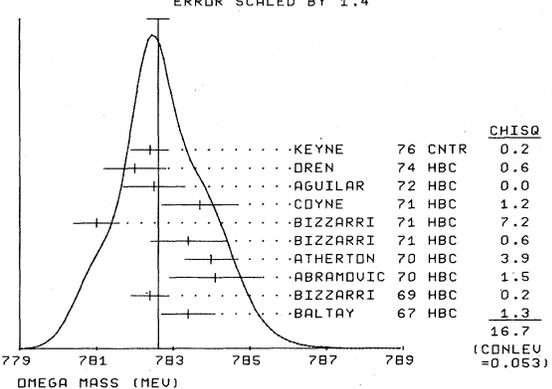


Table for OMEGA FULL WIDTH (MEV) listing various decay channels and their widths, including ARMENTERO, MILLER D, BARASH, UNFOLDED BY COYNE 71, ABRAMOVIC, ATHERTON, BIZZARRI, COYNE, AGUILAR, BENAKSAS, BROWN, BINNIE, KEVNE, etc.

Table for OMEGA PARTIAL DECAY MODES listing decay channels like OMEGA INTO PI+ PI- PI0, OMEGA INTO PI0 GAMMA, OMEGA INTO 2PI0 GAMMA, OMEGA INTO ETA GAMMA, OMEGA INTO E+ E-, OMEGA INTO MU+ MU-, OMEGA INTO ETA PI0, OMEGA INTO 3 GAMMA, OMEGA INTO PI0 MU+ ML-.

Table for FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS showing the matrix derived from the error matrix for the fitted partial decay mode branching fractions, P1, P2, P3.

Table for OMEGA BRANCHING RATIOS listing various decay channels and their branching ratios, including ARMENTERO, BUSHBECK, KRAEMER, ALFF-STEI, DI GIUGNO, FLATTE, JAMES, BALTAY, BARASH, AGUILAR, etc.

Table for OMEGA BRANCHING RATIOS (continued) listing various decay channels and their branching ratios, including ARMENTERO, BUSHBECK, KRAEMER, ALFF-STEI, DI GIUGNO, FLATTE, JAMES, BALTAY, BARASH, AGUILAR, etc.

Table for OMEGA BRANCHING RATIOS (continued) listing various decay channels and their branching ratios, including ARMENTERO, BUSHBECK, KRAEMER, ALFF-STEI, DI GIUGNO, FLATTE, JAMES, BALTAY, BARASH, AGUILAR, etc.

Mesons

$\omega(783)$, $M(940)$

Data Card Listings

For notation, see key at front of Listings.

Table listing meson properties for Omega and M(940) particles, including production methods, decay modes, and branching ratios.

REFERENCES FOR OMEGA

Table of references for Omega mesons, listing authors, journals, and years.

Table listing meson properties for M(940) particles, including production methods, decay modes, and branching ratios.

M(940) -> MM

66 M(940)

EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

66 M(940) MASS (MEV)

Table showing mass measurements for M(940) in MeV.

66 M(940) WIDTH (MEV)

Table showing width measurements for M(940) in MeV.

66 M(940) BRANCHING RATIOS

Table showing branching ratios for M(940).

REFERENCES FOR M(940)

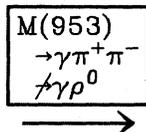
Table of references for M(940) mesons, listing authors, journals, and years.

Data Card Listings

For notation, see key at front of Listings.

Mesons

M(953), η'(958)



59 M(953, JPC= 0-+)
WHILE MASS AND WIDTH ARE CONSISTENT WITH ETA PRIME(958), THE PI+ PI- GAMMA DECAY SHOWS NO RHO SIGNAL. POSSIBLY SEEN IN MMS. OMITTED FROM TABLE.

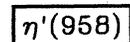
Table with 5 columns: M, M, M, M, M. Headers: 59 M MASS (MEV), 68 953.0, 2.0, AGUILAR 70 HBC, 3.9-4.6K-P, P K-M 1/71. Includes text: MISSING MASS SPECTRUM SHOWED THIS PEAK AT 953.4 INSTEAD OF ETA PRIME (958). PEAK LISTED UNDER M BECAUSE OF MASS COINCIDENCE. THE 1.5 MEV ERROR MAY BE UNDERESTIMATED BY A FACTOR OF 2 (SEE BRODY 72, TABLE III). OBSERVED PEAK COULD THEN WELL CORRESPOND TO ETA PRIME.

Table with 5 columns: W, M, M, M, M. Headers: 59 M WIDTH (MEV), 68 (10.0) OR LESS CL=.95, AGUILAR 70 HBC, 3.9-4.6K-P, P K-M 1/71. Includes text: (15.) OR LESS MAGLICH 71 MMS 3.8 P D, HE3 X0 2/72

Table with 5 columns: P1, P2, P4, P5, P5. Headers: 59 M PARTIAL DECAY MODES, DECAY MASSES. Includes text: M INTO PI+ PI- GAMMA 139+ 139+ 0, M INTO RHO GAMMA 773+ 0, M INTO PI+ PI- ETA 139+ 139+ 548, M INTO PIO ETA 134+ 548, M INTO PI+ PI- PIO 139+ 139+ 134

Table with 5 columns: R1, R2, R3, R4, R4. Headers: 59 M BRANCHING RATIOS, (P2)/(P1), (P1)/(P3), (P5), (P4N). Includes text: M INTO (RHO GAMMA)/(ALL PI+ PI- GAMMA) 70 HBC 3.9-4.6K-P, P K-M 1/71, M INTO (PI+ PI- GAMMA)/(PI+ PI- ETA NEUTR.) 70 HBC 3.9-4.6K-P, P K-M 1/71, M INTO (PI+ PI- PIO)/TOTAL 70 HBC 3.9-4.6K-P, P K-M 1/71, M INTO (PIO ETA NEUTR.)/TOTAL 70 HBC 3.9-4.6K-P, P K-M 1/71

Table with 5 columns: AGUILAR, MAGLICH, ROSNER, AGUILAR, BRODY, GRIGORIAN. Headers: REFERENCES FOR M. Includes text: AGUILAR-BENITEZ, BASSANO, SAMIOS, BARNES+(BNL), *OCSTENS, BRODY, CVIJANOVICH+(RUTG+PENN+UPNJ), J.L. ROSNER, E.W. COLGLAZIER (MINN+CIT), AGUILAR-BENITEZ, CHUNG, EISNER, SAMIOS (BNL), *GROVES, NOREM, CVIJANOVICH, *(PENN+RUTG+UPNJ), GRIGORIAN, LADAGE, MELLEMA, RUDNICK,+(UCLA)



2 ETA PRIME(958, JPC=0-+) I=0

Note on the J^P Assignment of η'(958)

From the Dalitz plot analyses of the η' → πππ and η' → π+π-γ decays, and from the observation of a η' → γγ decay mode, all assignments except J^{PC} = 0-+ and 2-+ are excluded. The Dalitz plot analyses favor spin 0, but cannot rule out spin 2. The indication of anisotropy in the decay of very forward produced η' (KALBFLEISCH 73) has not been confirmed by BALTAY 74, thus favoring strongly spin 0.

Table with 5 columns: M, M, M, M, M. Headers: 2 ETA PRIME MASS (MEV), ONLY EXPERIMENTS GIVING ERROR LESS THAN 2 MEV KEPT FOR AVERAGING 12/75*. Includes text: 85 (957.0) DAUBER 64 HBC 1.95 K-P 6/66, (958.0) (1.0) KALBFLEISCH 64 SUPERSEDED BY RITTENBERG 69 6/66, (957.0) (3.0) BADIER 65 HBC 3.0 K-P, (960.0) (2.0) TRILLING 65 HBC 3.65 PI+ P 12/75*, (955.0) (10.0) COMN 66 DBC 3.3 PI+D 6/66, (959.0) (3.0) LONDON 66 HBC 2.2 K-P 6/66, (960.0) (5.0) MOTT 69 HBC 4.1-5.5 K-P 7/69

Table with 5 columns: M, M, M, M, M. Headers: 957.0, 1.0, RITTENBERG 69 HBC, 1.7-2.7 K-P 9/69. Includes text: 0 (956.0) (2.0) AGUILAR 70 HBC 3.9-4.6K-P 12/75*, 3415 956.1 1.1 BASILEI 71 CNTR 1.6 PI- P,N XO 11/71, 535 957.4 1.4 BASILEI 71 CNTR 1.6 PI- P,N XO 11/71, 1414 958.2 0.5 DANBURG 73 HBC 2.2 K-P, LAM XO 2/74, 400 958. 1. JACOBS 73 HBC 2.9 K-P, LAM XO 1/74, 957.46 0.33 DUANE 74 MMS PI- P,N MM 1/74, AVG 957.57 0.25 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0), STUDENT 957.57 0.28 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Table with 5 columns: W, W, W, W, W. Headers: 2 ETA PRIME WIDTH (MEV), 85 (4.0) OR LESS DAUBER 64 HBC 1.95 K-P, 3415 (8.1) OR LESS CL=.90 BASILEI 71 CNTR 1.6 PI- P,N XO 11/71, 514 (4.7) OR LESS CL=.95 DANBURG 73 HBC 2.2 K-P, LAM XO 2/74, (0.8) OR LESS CL=.95 DUANE 74 MMS PI- P,N MM 1/74

Table with 5 columns: P1, P2, P3, P4, P10, P11, P12, P13, P14, P15, P16, P17, P18, P19, P20. Headers: 2 ETA PRIME PARTIAL DECAY MODES, DECAY MASSES. Includes text: ETA PRIME INTO PI+ PI- ETA 139+ 139+ 548, (N) ETAS DECAY INTO ALL NEUTRALS, (C) ETAS DECAY CHARGED, INTO PIO PIO ETA 134+ 134+ 548, (N) ETAS DECAY INTO ALL NEUTRALS, (C) ETAS DECAY CHARGED, INTO PI+ PI- GAMMA (INCLUDING RHO GAMMA) 139+ 139+ 0, INTO GAMMA GAMMA 0+ 0, INTO RHO GAMMA 0+ 773, INTO PI+ PI- E+ E- 139+ 139+ .5+ .5, INTO 2 PI 139+ 139, INTO 3 PI 139+ 139+ 134, INTO 4 PI 139+ 139+ 139+ 139, INTO 5 PI, INTO 6 PI 134+ .5+ .5, INTO ETA E+ E- (VIOLATES C IN BORN APPROX.) 548+ .5+ .5, INTO PIO RHO G O (VIOLATES C) 134+ 772, INTO PIO CMEGA (VIOLATES C) 134+ 783, INTO GAMMA CMEGA 0+ 783

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_i δP_i), while the off-diagonal elements are the normalized correlation coefficients (δP_i δP_j) / (δP_i δP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 5 columns: P 1, P 2, P 3, P 4. Includes text: P 1 .4352+- .0174, P 2 -.5783 .2407+- .0198, P 3 -.3541 -.5478 .3039+- .0171, P 4 .0314 -.1430 -.0218 .0202+- .0027

2 ETA PRIME BRANCHING RATIOS

Note on η'(958) Branching Fractions

In our calculation of the branching fractions of the η'(958), we use the decay modes ηππ (including ηπ⁰π⁰), ρ⁰γ, and γγ. It is assumed that the rate η → neutrals is 71.1%.

In the fit we do not use the constraint

R = Γ(η' → ηπ⁺π⁻) / Γ(η' → ηπ⁰π⁰) = 2

from I-spin conservation. The result of the fit is in agreement with it: R = 1.8 ± 0.2.

Table with 5 columns: R1, R1, R1, R1, R2, R2, R2, R2, R2. Headers: ETA PRIME INTO (PI+ PI-ETA (NEUTRAL DEC.))/TOTAL (PIN), KALBFLEZ 64 HBC 2.7 K-P 10/66, KALBFLEZ 64 SUPERSEDED BY RITTENBERG 69 10/66, 281 0.314 0.026 RITTENBERG 69 HBC 1.7-2.7 K-P 9/69, FIT 0.309 0.012 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0), INTO (PI+ PI- NEUTRALS) / TOTAL (PIN+P2C), 33 0.35 0.06 BADIER 65 HBC 3.0 K-P 10/66, 39 0.4 0.1 LONDON 66 HBC 2.2 K-P 10/66, AVG 0.363 0.051 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0), STUDENT 0.363 0.056 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT, FIT 0.379 0.010 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

Mesons
eta(958), delta(970)

Data Card Listings

For notation, see key at front of Listings.

Table of particle data for eta(958) and delta(970) mesons, including decay rates, branching ratios, and fit parameters.

Table of particle data for eta(970) meson, including decay rates, branching ratios, and fit parameters.

delta(970) -> eta pi pi

Under this entry, we list three types of I=1 peaks near KK threshold:
1) Missing-mass peaks, mostly controversial.
2) eta pi decays, peaking slightly below KK threshold. This defines I^G = 1^- and J^P = Normal.
3) Threshold enhancements in the (KK)^+ system with I=1. The Q value is low and J^P therefore probably 0^+.

Data Card Listings

For notation, see key at front of Listings.

Mesons

delta(970), H(990), S*(993)

In listing these types of peaks together under a common entry we do not imply that they are necessarily all related. However, the K-K bar threshold enhancement may be due to a virtual bound state that could also be responsible for the eta pi peaks (ASTIER 67). In a coupled-channel analysis of the most significant eta pi and K-K bar data, MORGAN 75 shows that the delta(970) probably is strongly coupled to both channels. It then constitutes the isovector member of the 0+ nonet. MORGAN 75 gets a width much larger than what is observed in the eta pi decays because of the strong K-K bar coupling.

36 DELTA(970) MASS (MEV)

Table with columns for mass (MEV), missing mass experiments, and references. Includes entries for K BAR ONLY, SCAT. LENGTH, and AVERAGE ERROR.

36 DELTA(970) WIDTH (MEV)

Table with columns for width (MEV), missing mass experiments, and references. Includes entries for K BAR ONLY, SCAT. LENGTH, and AVERAGE ERROR.

36 DELTA(970) PARTIAL DECAY MODES

Table with columns for decay modes and decay masses.

36 DELTA(970) BRANCHING RATIOS

Table with columns for branching ratios and references.

REFERENCES FOR DELTA(970)

Table of references for Delta(970) from various experiments and authors like TURKOT, ARMENTEROS, BARASH, etc.

H(990)

35 H(990, JPG=A -) I=0
66 HAS EVIDENCE OF BENSON 66 HAS DISAPPEARED AFTER RE-ANALYSIS (CHAUDHARY 70). NO SIGNIFICANT OTHER EVIDENCE HAS BEEN PUBLISHED. OMITTED FROM TABLE.

REFERENCES FOR H

Table of references for H(990) from various experiments and authors like BARTSCH, GOLDHABER, BENSON, etc.

S*(993)

3 S*(993, JPG=0++) I=0
UNDER THIS ENTRY WE LIST PARAMETERS OF THE POLE IN THE ISOSCALAR S-WAVE. FOR A MINI-REVIEW SEE UNDER EPSILON. POSSIBLE EVIDENCE OF D-WAVE PI-PI INTERACTIONS IN THIS REGION IS LISTED SEPARATELY UNDER ETA N(1080).

FOR EARLY WORK USING BREIT-WIGNER OR SCATTERING LENGTH PARAMETRIZATION IN FITS TO THE (K BAR) MASS SPECTRUM SEE REFERENCE SECTION AND OUR 1972 EDITION.

3 REAL PART OF THE S* POLE POSITION (MEV)

Table with columns for real part of the S* pole position (MEV) and references.

Mesons
S*(993), phi(1020)

Data Card Listings

For notation, see key at front of Listings.

3 NEGATIVE IMAG. PART OF THE S* PCLE POSITION (MEV)
CORRESPONDS TO HALF-WIDTH, NOT FULL WIDTH.
W H 40. 40. 60. HOANG 69 DSPK 4. PI-P,K,S KS N 1/73
W H 30. 30. 70. HOANG 69 DSPK 5. PI-P,K,S KS N 1/73
W B (13.) BEUSCH 70 DSPK 4.6 PI-P 1/73
W P 27. 8. PROTOPOPE 73 HBC 7. PI+ P 1/74
W P ANOTHER SOLUTION HAS 52 MEV AND NO EPSILON PCLE.
W S 24. 7. BINNIE 73 CNTR PI- P,S* N 1/74
W S SEE NOTE S ABOVE
W (5.) ESTABROOK 73 ASPK 17 PI-P,PI+PI-N 12/75*
W (16.) GRAYER 73 ASPK 17 PI-P,PI+PI-N 12/75*
W 15. (5.) HYAMS 73 ASPK 0 17 PI-P,N,PI+PI- 1/74
W A (19.) (3.) FUJII 75 RVUE 17 PI-P,PI+PI-N 12/75*
W A S* AMPLITUDES PARAMETRIZED IN TERMS OF POLE POSITIONS USING
W A HYAMS 73 PHASE SHIFTS
W AVG 20.0 3.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W STUDENT 20.1 4.2 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

3 S* WIDTH (MEV)
(180.) APPROX. MORGAN 75 RVUE 12/75*

REFERENCES FOR S*

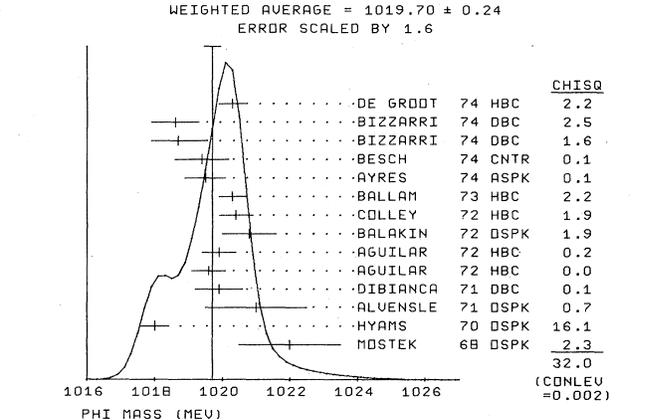
WANG 61 JETP 13 323 WANG TSU-TSENG, VEKSLER, VRANA,+ (JINR)
BIGI 62 CERN CONF 247 A BIGI'S BRANDT, R CARRARA + (CERN)
BINGHAM 62 CERN CONF 240 H BINGHAM, B BLOCH + (EPOL+CERN)
ERMIN 62 PR 163 1375 *HARDY+HESSE+KIRZMILLER (LRL)
ERMIN, MOYER, MARON, WALKER, WAMPLER (WISCONSIN)
BALTAY 64 DUBNA CONF 1 409 *BARNES, JENNELL, FLAMINIO, GOLDBERG,+ (BNL)
BARMIN 64 DUBNA CONF 1 433 BARMIN, DOLGICENKO, YEROFEEV, KRESTINI+ (ITEP)
CRENNELL 66 PRL 16 1025 CRENNELL, KALBFLEISCH, LAI, SCARR, SCHU+ (BNL)
HESS 66 PRL 17 1109 *CAHL+HARDY+KIRZ-MILLER (LRL)
BARLOW 67 NC 50A 701 *LILLESTOL+MONTANET+ (CERN+CDEF+IRAD+LIVP)
BEUSCH 67 PL 25 B 357 *FISCHER, GOBBI, ASTBURY+ (ETHZ+CERN)
DAHL 67 PR 163 1375 *HARDY+HESSE+KIRZMILLER (LRL)
ALITTI 68 PRL 21 1705 *BARNES, JENNELL, FLAMINIO, GOLDBERG,+ (BNL)
LAI 68 PHILAD. CONF. P-303 KWAN WU LAI (BNL)
PHELAN 68 THESIS JAMES J. PHELAN (ANL+ST. LOUIS UNIV)
ALSO 68 PRL 21 316 HOANG, EARTLY, PHELAN, ROBERTS+ (ANL+CHIC+NDAM)
AGUILAR- 69 PL 29 B 241 M. AGUILAR-BENITEZ, J. BARLOW,+ (CERN+CDEF)
ALSO 67 BARLOW
ALSO 69 NP 8 14 195 M. AGUILAR-BENITEZ, J. BARLOW,+ (CERN+CDEF)
HOANG 69 NC 61 A 325 T-F. HOANG (ANL)
HOANG 69 PR 184 1363 *EARTLY, PHELAN, ROBERTS,+ (ANL+ILLC)
BADIER 70 NP 8 22 512 *BCNNET, DREVILLON, BAUBILLIER,+ (EPOL+IPNP)
BATON 70 PL 33 B 528 *LAURENS, REIGNIER (SACLAY)
BEUSCH 70 PHILA. CONF. P-185 M. BEUSCH (ETHZ+CERN)
HYAMS 70 PHILA. CONF. P-41 *KOCH, BEUSCH,+ (CERN+MPIM+ETHZ+LDC+HAWAII)
ALSO 70 NP 8 22 189 HYAMS, KOCH, POTTER, VON LINDERN,+ (CERN+MPIM)
OH 70 PR D 1 2494 *GARRINKEL, MORSE, WALKER, PRENTICE+ (WISCONSIN)
ALSTON-G 71 PL 36 B 152 ALSTON-GARNJUST, BARBARO-GALTIERI,+ (LBL)
BASDEVAN 72 PL 41 B 178 BASDEVANT, FROGGATT, PETERSEN (CERN)
DAMERI 72 NC 9 A 1 *BORZATTA, GOUSSU,+ (GENO+MILA+SACL)
DUBOC 72 NP 8 46 429 *GOLDBERG, MAKOWSKI, DCNALD,+ (LNP+LIVP)
FLATTE 72 PL 38 B 232 *ALSTON-GARNJUST, BARBARO-GALTIERI,+ (LBL)
GRAYER 72 PHIL. CONF. PROC. 5 *HYAMS, JONES, SCHLEIN, BLUM, DIETL,+ (CERN+MPIM)
WILLIAMS 72 PR D 6 3178 P. K. WILLIAMS (FISU)
BINNIE 73 PRL 31 1534 *CARR, DEBENHAM, DUANE, GARBUTT,+ (LOIC+SHMP)
DIAMOND 73 PR D 7 1977 *BINKLEY,+ (WISC+DUKE+COLG+TNTD+OHIO)
ESTABROOK 73 TALLAHASSEE ESTABROOKS, MARTIN, GRAYER, HYAMS+ (CERN+MPIM)
FUJII 73 NC 13 A 311 Y. FUJII, N. KATO (TOKYO)
GRAYER 73 TALLAHASSEE *HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN+MPIM)
HYAMS 73 NP 8 64 134 *JONES, WEILHAMMER, BLUM, DIETL,+ (CERN+MPIM)
OGHS 73 THESIS W. OGHS (MPIM)
PROTOPOP 73 PR D 7 1280 PROTOPOPESCU, GARNJUST, GALTIERI, FLATTE+(LBL)
GRAYER 74 NP 8 76 375 *HYAMS, JONES, BLUM, DIETL (CERN+MPIM)
GRAYER 74 PR 8 75 189 *HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN+MPIM)
MORGAN 74 PL 518 71 D. MORGAN (RHEL)
FUJII 75 NP 885 179 Y. FUJII, N. FUKUGITA (TOKY)
MORGAN 75 PREPRINT RL75133 D. MORGAN (RHEL)
PANLICKI 75 PR D12 631 *AYRES-DIEBOLD, GREENE, KRAMER, WICKLUND (ANL)

REFERENCES FOR S*
WANG TSU-TSENG, VEKSLER, VRANA,+ (JINR)
A BIGI'S BRANDT, R CARRARA + (CERN)
H BINGHAM, B BLOCH + (EPOL+CERN)
ERMIN, MOYER, MARON, WALKER, WAMPLER (WISCONSIN)
*BARNES, JENNELL, FLAMINIO, GOLDBERG,+ (BNL)
BARMIN, DOLGICENKO, YEROFEEV, KRESTINI+ (ITEP)
CRENNELL, KALBFLEISCH, LAI, SCARR, SCHU+ (BNL)
*CAHL+HARDY+KIRZ-MILLER (LRL)
*LILLESTOL+MONTANET+ (CERN+CDEF+IRAD+LIVP)
*FISCHER, GOBBI, ASTBURY+ (ETHZ+CERN)
*HARDY+HESSE+KIRZMILLER (LRL)
*BARNES, JENNELL, FLAMINIO, GOLDBERG,+ (BNL)
KWAN WU LAI (BNL)
JAMES J. PHELAN (ANL+ST. LOUIS UNIV)
HOANG, EARTLY, PHELAN, ROBERTS+ (ANL+CHIC+NDAM)
M. AGUILAR-BENITEZ, J. BARLOW,+ (CERN+CDEF)
T-F. HOANG (ANL)
*EARTLY, PHELAN, ROBERTS,+ (ANL+ILLC)
*BCNNET, DREVILLON, BAUBILLIER,+ (EPOL+IPNP)
*LAURENS, REIGNIER (SACLAY)
M. BEUSCH (ETHZ+CERN)
*KOCH, BEUSCH,+ (CERN+MPIM+ETHZ+LDC+HAWAII)
HYAMS, KOCH, POTTER, VON LINDERN,+ (CERN+MPIM)
*GARRINKEL, MORSE, WALKER, PRENTICE+ (WISCONSIN)
ALSTON-GARNJUST, BARBARO-GALTIERI,+ (LBL)
BASDEVANT, FROGGATT, PETERSEN (CERN)
*BORZATTA, GOUSSU,+ (GENO+MILA+SACL)
*GOLDBERG, MAKOWSKI, DCNALD,+ (LNP+LIVP)
*ALSTON-GARNJUST, BARBARO-GALTIERI,+ (LBL)
*HYAMS, JONES, SCHLEIN, BLUM, DIETL,+ (CERN+MPIM)
P. K. WILLIAMS (FISU)
*CARR, DEBENHAM, DUANE, GARBUTT,+ (LOIC+SHMP)
*BINKLEY,+ (WISC+DUKE+COLG+TNTD+OHIO)
ESTABROOKS, MARTIN, GRAYER, HYAMS+ (CERN+MPIM)
Y. FUJII, N. KATO (TOKYO)
*HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN+MPIM)
*JONES, WEILHAMMER, BLUM, DIETL,+ (CERN+MPIM)
W. OGHS (MPIM)
PROTOPOPESCU, GARNJUST, GALTIERI, FLATTE+(LBL)
*HYAMS, JONES, BLUM, DIETL (CERN+MPIM)
*HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN+MPIM)
D. MORGAN (RHEL)
Y. FUJII, N. FUKUGITA (TOKY)
D. MORGAN (RHEL)
*AYRES-DIEBOLD, GREENE, KRAMER, WICKLUND (ANL)

phi(1020)

4 PHI(1020, JPG=1--) I=0

4 PHI MASS (MEV)
M S 18(1019.0) (2.0) SCHLEIN 63 HBC 2.0 K-P - 12/75*
M S 20(1018.6) (0.5) MILLER D 65 HBC 0.0 PBAR P 12/75*
M S 41(1020.0) (2.0) LONDON 66 HBC 2.2 K-P 12/75*
M S 46(1021.5) (0.8) ABRAMS 67 HBC 4.2 K-P 12/75*
M S 15(1019.) (2.1) BARLOW 67 HBC 1.2 PBAR P 12/75*
M S 32(1020.0) (4.0) DAHL 67 HBC 1.4 PI-P 12/75*
M 165 1022. 1.5 MOSTEK 68 DSPK 1.8 GAMMA + C 6/68
M 136 1018. 0.5 0.35 HYAMS 70 DSPK 11. PI-P 6/70
M 107 1021.0 1.5 ALVENSLE 71 DSPK GAMMA+K 1/72
M D 70 1019.9 0.7 DIBIANCA 71 DBC 4.93 K-N 12/75*
M 410(1019.9) (0.3) STOTTELY 71 HBC 2.9 K-P,Y K KEAR 11/71
M D 120 1019.6 0.5 AGUILAR 72 HBC 3.9,4.6 K-P 12/75*
M D 100 1019.9 0.5 AGUILAR 72 HBC 3.9,4.6 K-P 12/75*
M 87 1020.8 0.8 BALAKIN 72 DSPK E+ E- COLL. BEAMS 12/72
M 131 1020.4 0.5 COLLEY 72 HBC 10. K+ P,K+ P PHI 12/75*
M 100 1020.3 0.4 BALLAM 73 HBC 2.8 - 5.3 G P 1/74
M 500 1019.5 0.6 AYRES 74 ASPK 3-6PI/K-P,K+K- 12/75*
M 984 1019.4 0.8 BESCH 74 CNTR 2 GAMMA P,K+K- 12/75*
M 54 1018.7 0.8 BIZZARRI 74 DBC 0 PBAR N,K+K- PI 12/75*
M 71 1018.6 .7 BIZZARRI 74 DBC 0 PBAR N,K+K- PI 12/75*
M 170 1020.3 0.4 DE GROOT 74 HBC 4.2 K-P,L K+K- 12/75*
M AVG 20.0 3.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
M STUDENT1019.84 0.20 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
(OSEE IDEOGRAM BELOW)
M S INSIGNIFICANT DATA WITH SMALL STATISTICS NO LONGER AVERAGED
M D MASS ERRORS ENLARGED BY US TO WIDTH/SQRT(N)



4 PHI WIDTH (MEV)
W S 20 (3.5) (1.0) MILLER D 65 HBC 0.0 PBAR P 12/75*
W S 41 (6.0) (4.0) LONDON 66 HBC 2.2 K-P 12/75*
W S 46 (1.8) (3.0) ABRAMS 67 HBC 4.2 K-P 12/75*
W S 165 (4.5) (3.0) (1.5) (2.0) MOSTEK 68 DSPK 1.8 GAMMA + C 6/68
W 150 4.2 0.9 AUGUSTIN 69 DSPK E+ E- COLL. BEAMS 12/72
W 136 4.09 0.29 BIZOT 70 DSPK E+ E- COLL. BEAMS 12/72
W D 87 4.67 2.02 0.9 HYAMS 70 DSPK 11. PI-P 6/70
W D 70 5.5 2.6 DIBIANCA 71 DBC 4.93 K-N 12/75*
W 110 (4.5) (3.0) (4.0) STOTTELY 71 HBC 2.9 K-P,Y K KEAR 11/71
W D 120 4.6 1.7 AGUILAR 72 HBC 3.9,4.6 K-P 12/75*
W D 100 4.7 1.9 AGUILAR 72 HBC 3.9,4.6 K-P 12/75*
W D 131 5.0 1.8 COLLEY 72 HBC 10. K+ P,K+ P PHI 12/75*
W D 100 3.8 1.5 BALLAM 73 HBC 2.8 - 5.3 G P 12/75*
W 150 3.81 0.37 COSME 2 74 DSPK E+ E- COLL. BEAMS 2/74
W D 500 4.5 0.8 AYRES 74 ASPK 3-6PI/K-P,K+K- 12/75*
W D 984 4.4 0.6 BESCH 74 CNTR 2 GAMMA P,K+K- 12/75*
W S 54 (6.8) (3.1) (2.5) BIZZARRI 74 DBC 0 PBAR N,K+K- PI 12/75*
W S 71 (4.1) (3.8) (3.6) BIZZARRI 74 DBC 0 PBAR N,K+K- PI 12/75*
W D 170 4.2 1.3 DE GROOT 74 HBC 4.2 K-P,L K+K- 12/75*
W AVG 4.09 0.19 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W STUDENT 4.09 0.21 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
W S INSIGNIFICANT DATA WITH SMALL STATISTICS NO LONGER AVERAGED
W D WIDTH ERRORS ENLARGED BY US TO 4*WIDTH/SQRT(N)

4 PHI PARTIAL DECAY MODES
DECAY MASSES
P1 PHI INTO K+ K- 493+ 493
P2 PHI INTO KL KS 497+ 497
P3 PHI INTO PI+ PI- P10 (INCLUDING RHO PI) 139+ 139+ 134
P4 PHI INTO ETA GAMMA 548+ 0
P5 PHI INTO E+ E- .5+ .5
P6 PHI INTO MU+ MU- 105+ 105
P7 PHI INTO P10 GAMMA 134+ 0
P8 PHI INTO PI+ PI- (VICILATES G) 139+ 139
P9 PHI INTO PI+PI-GAMMA 139+ 139+ 0
P10 PHI INTO OMEGA GAMMA (VICILATES C) 783+ 0
P11 PHI INTO ETA P10 (VICILATES C) 548+ 134
P12 PHI INTO RHO GAMMA (VICILATES C) 773+ 0
P13 PHI INTO ETA NEUTRALS
P14 PHI INTO 5PI

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i +/- delta P_i, where delta P_i = sqrt(delta P_i^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 4 columns: P 1, P 2, P 3, P 4. Values include correlation coefficients like .4661+-.0232, -.7766, -.3503+-.0203, etc.

4 PHI BRANCHING RATIOS

R1 PHI INTO (K+ K-)/(K+K- + PI+ PI- P10) (P1)/(P1+P2+P3) 10/66
R1 B 27 (0.26) (0.06) BARBER 65 HBC 2.7 K-P 10/66
R1 C 252 (0.48) (0.04) LINDSEY 66 HBC 2.7 K-P 10/66
R1 C (0.493) (0.044) BIZOT 70 DSPK E+ E- COLL. BEAMS 11/71
R1 C SUPERSEDED BY CHATELUS 71 11/71
R1 0.540 0.026 BALAKIN 71 DSPK E+ E- COLL. BEAM 11/71
R1 0.486 0.044 CHATELUS 71 DSPK E+ E- COLL. BEAMS 11/71
R1 270 C.49 0.06 DE GROOT 74 HBC 4.2 K-P,L PHI 12/75*
R1 AVG 0.505 0.021 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R1 STUDENT 0.504 0.024 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R1 FIT 0.475 0.024 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)

Data Card Listings

For notation, see key at front of Listings.

Mesons

phi(1020), M(1033), B1(1040)

Table of particle data cards for phi(1020), M(1033), and B1(1040). Each entry includes particle name, mass, width, and various fit parameters and references.

Table of particle data cards for phi(1020), M(1033), and B1(1040). Each entry includes particle name, mass, width, and various fit parameters and references.

M(1033) -> MM

B1(1040) -> omega pi

REFERENCES FOR PHI

List of references for phi(1020) including authors like Bertanza, Gelfand, and Schleif, and their respective publications.

List of references for M(1033) and B1(1040) including authors like Garfinke, Binnie, and Butttram, and their respective publications.

Mesons

$B_1(1040)$, $\eta_N(1080)$, $A_1(1100)$

Data Card Listings

For notation, see key at front of Listings.

DEFIOX 73 PL 43 B 141 +GDBRZYNSKI,ESPIGAT,NASCIMENTO,+ (CDEF)
 DIAZ 74 PRL 32 260 +DIBIANCA,FICKINGER,ANDERSCN,+ (CASE+GARN)

REFERENCES FOR $B_1(1040)$

$\eta_N(1080)$
 $\rightarrow \pi\pi$

30 ETA N(1080, JPC=N +) I=0 J GREATER THAN 1
 SOME EXPERIMENTS SUGGEST J=2.
 OMITTED FROM TABLE

Note on $\pi^+\pi^-$ Peaks Called $\eta_N(1080)$

The $\eta_N(1080)$ is seen in $\pi^-p \rightarrow \pi^+\pi^-n$ predominantly at backward decay angles, $\cos\theta < -0.75$. OH 70 state that this "bump is almost certainly the result of P-D interference."

Note that the selection made in some bubble chamber experiments to reduce the background under the $\eta_N(1080)$ in the reaction $\pi^-p \rightarrow \pi^+\pi^-n$ may lead to a sample of events ambiguous with $\pi^-p \rightarrow p\pi^-\pi^0\pi^0$ (BATON 70 and private communications from G. Laurens).

30 ETA N MASS (MEV)						
M	1060.0	15.0	MILLER	68 HBC	4.0 PI- P	9/68
M	70 1085.0	10.0	WHITEHEAD	68 ASPK	3.1-3.6 PI-P	10/67
M	1120.0	100.0	OH	69 HBC	7.PI- P,PI+ D	9/69
M	1112.0	16.0	CLAYTON	70 HBC	2.5 PBAR P, 4 PI	1/71
M	(1080.0)		DIAZ	70 HBC	0. PBAR P, 4 PI	5/70
M	1070.0	20.0	REYNOLDS	70 HBC	2.26-2.36 PI- P	1/71
M				
M	AVG 1083.3	9.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)			
M	STUDENT 1083.0	8.3	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

30 ETA N WIDTH (MEV)						
W	(70.0)	DR LESS	MILLER	68 HBC	4.0 PI- P	9/68
W	(25.0)	DR LESS	WHITEHEAD	68 ASPK	3.1-3.6 PI-P	10/67
W	150.0	100.0	OH	69 HBC	7.PI- P,PI+ D	9/69
W	(80.0)		CLAYTON	70 HBC	2.5 PBAR P, 4 PI	1/71
W	(80.0)		DIAZ	70 HBC	0. PBAR P, 4 PI	5/70
W	85.0	35.0	REYNOLDS	70 HBC	2.26-2.36 PI- P	1/71
W				
W	AVG 98.0	31.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W	STUDENT 97.7	34.5	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

REFERENCES FOR ETA N

MILLER 68 PRL 21 1489 +GUTAY,JOHNSON,KENNEY+ (PURDUE+NDAM+SLAC)
 WHITEHEA 68 NC 53 A 817 C.WHITEHEAD + (AERE+SHMP+LOIC)
 OH 69 PRL 23 331 +WALKER,CARROLL,FIREBAUGH,+ (WISC+TNT0)
 BATON 70 PL 33 B 528 +LAURENS,ZEINIER (SACLAY)
 CLAYTON 70 NP 8 22 85 +NASON,MURHEAD,RIGOPOULOS,+ (LIVP+ATEN)
 DIAZ 70 NP 8 16 239 +GAVILLET,LABROSSE,MCNTANET,+ (CERN+CDEF)
 REYNOLDS 70 NP 8 21 77 +ALBRIGHT,BRADLEY,+ (OHIO+FSU+MNN+COLO)
 WHITEHEA 72 NP 8 48 365 WHITEHEAD,AULD,+ (AERE+RHEL+SHMP+LOUC)
 DIAMOND 73 PR D 7 1977 +BINKLEY,+ (WISC+DUKE+COLD+TNT0+OHIO)

$A_1(1100)$

10 A1(1100, JPC=1+-) I=1

The $A_1 \rightarrow \rho\pi$ bump has been mainly observed in the diffraction-like process $\pi N \rightarrow (\pi\pi\pi)N$ without quantum number exchange and at small momentum transfer. There are also observations of structure in the A_1 mass region in reactions where additional mesons are produced, in backward production (see however ABASHIAN 75), and in $\bar{p}p$ annihilations (see the Data Card Listings).

There is a special category of "diffractive mesons" consisting of statistically significant peaks like A_1 , A_3 , Q , and L (see the corresponding mini-reviews), which are not far above the respective $\rho\pi$, $f\pi$, $K^*(892)\pi$, and $K^*(1420)\pi$ thresholds. Because the behavior near threshold in these channels may be described by the Deck effect, a resonance interpretation is questionable. These peaks are included in the Meson Table, but we do not mean to imply that they are necessarily genuine resonances.

Other meson systems such as $K\omega$, $K\phi$, $K^*\bar{K}$ also exhibit low mass peaks when produced in reactions which may proceed without quantum number exchange (see for example DAVIS 72, CHUNG 74, THEOCHAROPOULOS 74, CARNEY 75, OTTER 75).

Partial-wave analyses of multi-meson systems in reactions like $\pi N \rightarrow (\pi\pi\pi)N$ are becoming available (ASCOLI 70). Several important assumptions are made in such analyses (see HANSEN 74, HERNDON 75 for detailed discussions), amongst which:

- (a) for a given t , the 3π vertex is independent of the NN vertex;
- (b) the 3π decay proceeds through quasi-two-body states ($\rho\pi, \epsilon\pi, \dots$) in the spirit of the isobar model (a partial-wave analysis with unitarity corrections to the isobar model is applied to the 3π system in ASCOLI 75; see also AARON 75, AITCHISON 75, and references therein).

The dominant effect in the A_1 mass region is a broad $J^P=1^+$ $\rho\pi$ S-wave enhancement, with a maximum intensity at ~ 1100 MeV and a width ~ 300 MeV (ANTIPOV1 73, KRUSE 74, OTTER 74, TABAK 74, THOMPSON 74, EMMS1 75). The phase of the $J^P=1^+$ wave shows little variation relative to various other "background" waves (see however the A_3 and Q mini-reviews). In contrast the A_2 peak has been confirmed as a resonance with a Breit-Wigner-like phase change of the $J^P=2^+$ partial wave (ASCOLI 70, ANTIPOV1 73, OTTER 74, TABAK 74, THOMPSON 74).

These results suggest that at most a small part of the A_1 enhancement corresponds to a $J^P=1^+$ resonance. Indeed ASCOLI 73 and ASCOLI 74 show that the partial-wave structure of the 3π

Data Card Listings

For notation, see key at front of Listings.

Mesons A₁(1100)

system is qualitatively reproduced by a Reggeized pion exchange model.

BOWLER 75 attempt to explain the lack of variation of the 1^+ phase by allowing for a phase difference between the Deck amplitude and a direct "A₁ resonance" production amplitude. Good fits to the data are obtained, but the A₁ is predicted to have $M \sim 1300$ MeV, $\Gamma \sim 250$ MeV.

Recent analyses of the $(3\pi)^0$ system produced by charge exchange find no evidence for A₁ production. WAGNER 75 study the reaction $\pi^+ p \rightarrow \pi^+ \pi^- \pi^0 \Delta^{++}$ at 7 GeV/c and are able to set an upper limit of 2 μ b on the A₁ production cross section if $\Gamma < 150$ MeV. Analyzing $\pi^+ n \rightarrow \pi^+ \pi^- \pi^0 p$ at 4 GeV/c, EMMS2 75 observe no A₁ with $M \sim 1100$ MeV, but cannot rule out resonance production above 1300 MeV.

References Not Included in the Data Card Listings

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- P. Theocharopoulos et al., Nucl. Phys. B83, 1 (1974).

10 A1 MASS (MEV)	
M	PRODUCED BY PI +
M	(1080.0)
M	(1080.0) APPROX.
M	(1040.0)
M	(1128.) (8.)
M	PRODUCED BY PI -
M	(1060.)
M	(1089.0) (12.0)
M	(1090.) APPROX.
M	(1055.0) (6.0)
M	(1119.) (30.)
M	S SHOULDER ON A2 ONLY
M	(1069.0) (7.0)
M	(1120.0)
M	F T (1150.)
M	T MASS AND WIDTH SEEN TO DEPEND ON T, UNIQUE DET. IMPOSSIBLE

M	PRODUCED BY PIONS, BACKWARDS SCATT.	ANDERSON	69 MMS	- 16 PI- P, BACKW9	8/69
M	(1115.0) (20.0)				
M	PRODUCED BY PBARS, SEE TYPED NOTE.	DANYSZ	67 HBC	+ 3.3.6 PBAR P	7/67
M	(1054.) (7.)	FRIDMAN	68 HBC	+ 5.7 PBAR P	6/68
M	(1042.) (21.)	ATHERTON	73 HBC	+ 5.7 PBAR P	1/74
M	A (1076.) (5.)				
M	A JP ANALYSIS GIVES SOME EVIDENCE FOR RHO PI D-WAVE				1/73
M	PRODUCED BY K-, SEE TYPED NOTE.	ALLISON	67 HBC	+ 6 K-P, LAM +5 PI	1/68
M	(111.) (10.)	ALLISON	67 HBC	+ 6 K-P, LAM +4 PI	1/68
M	(1117.) (35.)	JUHALA	67 HBC	0 4.6-5 K-P, 5BCDY	1/68
M	(1060.) (15.)				
M	PRODUCED BY K+, SEE TYPED NOTE.	ALEXANDER	69 HBC	+ 9 K+P	9/69
M	(1060.0) (20.0)	BERLINGHI	69 HBC	+ 0 12.7 K+ P	9/69
M	(1030.0) (20.0)				
M	K+ FOR CONTRADICTIONARY EVIDENCE SEE RABIN 70.				
M	F FROM A FIT TO JP=1+ RHO PI PARTIAL WAVE				
M	AVERAGING NOT MEANINGFUL				

10 A1 WIDTH (MEV)	
W	PRODUCED BY PIONS, RESONANCE INTERP. CONFUSED BY DECK EFFECT
W	PRODUCED BY PI +
W	(80.0)
W	(130.) APPROX.
W	(50.0) OR LES
W	F (300.) APPROX.
W	F (367.) (30.)
W	PRODUCED BY PI -
W	(140.0) (31.0)
W	(125.) APPROX.
W	(77.0) (17.0)
W	K (76.) (46.)
W	K SHOULDER ON A2 ONLY
W	(99.0) (15.0)
W	F T (300.)
W	T MASS AND WIDTH SEEN TO DEPEND ON T, UNIQUE DET. IMPOSSIBLE
W	PRODUCED BY PIONS, BACKWARDS SCATT.
W	(98.0) (45.0) (20.0)
W	PRODUCED BY PBARS, SEE TYPED NOTE.
W	(33.) (19.)
W	(130.) APPROX.
W	A (36.) (20.) (15.)
W	A JP ANALYSIS GIVES SOME EVIDENCE FOR RHO PI D-WAVE
W	PRODUCED BY K-, SEE TYPED NOTE.
W	(50.) (50.)
W	(50.) (25.)
W	(120.) (15.)
W	PRODUCED BY K+, SEE TYPED NOTE.
W	(160.0) (20.0)
W	B (120.0) (30.0)
W	K+ FOR CONTRADICTIONARY EVIDENCE SEE RABIN 70.
W	(130.0) (20.0)
W	F FROM A FIT TO JP=1+ RHO PI PARTIAL WAVE
W	AVERAGING NOT MEANINGFUL

10 A1 PARTIAL DECAY MODES		
P1	A1 INTO RHO PI	DECAY MASSES
P2	A1 INTO KBAR K	772+ 139
		493+ 497

10 A1 BRANCHING RATIOS		
R1	A1 INTO (KBAR K)/(RHO PI)	(P2)/(P1)
R1	(0.0025) OR LESS	DAHL 67 HBC - 4.0 PI- P .10/66

REFERENCES FOR A1

BELLINI 63 NC 29 896	BELLINI, FIORINI, HERZ, NEGRI, RATTI (MILAN)
ADERHOLZ 64 PL 10 226	AACH+BERL+BIHM+BONN+DESY+HAMBURG+LOIC+MPIM
GOLDBABE 64 PRL 12 336	GOLDBABER, BROWN, KADYK, SHEN+ (LRL+USC) JP
LANDER 64 PRL 13 346 A	LANDER, ABOLINS, CARMONY, HENDRICKS + (UCSD) JP
ABCLINS 65 ATHENS(OHIO)CONF.	+CARMONY, LANDER, XUONG, YAGER (LA JOLLA)I=1
ALITTI 65 PL 15 69	ALITTI, BATON, DELER, CRUSSARD+ (SACL+BGNA)
ALLARD 66 NC 46A 737	+DRIJARD+HENNESSY+ (ORSAY+MILAN+SACL+JUG)
DEUTSCHM 66 PL 20 82	DEUTSCHMANN, STEINBERG + (AACH+BERLIN+CERN)
HESS 66 UCRL-16832	R I FESS (THEISIS, BERKELEY) (LRL)
ALLISON 67 PL 25B 619	+CRUZ+ (OXF+MPIH+BIHM+RHEL+GLAS+LOIC)
DAHL 67 PR 163 1377	+HARDY+HESS+KIRZ+MILLER (LFL)
DANYSZ 67 NC 51 A 801	DANYSZ+FRENCH+SIMAK (CERN)
JUHALA 67 PRL 19 1355	+LEACOCK+RHODE+KOPPELMAN+ (TOWA+GOLO)
SLATTERY 67 NC 50A 377	+KRAYBILL+FORMAN+FERBEL (YALE+ROCH) JP
ARMENISE 68 PL 26 B 336	+FORING+CARTACCI+ (BARI+BGNA+FRIZ+ORSAY)
ASCOLI 68 PRL 21 113	+CRAWLEY, KRUSE, MORTARA, SCHAFFER+ (ILLINOIS)
BALLAM 68 PRL 21 934	+BRODY, CHADWICK, FRIES, GUIRAGOSSIAN+ (SLAC)JP
BOESEBEC 68 NP 8 501	BOESEBEC, DEUTSCHMANN+ (AACHEN+BERLIN+CERN)
CASO 68 NC 54 A 983	+CCNTE+CORDS+DIJAZ+ (GENOVA+HAMB+MILA+SACL)
CHUNG 68 PR 165 1491	S.U.CHUNG, O.DAHL, J.KIRZ, D.H.MILLER (LRL)
CNDPS 68 PRL 21 1609	+HOUGH, COHN, BUGG+ (BNL+ORNL+UCND+TENN+PENN)
FRIDMAN 68 PR 167 1268	+MAURER, MICHALOW, OUDET+ (HEID+STRASBOURG)
JUNKMANN 68 NP 88 471	+COCCENI+ (AACH+BERL+BGNN+CERN+WARS)
KEY 68 PR 166 1430	+PRENTICE+COOPER+MANNER+ (TNTD+ANL+WISCI)
ALEXANDE 69 PR 183 1168	G.ALEXANDER, A.FIRESTCNE, G.GOLDBABER (LRL)
ALLAY 69 PL 29B 198	+BINON+DIDDESS+DUTELL+KLOVNING+... (CERN)
ANDERSON 69 PRL 22 1390	+COLLINS+ (BNL+CARN)
BERLINGH 69 PRL 23 42	BERLINGHIERT, FARBET, + (ROCH)
DONALD 69 NP B 11 551	+EDWARDS, BURAN, BETTINI, + (LIVP+OSLO+PADO)

Mesons

A₁(1100), M(1150), A_{1,5}(1170), ε(1200)

Data Card Listings

For notation, see key at front of Listings.

Table listing meson properties and references. Includes entries for FAYOLLE, JUHALA, KENYON, ARMENISE, ASCOLI, BRANDENB, CASO, CRENNELL, CHIEN1, CHIEN2, GARELTICK, RABIN, SHIH, ASCOLI, BEMPRAD, BERGER, RINAUDO, BEREVNI, BLOODWDR, DIEBOLD, LAMSA, MORSE, ANTIPV1, ANTIPV2, ARNOLD, ASCOLI, ASCOLI, ATHERTON, READ, BOWLER, ASCOLI, KRUSE, LICHTMAN, OTTER, TABAK, THOMPSON, THOMPSON, ABASHIAN, BEUSCH, BOWLER, DIAZ, EMMS, EMMS, HORNE, KANE, WAGNER.

ε(1200) 14 PI PI S WAVE, CALLED EPSILON

S-Wave ππ Interactions in the Region 280-1800 MeV

In this note, we discuss information on the isoscalar ππ S-wave in terms of its phase shift δ₀⁰, from threshold to 1800 MeV.

The threshold behavior of elastic ππ scattering involves S and P waves which can be sufficiently well described by the scattering lengths a₀⁰, a₁¹, and a₀². In spite of many attempts (see PILKUHN 73, BONNIER 74, PASCUAL 74, RIESTER 75, SRINIVASAN 75), the determination of these parameters still meets with great difficulties (BASDEVANT 73, 75). The parameters a₀⁰ and a₀² are strongly correlated and must lie in a narrow band in the (a₀⁰, a₀²)-plane (MORGAN2 70). Thus if a₀⁰ is fixed, a₀² and a₁¹ are determined within small uncertainties (BASDEVANT 72, 75). However, all one knows is a region of finite extent within which a₀⁰ must lie. It is therefore not established whether, e.g., the Weinberg predictions are supported by experiment.

Near threshold the S wave shows no resonant behavior. The so-called ABC and DEF effects (BOOTH 63, HALL 69, BRODY 70, 72, BANAIGS 71, 73) occur only on nuclear targets (d, H³, He³) and move when kinematical conditions change. Thus they must be kinematical effects (DUBAL 71, BRODY2 72, RISSER 73, BAR-NIR 75, BARRY 75).

The region of elastic ππ scattering is known to extend from threshold to about 990 MeV, near the K₁⁻ threshold (BATON 70, CARROLL 72, PROTOPODESCU 73, HYAMS 73, OCHS 73).

Up to the ρ meson mass region, δ₀⁰ is (qualitatively) uniquely determined; it rises monotonically and reaches a value of 60° to 70° near 700 MeV (SONDEREGGER 69, BATON 70, BAILLON 72, CARROLL 72, FRENKIEL 72, GAIDOS 72, PROTOPODESCU 73, HYAMS 73, OCHS 73, ENGLER 74, ESTABROOKS 74, 75, GRAYER 74); see Fig. 1.

In the mass region of 700 to 900 MeV, all energy-independent phase-shift analyses using the constraint η₀⁰ = 1 find two solutions ("up-down ambiguity"). This ambiguity was resolved in favor of the "down" solution by the observation of a very rapid decrease in the S-wave amplitude

M(1150) → MM

68 M(1150) EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

Table with 7 columns: M, N, 65, 1148.3, 3.3, JACOBEL, 72, MMS, 0, 2.4, PI-, P, N, MM, 12/72, 12/75*

Table with 7 columns: M, 65, 15.0, 9.0, 11.7, JACOBEL, 72, MMS, 0, 2.4, PI-, P, N, MM, 12/72

REFERENCES FOR M(1150)

Table with 2 columns: JACOBEL, BUTTRAM, PRL, 29, 671, GARFINKEL, HOFFMAN, IOWA+PURD+ANL; PRL, 35, 970, CRAWLEY, DUKE, LAMB, LEEPER, PETERSON, ISU

A_{1,5}(1170) → 3π

44 A 1.5(1170, JPG = -) I=1 THIS ENTRY LISTS REFERENCES TO PEAKS OF LOW STATISTICAL SIGNIFICANCE IN THE 3 PI SYSTEM BETWEEN THE A1 AND THE A2. OMITTED FROM TABLE.

REFERENCES FOR A 1.5

Table with 2 columns: BUTTERWD, CASON, ASCOLI, DONALD, VCN, JUNKMANN, ARMENISE, GALLOWAY, MORSE, HEIDELB, CONF, P-28, REVIEW TALK ON MESONS AT HEIDELBERG CONF; PRL, 18, 880, DEBARD, GROVES, INOTREDAME; PRL, 21, 113, CRAWLEY, KRUSE, MORTARA, SCHAFFER, ILLINOIS; PL, 26, B, 327, FRODESEN, BETTINI, LIVERPOOL, OSLO, PADUA; PL, 27B, 253, MIYASHITA, KOPELMAN, MARSHALL, LIBBY, COLO; NP, 88, 471, ACCONTI, AACB+BERL+BOHN+CERN+MARS; PR D, 1, 3077, GHIDINI, FORINO, CARTACCI, BARI+BONA+FIRZ; NP, B, 43, 77, MOTT, ALYEA, LEE, MARTIN, PRICKETT, IND; NP, B, 43, 77, OH, WALKER, JOHNSTON, YDON, WISC+TNTD

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\epsilon(1200)$

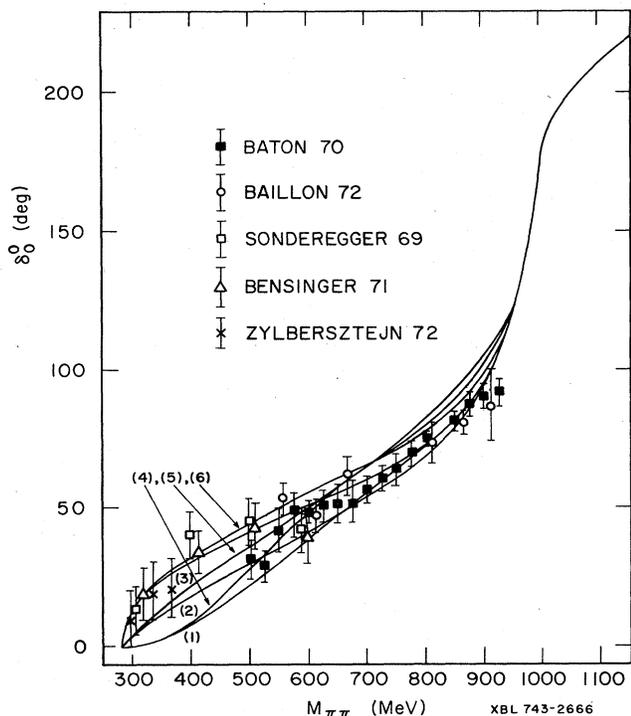


Fig. 1. The S-wave, $I=0$ $\pi\pi$ phase shift, as determined by several experiments. The curves are from BASDEVANT 73; solutions (1-3) were obtained by fitting the data of BATON 70, while (4-6) are from fits to the data of PROTOPOPESCU 73.

between 950 and 980 MeV (FLATTE 72, GAIDOS 72, HYAMS 73, BINNIE 73, ENGLER 74). The size of the observed drop corresponds to a change from nearly the unitarity limit to zero; see Fig. 2.

Independent evidence for the correctness of this "down" solution comes from studies of the $\pi^0\pi^0$ system (APEL 72, BRAUN 73, SKUJA 73, RIESTER 75). They observe a wide $\pi^0\pi^0$ enhancement of ~ 800 MeV which is much better described by the "down" solution than the "up" solution. Furthermore, indirect information from elastic $\pi\pi$ scattering in the crossed channel (ELVEKJAER 72, NIELSEN 70, 72) is compatible with the "down", but not with the "up", solution.

The ambiguities of the phase-shift solutions stem from the fact that there are more helicity amplitudes than observables. Thus in the absence of polarization measurements one is obliged to make some supplementary assumptions (see, e.g., DONOHUE 75). Analyzing the same data (GRAYER 74) by different methods, HYAMS 75 find four solutions in the region 1.0-1.8 GeV, ESTABROOKS 74 find eight solutions above 1.2 GeV and four solutions below, whereas FROGGATT 75 find arguments to favor one of these solutions [solution B of ESTABROOKS 74 (see Fig. 3), essentially the same as HYAMS 75 solution +--].

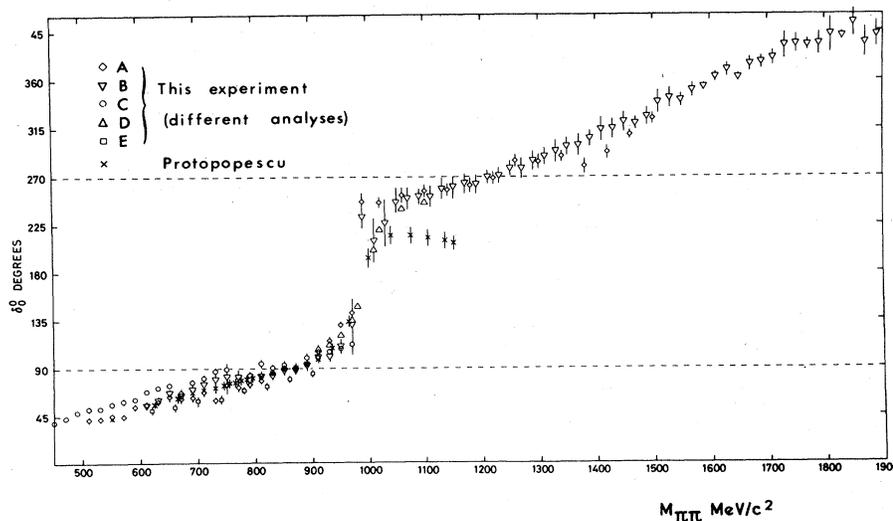


Fig. 2. The S-wave, $I=0$ $\pi\pi$ phase shift, as determined by various analyses of the data of GRAYER 74, compared with the previous results of PROTOPOPESCU 73. (Figure from GRAYER 74.)

Mesons

$\epsilon(1200)$

Data Card Listings

For notation, see key at front of Listings.

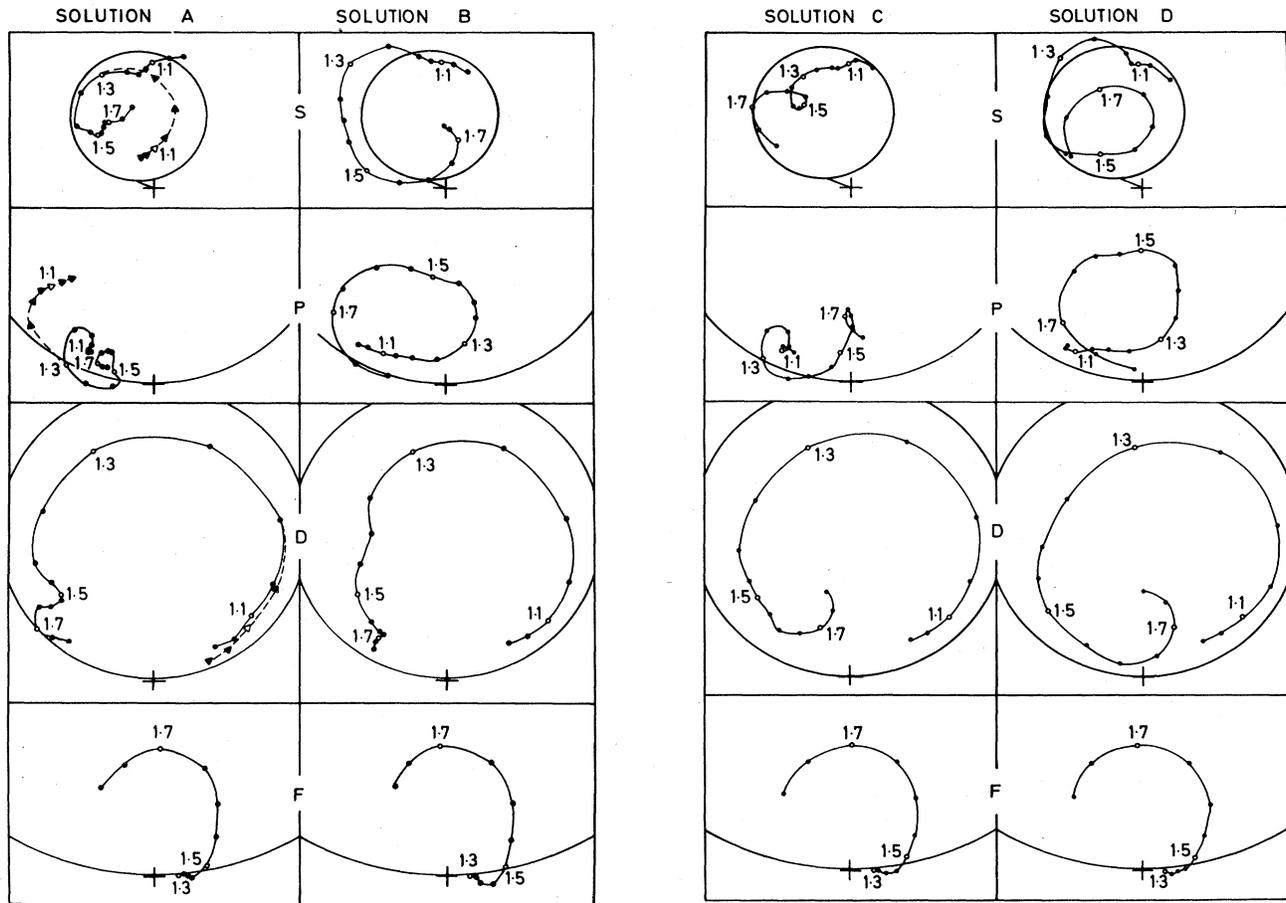


Fig. 3. Argand plots for $\pi\pi$ S, P, D, and F partial-wave amplitudes. Shown are solutions A-D of ESTABROOKS 75, obtained with the data of GRAYER 74. The points on the curves are at 50 MeV intervals. The diameters of the S-, P-, D-, and F-wave unitarity circles shown are given by $2/3$, $\sqrt{3}$, $(2/3)\sqrt{5}$, and $\sqrt{7}$, respectively. An elastic $I = 2$ S-wave phase shift of 25° independent of $M_{\pi\pi}$ has been assumed. (Figure from ESTABROOKS 75.)

Near 1.2 GeV all δ_0^0 solutions based on the GRAYER 74 data, as well as independent analyses of other experiments (CARROLL 72, 74, ENGLER 75), exhibit $\delta_0^0 = 270^\circ$.

It has now been made plausible (MORGAN 75) that all available data are compatible with the existence of just two poles: the $S^*(993)$ connected with the rapid variation of δ_0^0 near the $K\bar{K}$ threshold, and the $\epsilon(1200)$. The $S^*(993)$ is also responsible for the large $K\bar{K}$ $I=0$ S-wave scattering length.

The $\epsilon(1200)$ "replaces" the pole previously listed below 600 MeV. Note that although many analyses in the past found solutions with an

$\epsilon(600)$ pole, they also found solutions without such a pole (PROTOPODESCU 73, HYAMS 73).

Thus we have just the right number of isoscalar S-wave resonances to make up an SU(3) nonet, presumably together with the $\delta(970)$ and the $\kappa(1250)$ mesons (MORGAN 75).

Note that, although there is general consensus about the position of the S^* pole on the second sheet, it is not clear whether it is accompanied by a companion pole on the third sheet or not (FUJII 75, MORGAN 75). The contrast between the width of 180 MeV of MORGAN 75, and the value of 40 MeV obtained as twice the imaginary part, reflects just this uncertainty.

Data Card Listings

For notation, see key at front of Listings.

Mesons

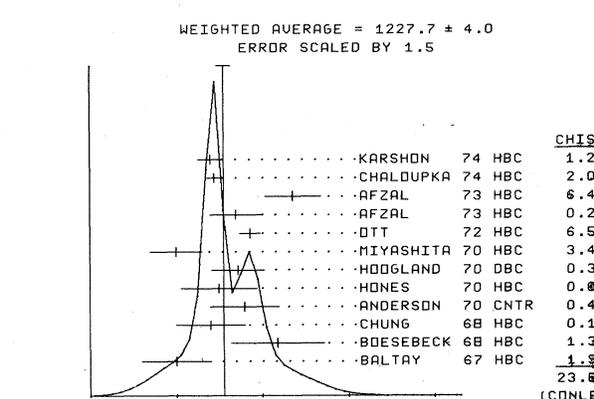
ε(1200), B(1235)

Table listing meson properties for ε(1200) and B(1235). Columns include particle name, mass (MEV), width (MEV), and references. Includes sub-sections for 'REFERENCES FOR EPSILON' and 'REFERENCES FOR B(1235)'.

Table listing meson properties for B(1235). Columns include particle name, mass (MEV), width (MEV), and references. Includes sub-sections for 'REFERENCES FOR B(1235)' and 'REFERENCES FOR B(1235)'. Includes a box labeled 'B(1235)'.

B(1235) 11 B(1235, JPC=1+-) I=1

Table listing meson properties for B(1235) with various quantum numbers and decay modes. Columns include particle name, mass (MEV), width (MEV), and references. Includes sub-sections for 'REFERENCES FOR B(1235)' and 'REFERENCES FOR B(1235)'.



Mesons

B(1235), rho'(1250), f(1270)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various decay parameters. Includes sub-sections for B WIDTH (MEV) and PARTIAL DECAY MODES.

Table listing experimental data for rho'(1250), including RHO PRIME(1250) MASS (MEV) and RHO PRIME(1250) WIDTH (MEV).

Table listing branching ratios for B(1235) into Omega pi, K K-bar pi, and other modes. Includes sub-sections for PARTIAL DECAY MODES and BRANCHING RATIOS.

Table listing experimental data for f(1270), including F(1270) MASS (MEV) and F(1270) WIDTH (MEV).

Table listing references for B(1235) and f(1270) decays, including authors and journal information.

Table listing references for rho'(1250) and f(1270) decays, including authors and journal information.

Data Card Listings

For notation, see key at front of Listings.

Mesons

f(1270)

Table with columns for particle name, mass, width, and other properties. Includes entries for SELOVE, JACOBS, RABIN, etc.

WEIGHTED AVERAGE = 181.1 ± 5.6
ERROR SCALED BY 1.7

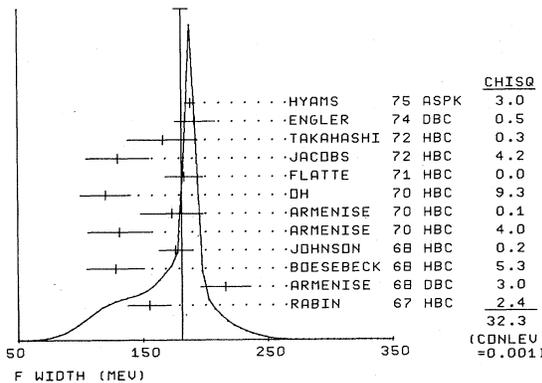


Table titled '5 F PARTIAL DECAY MODES' showing decay masses for various modes like F INTO PI PI, F INTO 2PI+ 2PI-, etc.

Table titled '5 F BRANCHING RATIOS' showing ratios for various decay channels and resonance parameters.

Table with columns for particle name, mass, width, and other properties. Includes entries for R3, R4, R5, R6, etc.

Table titled 'REFERENCES FOR F' listing various scientific references and authors associated with the f(1270) meson.

Mesons
D(1285), A2(1310)

Data Card Listings

For notation, see key at front of Listings.

D(1285)

8 D(1285, JP= +) I=0
(JP=0-+1+2- WITH 1+ FAVORED.)

Table with columns: M, (1290.0), APPROX., BARLOW, 67 HBC, 1.2 PBAR P, 4 PFS, 5/67. Includes student averages and missing mass spectrum.

Table with columns: W, R, (35.0), (10.0), DAHL, 67 HBC, 1.6-4.2 PI- P, 11/71. Includes student averages and resolution.

Table with columns: P1, D INTO K KBAR PI, 497+ 497+ 134. Includes decay masses and partial decay modes.

Table with columns: R1, D INTO (PI PI RHO) / (K KBAR PI), (P2)/(P1), 10/66. Includes branching ratios and partial decay modes.

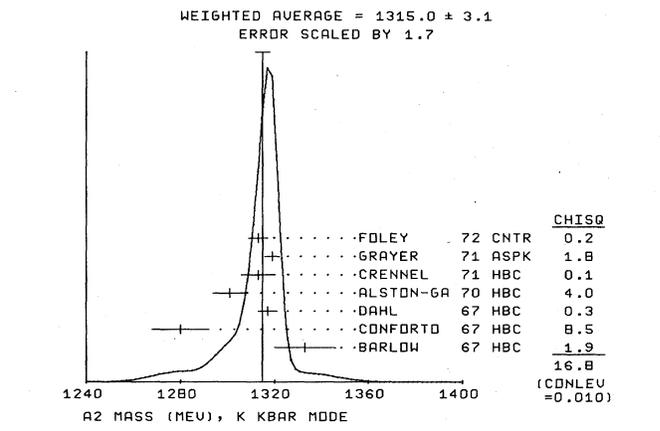
Table with columns: D-ANLAU, 65 PL 17 347, BARLOW, 45 PRL 14 1074. Includes references for D mesons.

A2(1310)

12 A2(1310, JP=2+-) I=1
WE LIST THE A2 AS AN ORDINARY BREIT-WIGNER RESONANCE. FOR DISCUSSION OF THE REPORTED SPLITTING SEE CUR APRIL 72 AND APRIL 73 EDITION.

Table with columns: M, (1320.0), (1301.0), ADERHOLZ, 64 HBC, 4.0 PI+P. Includes student averages and missing mass spectrum.

Table with columns: MK, 80(1317.0), (3.0), BARLOW, 67 HBC, 1.2 PBAR P, KK, 2/72. Includes branching ratios and partial decay modes.



Data Card Listings

For notation, see key at front of Listings.

Mesons

A₂(1310)

Table with columns for mass (MEV), error (ETA PI MODE), and various particle decay modes. Includes entries for ALSTON-GA, CASO, DZIERBA, JOHNSTON, ESPIGAT, KEY, CONFORTO, and MMS.

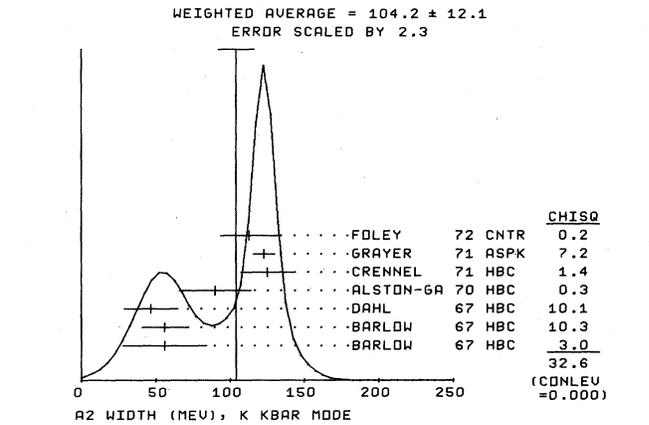
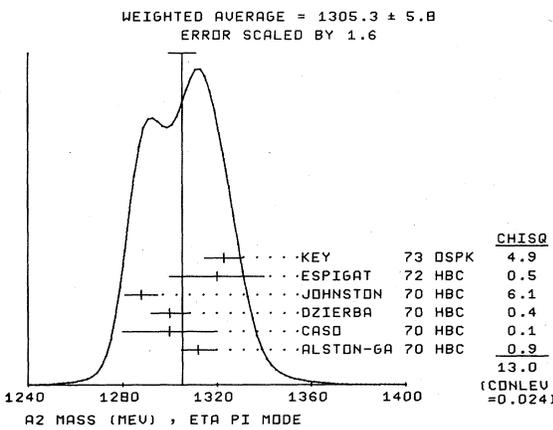


Table with columns for width (MEV), error (3PI MODE), and various particle decay modes. Includes entries for ADERHOLZ, GOLDHABER, LEBEVRES, BARNES, BENSON, LEVRAT, CHIKOVANI, ARMENISE, BOESEBECK, CHUNG, VON KROGH, JURKMANN, ANDERSON, ARMENISE, EISENBERG, ALSTON-GA, BOCKMANN, CASO, DIAZ, GARFINKEL, GORDON, BARNHAM, BINNIE, BINNIE, BOWEN, BOWEN, BOWEN, BLOODWORTH, ANTIPOVI, CHALUPKA, EMMS, WAGNER.

Table with columns for width (MEV), error (ETA PI MODE), and various particle decay modes. Includes entries for ALSTON-GA, CASO, DZIERBA, JOHNSTON, ESPIGAT, KEY, CONFORTO, MMS, and BARLOW.

Table with columns for partial decay modes and decay masses. Includes entries for A2 INTO RHO PI, A2 INTO K KBAR, A2 INTO ETA PI, A2 INTO OMEGA PI, A2 INTO PI+ PI- P, A2 INTO PI+ PI- PI- EXCL. RHO PI, A2 INTO PI GAMMA, A2 INTO ETA PRIME PI, and SMALL, NOT USED IN THE FIT.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS. Includes a matrix of branching fractions and a list of decay masses.

Table with columns for width (MEV), error (K KBAR MODE), and various particle decay modes. Includes entries for BARLOW, BEUSCH, CONFORTO, DAHL, DAHL, CRENNELL, ALSTON-GA, CRENNELL, GRAY, GRAY, FOLEY, and THE NEUTRAL MODE CAN INTERFERE WITH THE F MESON.

Table with columns for branching ratios and various particle decay modes. Includes entries for A2 (CHARGED ONLY) INTO (K KBAR)/(RHO PI), THE NEUTRAL MODE CAN INTERFERE WITH F, A2 INTO (ETA PI)/(RHO PI + K KBAR + ETA PI), and AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0).

Mesons

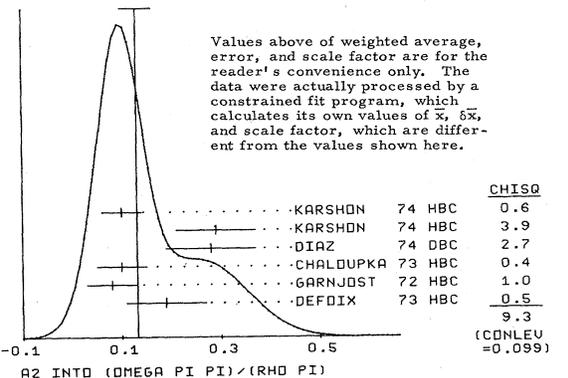
A₂(1310)

Data Card Listings

For notation, see key at front of Listings.

R3	A2 INTO (ETA PI) / (RHO PI)	(P3)/(P1)		
R3	0.3	0.2	ADERHOLZ 64 HBC	4.0 PI+P
R3	0.22	0.09	CONTE 67 HBC	11.0 PI-P
R3	0.25	0.08	ASGOLI 68 HBC	5 PI-P
R3	0.12	0.08	CHUNG 68 HBC	3.2 PI-P
R3	0.16	0.10	KEY 68 HBC	3 PI-P
R3	0.3	0.13	ABRAMVIC 70 HBC	3.93 PI-P
R3	0.18	0.06	VETLITSKY 69 HBC	3.3 PI-P
R3	0.25	0.09	BOCKMANN 70 HBC	5.0 PI+P
R3	0.34	0.17	BOCKMANN 70 HBC	5.0 PI+P
R3	0.39	0.07	DZIERBA 70 HBC	8. PI-P
R3	0.246	0.042	ALSTON-GA 71 HBC	7.0 PI+P
R3	0.211	0.044	CHALDUPKA 73 HBC	3.9 PI-P, P A2
R3	0.22	0.05	ANTIPOV 73 CNTR	40. PI-P, P A2-
R3	0.221	0.021	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R3	0.222	0.023	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R3	0.212	0.015	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R4	A2 INTO (ETA PRIME PI) / TOTAL	(P8)		
R4	0.14 OR LESS		CHUNG 65 HBC	3.2 PI-P
R4	0.02	0.57	BARNHAM 71 HBC	3.7 PI+P
R5	A2 INTO (ETA PRIME PI)/(RHO PI)	(P8)/(P1)		
R5	0.04	0.03	BOCKMANN 70 HBC	5.0 PI+P
R5	0.15	0.09	DZIERBA 70 HBC	8. PI-P
R5	0.04	0.04	ALSTON-GA 71 HBC	7.0 PI+P
R5	0.11	0.04	EISENSTEIN 73 HBC	5.1 PI-P, P 6P1
R6	A2 INTO (PI+ PI- P1) / (RHO PI)	(P5)/(P1)		
R6	0.17	0.08	BENSON 66 DBC	0.37 PI+D
R7	A2 INTO (ETA PI)/(K BEAR)	(P3)/(P2)		
R7	0.31	0.08	FOSTER 68 HBC	PBAR P, PBA REST.
R7	0.319	0.046	FROM FIT	
R8	A2 INTO (K KBAR)/(RHO PI + K KBAR + ETA PI)	(P2)/(P1+P2+P3)		
R8	0.06	0.03	BARNHAM 71 HBC	3.7 PI+P, KSK+P
R8	0.020	0.004	ESPIGAT 72 HBC	0. PBAR, P
R8	NOT AVERAGED BECAUSE OF DISCREPANCY BETWEEN MASSES			
R8	FROM (K KBAR) AND (RHO PI) MODES			
R8	0.03	0.02	DAMERI 72 HBC	11. PI-P
R8	0.05	0.02	TOET 73 HBC	5. PI+P, P K+ KO
R8	0.09	0.04	TOET 73 HBC	0.5 PI+P, P K K B
R8	0.048	0.012	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R8	0.048	0.014	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R8	0.0520	0.0058	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R9	A2 INTO (PI+ PI- PI-)/(RHO PI-)	(P6)/(P1C)		
R9	0.23	0.93	ABRAMVIC 70 HBC	3.93 PI-P
R11	A2 INTO (PI GAMMA)/(TOTAL)	(P7)		
R11	0.005	0.005	EISENBERG 72 HBC	PHOTOPRODUCTION
R11	PION EXCHANGE MODEL USED IN THIS ESTIMATION			
R12	A2 INTO (OMEGA PI PI)/(RHO PI)	(P4)/(P1)		
R12	0.19	0.08	DEFOIX 73 HBC	0.7 PBAR P, 7 P1
R12	DECAYS TO B1(1040) PI, B1 INTO OMEGA PI			
R12	D ERROR INCREASED TO ACCOUNT FOR POSSIBLE SYST. ERRORS			
R12	D OF COMPLICATED ANALYSIS.			
R12	0.08	0.05	GARNJOST 72 HBC	0.7 PI+P
R12	0.10	0.05	CHALDUPKA 73 HBC	3.9 PI-P, P A2
R12	0.28	0.09	DIAZ 74 DBC	0.6 PI+P, P(SPI)0
R12	0.29	0.08	KARSHON 74 HBC	4.9 PI+P, DEL+ A2
R12	0.10	0.04	KARSHON 74 HBC	4.9 PI+P, P A2
R12	K KARSHON 74 SUGGEST AN ADDITIONAL I=0 STATE, STRONGLY COUPLED			
R12	K TO OMEGA PI PI COULD EXPLAIN DISCREPANCIES IN BRANCHING RATIOS			
R12	K AND MASSES.			
R12	0.131	0.032	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
R12	0.127	0.027	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT	
R12	0.131	0.029	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)	
R12	(SEE IDEOGRAM BELOW)			

WEIGHTED AVERAGE = 0.131 ± 0.032
ERROR SCALED BY 1.4



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{x} , $\bar{\sigma}$, and scale factor, which are different from the values shown here.

ADERHCLZ 64	PL 10 248	(AACHEN+BERLIN+ERIM+BCNN+AMBURG+LOIC+MPIM)
CHUNG 64	PL 12 621	+DAHL+HARDY+HESS+KALBFLEISCH+KIRZ (LRL)
GOLDHABE 64	DUBNA CONF 1 480	G GOLDHABER'S GOLDHABER, D'ALLORAN, SHEN(LRL)
LANDER 64	PL 13 346	+ABOLINS, CARMONY, HENDRICKS, XUONG* (LA JOLLA)
ABOLINS 65	ATHENS(CHIO)CONF.	+CARMONY, LANDER, XUONG, YAGER (LA JOLLA)=1
ADERHOLZ 65	PR 138 B 897	(AACHEN+BERL+ERIM+BCNN+AMB+LOIC+MPIM)
ALITTI 65	PL 15 69	ALITTI, BATON, DELER, CRUSSARD* (SACLAY+BGNA) JP
CHUNG 65	PL 15 325	+CAHL, HARDY, HESS, JACOBS, KIRZ, MILLER (LRL)
FRINO 65	PL 19 68	+GESSAROLI+ (BGNA+BAR1+FRIZ+ORS+SACL)
LEFEVRE 65	PL 19 434	CERN MISSING MASS SPECTROMETER GROUP (CERN)
SEIDLITZ 65	PL 15 217	L SEIDLITZ, D I CAHL, D H MILLER (LRL)
BARNES 66	PL 16 41	BARNES, FOWLER, LAI, ORENSTEIN + (BNL+CUNY)
BENSON 66	MICH COO-1112-4	G.C. BENSON, THESIS (MICH)
ALSO 66	PL 16 1177	G BENSON, LOVELL, MARQUIT, ROE + (MICH)
EHRLICH 66	PL 152 1194	R. EHRLICH, W. SELOVE, H. YUTA (PENN)
FERBEL 66	PL 21 111	FERBEL (ROCHESTER)
LEVRAT 66	PL 22 714	CERN MISSING MASS SPECTROMETER GROUP (CERN)
ARMENISE 67	PL 258 53	ARMENISE, FORINO, + (BARI+BGNA+FRIZ+ORSAY)
BALTAY 67	PL 258 160	+KIRSCH+KUNG+YEN+RABIN (COLU+BNL+RUTGERS)
BARLOW 67	NC 50A 701	+ILLIESTOL+MCNANET* (CERN+COEF+IRAD+LVP)
BARTSCH 67	PL 258 48	+DEUTSCHMANN+GROTE+COCGNI+ (AACH+BERL+CERN)
BEUSCH 67	PL 25 8 357	+FISCHER, GOBBI, ASTBURY+ (ETHZ+CERN)
CASON 67	PL 18 880	+LAMS, BISWAS, DERADO, GROVES, + (NOTREDAME)
CHIKOVAN 67	PL 258 44	CERN MISSING MASS SPECTROMETER GROUP (CERN)
CHUNG 67	PL 18 100	+DAHL, HARDY, HESS, KIRZ, MILLER (LRL)
ALSO 66	UCRL-16832	RICHARD I HESS--THESIS, BERKELEY (LRL)
COHN 67	NP 81 57	+MCCULLOCH+BUGG+CONDO (ORNL+UNIV. TENN.)
CONFORTO 67	NP 83 469	+MARECHAL, MCNANET* (CERN+ANL+LVP)
CONTE 67	NC 51 A 175	+TOMASINI, CORDA+ (GENOVA+HAMB+MILAN+SACLAY)
DAHL 67	PL 163 1377	+HARDY+HESS+KIRZ+MILLER (LRL)
DANYSZ 67	NC 51 A 801	DANYSZ+FRENCH+SIMAK (CERN)
SLATTERY 67	NC 50A 377	+KRAYBILL+FRGMAN+FERBEL (YALE+ROCH) JP
ARMENISE 68	PL 268 336	ARMENISE, FORINO, + (BARI+BGNA+FRIZ+ORSAY)
ASCOLI 68	PL 20 1321	+CRAWLEY, MORTARA, SHAPIRO, BRIDGES+ (ILLINOIS) JP
BALLAM 68	PL 21 934	+BRODY, CHADWICK, FRIES, GUIRGOSSIAN+ (SLAC)
BENI 68	PL 28 B 233	CERN MISSING MASS SPECTROMETER GROUP (CERN)
BOESEBECK 68	NP 8 501	BOESEBECK, DEUTSCHMANN, + (AACHEN+BERLIN+CERN)
CASO 68	NC 54 A 983	+CONTE+CORDS+DIAZ+ (GENOVA+HAMB+MILAN+SACL)
CHUNG 68	PL 165 1491	S.U. CHUNG, O. DAHL, J. KIRZ, D. H. MILLER (LRL)
CRENNELL 68	PL 20 1318	+KARSHON+KWAN LAI, SCARR, SKILLICORN (BNL)
DONALD 68	PL 26 B 327	+FROESE+BETTINI+ (LIVERPOOL+OSLO+PAU)
FOSTER 68	NP 8 B 174	+GAVILLET, LABROSSE, MCNANET, + (CERN+COEF)
FRIDMAN 68	PL 167 1268	+MAURER, MICHALON, OUDET+ (HEID +STRASBOURG)
JUNKMANN 68	NP 88 471	+COCCONI, + (AACH+BERL+BCNN+CERN+WAR)
KEY 68	PL 166 1430	+PRENTICE+COOPER+MANNERS+ (INTO+ANL+LVP)
LAMSA 68	PL 166 1395	+CASON+BISWAS+DERADO+GROVES+ (NOTREDAME)
VGN KRGG 68	PL 27 B 253	+MIYASHITA, KOPPELMAN, MARSHALL LIBBY (COLO)
ADERHOLZ 69	NP 8 11 259	+BARTSCH, + (AACH+BERL+CERN+FRIZ+ORSAY)
AGUILAR 69	PL 29 B 62	+BARLOW, JACOBS, DELLA NEGRA+ (CERN+COEF+LVP)
AGUILAR 269	PL 29 B 241	M. AGUILAR-BENITEZ, J. BARLOW, + (CERN+COEF)
ANDERSON 69	PL 22 1390	+COLLINS, + (BNL+CERN)
ARMENISE 69	LNC 2 501	+G. IDINI, FORINO, CARTACCI+ (BARI+BGNA+FRIZ)
CHIKOVAN 69	PL 28 B 526	CERN MISSING MASS SPECTROMETER GROUP (CERN) JP
CRENNELL 69	PL 22 1327	+KARSHON, KWAN WU LAI, + (BNL) JP
DONALD 69	NP 8 12 325	+EDWARDS, FOSTER, MOORE (LIVERPOOL)
EISENBERG 69	PL 23 1322	EISENBERG, HABER, BALLAM, CHADWICK, + (REHO+SLAC)
VETLITSKY 69	SJNP 9 596	VETLITSKY, GRI GOREYEV, GRISHIN, + (ITEP)
ABRAMOVIC 70	NP 8 29 466	ABRAMOVICH, BLUMENFELD, BRUYANT, + (CERN) JP
ALSTON-GATO 70	PL 33 B 607	+BARBARO, BUEHL, DERENZO, EPPERSON, FLATTE+ (LRL)
ASCOLI 70	PL 25 962	+BRACKWAY, CRAWLEY, EISENSTEIN, HANFT, + (ILL) JP
BASILE 70	LNC 4 838	+CALPIAZ, FRABETTI, MASSAM, + (CERN+BGNA+STR)
BAUDI 70	PL 31B 397	CERN BOSON SPECTROMETER GROUP (CERN)
BAUD2 70	PHILAD. CONF. P. 311	CERN BOSON SPECTROMETER GROUP (CERN)
BAUD3 70	PL 31 B 401	CERN BOSON SPECTROMETER GROUP (CERN)
BOCKMANN 70	NP 8 16 221	+MAJOR, POLS, + (BNL+DURH+NIJ+EPOL+TORI)
BUTLER 70	UCRL 19845	THESIS (LRL)
CAROLL 70	PL 25 1393	+FIREBAUGH, GARFINKEL, MORSE, OH, + (WISC+INTO)
CASO 70	LNC 3 707	+CONTE, TOMASINI, CORDS+ (GENOVA+HAMB+MILAN+SACL)
DIAZ 70	NP 8 16 239	+GAVILLET, LABROSSE, MCNANET, + (CERN+COEF) JP
DZIERBA 70	PL 2 2 2544	+SHEPARD, BISWAS, CASO, JOHNSON, KENNEY (NDAM)
ALSO 68	LAMSA	
GARFINK 70	PL 33 B 536	GARFINKEL, AMMANN, CARMONY, YEN (PURD) JP
GORDON 70	COO 1195 179	THESIS, ILLINOIS (ILL)
JOHNSTON 70	NP 8 24 253	+KEY, PRENTICE, YOON, GARFINKEL, + (INTO+WISC)
KRUSE 70	PHILAD. CONF. P. 359	U. KRUSE, PARTIAL WAVE ANALYSIS (ILL) JP
NEF 70	THESIS+PRIV. COMM.	CERN BOSON SPECTROMETER GROUP (CERN)
SUTHERLA 70	PHILAD. CCF. P. 369	G. SUTHERLAND, INTERFERING RESONANCE (GLASGOW)
AGUILAR 71	PL D 4 2583	AGUILAR-BENITEZ, EISNER, KINSON (BNL)
ALSTON-GAT1 71	PL 34 B 156	+BARBARO, BUEHL, DERENZO, EPPERSON, FLATTE+ (LRL) JP
BARNHAM 71	PL 26 1494	+ABRAMS, BUTLER, COYNE, GOLDHABER, HALL, + (LRL)
BEKETOV 71	SJNP 4 765	+SOMBRONSKY, KONVALOV, KRUTSCHIN, + (ITEP)
BINNIE 71	PL 36 B 257	+CAMILLETTI, DUANE, FARUQI, BURTON, + (LOIC+SHMP)
BINNIE2 71	PL 36 B 537	+CAMILLETTI, DUANE, FARUQI, BURTON, + (LOIC+SHMP)
BOWEN 71	PL 26 1663	+FARLES, FAISSLER, BLIEDEN, + (NEAS+STON)
CRENNEL 71	PL 35 B 185	+GORDON, KWAN WU LAI, SCARR (BNL)
FARBER 71	NP 8 25 237	+DE PINTO, BISWAS, CASON, DEERY, KENNEY, + (NCAN)
FOLEY 71	PL 26 413	+LOVE, OZAKI, PLATNER, LINDENBAUM, + (BNL+CUNY)
GRAYNER 71	PL 34 B 333	+HYAMS, JONES, SCHLEIN, BLUM, DIETL, + (CERN+MPIM)
LYNCH 71	UCRL 20022 AND 71	G. LYNCH (LBL)
RINAUDO 71	NC 5 A 239	+BOECKMANN, MAJOR, + (TORI+BNL+DURH+NIJ+EPOL) JP
ANKENBRA 72	PL 29 1688	ANKENBRANDT, BRABSON, CRITTTENDEN, HEINZ, + (IND)
BERENYI 72	NP 8 37 621	+PRENTICE, STEENBERG, YOON, WALKER (INTO+WISC)
BLOODWORTH 72	NP 8 37 203	BLOODWORTH, JACKSON, PRENTICE, YOON (INTO)
DAMERI 72	NC 9 A 1	+BORIATTA, GOUSSU, + (GENOVA+MILAN+SACL)
DAMGAARD 72	UNPUBLISHED MEMO	+LECHANOINE, MARTIN (BOHR+GEVA)
DEIBOLD 72	BATAV. CONF.	R. DIEBOLD, RAPPORTEUR TALK (ANL)
EISENBER 72	PL D 5 15	EISENBERG, BALLAM, DAGAN, + (REHO+SLAC+TELA)
ESPIGAT 72	NP 8 36 29	+GHEZDELIER, L. ILLIESTOL, MONTANET (CERN+COEF)
FOLY 72	PL D 6 747	+LOVE, OZAKI, PLATNER, LINDENBAUM, + (BNL+CUNY)
GARNJOST 72	PL 29 1491	M. ALSTON-GARNJOST (LBL)
GETTNER 72	PREPRINT NUB2145	M. GETTNER (NEAS)
KIENZLE 72	UNPUBLISHED MEMO	W. KIENZLE (INDA)
LASSILA 72	PL 28 1491	LASSILA, YOUNG (INDA)
MORSE 72	NP 8 43 77	+OH, WALKER, JOHNSTON, YOON (WISC+INTO)

REFERENCES FOR A2

Data Card Listings

Mesons

For notation, see key at front of Listings. A2(1310), E(1420), X0(1430), X1(1440), f'(1514)

AMMANN 73 PR D 7 1345
ANKERMAN 73 PR D 8 2785
ANTIPOV 73 NP B 63 175
...
WAGNER 75 PL 588 201

E(1420) 6 E(1420, JPG=A) I=0
BAILLON 67 FAVOR JP=0-. DAHL 67 FAVOR 1+ BUT DO NOT
EXCLUDE 2-, 0-, LORSTAD 69 FIND 0- CR 1+.
VUILLEMI 75 FAVOUR 1+.

6 E MASS (MEV)
M 1425. 7. BAILLON 67 HBC 0. PBAR P 11/66
M 1420. 20. DAHL 67 HBC 1.6-4.2 PI- P 9/66
M 1423.0 10.0 FRENCH 67 HBC 3-4 PBAR P 6/67
...
M STUDENT 1415.5 4.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)
1416.0 4.3 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

6 E WIDTH (MEV)
W 80. 10. BAILLON 67 HBC 0. PBAR P 11/66
W 60.0 20.0 DAHL 67 HBC 1.6-4.2 PI- P 10/66
W 45. 20. FRENCH 67 HBC 3-4 PBAR P 6/67
...
W STUDENT 58.3 6.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
58.3 6.7 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

6 E PARTIAL DECAY MODES
P1 E INTO K K*(892)
P2 E INTO K KBAR PI
P3 E INTO PI PI RHO
P4 E INTO DELTA PI
P5 E INTO ETA PI PI

6 E BRANCHING RATIOS
R1 E INTO (KBAR K*(892) + C.C.)/(K KBAR PI) (P1)/(P2)
R1 .50 .10 BAILLON 67 HBC 0.0 PBAR P 11/66
R2 E INTO (PI PI RHO) / (K KBAR PI) (P3)/(P2)
R2 (2.0) OR LESS DAHL 67 HBC 0 1.6-4.2 PI- P .10/66
...
R4 E INTO (DELTA PI)/(ETA PI PI) (P4)/(P5)
R4 0.4 0.2 DEFCIX 72 HBC 0.7 PBAR P,7 PI 1/73

REFERENCES FOR E
BAILLON 67 NC 50A 393 +EDWARDS+D.ANDLAU+ASTIER+ (CERN+CDEF+IRAD)
BARASH 67 PR 156 1399 BARASH,KIRSCH,MILLER,TAN (COLUMBIA)
DAHL 67 PR 163 1377 +HARDY+ESS+KIRZ+MILLER (LRL II)
...
HANDLER 76 NP +PLANO,BRUCKER,KOLLER+ (RUTG+STEV+SETO)

X0(1430) -> KS KS, rho rho
29 X(1430, JPG=) I=0
EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.
PEAKS SEEN IN (KS KS) SPECTRA QUOTED UNDER
X(1440) (I=1) AS WELL.

29 X(1430) MASS (MEV)
M -----RHOD RHOD MODE-----
M (1410.0) BETTINI 66 DBC 0 0. PBAR N,5 PI 2/74
M B POSSIBLY SEEN ABRAMS 67 HBC 4.25 K- P 5/67
M B THE AUTHORS ASSOCIATE THE PEAK WITH THE F PRIME, BUT BACKGROUND
M B ESTIMATION IS DIFFICULT
...
M STUDENT 1437.7 5.9 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
1437.7 5.9 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

29 X(1430) WIDTH (MEV)
W -----RHOD RHOD MODE-----
W (90.0) BETTINI 66 DBC 0 0. PBAR N,5 PI 2/74
W B POSSIBLY SEEN ABRAMS 67 HBC 4.25 K- P 5/67
W B THE AUTHORS ASSOCIATE THE PEAK WITH THE F PRIME, BUT BACKGROUND
W B ESTIMATION IS DIFFICULT
...
W STUDENT 46.2 18.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
46.2 18.5 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

REFERENCES FOR X(1430)
BETTINI 66 NC 42A 695 +CRESTI,LIMENTANI,LORIA,PERUZZO+(PACO+PISA)
ABRAMS 67 PRL 18 620 +KEHOE,GLASSER,SECHI-ZORN,WOLSKY (MARYLAND)
BARLOW 67 NC 50 A 701 +MCNTANET,C-ANDLAU+ (CERN+CDEF+IRAD+LIVP)
BEUSCH 67 PL 25 B 357 +FISCHER,GOBBI,ASTBUR+ (ETHZ+CERN)
DONALD 69 NP 8 11 551 +EDWARDS,BURAN,BETTINI,+ (LIVP+OSLO+PADO)

X1(1440) -> KS KS
38 X(1440, JPG=) I=1
EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.
PEAKS SEEN IN (KS KS) SPECTRA QUOTED UNDER
X(1430) (I=0) AS WELL.

38 X(1440) MASS (MEV)
M B POSSIBLY SEEN ABRAMS 67 HBC 4.25 K- P 5/67
M B THE AUTHORS ASSOCIATE THE PEAK WITH THE F PRIME, BUT BACKGROUND
M B ESTIMATION IS DIFFICULT
...
M STUDENT 1437.7 5.9 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
1437.7 5.9 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

38 X(1440) WIDTH (MEV)
W 100. 70. BARLOW 67 HBC 1.2 PBAR P 5/67
W 43.0 17.0 18.0 BEUSCH 67 OSPK 5.7,12 PI-P 9/67
W (20.0) OR LESS FOLEY 71 CNTR - 20.3 PI- P,K- KS 2/71
...
W STUDENT 46.2 18.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
46.2 18.5 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

REFERENCES FOR X(1440)
ABRAMS 67 PRL 18 620 +KEHOE,GLASSER,SECHI-ZORN,WOLSKY (MARYLAND)
BARLOW 67 NC 50 A 701 +MCNTANET,D.ANDLAU+ (CERN+CDEF+IRAD+LIVP)
BEUSCH 67 PL 25 B 357 +FISCHER,GOBBI,ASTBUR+ (ETHZ+CERN)
FOLEY 71 PRL 26 413 +LOVE,OZAKI,PLATNER,LINDENBAUM,+ (BNL+CERN)
DEFOIX 73 PL 43 B 141 +DOBZYNSKI,ESPIGAT,NASCIMENTO,+ (CDEF)

f'(1514) 13 F PRIME(1514, JPG=2+) I=0
13 F PRIME MASS (MEV)
M 14(1480.0) CRENNELL 66 HBC 6.0 PI- P 8/66
M B 5(1460.) (10.) ABRAMS 67 HBC 4.25 K- P 5/67
M B BACKGROUND ESTIMATION DIFFICULT.
...
M STUDENT 1516.1 3.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
1516.1 3.1 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho'(1600)$, $A_3(1640)$

65 RHO PRIME MASS (MEV)

M C	(1600.)	APPROX.	BARBARINO 72	OSPCK	0 E+ E- TO 4 PI	1/73
M C	400	1430.	50.	BINGHAM 72	HBC	0 9.3 GAM P,P 4PI 12/72
M	1590.	20.	HYAMS 73	ASPCK	0 17 PI-P,N PI+PI-	1/74
M	1550.	60.	CONVERSI 74	OSPCK	0 E+ E- TO 4PI	12/75*
M	160	1550.	50.	SCHACHT 74	STRC	05.5-9 G P,P 4PI 12/75*
M	340	1450.	130.	SCHACHT 74	STRC	09-18 G P,P 4PI 12/75*
M	65	1570.	60.	ALEXANDER 75	HBC	07.5 GAM P,P 4PI 12/75*
M P	(1600.)	(50.)	FRÖGGATT 75	RVUE	17 PI-P,PI+PI-N	12/75*
M P	FROM PHASE SHIFT ANALYSIS OF HYAMS 73 DATA					
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)					

65 RHC PRIME WIDTH (MEV)

W	400	650.	100.	BINGHAM 72	HBC	0 9.3 GAM P,P 4PI 12/72
W	180.	50.	HYAMS 73	ASPCK	0 17 PI-P,N PI+PI-	12/75*
W	360.	100.	CONVERSI 74	OSPCK	0 E+ E- TO 4PI	12/75*
W E	160	400.	120.	SCHACHT 74	STRC	05.5-9 G P,P 4PI 12/75*
W E	340	850.	200.	SCHACHT 74	STRC	09-18 G P,P 4PI 12/75*
W E	65	340.	160.	ALEXANDER 75	HBC	07.5 GAM P,P 4PI 12/75*
W P	(220.)	(70.)	FRÖGGATT 75	RVUE	17 PI-P,PI+PI-N	12/75*
W P	FROM PHASE SHIFT ANALYSIS OF HYAMS 73 DATA					
W E	WIDTH ERRORS ENLARGED BY US TO 4*WIDTH/SORT(N)					
W	AVERAGE MEANINGLESS (SCALE FACTOR = 2.3)					

65 RHO PRIME PARTIAL DECAY MODES

P1	RHO PRIME INTO RHO 0 PI	PI	139* 139*	773
P2	NEUTRAL RHO PRIME INTO ALL 4 CHARGED PI MODES	139* 139* 139* 139*	773* 773*	
P3	RHO PRIME INTO RHO RHO	139* 139*		
P4	RHO PRIME INTO PI PI	492* 493		
P5	RHO PRIME INTO KBAR K	139* 783		
P6	RHO PRIME INTO PI 0 MEGA	139* 139*	134*	134*
P7	RHO PRIME INTO PI+ PI- PI0 PI0			

65 RHC PRIME BRANCHING RATIOS

R1	RHO PRIME INTO (RHO 0 PI+ PI-)/(4 PI, ALL CHARGED)	(P1)/(P2)	1/73	
R1 S	DOMINANT (1.80)	BINGHAM 72	HBC	0 E+ E- TO 4 PI 1/73
R1 S	500 0.7 0.1	SCHACHT 74	STRC	05.5-18 G P,P 4PI 12/75*
R1 S	THE PI PI SYSTEM IS IN S WAVE 1/73			
R2	RHO PRIME INTO (RHO 0 RHO 0)/(RHO 0 PI+ PI-)	(P3)/(P1)	1/73	
R2	NCNE (FORBIDDEN BY I=1)	BINGHAM 72	HBC	0 9.3 GAM P,P 4PI 1/73
R3	RHO PRIME INTO (PI+ PI-)/(4 PI, ALL CHARGED)	(P4)/(P2)	1/73	
R3	(0.14) OR LESS ESTIMATE DAVIER	73	STRC	0 6-18 G P,P 4PI 1/74
R4	RHO PRIME INTO (KBAR K)/(4 PI, ALL CHARGED)	(P5)/(P2)	1/73	
R4	(0.04) OR LESS 2 SIGMA	BINGHAM 72	HBC	0 9.3 GAM P 1/73
R5	RHO PRIME INTO (PI+PI-)/TOTAL	(P4)	1/74	
R5 E	(0.15) OR LESS	EISENBERG 73	HBC	0 5 PI+ P,DEL++2PI 1/74
R5 E	ESTIMATED USING OPE MODEL	HYAMS 73	ASPCK	0 17 PI-P,N PI+PI- 1/74
R5 P	(0.25) (0.10)	FRÖGGATT 75	RVUE	17 PI-P,PI+PI-N 12/75*
R5 P	FROM PHASE SHIFT ANALYSIS OF HYAMS 73 DATA			
R6	RHO PRIME INTO (PI+ PI- PI0 PI0)/(4PI, ALL CHARGED)	(P7)/(P2)	1/74	
R6	(1.) OR LESS ESTIMATE DAVIER	73	STRC	0 6-18 G P,P 4PI 1/74
R7	RHO PRIME INTO (PI+ PI- + NEUTRALS)/(4PI, ALL CHARGED)	(P7+...)/(P2)	12/75*	
R7 U	(2.6) (0.4)	BALLAM 74	HBC	0 9.3 GAMMA P 12/75*
R7 U	UPPER LIMIT. BACKGROUND NOT SUBTRACTED			

REFERENCES FOR RHO PRIME

ALVENSLE 71	PRL 26 273	ALVENSLEBEN, BECKER, BERTRAM, CHEN, + (DESEY+MIT) G
BAUON 71	NP 830 213	*FRIDMAN, GERBER, GIVERNAUD, + (STRASBOURG) G
BULOS 71	PR 26 149	*BUSZA, KEHOE, BENISTON, + (SLAC+UMD+IBM+LBL) G
BACCII 72	PL 388 551	*PENSO, SALVINI, STELLA, BALDINI, CE(ROMA+FRAS) JPC
BARBARINO 72	LNC 3 689	BARBARINO, CERADINI, + (FRAS+ROMA+PADO+UMD) IGJP
BARTOLI 72	PR D 6 2374	*FELICETTI, OGGREN, + (FRAS+ROMA+NAPL) IGJP
BINGHAM 72	PL 418 635	*RABIN, ROSENFIELD, SMOJJA, YCST+(LBL+UCB, SLAC) IGJP
BRAMON 72	LNC 3 693	*GRECO (THEORETICAL PAPER) (FRASCATI) J
DIEBOLD 72	BATAV, CONF.	R. DIEBOLD, RAPPORTEUR TALK (ANL)
EISENBERG 72	PR D 5 15	EISENBERG, BALLAM, DAGAN, + (REHO+SLAC+TELA)
LYSSAC 72	NC 10A 407	J. LAYSSAC, F. M. RENARD (MONT)
SMAJJA 72	PHIL. CONF. PROC349	*BINGHAM, FRETTER, BALLAM, CHADWICK+(LBL+SLAC)
CERADINI 73	PL 43 B 341	*CONVERSI, EKSTRAND, GRILLI, + (ROMA+FRAS+PADO) IGJP
DAVIER 73	NP 8 58 31	*DERADO, FRIES, LIU, MOZLEY, ODIAN, PARK, + (SLAC)
EISENBERG 73	PL 43 B 149	EISENBERG, KARSHCN, MIKENBERG, PITLUCK, + (REHO)
HYAMS 73	NP B 64 134	*JONES, WEILHAMMER, BLUM, DIETL, + (CERN+MPIIM)
KREUZER 73	PR D 8 1431	H. J. KREUZER, A. N. KAMAL (UNIV. OF ALBERTA)
OCHS 73	THESES	J. C. H. PARK (MPIIM)
PARK 73	NP 8 58 45	J. C. H. PARK (MPIIM) JP
BALLAM 74	NP 876 375	*CHADWICK, BINGHAM, FRETTER+ (SLAC+LBL+MPIIM)
BERNABEI 74	LNC 11 261	*D. ANGULO, SPILLANTINI, VALENTE (ROMA+FRAS)
CONVERSI 74	PL 528 493	*PAGLUZZI, CERADINI, GRILLI+ (ROMA+FRAS)
ESTABROO 74	NP 879 301	P. ESTABROOKS, A. C. MARTIN (DURH)
FERBEL 74	PR D9 824	T. FERBEL AND P. SLATTERY (ROCH)
GRAYER 74	NP 8 75 189	G. GRAYER, HYAMS, BLUM, DIETL, + (CERN+MPIIM)
HIRSHFEL 74	NP 874 211	A. C. HIRSHFELD, G. KRAMER (HAMB)
SCHACHT 74	NP 881 205	*DERADO, FRIES, PARK, YOUNG (MPIIM)
ALEXANDE 75	PL 578 487	ALEXANDER, BENARY, GANDSMAN, LISSAUER+ (TELA)
ALLES 75	NC 30A 136	ALLES-BORRELLI, BERNARDINI+ (CERN+BGNA+FRAS)
ESTABROO 75	NP 895 322	P. ESTABROOKS, A. C. MARTIN (DURH)
FRÖGGATT 75	NP 891 454	C. D. FRÖGGATT, J. L. PETERSEN (GLAS+NORD)
HYAMS 75	NP 8100 205	*JONES, WEILHAMMER, BLUM, DIETL+ (CERN+MPIIM)
LANG 75	PL 588 450	C. B. LANG, I. S. STEFANESCU (KARL)
LANGACKER 75	PREPRINT UPR00497	P. LANGACKER, G. SEGRE (PENN)
ROOS 75	NP 8 97 165	M. ROOS (HELS)

A₃(1640)

34 A3(1640, JPC=2-1) I = 1

The $A_3(1640)$ is seen as a bump in the diffractive-like process $\pi N \rightarrow (\pi\pi\pi)N$. The dominant effect is a ~ 300 MeV wide enhancement in the $J^P=2^-$ $f\pi$ S-wave system, starting from $f\pi$ threshold. Neither additional (narrower) structure in the 3π mass distribution, nor other decay modes, have been clearly established. The situation would appear to be similar to that of the A_1 , but there are certain additional complications:

- Experiments with incident π^- observe little variation of the $J^P=2^-$ $f\pi$ phase in the A_3 mass region (ANTIPOV1 73, ASCOLI1 73), whereas experiments with incident π^+ show evidence for variations consistent with a resonance interpretation (OTTER 74, THOMPSON 74).
- The A_3 region is not well described by the Deck-like model of Ascoli et al., although the agreement could probably be improved by a series of more or less well motivated adjustments to the model (ASCOLI2 73, ASCOLI 74). Nevertheless, relative phase variations of $\sim 40^\circ$ through the A_3 mass region are predicted for the 2^- $f\pi$ wave.
- Other partial waves contribute strongly in the A_3 region. The 2^- $\rho\pi P$ wave may exhibit a broad enhancement; no phase variation relative to 2^- $f\pi$ is observed (ANTIPOV1 73, ASCOLI1 73, OTTER 74, THOMPSON 74). ANTIPOV1 73 show some evidence for an enhancement in the 2^+ $f\pi$ P wave with $M \sim 1750$ MeV, $\Gamma \sim 200$ MeV. The relative phases are not inconsistent with a resonance interpretation.

34 A3 MASS (MEV)

M	30(1600.0)	FORING 65	DBC	04.5 PI+ D	10/66	
M	20(1630.0)	VETLITSKY 66	HBC	- 4.7 PI- P	12/75*	
M	(1630.)	BALTAY 68	HBC	+ 7. 8.5 PI+ P	12/75*	
M	(1600.0)	BARTSCH 68	HBC	+ 8. PI+ P, 3PI P	12/75*	
M	(1610.)	LAMSA 68	HBC	- 8.0 PI-P, PI-F	12/75*	
M	297(1673.0)	ARMENISE 69	DBC	+ 5.1 PI+D, 3PI++	12/75*	
M	(1680.)	CASO 69	HBC	- 11 PI- P	5/70	
M	(1660.0)	CASO 69	HBC	- 11 PI-P, PI-F	12/75*	
M	(1645.0)	CRENNELL 70	HBC	- 6. PI- P, PI-F	12/75*	
M	(1633.0)	MIYASHITA 70	HBC	- 6.7 PI-P, PI-F	1/71	
M	BACKGROUND SUBTRACTION DIFFICULT.					
M	(1672.0)	BEKETOV 71	HBC	- 4.45 PI- P	11/71	
M	(1600.)	PAER 71	DBC	+ 13. PI+ D, D13PI++	12/75*	
M	263(1660.)	CASO 72	HBC	+ 11.7 PI+ P	12/75*	
M	(1658.)	HARRISON 72	HBC	- 13..20. PI- P	12/72	
M	F FIT ASSUMES AN ADDITIONAL PEAK AT 1830 MEV.					
M	EVIDENCE FOR A SUBSTANTIAL DECAY INTO S PI CLAIMED					
M	(1650.)	ANTIPIV1 73	CNTR	- 25..40. PI- P	12/75*	
M	(1600.)	ASCOLI 1 73	HBC	- 5..25. PI- P, P A3	12/75*	
M	(1600.)	THOMPSON 74	HBC	+ 13. PI+ P, P A3	12/75*	
M	EVIDENCE FOR A ROTATION OF THE PHASE CLAIMED.					
M	575(1640.)	KALEKAR 75	HBC	+ 15 PI+P, P1F	12/75*	
M	FROM A FIT TO JP=2- F PI PARTIAL WAVE					
M	AVERAGING NOT MEANINGFUL					

Mesons

A₃(1640), ω(1675)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various decay parameters. Includes entries for A3(1640) and ω(1675) with detailed decay data.

Table titled '34 A3 PARTIAL DECAY MODES' showing decay channels like A3 INTO 3 PI and A3 INTO RHO PI with associated decay masses.

Table titled '34 A3 BRANCHING RATIOS' providing branching ratios for various decay modes of the A3 meson.

Table titled 'REFERENCES FOR A3' listing various scientific publications and sources related to the A3 meson.

Table listing references for the ω(1675) meson, including authors like ALEXANDE, ARNEMISE, and CASO.

***** REFERENCES FOR OMEGA(1675) *****

ω(1675) → ρ⁰π⁰ 45 OMEGA(1675, JPC=3--) I=0. THIS RESONANCE OVERLAPS IN ITS 3PI MODE WITH THE A3, BUT IN SOME EXPERIMENTS ONE CAN ESTABLISH THE DECAY MODE RHO PI0, THUS I=0, WAGNER 74 FINES JP=3- UNIQUELY. THE DECAYS INTO 5PI AND OMEGA PI+PI- NEED FURTHER CONFIRMATION (SEE ALSO XI16901).

Table titled '45 OMEGA(1675) MASS (MEV)' showing mass measurements from various experiments like ARNEMISE, KENYON, and MATTHEWS.

Table titled '45 OMEGA(1675) WIDTH (MEV)' showing width measurements from various experiments like ARNEMISE, KENYON, and MATTHEWS.

Table titled '45 OMEGA(1675) PARTIAL DECAY MODES' showing partial decay widths for various modes like OMEGA(1675) INTO 3 PI and OMEGA(1675) INTO 5 PI.

Table titled '45 OMEGA(1675) BRANCHING RATIOS' providing branching ratios for various decay channels of the ω(1675) meson.

Table titled 'REFERENCES FOR OMEGA(1675)' listing scientific references for the ω(1675) meson, including authors like ARNEMISE, KENYON, and MATTHEWS.

Data Card Listings
For notation, see key at front of Listings.

Mesons
g(1680)

g(1680)

15 G(1680.JPG = 3+-) 1=1

This entry contains the 2pi, 4pi, omega pi, KK-bar, and KK-bar pi peaks in the region of 1700 MeV. The spin-parity determination and the mass and width in the Meson Table come from the 2pi mode. An elasticity of 24% is found at resonance in the pi pi elastic partial-wave analysis (HYAMS 75); this is consistent with the assumption that at least some of the effects listed are due to g decay into various channels. On the other hand, some discrepancies in masses, widths, and branching ratios reported indicate that there may be more than one I_G=1+ meson in this region (BARNHAM 70, HOLMES 72, THOMPSON 74). Although we have collected all the data here under a common entry, we do not imply that they are necessarily all related.

15 G MASS (MEV)

Table with columns for mass (MEV), width, and various decay modes (e.g., HBC, DBC, ASPK, RVUE, ASK).

MASS ERRORS ENLARGED BY US TO WIDTH/SQRT(N)
FROM PHASE-SHIFT ANALYSIS
ERROR TAKES ACCOUNT OF SPREAD OF DIFFERENT PHASE-SHIFT SOLUTIONS
USES SAME DATA AS HYAMS 75

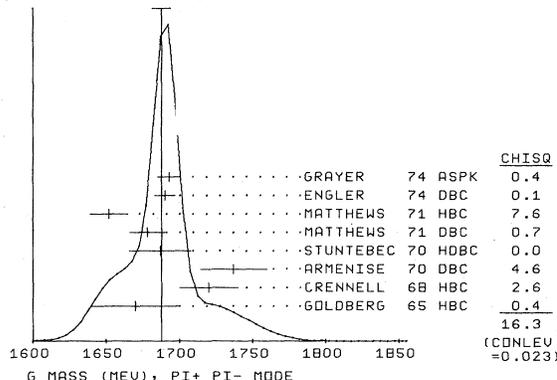
Table for (2PI)++ MODE with columns for mass, width, and decay modes.

Table for K KBAR + K KBAR PI MODE with columns for mass, width, and decay modes.

Table for (4PI)++ MODE with columns for mass, width, and decay modes.

FROM (RHOD-- RHOD) MODE
AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)
AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

WEIGHTED AVERAGE = 1687.8 +/- 6.4
ERROR SCALED BY 1.5



WEIGHTED AVERAGE = 1677.1 +/- 13.0
ERROR SCALED BY 2.0

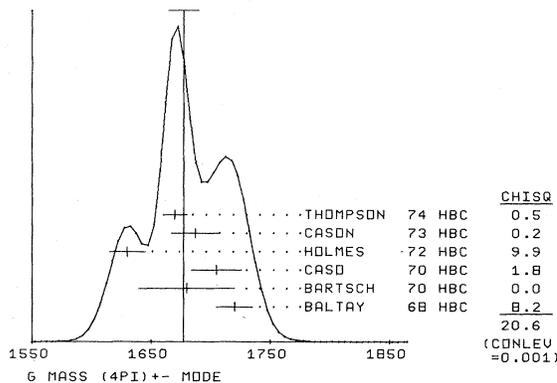


Table for RHOD RHOD MODE with columns for mass, width, and decay modes.

Table for OMEGA PI MODE with columns for mass, width, and decay modes.

R PEAKS FROM MMS. (FOR DIFFICULTIES WITH MMS EXPTS. SEE THE A2 MINI-REVIEW IN THE 1973 EDITION)
M NR1 (1632.) (15.) FOCACCI 66 MMS - 7-12 PI-P,P MMS 12/72
M NR2 (1700.) (15.) FOCACCI 66 MMS - 7-12 PI-P,P MMS 12/72
M NR3 (1748.) (15.) FOCACCI 66 MMS - 7-12 PI-P,P MMS 12/72
M N NCT SEEN BY BOWEN 72
M R (1700.) (47.0) ANDERSON 69 MMS - 16 PI- P, BACKW 8/69

15 G WIDTH (MEV)

Table for PI+ PI- MODE with columns for mass, width, and decay modes.

FROM PHASE-SHIFT ANALYSIS
ERROR TAKES ACCOUNT OF SPREAD OF DIFFERENT PHASE-SHIFT SOLUTIONS
USES SAME DATA AS HYAMS 75
AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Mesons
g(1680)

Data Card Listings

For notation, see key at front of Listings.

W	(2PI)++ MODE								
W	200.0	100.0	CRENNELL	68 HBC	- 6.0 PI- P	12/68			
W	122	180.0	BARTSCH	70 HBC	+ 8 PI+ P, 2 PI	5/70			
W	(42.1)	(20.1)	THOMPSON	74 HBC	+ 13 PI+ P	12/75*			
W	AVG	181.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
W	STUDENT	181.6	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT						
W	K KBAR + K KBAR PI MODE								
W	F	(120.)	APPROX.	FRENCH	67 HBC	(KO K+-) 3-4 PBAR P	11/69		
W	F	ABOVE VALUE ESTIMATED FROM FIG. 9 OF FRENCH 67							
W	79.0	70.0	25.0	CRENNELL	68 HBC	+ 6.0 PI- P, KBAR K	12/68		
W	112.0	60.0		ADERHOLZ	69 HBC	+ 8 PI+ P, KKBARPI	8/69		
W	205	20.0		BLUM	75 ASPK	018.4 PI- P, N K+-	11/75*		
W	AVG	179.6	33.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)					
W	STUDENT	182.6	23.1	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
W	(4PI)++ MODE								
W	J	100.	35.	BALTAY	68 HBC	+ 7, 8.5 PI+ P	6/68		
W	J	(90.0)	(20.0)	JOHNSTON	68 HBC	- 7.0 PI- P	6/68		
W	J	NOT SEPARATED FROM 2 PI DECAY							
W	144	135.0	30.0	BARTSCH	70 HBC	+ 8 PI+ P, 4 PI	4/71		
W	90	(180.0)	(30.0)	BARTSCH	70 HBC	+ 8 PI+ P, A2 PI	4/71		
W	F	102	(160.0)	(30.0)	BARTSCH	70 HBC	+ 8 PI+ P, 2 RHO	4/71	
W	F	(160.0)		CASO	70 HBC	- 11.2PI- P, RHO 2PI	5/70		
W	300	(200.)		ARMENISE	72 HBC	- 9.1 PI- P, P 4PI	12/72		
W	130.	30.		HCLMES	72 HBC	+ 10. -12. K+ P	1/73		
W	169.	70.	48.	CASON	73 HBC	- 8.1, 8.5 PI- P	1/74		
W	F	89.1	(83.1)	(35.1)	CASON	73 HBC	- 8.1, 8.5 PI- P	1/74	
W	F	66	(150.)		KLIGER	74 HBC	- 4.5 PI- P, P 4PI	12/75*	
W	F	106.	25.		THOMPSON	74 HBC	+ 13 PI+ P	12/75*	
W	F	FROM (RHO+- RHO0) MCDE							
W	AVG	120.6	14.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)					
W	STUDENT	120.6	15.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
W	OMEGA PI MODE								
W	F	130.	73.	43.	BARNHAM	70 HBC	+ 10 K+ P, OMEGA PI	6/70	
W	F	(60.0)			CASO	70 HBC	- 11.2PI- P, PI OMEG	5/70	
W	F	(194.1)	(94.1)	(60.)	CASON	73 HBC	- 8.1, 8.5 PI- P	1/74	
W	F	89.1	25.		THOMPSON	74 HBC	+ 13 PI+ P	12/75*	
W	AVG	95.4	23.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)					
W	STUDENT	95.3	25.0	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
W	R PEAKS FROM MMS. (FOR DIFFICULTIES WITH MMS EXPTS. SEE THE A2 MINI-REVIEW IN THE 1973 EDITION)								
W	NR1	(21.)	OR LESS	FOCACCI	66 MMS	- 7-12 PI- P, P MMS	12/72		
W	NR2	(30.)	OR LESS	FOCACCI	66 MMS	- 7-12 PI- P, P MMS	12/72		
W	NR3	(38.)	OR LESS	FOCACCI	66 MMS	- 7-12 PI- P, P MMS	12/72		
W	N	NOT SEEN BY BOWEN 72							
W	R	(195.0)		ANDERSON	69 MMS	- 16 PI- P, BACKK	8/69		

15 G PARTIAL DECAY MODES

P1	G INTO PI PI			DECAY MASSES
P2	G INTO 4PI			139+ 139
P3	G INTO 2 RHO			773+ 773
P4	G INTO PI PI RHO			139+ 139+ 773
P5	G INTO A2 PI			1310+ 135
P6	G INTO K KBAR			497+ 497
P7	G INTO OMEGA PI			139+ 783
P8	G INTO K KBAR PI			497+ 497+ 139
P9	G INTO PHI PI			1020+ 139
P10	G INTO ETA PI			548+ 139

15 G BRANCHING RATIOS

R1	G INTO (2PI)/TOTAL			(P1)			
R1 P	(0.4)		BARTSCH	70 HBC	+ 8. PI+ P	2/72	
R1 P	(0.22)	(0.04)	MATTHEWS	71 HBC	0 7. PI+ P, PI- P	2/72	
R1 P	OPE MODEL USED IN THIS ESTIMATION						
R1 G	(.245)	(.006)	ESTABROOK	75 RVUE	17 PI- P, PI+ P, N	12/75*	
R1 G	FROM PHASE-SHIFT ANALYSIS OF HYAMS 75 DATA						
R1 S	+.24	+.01	HYAMS	75 ASPK	0 17 PI- P, PI+ P, N	12/75*	
R1 S	ERROR TAKES ACCOUNT OF SPREAD OF 4 DIFFERENT PHASE-SHIFT SOLUTIONS						
R2	G INTO (PI+- P10) / (ALL PI+- PI+ PI- P10)			(P1)/(P2C)			
R2 D	(0.08) OR LESS		BALTAY	68 HBC	+ 7-8.5 PI+ P	6/68	
R2 D	USING DATA OF DEUTSCHMANN 65 ON PI+ P TO PI+ P10 P						
R2	0.8	0.2	JOHNSTON	68 HBC	- 7. PI+ P	2/72	
R2	0.8	0.15	BARTSCH	70 HBC	+ 8. PI+ P	2/72	
R2	(0.12) OR LESS		BALLAM	71 HBC	- 16. PI- P	2/72	
R2	(0.2) OR LESS		HOLMES	72 HBC	+ 10. -12. K+ P	1/73	
R2	0.35	0.11	CASON	73 HBC	- 8.1, 8.5 PI- P	1/74	
R2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)						
R2	STUDENT	0.56	0.16	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R2	STUDENT	0.57	0.12	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R3	G+- INTO (2RHO) / (ALL 4PI)			(P1)/(P3)			
R3 S	(0.48) OR LESS		BISHAS	68 HBC	- 8. PI- P	2/72	
R3 S	SUPERSEDED BY CASON 73						
R4	G+- INTO (K KBAR) / (2PI)			(P6)/(P1)			
R4	0.08	0.08	0.03	CRENNELL	68 HBC	+ 6.0 PI- P	12/68
R4	0.08	0.03		BARTSCH	70 HBC	+ 8. PI+ P	1/71
R4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
R4	STUDENT	0.080	0.028	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R5	G+- INTO (K KBAR PI) / (2PI)			(P8)/(P1)			
R5	0.10	0.03		BARTSCH	70 HBC	+ 8. PI+ P	2/72
R6	G+- INTO (RHO 2PI) / (ALL 4PI)			(P4)/(P2)			
R6	CONSISTENT WITH 1.		CASO	68 HBC	- 11 PI- P	6/68	
R6	0.18	0.15		BARTSCH	70 HBC	+ 8. PI+ P	2/72
R6	0.8	0.15		BALLAM	71 HBC	- 16. PI- P	2/72
R6	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
R6	STUDENT	0.94	0.11	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R6	STUDENT	0.94	0.12	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

R7	G+- INTO (2RHO) / (ALL 4PI)			(P3)/(P2)			
R7	0.72	0.15		BARTSCH	70 HBC	+ 8. PI+ P	2/72
R7	(0.92)			ARMENISE	72 HBC	+ 8.1 PI- P, P 4PI	12/72
R7 A	(0.78)	(0.33)		CASON	73 HBC	- 8.1, 8.5 PI- P	1/74
R7 A	ASSUMING (ALL 4PI)=(RHO RHO) + (OMEGA PI)						
R7	66	(.56)		KLIGER	74 HBC	- 4.5 PI- P, P 4PI	12/75*
R7 T	(0.13)	(0.09)		THOMPSON	74 HBC	+ 13 PI+ P	12/75*
R7 T	RHO RHO AND A2 PI MODES ARE INDISTINGUISHABLE						
R8	G+- INTO (2 RHO) / (ALL RHO 2PI)			(P3)/(P4)			
R8	0.48	0.16		CASO	68 HBC	- 11 PI- P	6/68
R8	(0.75) OR MORE			BISHAS	68 HBC	- 8. PI- P	2/72
R9	G+- INTO (PI+- A20) / (ALL 4PI)			(WITH A20 INTO (PI+ PI- P10))			
R9	0.40	0.20		BALTAY	68 HBC	+ 7, 8.5 PI+ P	6/68
R9	NOT SEEN			JOHNSTON	68 HBC	- 7. PI- P	6/68
R9	(0.6)	(0.15)		BARTSCH	70 HBC	+ 8. PI+ P	2/72
R9	NOT SEEN			CASON	73 HBC	- 8.1, 8.5 PI- P	1/74
R10	G+- INTO (PI OMEGA) / (ALL 4PI)			(P7)/(P2)			
R10	0.25	0.10		BALTAY	68 HBC	+ 7-8.5 PI+ P	5/68
R10	.25	0.10		JOHNSTON	68 HBC	- 7.0 PI- P	6/68
R10	0.12	0.07		BALLAM	71 HBC	- 16. PI- P	2/72
R10 A	(0.22)	(0.08)		CASON	73 HBC	- 8.1, 8.5 PI- P	1/74
R10 A	ASSUMING (ALL 4PI)=(RHO RHO) + (OMEGA PI)						
R10	(.09) OR LESS			KLIGER	74 HBC	- 4.5 PI- P, P 4PI	12/75*
R10	0.33	0.07		THOMPSON	74 HBC	+ 13 PI+ P	12/75*
R10	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)						
R10	AVG	0.233	0.050	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R10	STUDENT	0.236	0.051	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			
R11	G+- INTO (PI PHI) / (ALL 4PI)			(P9)/(P2)			
R11	(0.11) OR LESS			BALTAY	68 HBC	+ 7, 8.5 PI+ P	6/68
R12	G+- INTO (PI+- 2PI+ 2PI- P10) / (ALL PI+- PI+ PI- P10)			(P10)/(P2)			
R12	(0.15) OR LESS			BALTAY	68 HBC	+ 7, 8.5 PI+ P	6/68
R13	G+- INTO (PI ETA) / (ALL 4PI)			(P10)/(P2)			
R13	(0.02) OR LESS			THOMPSON	74 HBC	+ 13 PI+ P	12/75*
R14	GO INTO (2PI+ 2PI-1) / (2PI)			(P2C)/(P1)			
R14	GO INTO (2PI+ 2PI-1) FORBIDDEN IF 4PI MODE IS (RHO RHO)						
R14	0.44	0.20		KALELKAR	75 HBC	015 PI+ P	12/75*

REFERENCES FOR G

BELLINI	65 NC 40 A 948	BELLINI, DI CORATO, DUIMIO, FIORINI	(MILANO)
DEUTSCHMANN	65 PL 18 351	M. DEUTSCHMANN ET AL	(AACHEN+BERLIN+CEFN)
FORINO	65 PL 19 65	FCRINO, GESSAROLI +	(BOLOGNA+ORSAY+SACLAY)
GOLDBERG	65 PL 17 354	GOLDBERG+ (CERN+EPOL+ORSAY+MILANO+CEA-SACL)	
EHRlich	66 PR 152 1194	R. EHRlich, H. SELOVE, H. YUTA	(PENNSYLVANIA)
FOCACCI	66 PRL 17 890	CERN MISSING MASS SPECTROMETER GROUP	(CERN)
LEVRAT	66 PL 22 714	CERN MISSING MASS SPECTROMETER GROUP	(CERN)
SEGUINOT	66 PL 19 712	CERN MISSING MASS SPECTROMETER GROUP	(CERN)
ABRAMS	67 PRL 18 620	+KEHOE+GLASSER+SECHI-ZORN+WOLSKY	(MARYLAND)
DANYSZ	67 PL 248 309	+FRENCH+KINSON+SIMAK+ (CERN+LIVERPOOL)	
DUBAL	67 NP 83 435	+FOCACCI+KINZEL+LECHANDINE+LEVRAT+ (CERN)	
ALSO	68 THESIS 1456	L. DUBAL (GENEVE)	
FRENCH	67 NC 52A 442	+KINSON+MCDONALD+RIDDIFORD+ (CERN+BIFM)	
ARMENISE	68 NC 54 A 999	+FCRINO+CARTACCI+(PARI+BGNA + FIRENZE+ORSAY) I	
BALTAY	68 PRL 20 887	+KUNG+YEH+FERBEL+ (COLU+ROCH+RUTG+YALE) = 1	
BISWAS	68 PRL 21 50	+CASON, DZIERBA, GROVES, KENNEY, + (INDAN)	
BOESEBECK	68 NP 8 501	+BOESEBECK, DEUTSCHMANN, + (AACHEN+BERLIN+CEFN)	
CASO	68 NC 54 A 983	+CONTE+CORSOS+DIAZ+ (GENOVA+HAMB+MILA-SACL)	
CRENNELL	68 PL 28 B 136	+KARSHON, LAI, SCARR, SKILLI CORN (BNL)	
JOHNSTON	68 PRL 20 1414	+PRENTICE, STEENBERG, YGON (TORONTO+MISC) IJP	
ADERHOLZ	69 NP 8 B 11 259	+BARTSCH, + (AACH+BERL+CERN+JAGL+WARS)	
ANDERSON	69 PRL 22 1390	+COLLINS, BLIEDEN+ (BNL+CERN)	
BARISH	69 PR 184 1375	+SELOVE, BISWAS, CASON, + (PENN+NDAM+ROCH)	
CASO	69 NC 62 A 755	+CONTE, BENZ, + (GENO+DESY+HAMB+MILA-SACL)	
VELITSKY	69 SJMP 9 461	+GUZHAVIN, KLIGER, KILGANOV, LEBEDEV + (ITEP)	
ARMENISE	70 LNC 4 199	+GHIDINI, FORINO, CARTACCI, + (BARI+BGNA+FRIZ)	
BARNHAM	70 PRL 24 1083	+COLLEY, JONES, KENYON, PATHAK, RIDDIFORD (BRM)	
BARTSCH	70 NP 8 22 109	+KRAUS, TSANOS, GROTE, KOTZAN+ (AACH+BERL+CERN)	
CASO	70 LNC 3 707	+CONTE, TOMASINI, CORSOS+ (GENOVA+HAMB+MILA-SACL)	
KRAMER	70 PRL 25 396	+BARTON, GUTAY, LICHTMAN, MILLER, + (PURDUE)	
MAURER	70 THESIS NO. 588	G. MAURER (STRASBOURG)	
STUNTEBE	70 PL 32 B 351	+STUNTEBECK, KENNEY, DEERY, BISWAS, CASON+ (INDAN)	
BALLAM	71 PR D 3 2606	+CFADWICK, GUIRAGOSSIAN, JOHNSON, + (SLAC)	
BRAUN	71 NP 8 30 213	+FRIDMAN, GERBER, GIVERNAUD, KAHN, + (STRB)	
GRAYNER	71 PL 35 B 610	+HYAMS, JONES, SCHLEIN, BLUM, + (CERN+MPIM) JP3-	
MATTHEWS	71 NP 8 331	+PRENTICE, YOON, CARROLL, + (INTO+MISC) JP3-	
ARMENISE	72 LNC 4 205	+FORINO, CARTACCI, + (BARI+BGNA+FRIZ)	
ALSO	75 LNC 14 177	+FDGLI=MUC TACCI, FORINO+ (BARI+BGNA+FRIZ) JP	
ROWEN	72 PRL 29 890	+EARLES, FAISSLER, BLIECEN, + (NEAS+STCN)	
CLAYTON	72 NP 8 47 81	+MASON, MUIRHEAD, RIGOPoulos, + (LIVP+PATR)	
GRAYNER	72 PHIL. CNFN. PROC. 5	+HYAMS, JONES, SCHLEIN, BLUM, DIETL+ (CERN+MPIM)	
HOLMES	72 PR D 6 3336	+FERBEL, SLATTERY, WERNER (ROCH)	
ARNOLD	73 LNC 6 707	+ENGEL, ESCOBES, KURTZ, LLOP ET, PATY, + (STRB)	
CASON	73 PR D 7 1971	+BISHAS, KENNEY, MADDEN, SANDER, SHEPHERD (NDAM)	
CASCN 1	73 NP 8 64 14	+MADDEN, BISHOP, BISWAS, KENNEY, + (NDAM)	
HYAMS	73 NP 8 64 134	+JONES, WEILHAMMER, BLUM, DIETL, + (CERN+MPIM)	
ROBERTSON	73 PR D 7 2554	ROBERTSON, WALKER, DAVIS (DUKE+WISC)	
DUBOVIKO	74 SJMP 19 568	DUBOVIKOV, MATSYUK, MILCV, SKAZOLOV (ITEP)	

Data Card Listings

Mesons

For notation, see key at front of Listings.

X(1690), X⁻(1795), A₄(1900), S(1930)

X(1690)
→ ωTTT

64 X(1690)
THIS ENTRY CONTAINS (ΩMEGA PI PI) PEAKS AROUND 1690 MEV. EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.

64 X(1690) MASS (MEV)

M	N	(1689.)	(10.)	DANYSZ	67 HBC	0 3,3,6 PBAR P	2/74
M	N	NOT SEEN IN HIGH STATISTICS EXP. OF OREN 74		YOST	68 HBC	0 4,3 K-P, LMBD, 5PI	2/74
M	N	1670.0	18.0	BARNES	69 HBC	0 4,6 K-P, ΩMEG2PI	2/74
M	N	1695.0	20.0				
M	N	1681.2	13.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
M	N	STUDENT 1681.1	15.1	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

64 X(1690) WIDTH (MEV)

W	N	(38.1)	(18.)	DANYSZ	67 HBC	0 3,3,6 PBAR P	1/73
W	N	NOT SEEN IN HIGH STATISTICS EXP. OF OREN 74		YOST	68 HBC	0 4,3 K-P, LMBD, 5PI	1/73
W	N	59.0	15.0	BARNES	69 HBC	0 4,6 K-P, ΩMEG2PI	1/73
W	N	90.	20.				
W	N	64.4	19.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)			
W	N	STUDENT 63.8	14.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

REFERENCES FOR X(1690)

DANYSZ	67 NC 51 A 801	DANYSZ+FRENCH+SIMAK	(CERN)
YOST	68 UMD T-REPR87849	+YODH, EINSCHLAG, DAY, GLASSER	(UMD)
BARNES	69 PRL 23 142	+CHUNG, EISNER, FLAMINIO, +	(BNL)
OREN	74 NP 871 189	+COOPER, FIELDS, RHINES, WHITMORE, +	(ANL+OXF)

X⁻(1795)
→

63 X⁻(1795, JPG=) I=1
SEEN AS A (PBAR N) BOUND STATE IN PBAR D ANNIHILATIONS AT REST. NEEDS FURTHER CONFIRMATION. OMITTED FROM TABLE.

63 X⁻(1795) MASS (MEV)

M	D	1794.5	1.4	GRAY	71 DBC	- 0. PBAR D	1/72
M	D	DECAYS TO FOUR OR MORE PIONS					

63 X⁻(1795) WIDTH (MEV)

W	D	(8.)	OR LESS CL=.95	GRAY	71 DBC	- 0. PBAR D	1/72
W	D	DECAYS TO FOUR OR MORE PIONS.					

REFERENCES FOR X⁻(1795)

GRAY	71 PRL 26 1491	+HAGERT, KALOGEROPOULOS	(SYRA)
BOGDANOV	72 PRL 28 1418	+BOGDANOV, DALKAROV, SHAPIRO	(ITEP)
GRAY	73 PRL 30 1091	+PAPADOPOULOU, KARAGEROPOULOS, +	(ATEN+SYRA)

A₄(1900)
→

43 A₄(1900, JPG= -) I=1
OMITTED FROM TABLE.

THIS ENTRY CONTAINS THE DIFFRACTIVE-LIKE 3PI AND 5PI BUMPS IN THE REGION OF 1900 MEV, AS WELL AS VARIOUS PEAKS NEARBY. NOTE THAT THE EXISTENCE OF AN S-WAVE GPI THRESHOLD BUMP (AS A CONTINUATION TO A1 AND A3) IS NOT UNEXPECTED. OMITTED FROM TABLE.

43 A₄ MASS (MEV)

M	N	(1900.)		HUSON	68 HLBC	- 16. PI-A, A 5PI	2/74
M	N	(1830.)		SALZBERG	72 HBC	- 13, 20 PI-P, P 3PI	2/74
M	B	40(1960.)	(30.)	BASTIEN	73 CBC	- 15. PI-D, D 3PI	2/74
M	B	MARGINAL STATISTICAL SIGNIFICANCE.					
M	N	(1800.)		DEUTSCHM	75 HBC	+ 16 PI+P, P 3PI	12/75*
M	N	208(2080.)	(40.)	KALELKAR	75 HBC	+ 15 PI+P, P PI+G	12/75*
M	N	VARIOUS PEAKS					
M	K	(1820.)	(12.)	FRENCH	67 HBC	0 3,3,6 PBAR P	7/67
M	K	OBSERVED IN (KS KO PI0...) MODE (G-PARITY UNKNOWN)					

43 A₄ WIDTH (MEV)

W	N	(130.)		SALZBERG	72 HBC	- 13, 20 PI-P, P 3PI	2/74
W	B	40 (200.)	(80.)	BASTIEN	73 CBC	- 15. PI-D, D 3PI	2/74
W	N	208 (340.)	(80.)	KALELKAR	75 HBC	+ 15 PI+P, P PI+G	12/75*
W	N	VARIOUS PEAKS					
W	K	(50.)	(20.)	FRENCH	67 HBC	0 3-4 PBAR P	7/67

43 A₄ PARTIAL DECAY MODES

P1	A4 INTO 3PI	139+ 139+ 139
P2	A4 INTO RHD PI	77+ 139
P3	A4 INTO P PI	127+ 135
P4	A4 INTO G PI	1690+ 139

43 A₄ BRANCHING RATIOS
R1 A4 INTO (G PI)/(ALL 3PI) DOMINANT KALELKAR 75 HBC + (P4)/(P1) + 15 PI+P, P 3PI 12/75*

REFERENCES FOR A₄

DANYSZ	67 NC 51A 801	DANYSZ+FRENCH+SIMAK	(CERN)
FRENCH	67 NC 52A 442	+KINSON+MCDONALD+RIDDFORD+	(CERN+BIKM)
HUSON	68 PL 28 B 208	+LUBATTI, BELLINI, BINGHAM, +	(ORSA+MILA+LBL)
BEMPORAD	71 NP B 33 397	+DUFAY, COODLING, +	(CERN+ETHZ+LOIC+MLA)
CLAYTON	72 NP B 47 81	+MASON, MUIRHEAD, RIGGPOULOS, +	(LIVP+PATR)
HARRISON	72 PRL 28 775	+HEYDA, JOHNSON, KIM, LAN, MUELLER, +	(HARV)
SALZBERG	72 NP B 41 397	+HARRISON, FEYCA, JOHNSON, KIM, LAN, +	(HARV)
BASTIEN	73 UPPSALA CONF. 73	+DUNN, HARRIS, LUBATTI, BINGHAM, +	(SEAT+UCB)
OREN	74 NP 871 189	+COOPER, FIELDS, RHINES, WHITMORE, +	(ANL+OXF)
DEUTSCHM	75 NP 899 397	DEUTSCHMANN, +	(ABBCCH COLLABORATION)
KALELKAR	75 THESIS(NEVIS 207)	M.S.KALELKAR	(COLU)

S(1930) REGION
→

31 S(1930, JPG=)

This entry contains the structure observed in s-channel NN annihilations, as well as various peaks claimed in this region by production experiments. The resonant interpretation of the broad structure observed in pp backward elastic scattering (CLINE 70, D'ANDLAU 75) has been criticized (PINSKY 71, BIZZARRI 72). Although accurate measurements of both total and elastic scattering have produced evidence for a narrow bump in this region (CARROLL 74, MARZANO 76), the charge-exchange reaction shows no such structure (GARNJOST 75, MARZANO 76). Measurements of exclusive channel cross sections and of the spectator momentum distribution in pd annihilations have suggested more complicated structure (DEFOIX 75, KALOGEROPOULOS 75). We prefer to wait for further clarification before entering the S meson into the table of established resonances.

31 S MASS (MEV)

M	C	S CHANNEL NEAR N		CLINE	70 HBC	0 .25-.74 PBAR P	2/72
M	C	(1940.)	(8.)	BENVENUTI	71 HBC	0 .1 - .8 PBAR P	2/72
M	B	(1968.)		CARROLL	74 CNTR	S CHAN. PBAR P, D	12/75*
M	S	(1932.)	(2.)	D'ANDLAU	75 HBC	0 .175-.750 PBAR P	12/75*
M	C	(1942.)	(5.)	KALOGERO	75 DBC	- PBAR N ANNIH	12/75*
M	N	(1897.)	(1.)	KALOGERO	75 DBC	- PBAR N ANNIH	12/75*
M	N	(1934.4)	(2.6)	(1.4)	KALOGERO	75 HBC	OPBAR P, 5PI
M	A	(1940.)		DEFIOIX	75 HBC	OPBAR P, 5PI	12/75*
M	S	(1935-9)	(1.0)	MARZANO	76 HBC	OPBAR P TOT, ELAS	12/75*
M	N	PEAKS FROM PRODUCTION EXPERIMENTS					
M	N	226(1929.0)	(14.0)	CHIKOVANI	66 MMSP	- 12.0 PI-P	12/75*
M	N	100(1900.)	(40.)	BOESEBECK	68 HBC	+ 8 PI+ P, PI+ PI0	12/75*
M	N	30(1973.0)	(15.0)	CASO	70 HBC	- 11.2PI-P, RHD 2PI	12/75*
M	K	40(1975.0)	(12.0)	KRAMER	70 HBC	+ 13.1 PI+ P, 2PI	11/70
M	N	100(1898.)	(18.)	THOMPSON	74 HBC	+ 13 PI+ P, 2RHO	12/75*
M	N	50(1895.)	(19.)	ABASHIAN	75 STRC	- 8 PI-P, P 3PI	12/75*
M	A	FROM ENERGY DEPENDENCE OF 5PI CROSS-SECTION.					
M	B	SEEN AS A BUMP IN THE PBAR P - KS KL CROSS SECTION WITH JPC=1--.					
M	N	NOT SEEN BY CARSON 72 WITH EQUAL STATISTICS.					
M	C	FROM ENERGY DEPENDENCE OF FAR BACKWARD ELASTIC SCATTERING.					
M	N	SOME INDICATION OF ADDITIONAL STRUCTURE.					
M	K	2PI PEAK OF KRAMER NOT SEEN IN SAME EXP WITH MORE DATA(THOMPSON 74)					
M	N	NOT SEEN IN 3 CHARGED MODE, NOT SEEN BY BOWEN 73 WITH 6X STATISTICS.					
M	S	NARROW BUMP SEEN IN TOTAL PBAR P, D CROSS-SECTIONS. ISOSPIN UNCERTAIN					
M	N	NOT SEEN IN PBAR P CEX BY GARNJOST 75, MARZANO 76					

Mesons

Data Card Listings

For notation, see key at front of Listings.

T(2200), U(2360), $\bar{N}\bar{N}_{I=0}(2375)$, X(2500-3600)

ALEXANDE 72 NP B 45 29
 BERTANZA 72 CHEXBRES (CERN 72-10)
 BUGG 72 PR D 6 3047
 CLAYTON 72 NP B 47 81
 DIEBOLD 72 BATAV,CONF.
 DONALD 72 PL 40 B 586
 MING MA 72 NP B 51 77
 TAKAHASHI 72 PR D 6 1266

ALEXANDER, BAR-NIR, BEVARY, CARAN,+ (TEI A)
 L. BERTANZA, REVIEW TALK (PISA)
 +CONDO, HART, COHN, ENDERF,+ (TENN+ORNL+CINC)
 +MASON, MUIRHEAD, RIGOPoulos,+ (LIVP+PATR)
 R. DIEBOLD, RAPPORTEUR TALK (ANL)
 +GALLETLY, EDWARDS, DE BILLY,+ (LIVP+LPP)
 +EASTMAN, OH, PARKER, SMITH, SPRAFKA (MSU)
 TAKAHASHI, BARISH,+ (TG+O+PENN+NDAM+ANL)

ALSPECTOR, COHEN, CVIJANOVICH,+ (RUTG+UPNJ)
 +BUTTERWORTH, (RHEL+LIVP)
 +GARNJOST, BIGI,+ (PADO+LBL+PISA+TORI)
 +EARLES, FAISSLER, BLIEDEN,+ (NEAS+STON)
 +EDWARDS, GIBBINS, BRIAND, DUBOC,+ (LIVP+LPP)
 W. KIENZLE (CERN)

A. ASTBURY REVIEW AT PRAGUE 74 (RHEL)
 +BIGI, CASALI, LARICCIA,+ (PISA+PADO+TORI)

+JACQUES, JONES, PANDOLAS,+ (RUTG+STEV+ALBA)
 +GARNJOST, ROSS,+ (LBL+PADO+PISA+TORI)
 +DEMARZO, GUERRIERO,+ (CANB+PARI+BROW+MIT)

U(2360)
 REGION

33 U(2360, JPG=) I=1
 THIS ENTRY CONTAINS THE BROAD BUMP OBSERVED
 IN THE S CHANNEL NBAR N, AND VARIOUS OTHER PEAKS,
 MOSTLY CONTROVERSIAL, OMITTED FROM TABLE.
 FOR A REVIEW SEE ASTBURY 74.

33 U(2360) MASS (MEV)

M A	S CHANNEL NBAR N	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/73
M A	2350. (10.)				
M A	FOR I=1 NBAR N				
M N	(2360.0) (25.0)	OH	70 HBC	-OPBAR(P,N),K*2PI	1/73
M N	NO EVIDENCE FOR THIS BUMP SEEN IN THE PEAR, P DATA OF CHAPMAN 71				1/73
M N	NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.				12/75*
M I	(2359.1) (42.1)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
M I	ISOSPINS 0 AND 1 NOT SEPARATED				

M PEAKS FROM PRODUCTION EXPERIMENTS
 M M (2382.0) (24.0) CHIKOVANI 66 MMS - 12.0 PI-P 12/72
 M N NOT SEEN BY BOWEN 73.
 M M 2370. 17. ANDERSON 69 ASPK - 16 PI- BKSCAT 11/69
 M B 126 2340. 20. BALTAY 75 HBC + 15 PI+P, P5PI 12/75*

M B DOMINANT DECAY INTO RHO RHO PI
 M AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)

33 U(2360) WIDTH (MEV)

W	S CHANNEL NBAR N	ABRAMS	67 CNTR	S CHANNEL PBAR N	1/73
W N	(140.)				
W N	(60.0) OR LESS	OH	70 HBC	-OPBAR(P,N),K*2PI	11/71
W N	NO EVIDENCE FOR THIS BUMP SEEN IN THE PEAR, P DATA OF CHAPMAN 71				11/71
W N	NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.				12/75*
W I	(165.) (18.) (8.)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
W I	ISOSPINS 0 AND 1 NOT SEPARATED				

M PEAKS FROM PRODUCTION EXPERIMENTS
 W M (30.0) OR LESS CHIKOVANI 66 MMS - 12.0 PI-P 8/66
 W N NOT SEEN BY BOWEN 73.
 W M (57.) ANDERSON 69 ASPK - 16 PI- BKSCAT 11/69
 W B 126 180. 60. BALTAY 75 HBC + 15 PI+P, P5PI 12/75*

W B DOMINANT DECAY INTO RHO RHO PI

33 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON

CS A	(3.2)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
CS A	FOR I=1 NBAR N				
CS I	(2.1) (0.2) (0.1)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
CS I	ISOSPINS 0 AND 1 NOT SEPARATED				

REFERENCES FOR U(2360)

CHIKOVAN 66 PL 22 233 CERN MISSING MASS SPECTROMETER GROUP (CERN)
 FOCACCI 66 PRL 17 890 CERN MISSING MASS SPECTROMETER GROUP (CERN)

ABRAMS 67 PRL 18 1209 +COOL, GIACOMELLI, KYCIA, LEONTIC, LI,+ (BNL)
 ANDERSON 69 PRL 22 1390 +BLESER, BIRNBAUM, EDELSTEIN,+ (BNL+CERN)
 BRICMAN 69 PL 29 B 451 +FERRO-LUZZI, BIZARD,+ (CERN+CAEN+SACL)
 CASO 69 LNC 3 707 +CONTE, BENZ,+ (GENO+DESY+HAMB+MILA+SACL)

ABRAMS 70 PR D 1 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, LI,+ (BNL)
 OH 70 PRL 24 1257 +PARKER, EASTMAN, SMITH, SPRAFKA, MA (MSU)

CHAPMAN 71 PR D4 1275 +GREEN, LYS, MURPHY, RING,+ (MICH)
 FIELDS 71 PRL 27 1749 +COOPER, RHINES, ALL TSON (ANL+OXF)
 YOH 71 PRL 26 922 +BARISH, CAROLL, LOBKOVICZ+ (CIT+BNL+RGCH)

ASTBURY 72 CERN 72-10 A. ASTBURY REVIEW AT CHEXBRES 72 (RHEL)
 DIEBOLD 72 BATAV,CONF. R. DIEBOLD, RAPPORTEUR TALK (ANL)
 EASTMAN 72 NP B 51 29 +MING MA, OH, PARKER, SMITH, SPRAFKA (MSU)
 MING MA 72 NP B 51 77 +EASTMAN, OH, PARKER, SMITH, SPRAFKA (MSU)

ALSPECTOR 73 PRL 30 511 ALSPECTOR, COHEN, CVIJANOVICH,+ (RUTG+UPNJ)
 BOWEN 73 PRL 30 332 +EARLES, FAISSLER, BLIEDEN,+ (NEAS+STON)
 DONNACHI 73 LNC 7 285 A. DONNACHI, P. R. THOMAS (MANCHESTER)
 KIENZLE 73 PR D 7 3520 W. KIENZLE (CERN)
 OH 73 NP B 51 57 +EASTMAN, MING MA, PARKER, SMITH,+ (MSU)

ASTBURY 74 CERN 74-18 A. ASTBURY REVIEW AT PRAGUE 74 (RHEL)
 MING MA 74 NP B68 214 +MOUNTZ, ZEMANY, SMITH (MICH)

BALTAY 75 PRL 35 891 +CAUTIS, COHEN, KALELKA, PISELLO,+ (COLU+BING)
 KEMP 75 NC 27 A 155 +LOTTS, CONTRI, TEDDORC+ (DURH+GENO+MILA+LPP)

$\bar{N}\bar{N}_{I=0}(2375)$

41 N NBAR(2375, JPG=) I=0
 EVIDENCE FOR RESONANCE PRELIMINARY.
 OMITTED FROM TABLE.
 FOR A REVIEW SEE ASTBURY 74.

41 N NBAR(2375) MASS

M I	2375. (10.)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
M I	(2359.1) (2.)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
M I	ISOSPINS 0 AND 1 NOT SEPARATED				

41 N NBAR(2375) WIDTH

W I	(190.)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
W I	(165.) (18.) (8.)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
W I	ISOSPINS 0 AND 1 NOT SEPARATED				

41 N NBAR(2375) SIGMA (MB) FOR FORMATION BN

CS I	(2.5)	ABRAMS	70 CNTR	S CHANNEL PBAR P	1/71
CS I	(2.1) (0.2) (0.1)	ALSPECTOR	73 CNTR	S CHANNEL PBAR P	1/74
CS I	ISOSPINS 0 AND 1 NOT SEPARATED				

REFERENCES FOR N NBAR (2375)

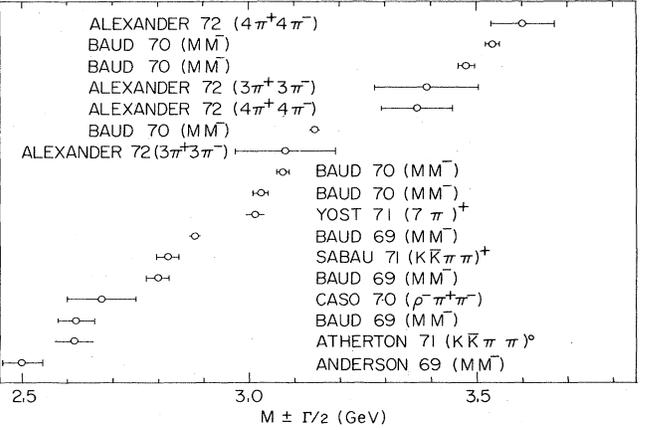
BRICMAN 69 PL 29 B 451 +FERRO-LUZZI, BIZARD,+ (CERN+CAEN+SACL)
 ABRAMS 70 PR D 1 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, LI,+ (BNL)
 EASTMAN 72 NP B 51 29 +MING MA, OH, PARKER, SMITH, SPRAFKA (MSU)
 MING MA 72 NP B 51 77 +EASTMAN, OH, PARKER, SMITH, SPRAFKA (MSU)
 ALSPECTO 73 PRL 30 511 ALSPECTOR, COHEN, CVIJANOVICH,+ (RUTG+UPNJ)

ASTBURY 74 CERN 74-18 A. ASTBURY REVIEW AT PRAGUE 74 (RHEL)
 MING MA 74 NP B68 214 +MOUNTZ, ZEMANY, SMITH (MICH)

X(2500-3600)

46 X(2500-3600)
 THIS ENTRY CONTAINS VARIOUS HIGH MASS NON-STRANGE
 PEAKS. OMITTED FROM TABLE.

The high mass region is covered nearly continuously by evidence for peaks of various widths and decay modes (see figure). As a satisfactory grouping into particles is not yet possible, we list all the Y=0 bumps with M > 2400 MeV together, ordered by increasing mass. Note that ANTIPOV 72 ($\pi^- p \rightarrow p \bar{M}$) at 25 and 40 GeV/c see no narrow bumps.



Masses and widths of reported enhancements with Y=0, M > 2400 MeV. (—o— indicates that upper limit only was reported for the width.)

Mesons

X(2500-3600), K±, K0, K*(892)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for mass, width, and references for X(2500-3600) mesons. Includes entries for ANDERSON, BAUD, CASO, SABAU, BRAUN, ALEXANDER, and YOST.

REFERENCES FOR X(2500-3600)
+COLLINS,+ (BNL+CERN)
CERN BOSON SPECTROMETER GROUP (CERN)
+BAR-NIR,DAGAN,GIDDI,GRUNHAUS+ (TEL-AVIV)
CERN BOSON SPECTROMETER GROUP (CERN)
+CENTE,TOMASINI,CORDS+(GENO+HAMB+MLA+SACL)
+URETSKY (BUCH+ANL)
+MORRIS,AL BRIGHT,BRUCKER,LANNUTTI (FSU)
ALEXANDER, BAR-NIR, BEVARY,DAGAN,+ (TEL-AVIV)
ANTIPQV 72 PL 408 147 +BRAU,BUSNELLO,DAMGAARD,+ (IHEP+CERN)
BRAUN 75 PREPRINT CRN 24 +BRICK,FRIDMAN,GERBER,JULLIOT,+ (STRB)
KALELKAR 75 THESIS(NEVIS 207) M.S.KALELKAR (COLU)

S=±1 MESON STATES

10 CHARGED K(494,JP=0-) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

K±

11 NEUTRAL K(498,JP=0-) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

K0

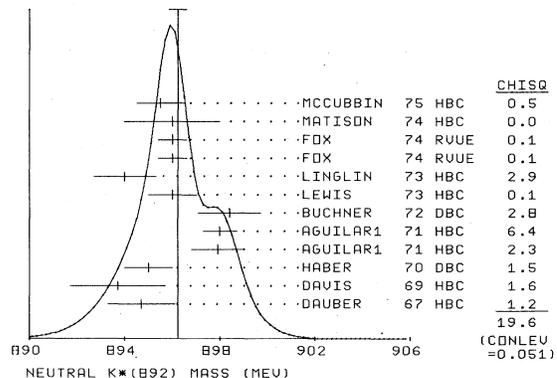
K*(892)

18 K*(892,JP=1-) I=1/2

Table with columns for charged only, neutral only, and student average for K*(892) mesons. Lists various experiments and their results.

A INCLUDED IN LINGLIN 73 WORLD K+P DST
C FROM POLE EXTRAPOLATION.
D ERRORS ENLARGED BY US TO GAMMA/SQRT(N). SEE TYPED NOTE.
I INCLUSIVE REACTION. COMPLICATED BACKGROUND AND PHASE-SPACE EFFECTS
S DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED
W NUMBER OF EVENTS IN PEAK REEVALUATED BY US

WEIGHTED AVERAGE = 896.23 ± 0.36
ERROR SCALED BY 1.3



Data Card Listings

For notation, see key at front of Listings.

Mesons

K*(892)

Note on K*(892) Masses and Mass Differences

Unrealistically small errors are reported by some experiments. We use simple "realistic" tests for the minimum errors on the determination of mass and width from a sample of N events:

$$\delta_{\min}(m) = \frac{\Gamma}{\sqrt{N}}, \quad \delta_{\min}(\Gamma) = 4 \frac{\Gamma}{\sqrt{N}}$$

(For detailed discussion see the April 1971 edition of this note.) We consistently increase unrealistic errors before averaging.

Although in the past we have argued against taking the mass difference $m(K^{*0}) - m(K^{*\pm})$ from the separate averages of $m(K^{*0})$ and $m(K^{*\pm})$, we no longer see any reason for such caution. In fact, the difference between the separate averages agrees with direct measurements of the mass difference, and it has a smaller statistical error.

18 K*(892) - K*(892) MASS DIFF. (MEV)

D	W	283	6.3	4.1	BARASH	67 HBC	0 PBAR P	12/75*
D	SD1400	(6.5)	(5.0)		FICENEC1	68 HBC	1.3 K-P	12/75*
D	SD1600	(9.5)	(5.0)		FICENEC2	68 HBC	2.7 K-P	12/75*
D	7338	5.7	1.7		AGUILARI	71 HBC	-0 3.9,4.6 K-P	11/71
D	AVG	5.8	1.6		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
D	STUDENT	5.8	1.7		AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

D ERRORS ENLARGED BY US TO GAMMA/SQRT(N). SEE TYPED NOTE.
S DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED
W NUMBER OF EVENTS IN PEAK REEVALUATED BY US

18 K*(892) WIDTH (MEV)

W	CHARGED ONLY. THIS IS WHAT APPEARS ON MESON TABLE	SD 100	(46.0)	(8.0)	CHADWICK	63 HBC	+ 1.5 K+P(K PI)	12/75*
W	D 1700	46.0	5.0		WOJCICKI	64 HBC	- 1.7 K-P(KO PI-)	12/75*
W	SD 200	(51.0)	(15.0)		ADELMAN	65 HBC	- 1.5 K-P	12/75*
W	SD 300	(47.0)	(11.0)		FERRO-LUZZI	65 HBC	+ 3.0 K+P	12/75*
W	SD 300	(50.0)	(12.0)		GELSEMA	65 HBC	- 1.5 K-P(KO PI-)	12/75*
W	SD 190	(50.1)	(15.0)		BOMSE	67 HBC	+ 2.3 K+P(KO PI+)	12/75*
W	D 620	56.0	9.0		DE BAERE	67 HBC	+ 3.5 K+P(KO PI+)	12/75*
W	SD 260	(53.1)	(13.0)		DE BAERE	67 HBC	+ 3.5 K+P(KO PI+)	12/75*
W	SD 70	(68.1)	(33.0)		SALLSTRM	67 HBC	+ 3. K+P (K+ P)	12/75*
W	SD 50	(47.1)	(27.0)		SALLSTRM	67 HBC	+ 3. K+P (K+ P)	12/75*
W	SD 210	(44.1)	(12.0)		BARLOW	67 HBC	+ 1.2 PBAR P	12/75*
W	D 720	43.1	9.0		BARLOW	67 HBC	+ 1.2 PBAR P(KO PI)	12/75*
W	D 600	53.1	9.0		BARLOW	67 HBC	+ 1.2 PBAR P(KO PI)	12/75*
W	SD 430	(58.1)	(12.0)		DE WIT	68 HBC	- 3. K- D	12/75*
W	D 540	44.0	8.0		DE WIT	68 HBC	- 3. K- D	12/75*
W	S	(58.1)	(16.0)		FICENEC1	68 HBC	- 1.3 K-P (K-PI0)	12/75*
W	S	(44.1)	(13.0)		FICENEC1	68 HBC	- 1.3 K-P (K-PI0)	12/75*
W	SD 115	(41.0)	(16.0)		SCHWEINGR	68 HBC	+ 4.1 K-P(KO PI-)	12/75*
W	SD 341	(47.0)	(10.0)		SCHWEINGR	68 HBC	- 5.5 K-P(KO PI-)	12/75*
W	S	(57.0)	(13.0)		FICENEC2	68 HBC	- 2.7 K- P(K-PI0)	12/75*
W	S	(48.0)	(9.0)		FICENEC2	68 HBC	- 2.7 K- P(KOPI-)	12/75*
W	SD 175	(52.0)	(16.0)		KANG	68 HBC	- 4.6 K- P	12/75*
W	D 2886	53.1	4.0		FRIEDMAN	69 HBC	- 2.1 K-P(KO PI-)	12/75*
W	D 728	49.1	7.3		FRIEDMAN	69 HBC	- 2.45K-P(KO PI-)	12/75*
W	D 3229	46.1	3.2		FRIEDMAN	69 HBC	- 2.6 K-P(KO PI-)	12/75*
W	D 1027	49.1	6.1		FRIEDMAN	69 HBC	- 2.7 K-P(KO PI-)	12/75*
W	SD 127	(50.1)	(18.0)		LIND	69 HBC	+ 9. K+ P	12/75*
W	D 4404	54.3	3.3		AGUILARI	71 HBC	- 3.9,4.6 K- P	12/75*
W	D 765	46.3	6.7		CLARK	73 HBC	- 3.13K-P P PI- KO	12/75*
W	W D1150	48.2	5.7		CLARK	73 HBC	- 3.3 K-P, P PI- KO	12/75*
W	I 9000	(52.1)	(2.2)		PALER	75 HBC	- 14.3 K-P, K- X	12/75*
W	AVG	49.5	1.5		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W	STUDENT	49.4	1.8		AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

W	NEUTRAL ONLY.	SD 70	(60.0)	(29.0)	COLLEY	62 HBC	0 2.0 PI-P	12/75*
W	S	310	(53.1)	(13.0)	BARLOW	67 HBC	0 1.2 PBAR P	12/75*
W	SD 120	(34.1)	(12.5)		BARLOW	67 HBC	0 1.2 PBAR P	12/75*
W	D 1040	44.1	5.5		DAUBER	67 HBC	0 2.0 K- P	12/75*
W	S	(52.1)	(12.0)		FICENEC1	68 HBC	0 1.3 K-P (K-PI+)	12/75*
W	S	(50.0)	(8.0)		FICENEC2	68 HBC	0 2.7 K- P(K-PI+)	12/75*
W	SD 200	(48.0)	(16.0)		KANG	68 HBC	0 4.6 K-P(K-PI+)	12/75*
W	SD 120	(51.0)	(19.0)		SCHWEINGR	68 HBC	0 5.5 K-P(K-PI+)	12/75*
W	SD 150	(53.0)	(17.5)		SCHWEINGR	68 HBC	0 4.1 K-P(K-PI+)	12/75*
W	D 10700	53.2	2.1		DAVIS	69 HBC	0 12. K+P(K+PI-)	12/75*
W	DAM2200	(58.0)	(5.0)		DE BAERE	69 HBC	0 5.0 K+P(K+PI-)	12/75*
W	W D4300	54.0	3.3		HABER	70 DBC	0 3. K-N (K-PI+)	12/75*
W	D 2924	55.8	4.2		AGUILARI	71 HBC	0 3.9,4.6 K- P	11/71
W	D 5362	48.5	2.7		AGUILARI	71 HBC	0 3.9,4.6 K- P	12/75*
W	D 1700	51.4	5.0		BUCHNER	72 DBC	0 4.6 K+ N, K+ P	12/72
W	D 3186	46.0	3.3		LEWIS	73 HBC	0 2.1-2.7 K+P	12/75*
W	C	(46.5)	(1.5)		LINGLIN	73 HBC	0 2-13 K+P, K+PI-	1/74
W	10600	47.1	2.1		FOX	74 RVUE	0 2 K+P, K+PI+	12/75*
W	C	51.1	2.1		FOX	74 RVUE	0 2 K+P, K+PI-	12/75*
W	C	(47.1)	(3.1)		MATISON	74 HBC	0 12 K+P, K+PI-	12/75*
W	3600	48.1	3.1		MCCUBBIN	75 HBC	0 3.6 K-P, K-PI+	12/75*
W	I 22K	(50.6)	(2.5)		PALER	75 HBC	0 14.3 K-P, K+XO	12/75*
W	AVG	50.0	1.0		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)			
W	STUDENT	49.9	1.1		AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

A INCLUDED IN LINGLIN 73 WORLD K+P DST
C FROM POLE EXTRAPOLATION.
D ERRORS ENLARGED BY US TO 4*GAMMA/SQRT(N). SEE TYPED NOTE.
I INCLUSIVE REACTION. COMPLICATED BACKGROUND AND PHASE-SPACE EFFECTS
S DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED
W NUMBER OF EVENTS IN PEAK REEVALUATED BY US

18 K*(892) PARTIAL DECAY MODES

P1	K*(892) INTO K PI	493+ 139
P2	K*(892) INTO K PI PI	493+ 139+ 139
P3	K*(892) INTO K GAMMA	493+ 0

18 K*(892) BRANCHING RATIOS

R1	K*(892) INTO (K PI PI)/(K PI)	(P2)/(P1)
R1	0 (0.002) OR LESS	WOJCICKI 2 64 HBC - 1.7 K-P
R2	K*(892) INTO (K GAMMA)/TOTAL	(UNITS 10**=-3) (P3)
R2	(1.6) OR LESS	BEMPRAD 72 CNTR + 10.-16. K+A, COUL 1/73
R2	1.5	0.7 CARITHERS 75 CNTR 0 8-16KOBAR A, COUL 12/75**

REFERENCES FOR K*(892)

ALSTON 61 PRL 6 300 ALSTON, ALVAREZ, EBERHARD, GOOD, GRAZIANO (LRL)
ALEXANDE 62 PRL 8 447 ALEXANDER, KALBFLEISCH, MILLER, G SMITH (LRL)
COLLEY 62 CERN CONF 315 COLLEY, N GELFAND + (COLUMBIA+RUTGERS)

CHADWICK 63 PL 6 309 CHADWICK, CRENNELL, DAVIES, BETTINI, FOX, PADO
GOLDHABE 63 ATHENS CONF 92 SULAMITH GOLDHABER (LRL)

WOJCICKI 64 PR 135 B 484 STANLEY G WOJCICKI (LRL)

ADELMAN 65 ATHENS 527 STUART LEE ADELMAN (CAVENDISH)
FERRO-LUZZI 65 NC 36 1101 FERRO-LUZZI, GEORGE, HENRI, JONGEJANS (CERN)
FERRO-LUZZI 65 NC 39 417 FERRO-LUZZI, GEORGE, GOLDSCHMIDT-CLER + (CERN)
GELSEMA 65 THESIS E. S. GELSEMA (SEE ALSO PL 10 341) (AMSTERDAM)
WANGLER 65 PR 137 B 414 WANGLER, ERWIN, WALKER (MISCOSIN)

BARASH 67 PR 156 1399 BARASH, KIRSCH, MILLER, TAN (COLUMBIA)
BARLOW 67 NC 50 A 701 +MONTANET, D. AND LAU+ (CERN+CDEF+IRAD+LIVP)
BOMSE 67 PR 158 1298 +BORNSTEIN, COLE+GILLESPIE+ (JOHN HOPKINS)
CONFORTO 67 NP 83 465 +MARECHAL, MONTANET+CERN+CDEF+IPNL+IVERPOOL
DAUBER 67 PR 153 1403 +SCHLEIN, SLATER, TICH (UCLA)
DE BAERE 67 NC 51 A 401 +GOLDSCHMIDT-CLERMONT, HENRI+ (BRUX+CERN)
FRENCH 67 NC 42A 442 +KINSON+MCDONALD+RIDDIFORD+ (CERN+BIPM)
GEORGE 67 NC 49A 9 +GOLDSCHMIDT-CLERMONT+HENRI+ (CERN+BRUX)
SALLSTRG 67 NC 49A 348 SALLSTRM+OTTER+EKSPONG (STOCKHOLM)

DE WIT 68 THESIS S. DE WIT (AMSTERDAM)
FICENEC1 68 PR 169 1034 +MULLISER+SHANSON+TROWER (ILL)
FICENEC2 68 PR 175 1725 FICENEC, GORDON, TROWER (ILLINOIS)
KANG 68 PR 176 1587 Y. W. KANG (IOWA)
SCHWEINGR 68 PR 166 1317 SCHWEINGRUBER, DERRICK, FIELDS+ (ANL+MESH)

CRENNELL 69 PRL 22 487 +KARSHON, LAI, ONEALL, SCARR (BNL)
DAVIS 69 PRL 23 1071 +DERENZO, FLATTE, ALSTON, LYNCH, SOLMITZ (LRL)
DE BAERE 69 NC 61 A 397 +GOLDSCHMIDT-CLERMONT, HENRI, + (BELG+CERN)
FRIEDMAN 69 UCRL-18860 J. FRIEDMAN, PH. D., THESIS (LRL)
JUHALLA 69 PR 184 1461 +LEACOCK, RHODE, KOPELMAN, LIBBY, + (ISUC-COLO)
LIND 69 NP 8 B 14 1 +ALEXANDER, FIRESTONE, FU, GOLDHABER (LRL)

ATHERTON 70 NP 8 B 16 416 +FRANEK, FRENCH, FRISK, BEDNAR+ (CERN+PRAG)
DE BAERE 70 CERN PHYS TO 41 +DEBAISIEUX, DE WOLF, DUFOUR, + (BELG+CERN)
HABER 70 NP 8 B 17 289 +SHAPIRA, ALEXANDER+ (REHO+SACL+BGNA+EPOL)

AGUILARI 71 PRL 26 466 +BARNES, BASSANO, EISNER, KINSON, SAMIOS (BNL)
AGUILARI 71 PR D 4 2583 +EISNER, KINSON (BNL)
BARNHAM 71 NP 8 B 21 171 +COLLEY, JONES, GRIFFITHS, HUGHES, + (BIRM+GLAC)
BUCHNER 71 NP 8 B 29 381 +DEHM, GOEBEL, GOLDSCHMIDT, + (MPI+CERN+BELG)
CORDS 71 PR D 4 1974 +CARMONY, ERWIN, MEIERE, + (PURD+UCD+IUPUI)
MERCER 71 NP 832 381 +ANTICH, CALLAHAN, CHIEN, COX, + (JOHN HOPKINS)
YUTA 71 PRL 26 1502 +DERRICK, ENGELMANN, MUSGRAVE (ANL+EFF1)

ABRAMOVICH 72 NP 8 B 39 189 +ABRAMOVICH, CHALOUPIKA, CHUNG, HILPERT, + (CERN)
BINGHAM 72 NP 8 B 41 1 + (INTERNATIONAL K+ COLLABORATION)
BEMPRAD 72 NP 8 B 51 171 +BEUSCH, FROEDERICH, + (CERN+ETHZ+GLAC)
BRUNET 72 NP 8 B 37 114 +DANYSZ, GOLDSACK, + (CDEF+SACL+LOIC+LDMC)
BUCHNER 72 NP 8 B 45 333 +DEHM, CHARRIERE, CORNET, + (MPI+CERN+BRUX)
CHARRIERE 72 NP 8 B 51 317 +CHARRIERE, DRIJARD, DE BAERE, + (CERN+BRUX)
CRENNELL 72 PR D 6 1220 +GORDON, KWAN, WU LAI, SCARR (BNL)
DEUTSCHMANN 72 NP 8 B 36 373 +DEUTSCHMANN, + (ABCLV COLLABORATION)
ENGELMANN 72 PR D 5 2162 ENGELMANN, MUSGRAVE, FORMAN, + (ANL+EFF1)
ROUGE 72 NP 8 B 46 29 +VIDEAU, VOLTE, DE BRION, + (EPCL+SACL)
TIECKE 72 NP 8 B 39 596 +GRUJNS, HEINEN, DE GRIGT, + (NIJM+ZEEB)

Mesons

$K^*(892)$, $\kappa(1250)$

BERTHCN 73 NP B 63 54	+MONTANET, PAUL, BERTRANET, * (CERN+SACL)
CLARK 73 NP B 54 432	+LYONS, RADOJICIC (OXFORD)
LEWIS 73 NP B 60 283	+ALLEN, JACOBS, DANYSZ, BORG, + (LOWC+LOIC+CDEF)
LINGLIN 73 NP B 55 408	D. LINGLIN (CERN)
WALUCH 73 PR D 8 2837	+FLATTE, FRIEDMAN (LBL)
FOX 74 NP B80 403	G.C. FOX, M. L. GRIS (CIT)
MATISCN 74 PR D9 1872	+GALTIERI, GARNJUST, FLATTE, FRIEDMAN, + (LBL)
BRANDENB 75 PL 59 B 405	BRANDENBURG, CARNEGIE, CASHMERE, DAVIER+ (SLAC)
CARITHER 75 PRL 35 349	CARITHERS, MUHLEMANN, UNDERWOOD, + (ROCH+MCGI)
MCCUBBIN 75 NP B86 13	N. A. MCCUBBIN, L. LYONS (OXF)
PALER 75 NP B96 1	+TOVEY, SHAH, SPIRO, CHAURAND+ (RHEL+SACL+EPCL)

$\kappa(1250)$

19 K PI S WAVE, CALLED KAPPA

S-Wave $K\pi$ Interactions in the Region 750-1700 MeV

$K\pi$ interactions in the $I(J^P)=1/2(0^+)$ wave can be described by the elastic phase shift δ_0^1 from the $K\pi$ threshold (~ 630 MeV) up to at least 1100 MeV (BINGHAM 72). The first inelastic S-wave thresholds are $K\pi\pi$ and $K\eta$, neither of which is known to be important below 1400 MeV. Apart from the inelastic thresholds, the S-wave $\pi\pi$ and $K\pi$ interactions are reminiscent of each other.

All phase-shift solutions (MERCER 71, BINGHAM 72, FIRESTONE 71, 72, MATISON 72, 74, BAKER 73, GALTIERI 73, YUTA 73, FOX 74, LAUSCHER 75) share the following intrinsic ambiguities:

- 1) The standard modulo-180° ambiguity.
- 2) If one amplitude is dominant [the P wave near $K^*(892)$ or the D wave near $K^*(1420)$], then the observed S-P or S-D interference can be explained by two ambiguous S-wave solutions, known as "up" and "down". For an illustration see our 1972 edition.

The resulting "up-down" ambiguity of the S-wave phase shift in the $K^*(892)$ region (BINGHAM 72) was restricted by an analysis of $12 \text{ GeV}/c \text{ } K^+p \rightarrow K^+\pi^-\Delta^{++}$ (MATISON 72, 74, GALTIERI 73) to only two points at 890 and 900 MeV, where the P wave goes through 90°. The "up" solution is therefore excluded, except for a possibility of a very narrow ($\Gamma < 7$ MeV) S-wave resonance. Moreover, CHUNG 72 imposes positivity on physical-region $K\pi$ moments, and finds a narrow resonance most unlikely.

FIRESTONE 71 and 72 have continued $K\pi$ partial-wave analysis up to 1700 MeV. They find that δ_0^1 crosses 90° near 1300 MeV, and indeed shows the "up-down" ambiguity near the $K^*(1420)$.

Meanwhile several groups have attempted to clarify the situation around 1300 MeV. CORDS 72, FRATI 72, ROUGE 72, and LAUSCHER 75

Data Card Listings

For notation, see key at front of Listings.

give support to the resonant S-wave interpretation of FIRESTONE 71. The other groups (AGUILAR 72, BUCHNER 72, CRENNELL 72, ENGELMANN 72, BAKER 73) agree that the S wave is important but not necessarily resonant.

In a review of the status of the 0^+ nonet, MORGAN 75 shows that the resonant interpretation, $\kappa(1250)$, is consistent with the information on the other S-wave systems, $\pi\pi$, $K\bar{K}$, and $\eta\pi$ (see the corresponding mini-reviews), suggesting a non-ideally mixed nonet.

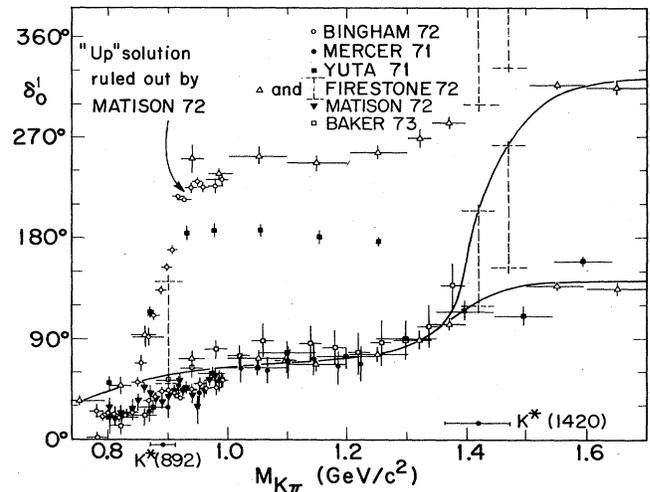


Fig. 1. The S-wave $K\pi$ phase shift. An "up-down" ambiguity occurs at the mass of $K^*(892)$, $K^*(1420)$, ..., which can be resolved by precise measurement of $\sigma(K\pi)$. At $K^*(892)$, the "up" solution was ruled out by MATISON 72; earlier "up" solutions are plotted only to show historical progress.

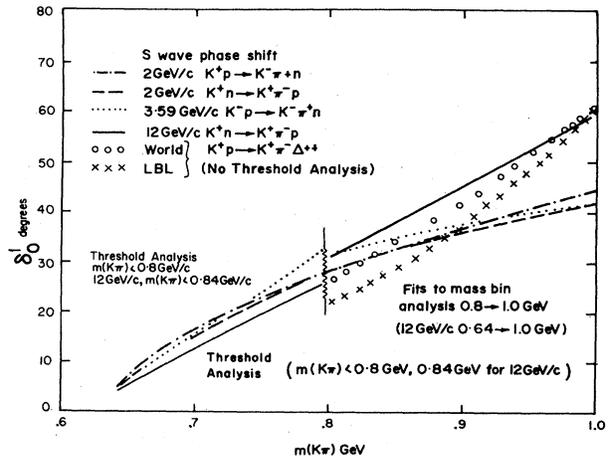


Fig. 2. Summary of six different S-wave $K\pi$ phase-shift analyses, including four solutions of FOX 74. (Figure from FOX 74.)

Data Card Listings

For notation, see key at front of Listings.

Mesons
 $\kappa(1250), Q$

19 KAPPA MASS (MEV)				
M	1250.	100.	MORGAN	75 RVUE

19 KAPPA WIDTH (MEV)				
M	(450.)	APPROX.	MORGAN	75 RVUE

REFERENCES FOR KAPPA				
TRIPPE	68 PL 28 B 203		*CHIEN, MALAMUD, NELLEMA, SCHLEIN, +	(UGLA)
CRENNELL	69 PRL 22 487		*KARSHON, LAI, O. NEALL, SCARR	(BNL)
DODD	69 PR 177 1994		*JOLDERSMA, PALMER, SAMIOS	(BNL)
GOLDBERG	69 PL 30 B 434		SABRE COLLABOR. (SACL+AMST+BGNA+REHO+EPOL)	
SCHLEIN	69 ARGONNE CONF. 446		P. SCHLEIN	(UGLA)
FIRESTONE	71 PRL 26 1460		A. FIRESTONE, G. GOLDBERGER, D. LISSAUER	(LRL)
MERCER	71 NP B32 381		*ANTICH, CALLAHAN, CHIEN, COX, + (JOHN HOPKINS)	
YUTA	71 PRL 26 1502		*DERRICK, ENGELMANN, MUSGRAVE	(ANL+EFI)
AGUILAR	72 PR D 6 11		AGUILAR-BENITEZ, CHUNG, ETSNER	(BNL)
BINGHAM	72 NP B 41 1		+ (INTERNATIONAL K+ COLLABORATION)	
BUCHNER	72 NP B 45 333		*DEHM, CHARRIERE, CORNET, +	(MPIN+CERN+BRUX)
CHUNG	72 PRL 29 1570		*ETSNER, AGUILAR-BENITEZ	(BNL)
CRENNELL	72 PR D 6 1220		*GORDON, KWAN, WU LAI, SCARR	(BNL)
DIEBOLD	72 BATAV. CONF.		R. DIEBOLD RAPPORTEUR TALK	(ANL)
ENGELMANN	72 PR D 5 2162		ENGELMANN, MUSGRAVE, FORMAN, +	(ANL+EFI)
FIRESTONE	72 PR D 5 2188		*GOLDBERGER, LISSAUER, TRILLING	(LBL+PMA)
FRATI	72 PR D 6 2361		*HALPERN, HARGIS, SHAPIRO, CARNAHAN, + (PENNSYLVANIA)	
ROUGE	72 NP B 46 29		*VIDEAU, VOLTE, DE BRION, +	(EPOL+SACL)
MATISON	72 LBL 1537 (THESIS)		REVISED VERSION WILL GO TO PHYS. REV.	LBL
BAKER	73 IC/HEP/73/12		*BANERJEE, CAMPBELL, HALL, ISLAM, + (LOIC+LOWC)	
CORDS	73 NP B 54 109		*CARMONY, LANDER, MEIERE, + (PURDUE+IUPUI)	(LBL)
GALTIERI	73 PREP. LBL 1772		*MATISON, GARNJOST, FLATTE, FRIEDMAN, +	(LBL)
LINGLIN	73 NP B 55 408		D. LINGLIN	(CERN)
YUTA	73 NP B 52 70		*ENGELMANN, MUSGRAVE, FORMAN, +	(ANL+EFI)
FOX	74 NP B80 403		G.C. FOX, M.L. GRISS	(CIT)
MATISON	74 PR D9 1872		*GALTIERI, GARNJOST, FLATTE, FRIEDMAN, +	(LBL)
ALSO	72 MATISON, 73 GALTIERI			
MORGAN	74 PL 518 71		D. MORGAN	(RHEL)
BAKER	75 NP B99 211		*BANERJEE, CAMPBELL, ALLEN, MARCH, + (LOIC+LOWC)	
LAUSCHER	75 NP B86 189		*OTTER, WIECZOREK, + (ABCLV COLLABORATION)	
MORGAN	75 PREPRINT RL75133		D. MORGAN	(RHEL)

Q REGION, $K\pi\pi(1240-1400)$

28 Q REGION (1200-1400) MEV I=1/2

The main effect in the Q region is a broad bump in the $K\pi\pi$ spectrum between 1200 and 1400 MeV (not far above the $K^*(892)\pi$ threshold), produced by K beams without charge exchange. In particular, it has been observed in coherent K^+d interactions (FIRESTONE 72) and in coherent interactions on heavy nuclei (BINGHAM 73). Throughout the entire region, $J^P = 1^+$ and $I = 1/2$.

Evidence for narrower states in the Q region has been reported from reactions with incident π^- and \bar{p} (ASTIER 69, CRENNELL 67, 72, DAVIDSON 74, DORE 75). FIRESTONE 72 observe a bump in the backward direction with a shape similar to that of the Q. WERNER 73 find no evidence for narrow states in the charge-exchange reaction $K^-p \rightarrow (K\pi\pi)^0 n$, but the data can accommodate some broad Q production. The $(K\pi\pi)^0$ system appears to have an important $J^P=1^+$ contribution (OTTER 75).

The broad Q peak does not have a simple Breit-Wigner shape. It can be fitted at all

energies by a superposition of two Breit-Wigner amplitudes (FIRESTONE 70, BARNHAM 71, BOWLER 71, BARLOUTAUD 73).

Dalitz plot analyses of the interference between the $K^*\pi$ and $K\rho$ modes show the relative magnitude and relative phase of the two decay amplitudes varying with $K\pi\pi$ mass. The $K\rho$ mode has a maximum intensity below that of $K^*\pi$. This is suggestive of the presence of two $J^P=1^+$ states, possibly mixtures of the strange members of the $J^P=1^{++}$ (" A_1 ") and 1^{+-} (B) nonets (GOLDBERGER 67, BARNHAM 71, BOWLER 71, GARFINKEL 71, BINGHAM 72, FIRESTONE 72, BOWLER 74). In addition to the dominant modes $K^*\pi$ and $K\rho$, there is some evidence for a $K\pi\pi$ mode, with the $\pi\pi$ or the $K\pi$ system in an S wave (ALEXANDER 69, BARNHAM 71, DAVIS 72, BARLOUTAUD 73).

Recent partial-wave analyses have confirmed the rather complex situation in the Q region (DEUTSCHMANN 74, ANTIPOV 75, OTTER2 75, OTTER3 75, TOVEY 75, BRANDENBURG 76). Although $J^P=1^+$ is the dominant contribution, other spin-parity states are necessary to describe the whole Q enhancement. Moreover, the $K^*\pi$ and $K\rho$ modes are not produced coherently and have different polarization properties (OTTER2 75, OTTER3 75, TOVEY 75, BRANDENBURG 76). Whereas the $K\rho$ mode approximately conserves s-channel helicity, the $K^*\pi$ mode is approximately t-channel helicity-conserving.

ANTIPOV 75 do not require the $K\rho$ decay mode to be present at 40 GeV/c.

ANTIPOV 75 and OTTER2 75 find no variation in the phase of the $1^+ K^*\pi$ wave. On the other hand BRANDENBURG 76, with high-statistics spectrometer data, observe sufficient phase variation of both the $1^+ K\rho$ and $K^*\pi$ waves to warrant proposing the existence of two 1^+ mesons, superposed on a large 1^+ , mainly $K^*\pi$, Deck background. The first state has $M \sim 1300$ MeV, $\Gamma \sim 200$ MeV and decays mainly to $K\rho$; the second has $M \sim 1400$ MeV, $\Gamma \sim 160$ MeV and decays mainly to $K^*\pi$.

Mesons

Q

Data Card Listings

For notation, see key at front of Listings.

28 Q REGION MASS (MEV)

M A	PRODUCED BY BEAMS OTHER THAN K MESONS	1242.0	9.0	10.0	ASTIER	69 HBC	0 PBAR P	9/69
M	THIS IS THE C MESON.	45(1300.)			CRENNELL	67 HBC	0 6 PI-P, LK2PI	7/67
M		40(1300.)			CRENNELL	72 HBC	0 4.5PI-P, LK2PI	12/72
M		40(1278.)	(5.)		DAVIDSON	74 HBC	+- 1.6-2.2 PBAR P	12/75*
M		43(1235.)	(10.)		DORE	75 OSPK	06.2 PI-P, L MM	12/75*

PRODUCED BY K BEAMS

M		12(1320.0)	(25.0)		ALMEIDA	65 HBC	+ 3-5 K+ P	12/72
M		(1230.0)	(15.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67
M		35(1280.0)	(10.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67
M		(1320.0)	(15.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67

SPLIT THE Q REGION INTO 3 BUMPS

M		(1270.)	APPROX.		DE BAERE	67 HBC	+ 3.5 K+ P	7/67
M		(1335.0)	(6.0)		BARTSCH	68 HBC	10. K-P, K NP1	12/75*
M		(1300.)	APPROX.		BARBARO	69 HBC	+ 12. K+ P (K 2PI)	9/69
M		45(1301.0)	(10.0)		BISHOP	69 HBC	+ 3.5 K+P(K* PI)	12/75*
M		21(1300.0)	(10.0)		ERWIN	69 HBC	0 3.5 K+P(K* PI)	12/75*
M		(1281.)	(7.)		FRIEDMAN	69 HBC	- 2.6, 2.7 K- P	12/75*
M		(1300.0)	(10.0)		ABRAMS	70 HBC	+ 2.5-3.2 K+ P	12/75*
M		(1260.)	(20.)		FARBER	70 HBC	+ 12.7 K+ P	12/75*
M		(1325.0)	(7.)		DENEGR	71 DBC	- 12.6 K-D, K 2PI D	12/75*
M		(1256.)	(5.)		BARLOUTAU	73 HBC	- 14.3 K-P, P K-2PI	12/75*
M		(1283.)	(6.)		BARLOUTAU	73 HBC	- 14.3 K-P, P K-2PI	12/75*
M		(1315.)	(7.)		BINGHAM	73 HBC	- 5.5-12.7 CDH K-A	12/75*
M		(1260.)	(10.)		LEWIS	73 HBC	+ 2.1-2.7 K+ P	12/75*
M		(1260.)	(5.)		LEWIS	73 HBC	+ 2.1-2.7 K+ P	12/75*

FROM NON-PERIPHERAL EVENTS (T GT 0.8)

AVERAGING NOT MEANINGFUL

28 Q LCW (CA) MASS (MEV)

ML F	FROM EXPERIMENTS SPLITTING Q REGION INTO TWO PEAKS							
ML		(1280.)			SHEN	66 HBC	+ 0 4.6 K+P, 5 BODY	12/72
ML		(1260.0)	(10.0)		ALEXANDER	69 HBC	9.0 K+ P	12/75*
ML		(1240.0)	(5.0)		BARNHAM	70 HBC	+ 10.0 K+P, K 2PI	12/75*
ML		(1243.)	(8.)		GARFINKEL	71 DBC	+ 9. K+ D	12/75*
ML		(1228.)	(14.)		ANDERSON	72 DBC	- 7.3 K- D	12/75*
ML		(1240.)	(12.)		DAVIS	72 HBC	+ 12. K+ P	12/72
ML		(1234.)	(12.)		FIRESTONE	72 DBC	+ 12. K+ D	2/73
ML C		(1300.)	APPROX.		BRANDENB	76 ASPK	+- 13 K+-P, (KPIPIP)	12/75*

COUPLES MAINLY TO RHO K

AVERAGING NOT MEANINGFUL

28 Q HIGH (QB) MASS (MEV)

MH F	FROM EXPERIMENTS SPLITTING Q REGION INTO TWO PEAKS							
MH		70(1320.0)	(10.0)		SHEN	66 HBC	+ 4.6 K+ P	12/75*
MH		(1380.0)	(20.0)		ALEXANDER	69 HBC	9.0 K+ P	12/75*
MH		(1420.0)	(5.0)		BARNHAM	70 HBC	+ 10.0 K+P, K 2PI	12/75*
MH		(1344.)	(8.)		GARFINKEL	71 DBC	+ 9. K+ D	12/75*
MH		(1414.)	(15.)		ANDERSON	72 DBC	- 7.3 K- D	12/75*
MH		(1420.)	(12.)		DAVIS	72 HBC	+ 12. K+ P	12/72
MH		1368.	18.		FIRESTONE	72 DBC	+ 12. K+ D	12/75*
MH D		(1400.)	APPROX.		BRANDENB	76 ASPK	+- 13 K+-P, (KPIPIP)	12/75*

COUPLES MAINLY TO K* PI

AVERAGING NOT MEANINGFUL

28 Q REGION WIDTH (MEV)

W	PRODUCED BY BEAMS OTHER THAN K MESONS	127.0	7.0	25.0	ASTIER	69 HBC	0 PBAR P	9/69
W		45 (60.)			CRENNELL	67 HBC	0 6 PI-P	7/67
W		40 (60.)			CRENNELL	72 HBC	0 4.5PI-P, LK2PI	12/72
W D		40 (25.)	(15.)		DAVIDSON	74 HBC	+- 1.6-2.2 PBAR P	12/75*
W D	ERROR INCREASED BY US. SEE K* TYPED NOTE.	43 (30.)	(25.)	(18.)	DORE	75 OSPK	06.2 PI-P, L MM	12/75*

PRODUCED BY K BEAMS

W		12 (60.0)	(20.0)		ALMEIDA	65 HBC	+ 3-5 K+P	12/72
W		(60.0)	(20.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67
W		35 (80.0)	(20.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67
W		(60.0)	(20.0)		BASSOMPIE	67 HBC	+ 5. K+ P	11/67

SPLIT THE Q REGION INTO 3 BUMPS

W		(200.)	APPROX.		DE BAERE	67 HBC	+ 3.5 K+ P	7/67
W		(190.0)	(15.0)		BARTSCH	68 HBC	10. K+P, K NP1	12/75*
W		250.	APPROX.		BARBARO	69 HBC	+ 12. K+ P (K 2PI)	9/69

NO BACKGROUND SUBTRACTION.

W		45 (40.0)	(10.0)		BISHOP	69 HBC	+ 3.5 K+P(K* PI)	12/75*
W		(180.0)	(28.)		ERWIN	69 HBC	0 3.5 K+P(K* PI)	12/75*
W		(51.)	(22.)		FRIEDMAN	69 HBC	- 2.6, 2.7 K- P	12/75*
W		(80.0)	(20.0)		ABRAMS	70 HBC	+ 2.5-3.2 K+ P	12/75*
W		(180.)	(28.)		FARBER	70 HBC	+ 12.7 K+ P	12/75*
W		(130.0)	(17.)		DENEGR	71 DBC	- 12.6 K-D, K 2PI D	5/71
W		(266.)	(21.)		BARLOUTAU	73 HBC	- 14.3 K-P, P K-2PI	12/75*
W		(150.)	(70.)		LEWIS	73 HBC	+ 2.1-2.7 K+ P	12/75*
W		(47.)	(18.)		LEWIS	73 HBC	+ 2.1-2.7 K+ P	12/75*

FROM NON-PERIPHERAL EVENTS (T GT 0.8)

AVERAGING NOT MEANINGFUL

28 Q LCW (CA) WIDTH (MEV)

WL F	FROM EXPERIMENTS SPLITTING Q REGION INTO TWO PEAKS							
WL		(100.0)	(20.0)		SHEN	66 HBC	+ 0 4.6 K+P, 5 BODY	12/75*
WL		(40.0)	(10.0)		ALEXANDER	69 HBC	9.0 K+ P	12/75*
WL		(110.0)	(15.0)		BARNHAM	70 HBC	+ 10.0 K+P, K 2PI	12/75*
WL		(270.)	(26.)	(18.)	GARFINKEL	71 DBC	+ 9. K+ D	12/75*
WL		(111.)	(33.)		ANDERSON	72 DBC	- 7.3 K- D	12/75*
WL		(120.)			CAVIS	72 HBC	+ 12. K+ P	12/72
WL		(188.)	(21.)		FIRESTONE	72 DBC	+ 12. K+ D	12/75*
WL C		(200.)	APPROX.		BRANDENB	76 ASPK	+- 13 K+-P, (KPIPIP)	12/75*

COUPLES MAINLY TO RHO K

AVERAGING NOT MEANINGFUL

28 Q HIGH (QB) WIDTH (MEV)

WH F	FROM EXPERIMENTS SPLITTING Q REGION INTO TWO PEAKS							
WH		70 (80.0)	(20.0)		SHEN	66 HBC	+ 4.6 K+ P	12/75*
WH		(120.0)	(20.0)		ALEXANDER	69 HBC	9.0 K+ P	12/75*
WH		(120.0)	(15.0)		BARNHAM	70 HBC	+ 10.0 K+P, K 2PI	12/75*
WH		(60.)	OR LESS		GARFINKEL	71 DBC	+ 9. K+ D	12/72
WH		(89.)	(24.)		ANDERSON	72 DBC	- 7.3 K- D	12/75*
WH		(80.)			DAVIS	72 HBC	+ 12. K+ P	12/72
WH		(241.)	(30.)		FIRESTONE	72 DBC	+ 12. K+ D	12/75*
WH D		(160.)	APPROX.		BRANDENB	76 ASPK	+- 13 K+-P, (KPIPIP)	12/75*

COUPLES MAINLY TO K* PI

AVERAGING NOT MEANINGFUL

28 Q REGION PARTIAL DECAY MODES

P1	Q REGION INTO K*(892) PI							DECAY MASSES
P2	Q REGION INTO K RHO	75.0			ARMENTEROS	64 HBC		892+ 139
P3	Q REGION INTO K PI							497+ 773
P4	Q REGION INTO K ETA							497+ 139
P5	Q REGION INTO K CMEGA							497+ 548
P6	Q REGION INTO K PI PI							497+ 783
								497+ 139+ 139

28 Q REGION BRANCHING RATIOS

R1	PRODUCED BY BEAMS OTHER THAN K MESONS							
R1	Q REGION INTO (K RHO)/TOTAL (UNITS OF 10**=-2)	75.0			ARMENTEROS	64 HBC	0.0 PBAR P	6/66
R1	DOMINANT	10.0			CRENNELL	72 HBC	0 4.5PI-P, LK2PI	12/72
R2	Q REGION INTO (K* PI)/TOTAL (UNITS OF 10**=-2)	25.0			ARMENTEROS	64 HBC	0.0 PBAR P	6/66
R3	Q REGION INTO (K+ PI-) / (K+0 PI+0 PI-)	(0.2)	OR LESS	CL=.90	CRENNELL	67 HBC	0 6.0 PI-P	7/67
R4	Q REGION INTO (K0 PI+ PI-) / (K+0 PI+0 PI-)	(0.1)	OR LESS	CL=.90	CRENNELL	67 HBC	0 6.0 PI-P	7/67

PRODUCED BY K BEAMS

R10	Q REGION INTO (K PI) / (K*(892) PI)	(0.8)	OR LESS		SHEN	66 HBC	4.6 K+P, 5 BODY	11/67
R10	Q REGION INTO (K*(892) PI AND K RHO (OVERLAPPING BANDS))	70 (1.0)			SHEN	66 HBC	+ 4.6 K+P	8/66
R11	Q REGION INTO (K OMEGA)/(K*(892) PI)	(0.1)	OR LESS		SHEN	66 HBC	+ 4.6 K+P	10/66
R12	Q REGION INTO (K PI) / (K*(892) PI)	(0.30)	OR LESS		SHEN	66 HBC	+ 4.6 K+P	10/66

Q REGION INTO K*(892) PI AND K RHO (OVERLAPPING BANDS)

R13		200 (1.0)			BERLINGHI	67 HBC	+ 12.7 K+ P	7/67
R14	Q REGION INTO (K PI) / TOTAL	(0.02)	OR LESS	CL=.95	BERLINGHI	67 HBC	+ 12.7 K+ P	11/67
R14		(0.02)	OR LESS	CL=.95	BARTSCH	68 HBC	- 10.0 K- P	

Q REGION INTO (K PI) / (K*(892) PI)

R16	Q REGION INTO (K OMEGA) / TCTAL	(0.02)	OR LESS		BERLINGHI	67 HBC	+ 12.7 K+ P	11/67
R16		12 (0.01)	0.005		BARTSCH	68 HBC	- 10.0 K- P	9/68
R17	Q REGION INTO (K RHO) / (K*(892) PI)	0.91	0.25		BERLINGHI	67 HBC	+ 12.7 K+ P	11/67
R17		701 (0.4)	0.1		BARTSCH	68 HBC	- 10.0 K- P	9/68
R17								
R17 AVG		0.47	0.18		AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)			
R17 STUDENT		0.46	0.11		AVERAGE USING STUDENT (H/1.11) -- SEE TEXT			

Q REGION INTO (K PI) / (K*(892) PI)

R18	Q REGION INTO (K PI) / (K*(892) PI)	(0.21)	OR LESS		DE BAERE	67 HBC	+ 3.5 K+ P	11/66
R19	Q REGION INTO (K PI PI) / TOTAL	201 (0.22)	0.08		BARTSCH	68 HBC	- 10.0 K- P	9/68
R19 S	POSSIBLY SEEN				ALEXANDER	69 HBC	9.0 K+ P	2/73
R19 S	POSSIBLY SEEN				DAVIS	72 HBC	+ 12. K+ P	1/73
R19 S	WITH THE (PI PI) SYSTEM IN S-WAVE							1/73

REFERENCES FOR C REGION

PRODUCED BY BEAMS OTHER THAN K MESONS

ARMENTEROS	64 DUBNA CONF 1 577	ARMENTEROS, EDWARDS, D-ANDLAU + (CERN+CDEF)
	ALSO 64 DUBNA CONF 1 617	R ARMENTEROS (RAPORTEUR)
	ALSO 66 PR 145 1095	BARASH, KIRSCH, MILLER, TAN (COLUMBIA)
CRENNELL	67 PRL 19 44	*KALBFLEISCH, LAI, SCARR, SCHUMANN (BNL)
ASTIER	69 NP 8 10 65	*MARECHAL, MONTANET, + (CDEF+CERN+INP+IIV)JP
BETTINI	69 NC 62 A 1038	*CRESTILLI, MONTANET, BERTAUZA, BIGI+(PADO+PISAI)
CRENNELL	72 PR D 6 1220	*GORDON, KWAN, WU, LAI, SCARR (BNL)
DAVIDSON	74 PR D 9 77	*CHAPMAN, GREEN, LYS, ROE (MICH)
DORE	75 LNC 13 265	*GUIDONI, LAASKO, MARINI, CONFORTO+(RGM+RHEL)

PRODUCED BY K BEAMS

ALMEIDA	65 PL 16 184	ALMEIDA, ATHERTON, BYER, DORNAN, FORSON+ (CAVE)
SHEN	66 PRL 17 726	*BUTTERWORTH, FUJ, GOLDBABERS, TRILLING (LRL)
	ALSO 66 (PRIVATE COMMUNICATION)	GOLDBABER (LRL)
BASSOMPIE	67 PL 268 30	BASSOMPIERRE, GOLDSCHMIDT+ (CERN+BRUX+BIRM)JP
BERLINGHI	67 PRL 18 1087	BERLINGHI, ERBER+FERSEL, FORNAN (ROCHIIJP)
DE BAERE	67 NC 49A 374	*DEBAISIEUX+FAST+FILIPPAS+ (CERN+BRUX)
		ALSO PRIVATE COMMUNICATION BY B. JONGEJAANS (BNL)
GOLDBABER	67 PRL 19 976	G. GOLDBABER (LBL)
BARTSCH	68 NP 88 9	*COCCONI, + (AACH+BERL+CERN+LOIC+VIEN)
BOHME	68 PRL 20 1519	*BORENSTEIN, CALLAHAN, COLE, COX, + (JOHNHOPK) 1+
DENEGR	68 PRL 20 1194	*CALLAHAN+ETTLINGER+GILLESPIE+ (JOHNHOPK) 1+
	ALSO 70 ANTICH	

Data Card Listings

For notation, see key at front of Listings.

Mesons Q, K*(1420)

Table listing meson data cards including names (ALEXANDER, ANDREWS, BARBARO, etc.), PRL numbers, and various physical parameters like mass and width.

Table with 22 K*(1420) WIDTH (MEV) header, listing charged and neutral decay modes with branching ratios and scale factors.

Table with 22 K*(1420) PARTIAL DECAY MODES header, listing partial decay modes (P1, P2, P3, P4, P5) and their corresponding masses.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS. The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i +/- delta P_i, where delta P_i = sqrt(delta P_i^2 + delta P_j^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 22 K*(1420) MASS (MEV) header, listing mass measurements from various experiments (BASSANO, DE BAERE, SCHWEINGR, etc.) with error bars.

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R1, R2, R3, R4, R5) and their fits.

Table with 22 K*(1420) PARTIAL DECAY MODES header, listing partial decay modes (R1, R2, R3, R4, R5) and their fits.

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios (R1, R2, R3, R4, R5) and their fits.

Table with 22 K*(1420) PARTIAL DECAY MODES header, listing partial decay modes (R6, R7, R8, R9, R10) and their fits.

Mesons

K*(1420), KN(1700), L(1770)

Data Card Listings

For notation, see key at front of Listings.

R7 K*(1420) INTO (K OMEGA) / K PI (P4)/(P1)
R7 (0.08) OR LESS SHEN 66 HBC 4.6 K+P 8/66
R7 (0.21) OR LESS BASSOMPIE 69 HBC + 5 K+P 9/69
R7 0.13 0.07 BASSOMPIE 69 HBC 0 5 K+P 9/69
R7 0.05 0.04 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R7 (0.21) OR LESS CL=.95 CHUNG 74 HBC - 7.3 K-P, K- P 12/75*

R8 K*(1420) INTO (K RHO) / (K PI) (P3)/(P1)
R8 (0.09) OR LESS CHUNG 65 HBC + 0 3.9-4.2 PI- P 8/66
R8 0.26 0.16 SCHWEINGR 68 HBC 0 4.1+5.5 K- P 10/67
R8 (0.23) OR LESS BASSOMPIE 69 HBC + 5 K+P 9/69
R8 (0.31) OR LESS BASSOMPIE 69 HBC 0 5 K+P 9/69
R8 Q 15 (0.11) (0.06) BISHOP 69 HBC 3.5 K+P 9/69
R8 0.16 0.05 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R8 0.02 0.10 0.02 DEHM 74 DBC 0 4.6 K+ N 12/75*

R9 K*(1420) INTO (K RHO) / (K*(892) PI) (P3)/(P2)
R9 (0.39) OR LESS BASSOMPIE 67 HBC + 5 K+ P 9/67
R9 (0.40) OR LESS FIELD 67 HBC - 3.8 K- P 6/67
R9 P 130 .13 .09 OTTER 75 HBC 08.10.16 K-P, K* N 12/75*

R10 K*(1420) INTO (K OMEGA) / (K*(892) PI) (P4)/(P2)
R10 Q (0.10) (0.04) FIELD 67 HBC - 3.8 K- P 6/67
R10 FIT .145 .059 FROM FIT

R11 K*(1420) INTO (K ETA) / (K*(892) PI) (P5)/(P2)
R11 Q (0.07) (0.04) FIELD 67 HBC - 3.8 K- P 6/67
R11 FIT .064 .066 FROM FIT

R12 K*(1420) INTO (K ETA) / (K PI) (P5)/(P1)
R12 (0.02) OR LESS BISHOP 69 HBC 3.5 K+ P 9/69
R12 (0.04) OR LESS CL=.95 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R12 FIT .035 .036 FROM FIT

R Q FOLLOWING SUGGESTION BY AGUILAR 70, WE DO NOT MAKE USE OF MEASURE-
MENTS WHERE THE (K PI PI) BACKGROUND SUBTRACTION IS DIFFICULT DUE
TO THE NEARBY C REGION.

REFERENCES FOR K*(1420)

BADIER 65 PL 19 612 BADIER, DEMOULIN, GOLDBERG+ (EPOL+SACL+ZEEHAN)
CHUNG 65 PRL 15 325 +DAHL, HARDY, HESS, JACOBS, KIRZ, MILLER (LRL)
FOCARDI 65 PL 16 351 FOCARDI, MINGUZZI, RANZI, SERRA+ (BOLOGNA+SACL)
SHEN 66 PRL 17 726 +BUTTERWORTH, FU, GOLDBERG, TRILLING (LRL)
ALSO 66 (PRIVATE COMMUN) GERSON GOLDHABER (LRL)
BASSANO 67 PRL 19 968 +GOLDBERG, GDZ, BARNES, LEITNER+ (BNL+SYRACUSE)
BASSOMPIE 67 PL 268 30 BASSOMPIERE, GOLDSCHMIDT+ (CERN+BRUX+BIRM) JIP
CRENELL 67 PRL 19 66 +KALBFLEISCH, LAI, SCARR, SCHUMANN (BNL)
DAHL 67 PR 163 1377 +HARDY+HESS+KIRZ+MILLER (LRL)
ALSO 65 PRL 14 401 HARDY, CHUNG, DAHL, HESS, KIRZ, MILLER (LRL)
DE BAERE 67 NC 51 A 401 +GOLDSCHMIDT-CLERMONT, HENRI+ (BRUX+CERN)
FIELD 67 PL 248 638 +HENDRICKS+PICCIONI+YAGER (LAJOLLA)
GOLDBERG 67 PRL 19 972 G. GOLDBERGER, FIRESTONE, SHEN (LRL)
ADERHOLZ 68 NP B 5 567 +DEUTSCHMANN+ (AACH+BERL+CERN+LOIC+VIENNA)
ALSO 66 PL 22 357 BARTSCH, DEUTSCHMANN, MORRISON+ (ABCL+ICIV)
ANTICH 68 PRL 21 1862 +CALLAHAN, CARSON, COX, DENEGRI+ (JHU)
DUBAL 68 THESIS 1456 L. DUBAL (GENEVE)
KANG 68 PR 176 1587 Y. W. KANG (IOWA)
SCHWEINGR 68 PR 166 1317 SCHWEINGRUBER, DERRICK, FIELDS+ (ANL+NWES)
ALSO 67 THESIS F. L. SCHWEINGRUBER (NORTH-WESTERN, EVANSTON)
BASSOMPIE 69 NP B 13 189 BASSOMPIERE, GOLDSCHMIDT-CLERM.+ (CERN+BRUX) JIP
BISHOP 69 NP B 9 403 +GOSHAW, ERWIN, WALKER (WISC)
CRENELL 69 PRL 22 487 +KARSHON, LAI, ONEALL, SCARR (BNL)
DAVIS 69 PRL 23 1071 +DERENZO, FLATTE, ALSTON, LYNCH, SOLMITZ (LRL)
DE BAERE 69 NC 61 A 397 +GOLDSCHMIDT-CLERMONT, HENRI+ (BELG+CERN)
FRIEDMAN 69 UCRL-18860 J. FRIEDMAN, PH. D. THESIS (LRL)
LIND 69 NP B 14 1 +ALEXANDER, FIRESTONE, FU, GOLDHABER (LRL) JIP
ABRAMS 70 PR D 1 2433 +EISENBERG, KIM, MARSHALL, C. HALLORAN, + (ILL)
AGUILAR 70 PRL 25 1362 AGUILAR-BENITZ, BASSANO, EISNER, + (BNL+GLAS)
BIRMINGH 70 KIEV CONF. ASTIER RAPPORTEURS TALK (BIRM+GLAS+OXF)
AGUILAR 71 PR D 4 2583 +EISNER, KINSCH (BNL)
BARNHAM 71 NP B 28 171 +COLLEY, JOBS, GRIFFITHS, HUGHES, + (BIRM+GLAS)
CORDS 71 PR D 4 1974 +CARMONY, ERWIN, MEIERE, + (PURD+UCD+IUPUI)
SLATTERY 71 UR-875-332 (PREP) P. SLATTERY: A REVIEW OF STRANGE MESONS (ROCH)
BUCHNER 72 NP B 45 333 +DEHM, CHARRIERE, CORNET, + (MPI+M+CERN+BRUX)
CRENELL 72 PR D 6 1220 +GORDON, KWAN-WU LAI, SCARR (BNL)
DEUTSCHM 72 NP B 36 373 DEUTSCHMANN, + (ABCLV COLLABORATION)
ENGELMAN 72 PR D 5 2162 ENGELMANN, MUSGRAVE, FORMAN, + (ANL+EFI)
FRATI 72 NP B 5 2361 +HALPERN, HARCIS, SWAPS, CARNAHAN, + (PENN+CINC)
ROUGE 72 NP B 46 29 +VIDEAU, VOLTE, DE BRION, + (EPCL+SACL)
TIECKE 72 NP B 39 596 +GRIJNS, HEINEN, DE GROOT, + (NIJN+ZEEM)
CHARRIERE 73 NP B 51 317 CHARRIERE, ORJARD, DE BAERE, + (CERN+BELG)
ALSO 73 (PRIVATE COMMUNICAT) GOLDSCHMIDT-CLERMONT (CERN)
CLARK 73 NP B 54 432 +LYONS, RADJICIC (OXFORD)
DE JONGH 73 NP B 58 110 +CORNET, CHARRIERE, + (BRUX+MONS+CERN+MPI)
LINGLIN 73 NP B 5 408 D. LINGLIN (CERN)
WALUCH 73 PR D 8 2837 +FLATTE, FRIEDMAN (LRL)
DEHM 74 NP B 75 47 +GOEBEL, WITTEK, WOLF, + (MPI+BRUX+MONS+CERN)
CHUNG 74 PL 518 413 +EISNER, PROTOPODESCU, SAMIOS, STRAND (BNL)
ANTIPOV 75 NP B 86 381 +ASCOLI, BUSNELLO, KIENZLE+ (SERP+CERN+ILL)
LAUSCHER 75 NP B 86 189 +OTTER, WIECZOREK, + (ABCLV COLLABORATION) JIP
MCUBBIN 75 NP B 86 13 N. A. MCCUBBIN, L. LYONS (OXF)
OTTER 75 NP B 84 333 + (AACH+BERL+CERN+LOIC+VIEN+ATHU+AATEN+LIVP)
BRANDENB 76 PL BRANDENBURG, CARNEGIE, CASHMORE, DAVERI+ (SLAC) JIP

KN(1700) 27 KN(1700, JP= 1) I = 1/2
THIS ENTRY CONTAINS VARIOUS PEAKS IN STRANGE MESON
SYSTEMS REPORTED IN THE 1700 MEV REGION.
EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

27 KN(1700) MASS (MEV)
M (1660.0) CARMONY 67 HBC + 3.8 K-P, OMEGA K 11/67
M J (1660.0) (10.0) JOBS 67 HBC + 5 K+ P 12/75*
M J CLAIMED BY JOBS IN (K PI), (K*(892) PI), AND (K*(1420) PI)
M J MODES, K PI BUMP INTERFERES MOSTLY WITH DELTA(1236).
M (1660.) CHARRIERE 73 HBC 0 5 K+ P, K- P 3PI 1/73
M 60(1710.) (15.) CHUNG 74 HBC - 7.3K-P, K- OMEGA P 12/75*

27 KN(1700) WIDTH (MEV)
W (60.0) (20.0) JOBS 67 HBC + 5 K+ P 12/75*
W (60.) CHARRIERE 73 HBC 0 5 K+ P, K- P 3PI 1/73
W 60 (110.) (50.) CHUNG 74 HBC - 7.3K-P, K- OMEGA P 12/75*

27 KN(1700) PARTIAL DECAY MODES
P1 KN(1700) INTO K PI 493+ 139
P2 KN(1700) INTO K PI PI 493+ 139+ 139
P3 KN(1700) INTO K*(892) PI 892+ 139
P4 KN(1700) INTO K RHO 493+ 773
P5 KN(1700) INTO K*(1420) PI 1421+ 139
P6 KN(1700) INTO K OMEGA 493+ 783

27 KN(1700) BRANCHING RATIOS
R1 KN(1700) INTO (K PI)/(K OMEGA) (P1)/(P6)
R1 N (0.5) (0.5) CHUNG 74 HBC - 7.3 K-P, K- P 12/75*
R1 N NO K PI SIGNAL SEEN IN THIS EXPERIMENT

REFERENCES FOR KN(1700)

CARMONY 67 PRL 18 615 D. CARMONY, T. HENDRICKS, L. LANDER (LA JOLLA)
JOBS 67 PL 268 49 +BASSOMPIERE, DE BAERE + (BIRM+CERN+BRUX)
CHARRIERE 73 NP B 51 317 CHARRIERE, ORJARD, DE BAERE, + (CERN+BELG)
CHUNG 74 PL 518 412 +EISNER, PROTOPODESCU, SAMIOS, STRAND (BNL)

L(1770) 23 L(1770, JP= 1) I = 1/2

The L(1770) is seen as a bump in the diffractive-like process KN -> (KPI)N. BARBARO 69 and FIRESTONE 72 find the decay is consistent with being entirely K*(1420)pi, whereas AGUILAR 70, BARTSCH 70, COLLEY 71, and DENEGRI 71 present evidence for alternate decay modes. For a review see EISNER 74.

Recent partial-wave analyses (DEUTSCHMANN 74, ANTIPOV 75, OTTER 75) have shown that the situation in the L region is complicated, with many waves contributing. The 2- K*(1420)pi S wave is important, but cannot explain the whole L enhancement (DEUTSCHMANN 74). The phase variation of the 2- wave shows no evidence for any significant resonance contribution with Gamma < 300 MeV (OTTER 75).

23 L MASS (MEV)
M 20(1780.) BERLINGHI 67 HBC + 12.7 K+P 7/67
M (1760.0) (15.0) JOBS 67 HBC + 5 K+ P 1/73
M 1745.0 20.0 AGUILAR 70 HBC - 4.6 K- P 6/70
M 1780.0 15.0 BARTSCH 70 HBC - 10.1 K- P 1/71
M (1760.0) (15.0) LUDLAM 70 HBC - 12.6 K- P 1/73
M X 1765.0 40.0 COLLEY 71 HBC + 10.4 K+P, K 2PI 1/73
M X SYSTEMATIC ERRORS ADDED CORRESP. TO SPREAD OF DIFFERENT FITS.
M (1740.0) DENEGRI 71 DBC - 12.6 K-D, K 2PI D 5/71
M 1767. 6. BLIEDEN 72 MMS - 11.-16. K- P 12/72
M P 306 1730. 20. FIRESTONE 72 DBC + 12. K+ D 1/73
M P PRODUCED IN CONJUNCTION WITH D*
M AVG 1764.6 6.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
M STUDENT1765.0 5.8 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Data Card Listings

For notation, see key at front of Listings.

Mesons

L(1770), K_N(1800), K*(2200)

Table with 5 columns: W, L WIDTH (MEV), BERLINGHI, JONES, AGUILAR, BARTSCH, LUDLAM, COLLEY, DENEGRI, FIREFSTONE. Includes average values and error scales.

Table with 3 columns: P1, L INTO K PI PI, DECAY MASSES. Lists partial decay modes and associated masses.

Table with 3 columns: R1, L INTO K*(1420) PI / (K PI PI), (P2)/(P1). Lists branching ratios and experimental evidence.

Table with 3 columns: Author, PRL, PL, REFERENCE. Lists references for L(1770) from various experiments.

K_N(1800) 60 KN(1800, JP=3-)

This entry includes peaks of low statistical significance seen in the K_N and K_N spectra in the region of 1800 MeV (CARMONY 71, AGUILAR 73, SPIRO 76). FIRESTONE 71 and BRANDENBURG 76 observe structure in the K_N angular distribution, the simplest explanation for which is a rapid rise of the F-wave amplitude around 1800 MeV, interfering strongly with other waves. BRANDENBURG 76 propose the existence of a J^P=3⁻ resonance, but we prefer to wait for further confirmation before including this entry in the Table.

Table with 5 columns: M, C, DISAGREEMENT BETWEEN THE FIT AND DATA ON BOTH SIDES OF THE SIGNAL. Lists experimental data for 60 KN(1800) MASS (MEV).

Table with 5 columns: W, C, DISAGREEMENT BETWEEN THE FIT AND DATA ON BOTH SIDES OF THE SIGNAL. Lists experimental data for 60 KN(1800) WIDTH (MEV).

Table with 3 columns: P1, KN(1800) INTO K PI, DECAY MASSES. Lists partial decay modes and masses for 60 KN(1800) PARTIAL DECAY MODES.

Table with 3 columns: R1, KN(1800) INTO (K PI)/(K*(892) PI + K RHO), (P1)/(P2+P3). Lists branching ratios for 60 KN(1800) BRANCHING RATIOS.

Table with 3 columns: Author, PRL, PL, REFERENCE. Lists references for KN(1800) from various experiments.

K*(2200) 40 K*(2200, JP=)

Table with 3 columns: M, C, COMPILATION OF (ANTIHYPERON-NUCLEON) MASS IN K* P 8.-13. GEV/C. Lists experimental data for 40 K*(2200) MASS (MEV).

Table with 3 columns: W, C, COMPILATION OF (ANTIHYPERON-NUCLEON) MASS IN K* P 8.-13. GEV/C. Lists experimental data for 40 K*(2200) WIDTH (MEV).

Table with 3 columns: Author, PRL, PL, REFERENCE. Lists references for K*(2200) from various experiments.

Mesons

EXOTICS, NEW HEAVY MESONS

Data Card Listings

For notation, see key at front of Listings.

EXOTIC MESONS

EXOTICS

70 EXOTICS

THE PURPOSE OF THIS ENTRY IS TO PROVIDE A LIST OF REFERENCES FOR EXOTIC MESON SEARCHES (SEE MAIN TEXT, SEC. III AND TABLE I), AS WELL AS THEORETICALLY BASED SUGGESTIONS FOR EXPERIMENTS. NOTE THAT LIPKIN 73 PROPOSES EXPERIMENTS WHICH ARE CONCLUSIVE EVEN IF NEGATIVE RESULTS ARE OBTAINED.

REFERENCES FOR EXOTICS

REPORTS ON SEARCHES

ROSENFEL	68 PHILA.CONF.P.455	A.H.ROSENFELD	(LRL)
ODD	69 PR 177 1991	+JGLDERSMA, PALMER, SAMIOS	(BNL)
CHO	70 PL 32 B 409	+DERRICK, JOHNSON, MUSGRAVE, +	(ANL+NWES+KANS)
GIACOMEL	70 PL 33 B 373	G.GIACOMELLI +	(BGNA+SACL+ZEEP+REHO+EPOL)
LYS	70 PR D 2 2525	J.LYS +	(MICH)
ROSNER	70 EXP.MESON SPECTROSCOPY, ED. C. BALTAY AND A.H.ROSENFELD, P.499		
BUHL	72 NP B 37 421	+CLINE, TERRELL	(WISCONSIN)
COHEN	73 NP B 53 1	+FERBEL, SLATTERY, WERNER	(ROCHESTER)
ALAM	74 PL 53 B 207	+BRABSON, GALLOWAY, +	(IND+PURD+SLAC+VAND)
COHEN	74 BOSTON	D.COHEN REVIEW TALK	(COLU)
OREN	74 NP B 71 189	+COOPER, FIELDS, RHINES, WHITMORE, +	(ANL+OXF)
BALTAY	75 PL 57 B 293	+CAUTIS, COHEN, KALEL KAR, PISELLO, +	(COLU+BING)
DAVIS	75 NP B 96 426	+AMAR, KRUPAC, YARGER, +	(KANS+CCAC+ANL)

SUGGESTIONS FOR SEARCHES

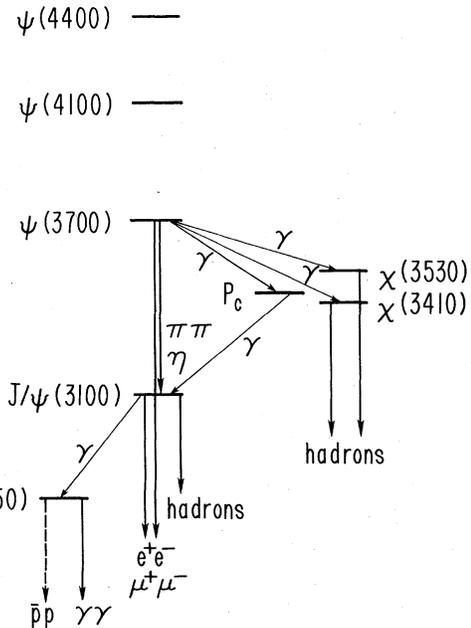
ROSNER	68 PRL 21 950+1468	J.L.ROSNER	(TEL-AVIV)
ROSNER	70 EXP.MESON SPECTROSCOPY, ED. C. BALTAY AND A.H.ROSENFELD, P.499		
FATMAN	73 PL 43 B 307	D.FATMAN, G.GOLDBABER, Y.ZARMI	(CERN)
LIPKIN	73 PR D 7 2262	H.J.LIPKIN	(ARGONNE+FNAL)

NEW HEAVY MESONS

The New Heavy Mesons

Since the discovery of the $J/\psi(3100)$ particle in p-Be collisions at the AGS (AUBERT 74), and in e^+e^- collisions at SPEAR (AUGUSTIN 74), the size of the family of "new particles" has increased to at least 8 members, with a complex set of decay modes and mutual transitions; the main features of the experimental situation are illustrated in the figure below.

In the following entries we compile the experimental data as of January 1976, with only a minimum of mini-reviews and footnotes. For detailed discussions of the implications and possible interpretations of the properties of the new particles, see, e.g., the Proceedings of the 1975 International Symposium on Lepton and Photon Interactions at High Energies, Stanford, 1975 [W. T. Kirk (SLAC), editor].



XBL 763 2361

The New Heavy Mesons

Extracting Resonance Widths from e^+e^- Colliding Beam Formation Experiments

In an e^+e^- colliding beam formation experiment, the true shape of an observed resonance is distorted primarily by the effects of (1) soft-photon processes, and (2) beam energy spread due to processes such as quantum fluctuations in emission of synchrotron radiation. The spread in energy due to (2) may usually be approximated by a Gaussian distribution, the effect of which vanishes rapidly at energies sufficiently removed from resonance. The major effect of (1) is a decrease in the effective c.m. energy for some fraction of the collisions, because of the emission of bremsstrahlung by the electron or positron before annihilation. Hence, though the nominal beam energy may be well above the resonance region, a certain fraction of the collisions occur at or near resonance. This gives rise to the well-known high-mass radiative tails of the $J/\psi(3100)$ and $\psi(3700)$ resonances.

Data Card Listings

For notation, see key at front of Listings.

Mesons

NEW HEAVY MESONS, J/ψ(3100)

Because of these effects, the most reliable means for determining resonance widths is to use a method based on the area under the line shape [see, e.g., "Notes from the SLAC Theory Workshop on the ψ", ed. R. Pearson, SLAC-PUB-1515 (1974)]. This method, familiar in nuclear physics, minimizes the sensitivity to the details of the beam energy spread. Corrections for the radiative processes, which depend on the limits of integration of the areas, still need to be made. This discussion assumes the resolution is adequate to allow a reasonable separation of signal from background (which itself is subject to radiative processes).

For formation of a resonance of mass M in e⁺e⁻ collisions, with subsequent decay via channel i,

$$\sigma_i(W) = \sigma_0 \frac{\Gamma_e \Gamma_i / 4}{(M-W)^2 + \Gamma^2 / 4}$$

where W is the total center of mass energy; Γ, Γ_e, and Γ_i are the total width and partial widths for coupling to e⁺e⁻ and channel i, respectively; and a Breit-Wigner line shape with energy-independent partial widths is assumed. The quantity σ₀ is given by

$$\sigma_0(W) = \frac{4\pi(2J+1)}{W^2}$$

where J is the spin of the resonance. For a narrow resonance, the area under the resonant line is given by

$$A_i = \frac{\pi}{2} \frac{\Gamma_e \Gamma_i}{\Gamma} \sigma_0(M)$$

independent of the energy resolution of the apparatus. Determination of the mass, spin, and integrated channel cross sections, A_i, allows, then, determination of the quantities Γ_eΓ_i/Γ, and, assuming there are no undetected decay modes (i.e., in our case Γ = Γ_{e⁺e⁻} + Γ_{μ⁺μ⁻} + Γ_{hadrons}), determination of the total and partial widths.

In the Data Card Listings, we tabulate and average the quantity Γ_eΓ_i/Γ only when it was not used to determine the partial widths and/or branching ratios.

J/ψ(3100)

70 J/PSI(3100,JPG=1--1 1=0

The J/ψ(3100) was discovered in p-Be collisions at the AGS (AUBERT 74), and in e⁺e⁻ collisions at SPEAR (AUGUSTIN 74), and since then has been extensively investigated in e⁺e⁻ collisions at SPEAR, DORIS, and ADONE, as well as at conventional accelerators with photon and hadron beams.

Evidence bearing on the hadronic nature of the J/ψ(3100) comes from the observation of a large forward photoproduction cross section. Using vector dominance arguments one derives a total cross section for J/ψ(3100) on nucleons over a large energy range of the order of 1 mb (ANDREWS 75, DAKIN 75, KNAPP1 75, MARTIN 75). Although this cross section is an order of magnitude smaller than the corresponding cross sections of the well known vector mesons, it is large enough to suggest that the J/ψ(3100) is probably a hadron. The hadronic interpretation of the J/ψ(3100) is also supported by the apparent conservation of isospin and G-parity in direct hadronic decays (JEAN-MARIE 75). For a detailed discussion of the hadronic nature of J/ψ(3100), see HARARI 75.

The J^{PC} = 1⁻⁻ assignment, suggested by production in s-channel e⁺e⁻ collisions, is confirmed by the observation of an interference between the resonant and QED amplitudes, and by the angular distribution of the lepton pairs in the final state (BOYARSKI 75). The I^G = 0⁻ assignment was determined by a study of multipion decays (JEAN-MARIE 75). These quantum numbers are compatible with all the data on partial decay widths (non-observation of the γγ mode, observation of the ΛΛ̄ mode, etc.).

70 J/PSI(3100) MASS (MEV)

M	(3100.)		AUBERT	74	SPEC	28.	PP(E+E-)	2/75*
M	L (3105.)	(3.)	AUGUSTIN	74	SPEAR		E+E-	2/75*
M	D 3095.	4.	BOYARSKI	75	SPEAR		E+E-	3/75*
M	S (3089.5)		CRIGGEE	75	DORIS		E+E-	2/75*
M	3098.	6.	PREPOST	75	SPEC	13.-21.	GAMMA D	1/76*
M	3103.	6.	BEMPORAD	75	ADONE		E+E-	1/76*
M	L		BOYARSKI 75 IS A REEVALUATION OF AUGUSTIN 74 BASED ON A RECALIBRATION OF THE SPEAR BEAM ENERGY.					
M	D		MASS, WIDTH, PARTIAL WIDTHS, AND BRANCHING RATIOS ALL OBTAINED FROM ONE OVERALL FIT TO DATA OF THIS EXPERIMENT.					
M	S		ERROR OF ABOUT 1 PER CENT FROM THE UNCERTAINTY IN CALIBRATION OF THE BEAM ENERGY.					
M	AVG	3097.6	2.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)				3/75*
M	STUDENT	3097.6	3.3	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT				2/75*

70 J/PSI(3100) WIDTH (KEV)

SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.									
W	69.	15.	BOYARSKI	75	SMAG	E+E-	3/75*		
W	68.	26.	BALDINI1	75	FRAG	E+E-	1/76*		
W	60.	25.	ESPOSITO	75	FRAM	E+E-	1/76*		
W	66.9	11.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
W	STUDENT	66.9	12.4	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					

Mesons

J/ψ(3100)

Data Card Listings

For notation, see key at front of Listings.

70 J/PSI(3100) PARTIAL DECAY MODES

		DECAY MASSES		
P1	J/PSI(3100) INTO E+ E-	.5+ .5		
P2	J/PSI(3100) INTO MU+ MU-	105+ 105		
P3	J/PSI(3100) INTO HADRONS			
P4	J/PSI(3100) INTO GAMMA INTO HADRONS			
HADRONIC DECAYS				
P				
P11	J/PSI(3100) INTO PI+ PI-			
P12	J/PSI(3100) INTO PI+ PI- P10			
P13	J/PSI(3100) INTO 2(PI+ PI-)			
P14	J/PSI(3100) INTO 2(PI+ PI-) P10			
P15	J/PSI(3100) INTO 3(PI+ PI-)			
P16	J/PSI(3100) INTO 3(PI+ PI-) P10			
P17	J/PSI(3100) INTO 4(PI+ PI-)			
P18	J/PSI(3100) INTO 4(PI+ PI-) P10			
P19	J/PSI(3100) INTO K0S K0L			
P20	J/PSI(3100) INTO K+ K-			
P21	J/PSI(3100) INTO PI+ PI- K+ K-			
P22	J/PSI(3100) INTO 2(PI+ PI-) K+ K-			
P23	J/PSI(3100) INTO RHO PI			
P24	J/PSI(3100) INTO RHO PI PI			
P25	J/PSI(3100) INTO OMEGA PI PI			
P26	J/PSI(3100) INTO PHI PI PI			
P27	J/PSI(3100) INTO K K*(892)			
P28	J/PSI(3100) INTO K K*(1420)			
P29	J/PSI(3100) INTO K*(892) K*(892)			
P30	J/PSI(3100) INTO K*(1420) K*(1420)			
P31	J/PSI(3100) INTO K*(892) K*(1420)			
P32	J/PSI(3100) INTO PEAR P			
P33	J/PSI(3100) INTO LAMBDA ANTILAMBDA			
P34	J/PSI(3100) INTO NUCLEON ANTINUCLEON PI			
P35	J/PSI(3100) INTO P PEAR PI+ PI-			
P36	J/PSI(3100) INTO P PBAR PI+ PI- NEUTRALS			
RADIATIVE DECAYS				
P				
P51	J/PSI(3100) INTO GAMMA GAMMA			
P52	J/PSI(3100) INTO P10 GAMMA			
P53	J/PSI(3100) INTO ETA GAMMA			
P54	J/PSI(3100) INTO ETA PRIME GAMMA			
P55	J/PSI(3100) INTO X(2750) GAMMA	2750+ 0		

70 J/PSI(3100) PARTIAL WIDTHS (KEV)

SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.

W1	J/PSI(3100) INTO E+ E-	(G1)		
W1	4.8	0.6	BOYARSKI 75 SMAG	E+E- 2/75*
W1 B	(4.6)	(.8)	BALDINI 75 FRAG	E+E- 3/75*
W1 B	ASSUMING EQUAL PARTIAL WIDTHS FOR (E+E-) AND (MU+MU-)			
W1	4.6	1.0	ESPOSITO 75 FRAM	E+E- 1/76*
W1	STUDENT			
W1 AVG	4.75	0.51	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
W1 STUDENT	4.75	0.55	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
HADRONIC DECAYS (G2)				
W2	J/PSI(3100) INTO MU+ MU-			
W2	4.8	0.6	BOYARSKI 75 SMAG	E+E- 2/75*
W2	5.0	1.0	ESPOSITO 75 FRAM	E+E- 3/75*
W2	STUDENT			
W2 AVG	4.85	0.51	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
W2 STUDENT	4.85	0.55	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
HADRONIC DECAYS (G3)				
W3	J/PSI(3100) INTO HADRONS			
W3	59	14	BOYARSKI 75 SMAG	E+E- 2/75*
W3	59	24	BALDINI 75 FRAG	E+E- 3/75*
W3	50	25	ESPOSITO 75 FRAM	E+E- 1/76*
W3	STUDENT			
W3 AVG	57.3	10.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
W3 STUDENT	57.3	11.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
HADRONIC DECAYS (G4)				
W4	J/PSI(3100) INTO GAMMA INTO HADRONS			
W4	12.	2.	BOYARSKI 75 SMAG	E+E- 1/76*

70 J/PSI(3100) BRANCHING RATIOS

FOR THE BRANCHING RATIOS R1 - R4, SEE ALSO THE PARTIAL WIDTHS ABOVE, AND (PARTIAL WIDTHS)*R1 BELOW.

R1	J/PSI(3100) INTO (E+ E-)/TOTAL	(P1)		
R1	0.069	0.009	BOYARSKI 75 SMAG	E+E- 3/75*
R2	J/PSI(3100) INTO (MU+ MU-)/TOTAL			
R2	0.069	0.009	BOYARSKI 75 SMAG	E+E- 3/75*
R3	J/PSI(3100) INTO (HADRONS)/TOTAL			
R3	0.86	0.02	BOYARSKI 75 SMAG	E+E- 3/75*
R4	J/PSI(3100) INTO (E+ E-)/(MU+ MU-)	(P1)/(P2)		
R4	1.00	0.05	BOYARSKI 75 SMAG	E+E- 2/75*
R4	0.93	0.10	FORD 75 SPEC	E+E- 2/75*
R4	.91	.15	ESPOSITO 75 FRAM	E+E- 1/76*
R4	STUDENT			
R4 AVG	0.980	0.043	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R4 STUDENT	0.980	0.047	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R5	J/PSI(3100) INTO (GAMMA INTO HADRONS)/TOTAL			
R5	.17	.02	BOYARSKI 75 SMAG	E+E- 1/76*
HADRONIC DECAYS				
R				
R8	J/PSI(3100) INTO (PI+ PI-)/TOTAL			
R8	(.0003)OR LESS CL=.90	WIJK 75 DASP		E+E- 1/76*
R9	J/PSI(3100) INTO (2(PI+ PI-))/TOTAL			
R9	76	.004	JEAN-MARI 75 SMAG	E+E- 1/76*
R10	J/PSI(3100) INTO (2(PI+ PI-) P10)/TOTAL			
R10	675	.04	JEAN-MARI 75 SMAG	E+E- 1/76*
R11	J/PSI(3100) INTO (3(PI+ PI-))/TOTAL			
R11	32	.004	JEAN-MARI 75 SMAG	E+E- 1/76*

R12	J/PSI(3100) INTO (3(PI+ PI-) P10)/TOTAL			
R12	181	.029	JEAN-MARI 75 SMAG	E+E- 1/76*
R13	J/PSI(3100) INTO (4(PI+ PI-) P10)/TOTAL			
R13	13	.009	JEAN-MARI 75 SMAG	E+E- 1/76*
R14	J/PSI(3100) INTO (PI+ PI- K+ K-)/TOTAL			
R14 A	83	.005	ABRAMS4 75 SMAG	E+E- 1/76*
R14 A	INCLUDING CONTRIBUTION FROM K*(892)K*(1420)			
R15	J/PSI(3100) INTO (2(PI+ PI-) K+ K-)/TOTAL			
R15	.003	.001	ABRAMS4 75 SMAG	E+E- 1/76*
R16	J/PSI(3100) INTO (RHO PI)/(PI+ PI- P10)			
R16	(.7)	OR MORE	JEAN-MARI 75 SMAG	E+E- 1/76*
R17	J/PSI(3100) INTO (RHOO P10)/(RHOO+ PI+)			
R17	.59	.17	JEAN-MARI 75 SMAG	E+E- 1/76*
R18	J/PSI(3100) INTO (RHO PI)/TOTAL			
R18	153	.013	JEAN-MARI 75 SMAG	E+E- 1/76*
R19	J/PSI(3100) INTO (OMEGA PI PI)/(2(PI+ PI-) P10)			
R19 J	(.2)		JEAN-MARI 75 SMAG	E+E- 1/76*
R20	J/PSI(3100) INTO (RHO PI PI PI)/(2 (PI+ PI-) P10)			
R20 J	(.3)		JEAN-MARI 75 SMAG	E+E- 1/76*
R20 J	FINAL STATE 2(PI+PI-)P10			
R21	J/PSI(3100) INTO (PHI PI+ PI-)/(OMEGA PI+ PI-)			
R21	.2	.1	ABRAMS4 75 SMAG	E+E- 1/76*
R22	J/PSI(3100) INTO (K0S K0L)/TOTAL			
R22	(.0002)OR LESS CL=.90	ABRAMS4 75 SMAG		E+E- 1/76*
R23	J/PSI(3100) INTO (K+ K-)/TOTAL			
R23	(.0006)OR LESS CL=.90	WIJK 75 DASP		E+E- 1/76*
R24	J/PSI(3100) INTO (K0 K*(892))/TOTAL			
R24	57	.0024	ABRAMS4 75 SMAG	E+E- 1/76*
R25	J/PSI(3100) INTO (K+ K*(892)+)/TOTAL			
R25	87	.0031	ABRAMS4 75 SMAG	E+E- 1/76*
R25	POSSIBLY SEEN			
R25			WIJK 75 DASP	E+E- 1/76*
R26	J/PSI(3100) INTO (K0 K*(1420))/TOTAL			
R26	(.0019)OR LESS CL=.90	ABRAMS4 75 SMAG		E+E- 1/76*
R27	J/PSI(3100) INTO (K+ K*(1420+))/TOTAL			
R27	(.0019)OR LESS CL=.90	ABRAMS4 75 SMAG		E+E- 1/76*
R28	J/PSI(3100) INTO (K*(892)0 K*(892)0)/TOTAL			
R28	(.0006)OR LESS CL=.90	ABRAMS4 75 SMAG		E+E- 1/76*
R29	J/PSI(3100) INTO (K*(1420)0 K*(1420)0)/TOTAL			
R29	(.0018)OR LESS CL=.90	ABRAMS4 75 SMAG		E+E- 1/76*
R30	J/PSI(3100) INTO (K*(892)0 K*(1420)0)/TOTAL			
R30	30	.0037	ABRAMS4 75 SMAG	E+E- 1/76*
R31	J/PSI(3100) INTO (PEAR PI)/TOTAL			
R31 A	105	.0021	ABRAMS4 75 SMAG	E+ E- 1/76*
R31 A	40	.0023	WIJK 75 DASP	E+ E- 1/76*
R31 A	ASSUMING ANGULAR DISTRIBUTION (1.+COS(THETA)**2)			
R31	STUDENT			
R31 AVG	0.00216	0.00033	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R31 STUDENT	0.00216	0.00036	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R32	J/PSI(3100) INTO (PEAR PI)/(MU+ MU-)			
R32 A	20	.051	CRTEGEE 75 PLUT	E+ E- 1/76*
R32 A	ASSUMING ANGULAR DISTRIBUTION (1.+COS(THETA)**2)			
R33	J/PSI(3100) INTO (LAMBDA ANTILAMBDA)/TOTAL			
R33	19	.0016	ABRAMS4 75 SMAG	E+ E- 1/76*
R34	J/PSI(3100) INTO (NUCLEON ANTINUCLEON PI)/TOTAL			
R34 A	87	.0037	ABRAMS4 75 SMAG	E+ E- 1/76*
R34 A	INCLUDES P PBAR P10, NBAR P PI- AND N PBAR PI+			
R35	J/PSI(3100) INTO (P PBAR PI+PI-)/TOTAL			
R35 G	125	SEEN	GOLDHABER 75 SMAG	E+E- 2/76*
R35 G	INCLUDES LAMBDA ANTILAMBDA			
R36	J/PSI(3100) INTO (P PEAR PI+ PI- NEUTRALS)/TOTAL			
R36 I	91	SEEN	GOLDHABER 75 SMAG	E+E- 2/76*
R36 I	INCLUDES SIGMA ANTISIGMA ETC.			
R37	J/PSI(3100) INTO (LAMBDA ANTISIGMA)/(LAMBDA ANTILAMBDA)			
R37	(.22) OR LESS	GOLDHABER 75 SMAG		E+E- 2/76*
RADIATIVE DECAYS				
R				
R51	J/PSI(3100) INTO (2 GAMMA)/(E+ E-)			
R51	(.050)OR LESS CL=.90	BRAUNSCH 75 DASP		(P4)/(P1) E+E- 2/75*
R52	J/PSI(3100) INTO (PI0 GAMMA)/(E+ E-)			
R52	(.0.13) OR LESS CL=.90	BRAUNSCH 75 DASP		(P5)/(P1) E+E- 2/75*
R52 B	(.06) OR LESS CL=.90	BACCI 75 FRAG		E+E- 1/76*
R52 B	RE-STAT'ED BY US USING (E+E-)/HADRONS = .08			
R53	J/PSI(3100) INTO (ETA GAMMA)/TOTAL			
R53 U	.0014	.0005	WIJK 75 DASP	E+ E- 1/76*
R53 U	USING TOTAL WIDTH (69+-15) KEV			
R53	(.016)OR LESS CL=.90	BACCI 75 FRAG		E+E- 1/76*
R54	J/PSI(3100) INTO (ETA PRIME GAMMA)/(ETA GAMMA)			
R54	(5.) OR LESS CL=.90	WIJK 75 DASP		E+E- 1/76*
R54	4.0	2.5	HEINTZE 75 DESY	E+E- 1/76*
R54 B	(10.) OR LESS CL=.90	BALDINI2 75 FRAG		E+E- 1/76*
R54 B	RE-STAT'ED BY US USING (ETA GAMMA)/TOTAL = .0014			
R56	J/PSI(3100) INTO (X(2750) GAMMA)/(ETA GAMMA) FINAL STATE (3 GAMMA)			
R56	16	(1.)	HEINTZE 75 DESY	E+E- 1/76*

Data Card Listings

For notation, see key at front of Listings.

Mesons

J/ψ(3100), ψ(3700)

Table with columns for particle name, G-parity, JPC, and various decay widths. Includes entries for J/ψ(3100) and ψ(3700) with detailed decay data and branching ratios.

ψ(3700) 71 PSI(3700, JPC=1--) I=0
The ψ(3700) was discovered in e+e- collisions at SPEAR (ABRAMS 74), and since then has been extensively investigated in e+e- collisions at SPEAR and DORIS, as well as at conventional accelerators.
The hadronic interpretation of the ψ(3700) is closely related to that of J/ψ(3100), due to the decays ψ(3700) → J/ψ(3100)ππ and ψ(3700) → J/ψ(3100)η. These decays also determine the IG = 0- assignment for the ψ(3700) (TANENBAUM 75).
The JPC assignment, suggested by production in s-channel e+e- collisions, is confirmed by the observation of an interference between the reso-

nant and QED amplitudes, and by the angular distribution of the lepton pairs in the final state (LÜTH 75). These quantum numbers are compatible with all the data on partial decay widths (the multiplication decays mentioned above, the non-observation of J/ψ(3100)π0 and γγ modes, etc.).

71 PSI(3700) MASS (MEV)
Table listing mass measurements from various experiments like ABRAMS 74, CRIEGEE 75, LUTH 75, and PREPOST 75.

PSI(3700) - J/PSI(3100) MASS DIFFERENCE (MEV)
Table showing mass differences between PSI(3700) and J/PSI(3100) from experiments like DM 588.7 and LUTH 75.

71 PSI(3700) WIDTH (KEV)
Table listing partial widths for PSI(3700) from experiments like W 228 and LUTH 75.

71 PSI(3700) PARTIAL DECAY MODES
Table listing various decay channels for PSI(3700) such as PSI(3700) INTO E+ E-, PSI(3700) INTO MU+ MU-, and PSI(3700) INTO GAMMA INTO HADRONS.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS
The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, Pij, as follows: The diagonal elements are Pij ± δPij, where δPij = sqrt(δPij δPij), while the off-diagonal elements are the normalized correlation coefficients (δPij δPkl) / (δPij δPkl). For the definitions of the individual Pij, see the listings above; only those Pij appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.
Table with columns P13, P14, P15, P58, POTHER and corresponding values.

71 PSI(3700) PARTIAL WIDTHS (KEV)
Table listing partial widths for PSI(3700) into E+ E- and PSI(3700) into HAERONS from experiments like W1, W3, and W3.

71 PSI(3700) BRANCHING RATIOS
Table listing branching ratios for PSI(3700) into E+ E- and PSI(3700) into HAERONS from experiments like R1, R1 L, and R1 L.

Mesons

$\psi(3700)$

Data Card Listings

For notation, see key at front of Listings.

R2	PSI(3700) INTO (MU+ MU-)/TOTAL								
R2 H	.0084 .0018 HILGER 75 SPEC	E+E-	1/76*						
R2 H	RE-STAT'D BY US USING (J/PSI(3100)+ANYTHING)/TOTAL = .6								
R3	PSI(3700) INTO (HADRONS)/TOTAL								
R3	.981 .003 LUTH 75 SMAG	E+E-	1/76*						
R4	PSI(3700) INTO (MU+ MU-)/(E+ E-)								
R4	.89 .16 LUTH1 75 SMAG	E+E-	1/76*						
R5	PSI(3700) INTO (GAMMA INTO HADRONS)/TOTAL								
R5	.029 .004 LUTH 75 SMAG	E+E-	1/76*						
R	DECAYS INTO J/PSI(3100) + ANYTHING								
R10	PSI(3700) INTO (J/PSI(3100) + ANYTHING)/TOTAL								
R10	.57 .08 ABRAMS4 75 SMAG	E+E-	1/76*						
R10	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R10	0.581 0.045								
R11	PSI(3700) INTO (J/PSI(3100) + NEUTRALS)/(J/PSI(3100) + ANYTHING)								
R11 S	(.44) (.03) ABRAMS1 75 SMAG	E+E-	1/76*						
R11 S	SUPERSEDED BY TANENBAUM 76 (FOOTNOTE 13 OF TANENBAUM 75, AND R11 S PRIV. COMMUNICATION).								
R11	.41 .02 TANENBAUM 76 SMAG	E+E-	2/76*						
R11	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R11	0.410 0.019								
R12	PSI(3700) INTO (J/PSI(3100) P+ P-)/TOTAL								
R12	.32 .04 TANENBAUM 75 SMAG	E+E-	1/76*						
R12	.36 .06 WIJK 75 DASP	E+E-	1/76*						
R12	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R12	0.332 0.033								
R12	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT								
R12	0.331 0.027								
R12	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R13	PSI(3700) INTO (J/PSI(3100) P0 P0)/TOTAL								
R13	.18 .06 WIJK 75 DASP	E+E-	1/76*						
R13	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R13	0.172 0.024								
R14	PSI(3700) INTO (J/PSI(3100) P0 P0)/(J/PSI(3100) P+ P-)								
R14 H	(.64) (.15) HILGER 75 SPEC	E+E-	1/76*						
R14 H	IGNORING THE (J/PSI ETA) AND (J/PSI GAMMA GAMMA) DECAYS								
R15	PSI(3700) INTO (J/PSI(3100) ETA)/TOTAL								
R15	.043 .008 TANENBAUM 75 SMAG	E+E-	1/76*						
R15	.037 .015 WIJK 75 DASP	E+E-	1/76*						
R15	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R15	0.0417 0.0071								
R15	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT								
R15	0.0416 0.0070								
R15	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R16	PSI(3700) INTO (J/PSI(3100) GAMMA OR J/PSI(3100) P0)/TOTAL								
R16	(.00151)OR LESS CL=.90 TANENBAUM 75 SMAG	E+E-	2/76*						
R50	PSI(3700) INTO (J/PSI(3100) GAMMA GAMMA)/TOTAL								
R50 A	10 SEEN CASP3 75 DASP	E+E-	1/76*						
R50 B	5 SEEN DASP3 75 DASP	E+E-	1/76*						
R50 A	15 SEEN HEINTZE 75 SPEC	E+E-	1/76*						
R50 B	150 SEEN TANENBAUM 75 SMAG	E+E-	1/76*						
R50 AB	DECAY PROCEEDS VIA INTERMEDIATE STATE PC (SEE R51)								
R50 A	WITH J/PSI(3100) INTO (E+ E-)								
R50 B	WITH J/PSI(3100) INTO (MU+ MU-)								
R	HADRONIC DECAYS								
R20	PSI(3700) INTO (P+ P-)/TOTAL								
R20	(.00091)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R20	(.00019)OR LESS CL=.90 FELDMAN 75 SMAG	E+E-	1/76*						
R21	PSI(3700) INTO (RH00 P0)/TOTAL								
R21	(.0011)OR LESS CL=.90 ABRAMS4 75 SMAG	E+E-	1/76*						
R22	PSI(3700) INTO (2(P+ P-) P0)/TOTAL								
R22	.0035 .0015 ABRAMS4 75 SMAG	E+E-	1/76*						
R23	PSI(3700) INTO (K+ K-)/TOTAL								
R23	(.0016)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R23	(.00023)OR LESS CL=.90 FELDMAN 75 SMAG	E+E-	1/76*						
R24	PSI(3700) INTO (P+ P- K+ K-)/TOTAL								
R24	(.0005) ABRAMS4 75 SMAG	E+E-	1/76*						
R25	PSI(3700) INTO (PBAR P)/TOTAL								
R25	(.0011)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R25	.0004 .0002 ABRAMS4 75 SMAG	E+E-	1/76*						
R	RADIATIVE DECAYS								
R41	PSI(3700) INTO (GAMMA GAMMA)/TOTAL								
R41	(.0081)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R41 U	(.0051)OR LESS CL=.95 HUGHES 75 SPEC	E+E-	1/76*						
R42	PSI(3700) INTO (P0 GAMMA)/TOTAL								
R42	(.01) OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R42 U	(.0071)OR LESS CL=.95 HUGHES 75 SPEC	E+E-	1/76*						
R43	PSI(3700) INTO (ETA GAMMA)/TOTAL								
R43	(.0013)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R43 U	(.024)OR LESS CL=.95 HUGHES 75 SPEC	E+E-	1/76*						
R43 U	RE-STAT'D BY US USING (MU+MU-)/TOTAL = .01								
R44	PSI(3700) INTO (ETA PRIME GAMMA)/TOTAL								
R44	(.014)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R51	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (J/PSI(3100) GAMMA)								
R51	.04 .02 WIJK 75 DASP	E+E-	1/76*						
R51	.036 .007 TANENBAUM 75 SMAG	E+E-	1/76*						
R51	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R51	0.0364 0.0066								
R51	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT								
R51	0.0363 0.0066								
R51	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)								
R52	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (GAMMA GAMMA)								
R52	(.000013)OR LESS CL=.90 WIJK 75 CASP	E+E-	1/76*						
R53	PSI(3700) INTO (X(2750) GAMMA)/TOTAL, X(2750) INTO (GAMMA GAMMA)								
R53	(.00037)OR LESS CL=.90 WIJK 75 DASP	E+E-	1/76*						
R53	(.0031)OR LESS CL=.95 HUGHES 75 SPEC	E+E-	1/76*						

R54	PSI(3700) INTO (CHI(3410) GAMMA)/TOTAL *10**2								
R54	FOR DECAY MODES OF THE CHI INDICATED								
R54 F	(.13) (.05) FELDMAN 75 CHI INTO (P+P-1)OR(K+K								
R54 F	(.07) FELDMAN 75 CHI INTO (P+P-1)K+K-								
R54 F	(.14) (.07) FELDMAN 75 CHI INTO (2(P+P-1))								
R54 F	(.1) FELDMAN 75 CHI INTO (3(P+P-1))								
R54 F	(.5) OR LESS CL=.90 FELDMAN 75 CHI INTO (J/PSI GAMMA)								
R54 F	FOR DECAY MODES OF THE CHI INDICATED								
R55	PSI(3700) INTO (CHI(3530) GAMMA)/TOTAL *10**2								
R55 F	FOR DECAY MODES OF THE CHI INDICATED								
R55 F	(.027)OR LESS CL=.90 FELDMAN 75 CHI INTO (P+P-1)OR(K+K								
R55 F	(.2) (.1) FELDMAN 75 CHI INTO (2(P+P-1))								
R55 F	(.2) FELDMAN 75 CHI INTO (3(P+P-1))								
R55 F	FOR DECAY MODES OF THE CHI INDICATED								
R56	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (P+P-1) OR (K+K-)								
R56	(.00027)OR LESS CL=.90 FELDMAN 75 SMAG	E+E-	1/76*						
71 PSI(3700) G(I)*G(E+E-)/G(TOTAL) (KEV)									
THIS COMBINATION OF A PARTIAL WIDTH WITH THE PARTIAL WIDTH INTO E+E- AND WITH THE TOTAL WIDTH IS OBTAINED FROM THE INTEGRATED CROSS-SECTION INTO CHANNEL (I) IN THE E+E- ANNIHILATION. SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.									
G3	G(HADRONIC)*G(E+E-)/G(TOTAL)								
G3	2.2 .4 ABRAMS4 75 SMAG	E+E-	1/76*						
***** REFERENCES FOR PSI(3700) *****									
ABRAMS	74 PRL 33 1453	+BRIGGS,AUGUSTIN,BOYARSKI+	(LBL+SLAC)	V71					
ABRAMS1	75 PRL 34 1181	+BRIGGS,CHINCENSKY,FRIEDBERG,+	(LBL+SLAC)	V71					
ABRAMS4	75 STANFORD SYMP.	G.S.ABRAMS	(LBL)	V71					
AUBERT	75 PRL 33 1624	+BECKER,BIGGS,BURGER,GLENN+	(MIT+BNL)	V71					
CAMERINI	75 PRL 35 483	+EARNED,PREFPOST,ASH,ANDERSON,+	(WISC+SLAC)	V71					
CRIGEE	75 PL 33B 489	+DEHNE,FRANKE,HORLITZ,KRECHLOCK+	(DESY)	V71					
DASP3	75 PL 57B 407	BRAUNSCHWEIG,KONIGS,+	(AACH+DESY+MIM+TKY)	V71					
FELDMAN	75 PRL 35 821	+JEAN-MARIE,SADOULET,VANNUCCI,+	(LBL+SLAC)	V71					
FELDMAN	75 STANFORD SYMP.	G.J.FELDMAN	(SLAC)	V71					
GREGO	75 PL 56B 367	+PANCHER-SRIVASTAVA,SRIVASTAVA	(FRAS)	V71					
HEINTZE	75 STANFORD SYMP.	J.FEINTZE	(HEIDELBERG)	V71					
JACKSON	75 NIM 128 13	J.D.JACKSON,D.SCHARRE	(LBL)	V71					
HILGER	75 PRL 35 625	+BERON,FORD,HOFSTADTER,HOWELL,+	(STAN+PENN)	V71					
HUGHES	75 PREP-HEPL 765	+BERON,CARRINGTON,FORD,HILGER,+	(STAN+PENN)	V71					
LUTH	75 PRL 35 1124	+BOYARSKI,LYNCH,BREIBENBACH,+	(SLAC+LBL+JPC)	V71					
LUTH1	75 SLAC-PUB-1599	V.LUTH, PALERMC CNF.PROC.+	(SLAC+LBL)	V71					
LIBERMAN	75 STANFORD SYMP.	A.D.LIBERMAN	(STANFORD)	V71					
PREFPOST	75 STANFORD SYMP.	R.PREFPOST	(WISCONSIN)	V71					
SIMPSON	75 PRL 35 699	+BERON,FORD,HILGER,HOFSTADTER,+	(STAN+PENN)	V71					
TANENBAU	75 PRL 35 1323	TANENBAUM,WHITAKER,ABRAMS,+	(LBL+SLAC)	V71					
TANENBAUM	75 PRL SLAC-PUB1696	TANENBAUM,ABRAMS,BOYARSKI,+	(SLAC+LBL)	V71					
WIJK	75 STANFORD SYMP.	B.+WIJK	(DESY)	V71					

Additional States in S-channel e+e- Collisions

No evidence for narrow states has been found in the mass intervals 1910-2545 MeV and 2970-3090 MeV (BACCI 75, ESPOSITO 75). Except for the $\psi(3700)$, no additional narrow states have been found in the mass interval 3200-7600 MeV (BOYARSKI 75, SCHWITTERS 75).

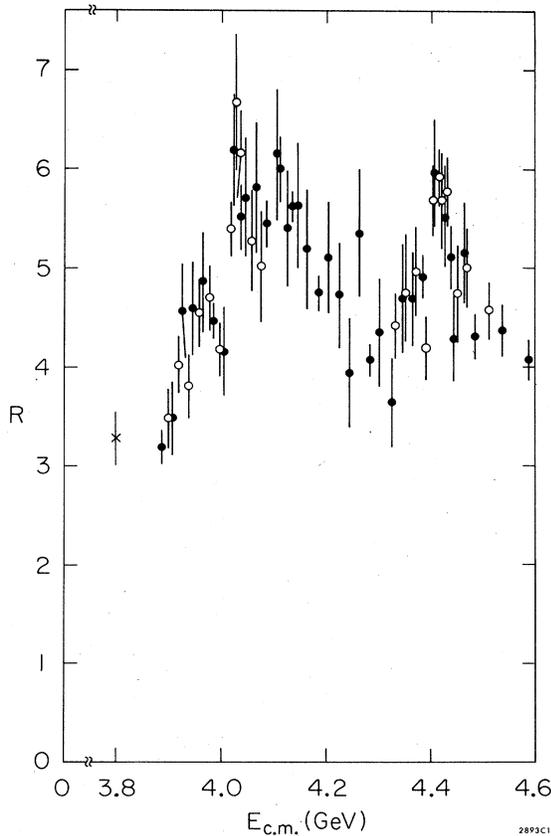
The broad enhancement observed in the e^+e^- total cross section at SPEAR (AUGUSTIN 75) is now resolved into two distinct objects at 4100 and 4400 MeV (SIEGRIST 76), with indications of even more detailed substructure especially at 4100 MeV (see figure). Due to there being very limited information available on the properties of $\psi(4100)$ and $\psi(4400)$, their resonant interpretation is not yet established. If in fact they are resonances, then their relatively large widths suggest that they are hadronic, with $J^{PC} = 1^{--}$.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\psi(4100)$, $\psi(4400)$, $X(2750)$



The ratio of cross sections $\frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$ in the $\psi(4100)$ - $\psi(4400)$ region (from SIEGRIST 76).

$\psi(4100)$

72 PSI(4100, JPG=) I=

72 PSI(4100) MASS (MEV)				
M	(4150.)	AUGUSTIN 75 SPEAR	2.4-5.0 E+E-	2/75*
M	(4100.)	SCHWITTER 75 SMAG	E+E-	1/76*

REFERENCES FOR PSI(4100)

AUGUSTIN 75 PRL 34 764	+BOYARSKI, ABRAMS, BRIGGS+	(SLAC+LBL)	V72
BACCI 75 PL 58B 481	+BIDDLI, PENSO, STELLA, +	(ROMA+FRAS)	V72
BOYARSKI 75 PRL 34 762	+BREIDENBACH, ABRAMS, BRIGGS, +	(SLAC+LBL)	V72
ESPOSITO 75 PL 58B 478	+FELICETTI, PERUZZI, +	(FRAS+NAPO+PADO+ROMA)	V72
SCHWITTE 75 STANFORD SYMP.	R.F.SCHWITTERS	(SLAC)	V72
SIEGRIST 76 SUBM. TO PRL	+ABRAMS, BOYARSKI, BREIDENBACH, +	(LBL+SLAC)	V72

$\psi(4400)$

73 PSI(4400, JPG=) I=

73 PSI(4400) MASS (MEV)				
M	4414.	7.	SIEGRIST 76 SMAG	E+E-
				2/76*

73 PSI(4400) WIDTH (MEV)				
M	33.	10.	SIEGRIST 76 SMAG	E+E-
				2/76*

73 PSI(4400) BRANCHING RATIOS				
R1	PSI(4400) INTO (E+ E-)/TOTAL	*10**5		
R1	1.3	.3	SIEGRIST 76 SMAG	E+E-
				2/76*

REFERENCES FOR PSI(4400)
 SCHWITTE 75 STANFORD SYMP. R.F.SCHWITTERS (SLAC) V73
 SIEGRIST 76 SUBM. TO PRL +ABRAMS, BOYARSKI, BREIDENBACH, + (LBL+SLAC) V73

X(2750)

54 X(2750, JPG=) I=

OBSERVED IN THE SEQUENTIAL RADIATIVE DECAY OF THE J/PSI(3100) INTO X(2750) GAMMA, X(2750) INTO GAMMA GAMMA, AND POSSIBLY PBAR P (HEINTZE 75). THIS SUGGESTS QUANTUM NUMBER ASSIGNMENTS C=+, IG=0+ OR 1-. NEEDS CONFIRMATION. OMITTED FROM TABLE.

54 X(2750) MASS (MEV)				
M	(2750.)	HEINTZE 75 DESY	E+E-	1/76*

54 X(2750) PARTIAL DECAY MODES				
P1	X(2750) INTO GAMMA GAMMA	DECAY MASSES		
P2	X(2750) INTO PBAR P			

54 X(2750) BRANCHING RATIOS				
R1	X(2750) INTO (PBAR P)/TOTAL	WIJK 75 DASP	E+E-	1/76*
R1	2 POSSIBLY SEEN			

REFERENCES FOR X(2750)

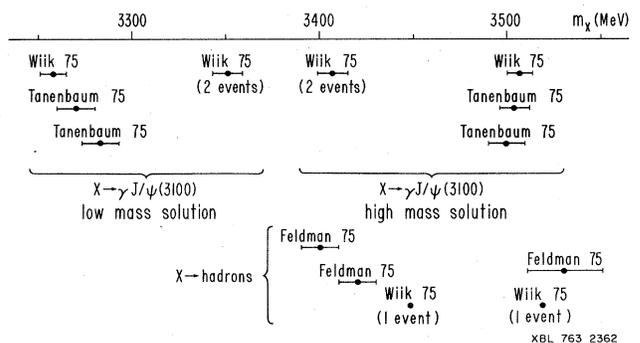
HEINTZE 75 STANFORD SYMP.	J.FEINTZE	(HEIDELBERG)	V54
WIJK 75 STANFORD SYMP.	B.H.WIJK	(DESY)	V54

States Observed in Radiative Decays of $\psi(3700)$

The properties, and even the number, of intermediate states in radiative decay of $\psi(3700)$ are not well known. The main problems are:

- 1) small statistics in most experiments;
- 2) ambiguity in the determination of the mass of P_c ;
- 3) it is not known whether the $X(3530)$ is a single state, or whether the relatively large observed width is due to a superposition of several narrow states.

The data seem to require the presence of at least 3 states in the mass interval 3250-3550 MeV; the figure summarizes all the present claims.



Claims for intermediate states X in the radiative decay $\psi(3700) \rightarrow \gamma X$.

Mesons

P_c(3300 or 3500), χ(3410), χ(3530)

Data Card Listings

For notation, see key at front of Listings.

P_c(3300 or 3500)

55 PC(3300 OR 3500, JPC=) I=

OBSERVED IN THE RADIATIVE SEQUENTIAL DECAY OF THE PSI(3700) INTO PC GAMMA, PC INTO J/PSI(3100) GAMMA (BRAUNSWIG 75, CONFIRMED BY TANENBAUM 75), THEREFORE C=+. MASS DETERMINATION AMBIGUOUS DUE TO TWO POSSIBLE (GAMMA J/PSI) COMBINATIONS IN THE FINAL STATE (J/PSI GAMMA GAMMA). NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

55 PC MASS (MEV)

Table with columns for mass solutions (LW, HIGH MASS), particle names (DASP3, TANENBAUM, WIJK), and branching ratios (E+E-, 1/76*).

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Table with columns for high mass solution, particle names, and branching ratios.

AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0) AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

55 PC PARTIAL DECAY MODES

Table with columns for decay modes (PC INTO J/PSI(3100) GAMMA, etc.) and decay masses.

55 PC BRANCHING RATIOS

Table with columns for branching ratios (DOMINANT, NOT SEEN, etc.) and branching ratios.

REFERENCES FOR PC

Table listing references for PC (BRAUNSCHWEIG, FELDMAN, HEINTZE, SIMPSON, TANENBAUM, WIJK).

χ(3410)

56 CHI(3410, JPC=) I=

OBSERVED IN THE RADIATIVE DECAY OF PSI(3700) INTO CHI(3410) GAMMA (FELDMAN 75), THEREFORE C=+. THE OBSERVED DECAY INTO (PI+PI-) OR (K+K-) IMPLIES G=+, JP=0+, 2+, CONFIRMATION OF BOTH (PI+PI-) AND (K+K-) MODES WOULD ESTABLISH I=0. NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

56 CHI(3410) MASS (MEV)

Table with columns for mass solutions (LW, HIGH MASS), particle names (FELDMAN, STUDENT), and branching ratios (E+E-, 1/76*).

56 CHI(3410) PARTIAL DECAY MODES

Table with columns for decay modes (CHI(3410) INTO PI+ PI-, etc.) and decay masses.

56 CHI(3410) BRANCHING RATIOS

SEE BRANCHING RATIOS R54 OF PSI(3700)

REFERENCES FOR CHI(3410)

Table listing references for CHI(3410) (FELDMAN, TANENBAUM).

χ(3530)

57 CHI(3530, JPC=) I=

OBSERVED IN RADIATIVE DECAY OF PSI(3700) INTO CHI(3530) GAMMA (FELDMAN 75), THEREFORE C=+. CAN BE INTERPRETED AS A SINGLE (BROAD) STATE, OR AS A SUPERPOSITION OF SEVERAL NARROW STATES. NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

57 CHI(3530) MASS (MEV)

Table with columns for mass solutions (LW, HIGH MASS), particle names (FELDMAN), and branching ratios (E+E-, 1/76*).

57 CHI(3530) PARTIAL DECAY MODES

Table with columns for decay modes (CHI(3530) INTO PI+ PI-, etc.) and decay masses.

57 CHI(3530) BRANCHING RATIOS

SEE BRANCHING RATIOS R55 OF PSI(3700)

REFERENCES FOR CHI(3530)

Table listing references for CHI(3530) (FELDMAN, TANENBAUM).

Data Card Listings

For notation, see key at front of Listings.

Baryons
N's and Δ 'sNote on N's and Δ 'sI. Determination of Resonance Parameters

Values of masses, widths, and branching ratios are obtained mainly from phase-shift analyses. In production experiments, in fact, it is seldom clear which of the many states at similar masses is being observed. In addition to a few complete phase-shift analyses, we have other analyses, done by using somewhat incomplete data, by various different groups, but we are quite far from having reliable masses and widths derived therefrom.

There are essentially two problems in obtaining reliable resonance parameters. First there is often disagreement as to just what the values of the partial-wave amplitudes are. This problem is obviously related to the quality and quantity of the data and to the procedures used to determine the amplitudes. Secondly, even if smooth curves were available for the amplitudes, there would still be some parametrization-dependent ambiguity in deciding what the resonance parameters should be. From a theoretical standpoint the most unambiguously defined resonance parameters are the pole position and residue, and it has been found in practice that, given sufficiently precise partial-wave amplitudes, these quantities can be extracted in a stable and parametrization-independent way, in spite of the fact that they require an extrapolation away from the physical region. This point has been discussed in detail with regard to the $\Delta(1232)$ in previous editions of this review^{1,2}. We list available pole parameter determinations in the Data Card Listings, and many further discussions can be found in the corresponding references, e.g., NOGOVA 73, SPEARMAN 74, BALL 75, LICHTENBERG 75, LONGACRE 75, and VASAN 75.

At the beginning of the Data Card Listings for N's and Δ 's, we present a table giving our evaluation of the N and Δ resonances based on information contained in the Listings. In the Table of Particle Properties, we do not quote values and errors for most parameters, but give only ranges for masses and widths in order to

emphasize that in some cases these parameters are quite poorly determined.

References for Section I

1. Particle Data Group, Rev. Mod. Phys. **43**, S114 (1971).
2. Particle Data Group, Phys. Lett. **39B**, 103 (1972).

For other references see the Data Card Listings.

II. $\pi N \rightarrow \pi N, \eta N, K\Sigma, K\Lambda$

The most recent available $\pi N \rightarrow \pi N$ amplitudes are from the analysis of the Saclay group, AYED 74. Preliminary results of more recent analyses have been discussed at meetings^{1,2}, but resonance parameters are not yet available. The results of AYED 74 are shown in Figs. II.1-II.6. The figure captions summarize the known resonances in the various partial waves. We restrict our discussion here to mentioning the possible new effects seen in recent analyses. All resonance parameters quoted under AYED 74 derive from energy-dependent fits (Breit-Wigner with background) to the partial-wave amplitudes which were determined in an energy-independent shortest-path analysis.

S₃₁^{''}: In Fig. II.1, the energy-dependent fit of AYED 74 suggests that the broad structure around 2000 MeV is associated with a second resonance with

$$M = 2001 \text{ MeV}, \Gamma_{\text{tot}} = 307 \text{ MeV}, x_{e1} = 0.08$$

There is some additional evidence for such an effect from LANGBEIN 73 and DEANS 75, which are energy-independent and energy-dependent analyses, respectively, of $K\Sigma$ associated production.

P₃₃^{''}: The small dip in the imaginary part of P_{33} and the zero in the real part around 1900 MeV are interpreted in the energy-dependent fit of AYED 74 as a second P_{33} resonance,

$$M = 1904 \text{ MeV}, \Gamma_{\text{tot}} = 204 \text{ MeV}, x_{e1} = 0.19$$

ALMEHED 72 found two effects, one at ~ 1680 MeV and another at ~ 2150 MeV, as did DEANS 75.

However, LANGBEIN 73 found evidence for the higher mass effect only. The existence of some sort of effect at ~ 2200 MeV seems definitely established

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

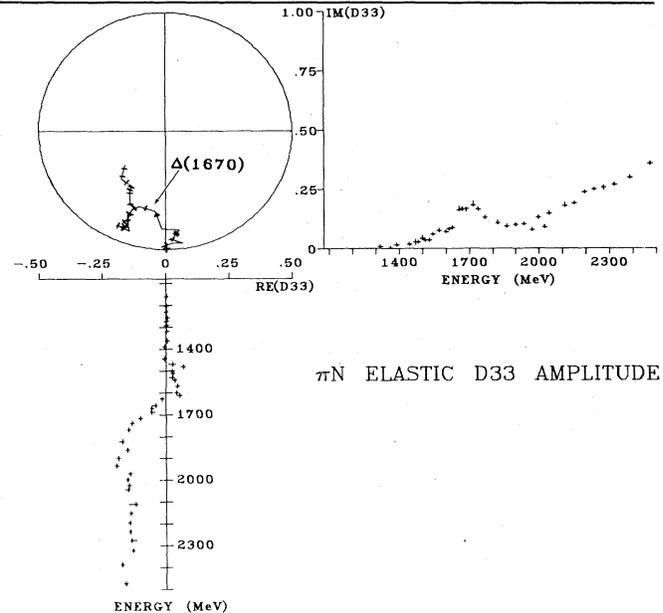
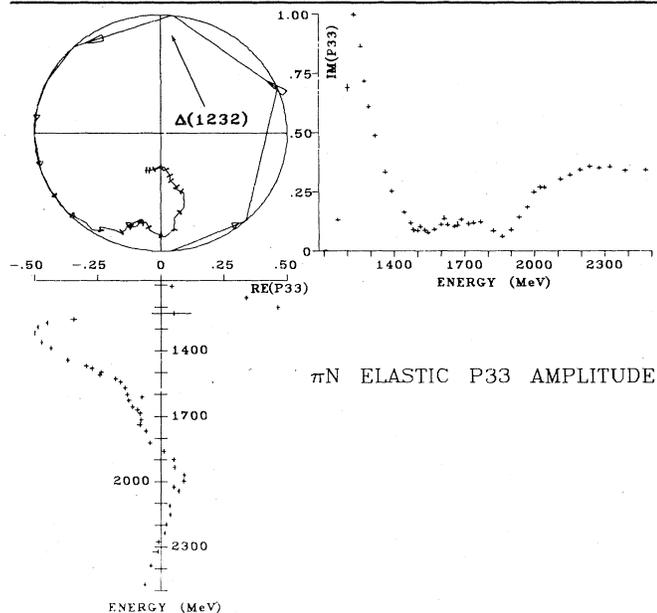
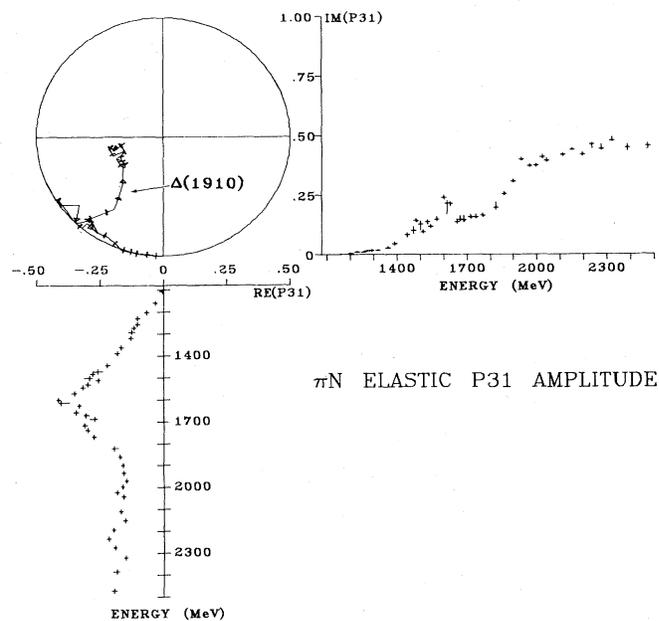
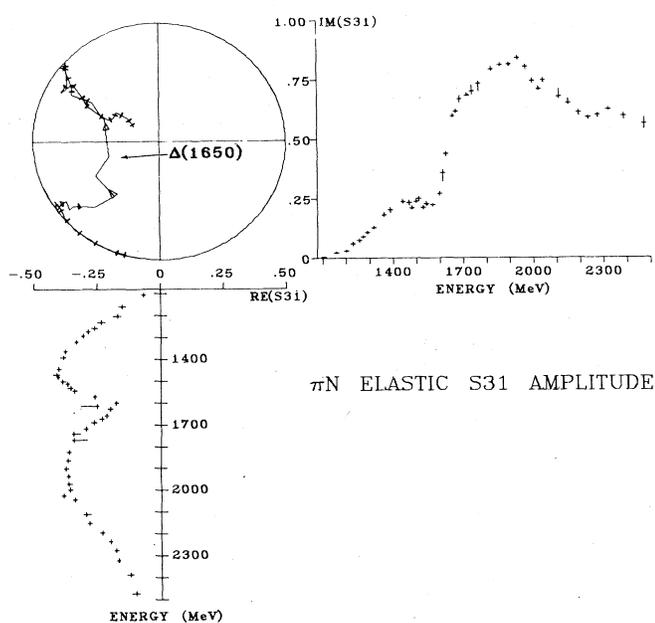


Fig. II.1. Amplitudes for $I = 3/2$ π N elastic scattering in the $J = 1/2$ and $J = 3/2$ waves from AYED '74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $\Delta(1650)$, $\Delta(1910)$, $\Delta(1232)$, and $\Delta(1670)$ in the S_{31} , P_{31} , P_{33} , and D_{33} waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for other possible resonances in the S_{31} and P_{33} waves.

Data Card Listings

For notation, see key at front of Listings.

Baryons

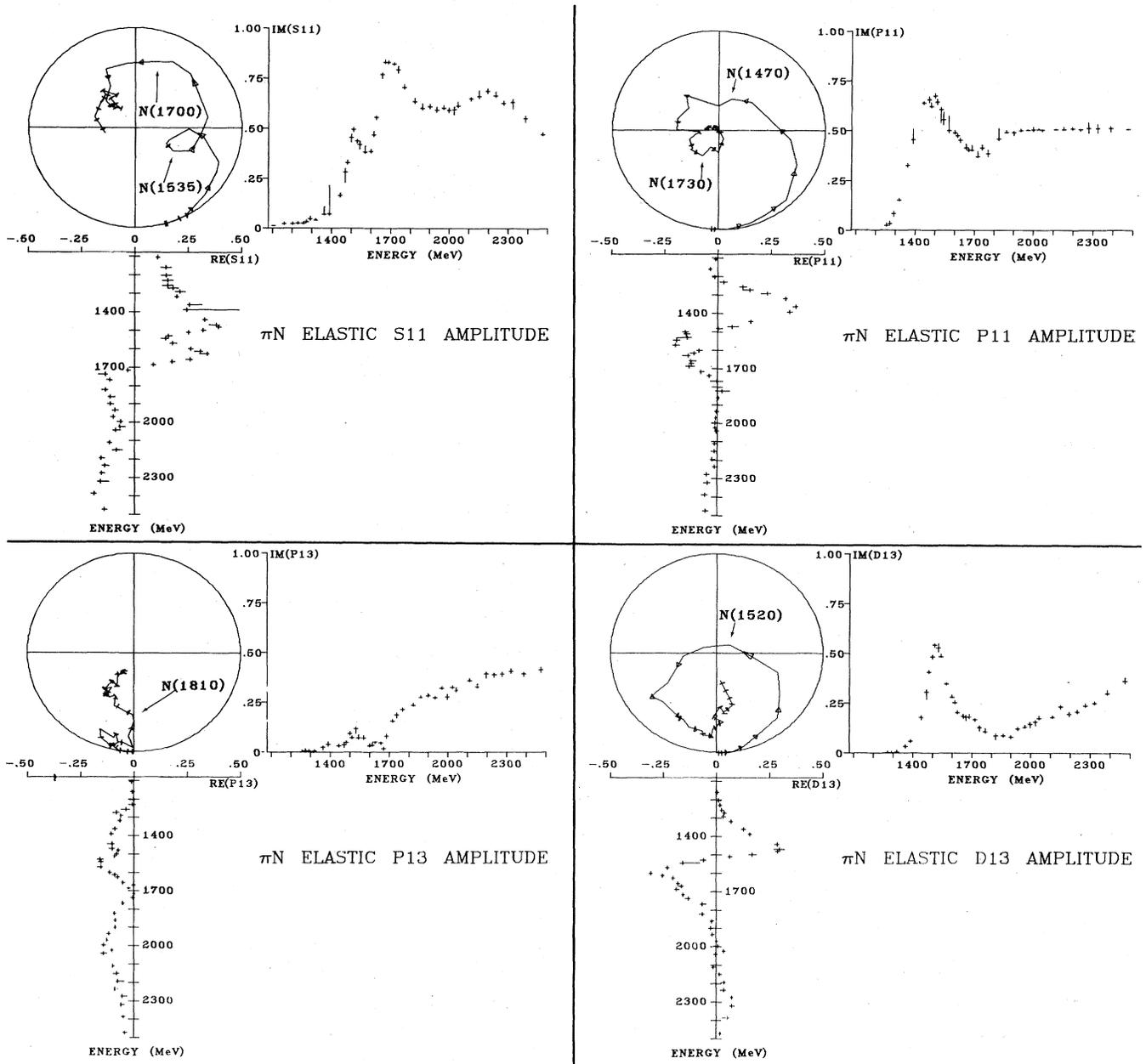
N's and Δ 's

Fig. II.2. Amplitudes for $I = 1/2$ πN elastic scattering in the $J = 1/2$ and $J = 3/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $N(1535)$ and $N(1700)$ in the S_{11} wave, the $N(1470)$ and $N(1730)$ in the P_{11} wave, the $N(1810)$ in the P_{13} wave, and the $N(1520)$ in the D_{13} wave; these are indicated on the above Argand plots. The P_{11} wave also contains the nucleon pole, 138 MeV below threshold. See the Data Card Listings and the accompanying mini-review for other possible resonances in these waves.

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

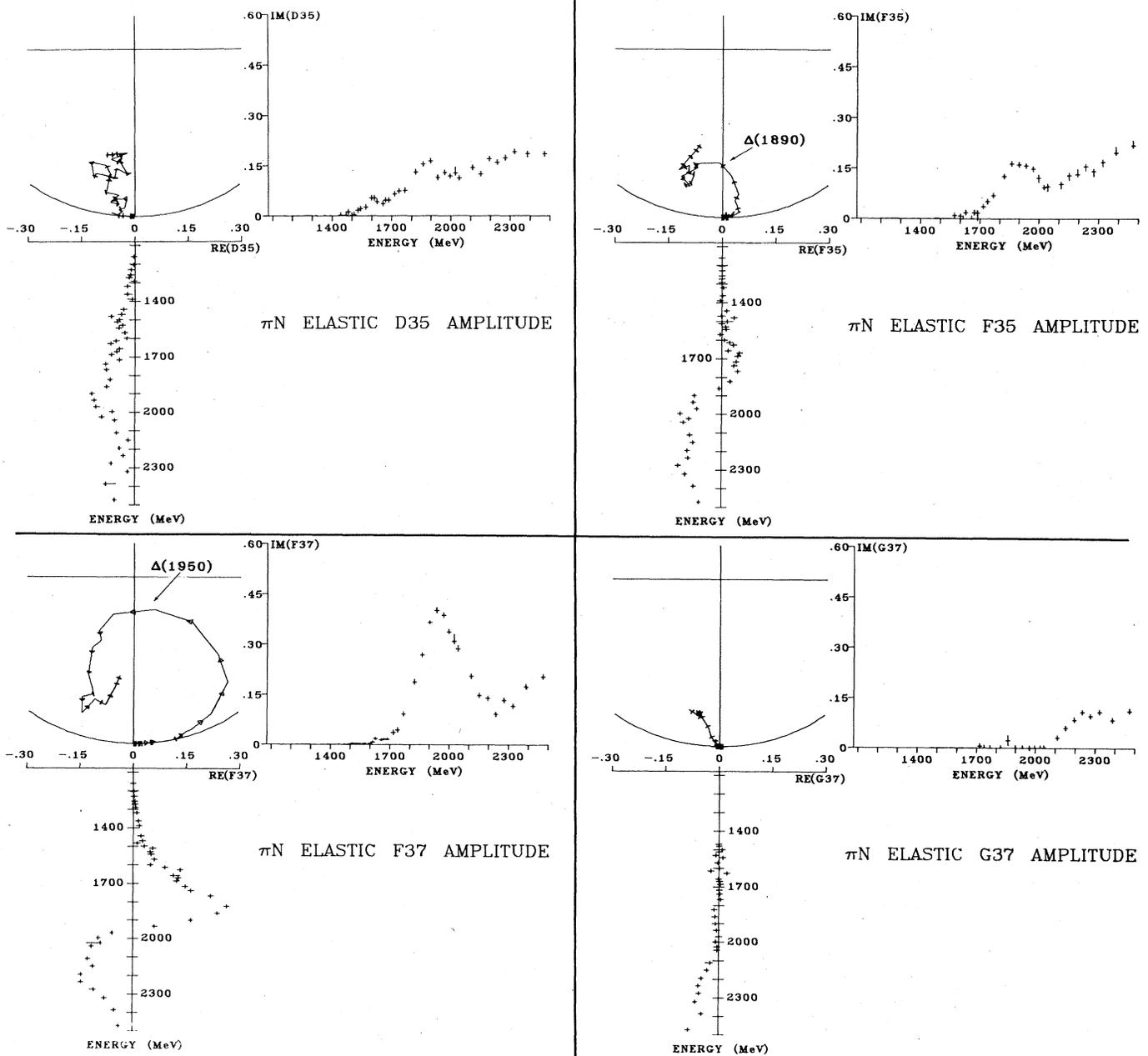


Fig. II.3. Amplitudes for $I = 3/2$ π N elastic scattering in the $J = 5/2$ and $J = 7/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $\Delta(1890)$ and $\Delta(1950)$ in the F35 and F37 waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for another possible resonance in the D35 wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

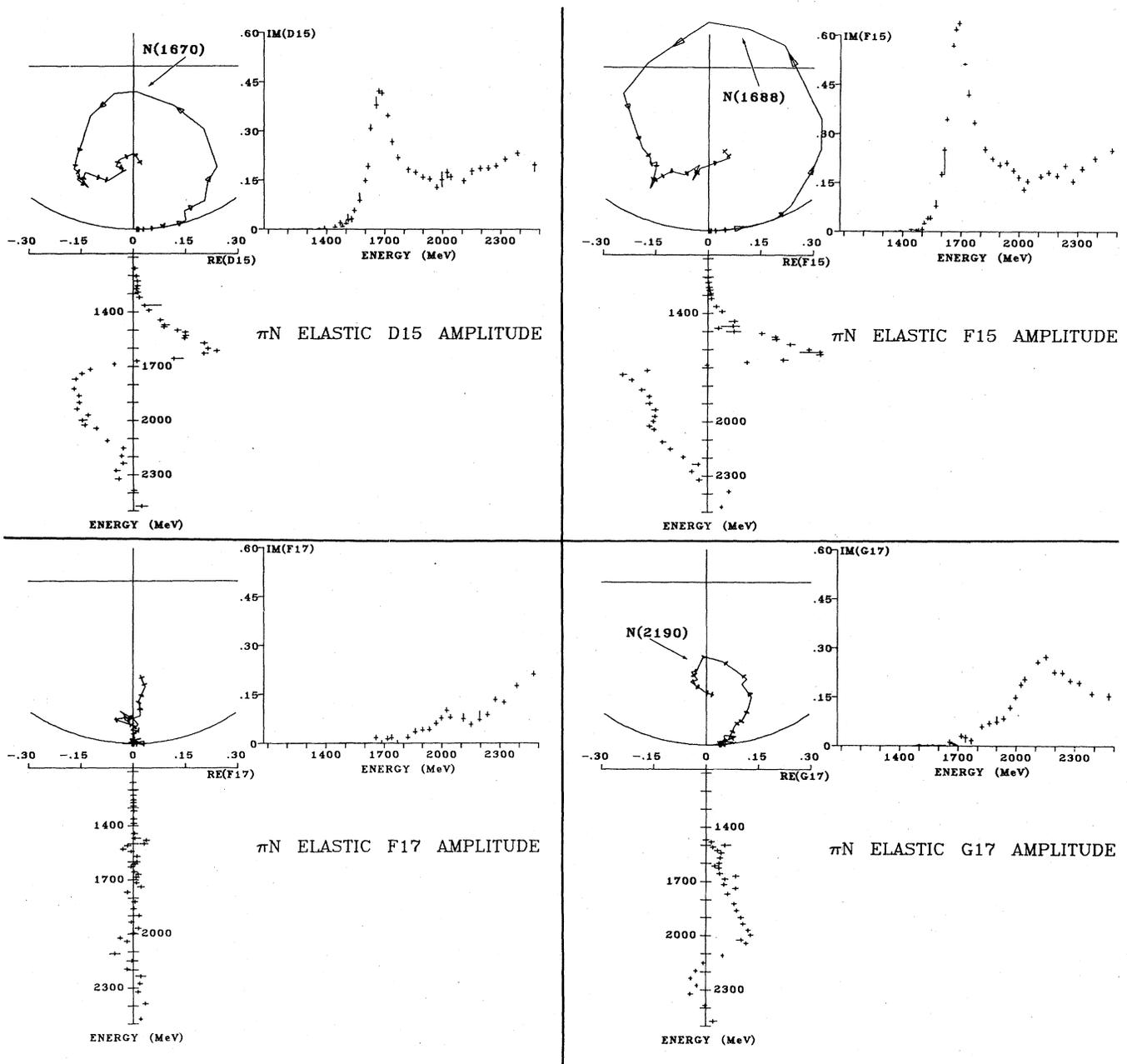
N's and Δ 's

Fig. II.4. Amplitudes for $I = 1/2$ πN elastic scattering in the $J = 5/2$ and $J = 7/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the N(1670), N(1688), and N(2190) in the D₁₅, F₁₅, and G₁₇ waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for other possible resonances in the D₁₅, F₁₅, and F₁₇ waves.

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

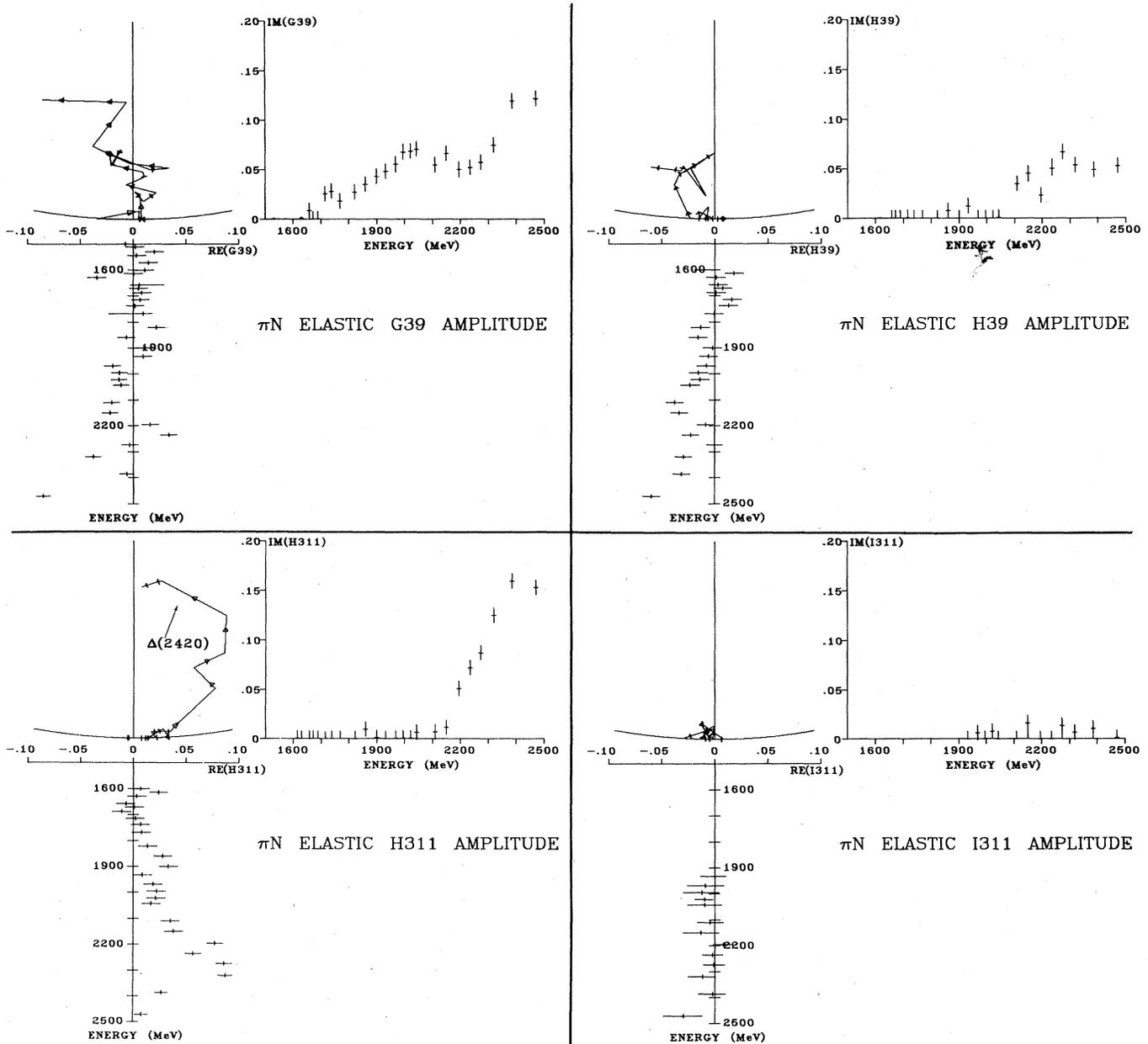


Fig. II.5. Amplitudes for $I = 3/2$ πN elastic scattering in the $J = 9/2$ and $J = 11/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All energy axes run from 1500 to 2500 MeV. The only established resonance in these waves is the $\Delta(2420)$ in the $H_3 11$ wave; it is indicated on the above $H_3 11$ Argand plot. See the Data Card Listings and the accompanying mini-review for another possible resonance in the G_{39} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

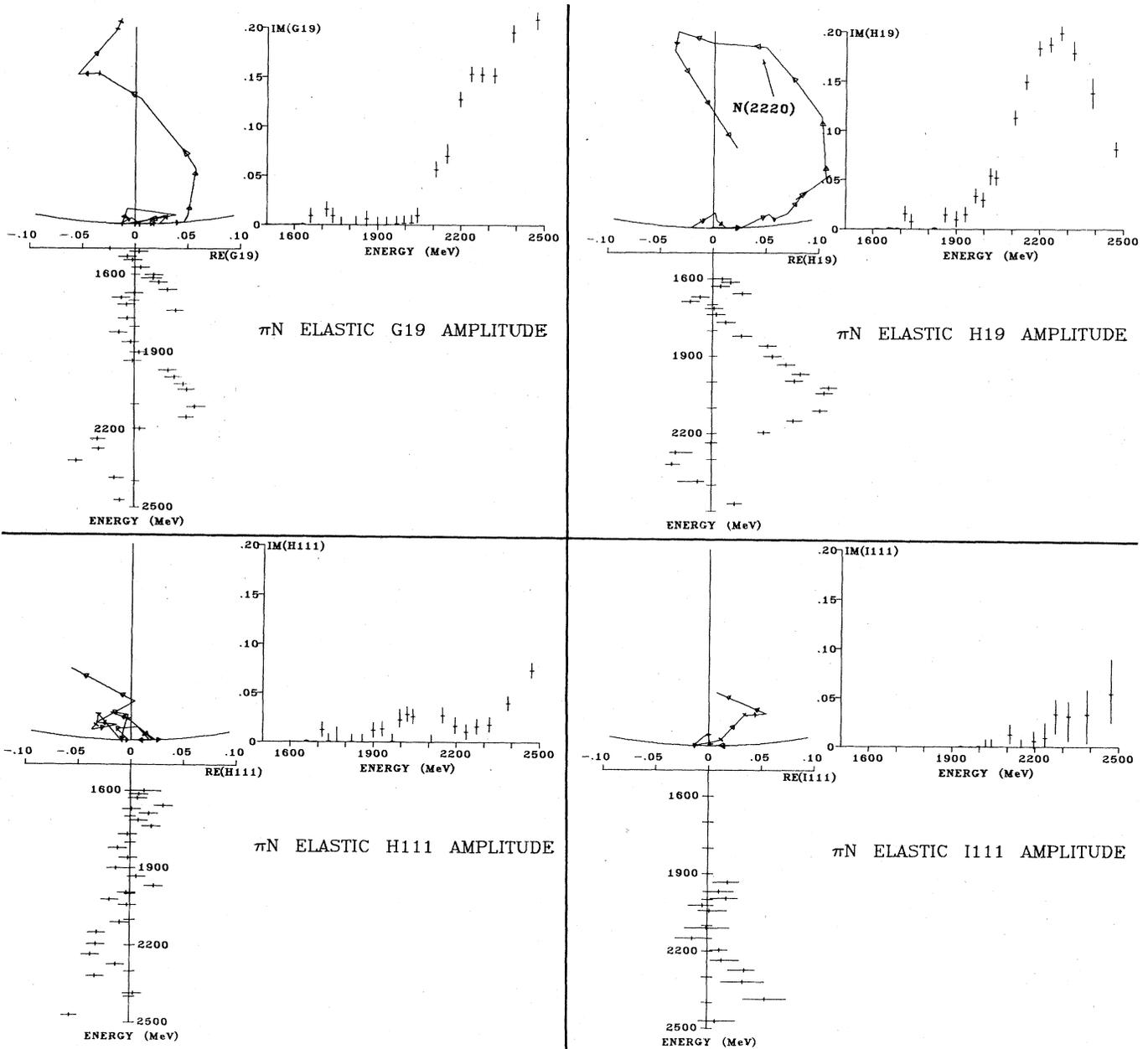
N's and Δ 's

Fig. II.6. Amplitudes for $I = 1/2$ πN elastic scattering in the $J = 9/2$ and $J = 11/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from 1500 to 2500 MeV. The only established resonance in these waves is the N(2220) in the H₁₉ wave; it is indicated on the above H₁₉ Argand plot. See the Data Card Listings and the accompanying mini-review for another possible resonance in the G₁₉ wave.

Baryons

N's and Δ 's

by the backward π^+p scattering experiment of REY 74, but the quantum numbers are not yet established [see the Listings for $\Delta(2160)$].

P_{31} : The P_{31} resonance mass quoted by AYED 74 ($M = 1786$ MeV) is considerably lower than in previous analyses (except AYED 70) which find M in the region 1900 to 1950 MeV.

D_{35} : AYED 74, LANGBEIN 73, and DEANS 75 all find evidence for a D_{35} state somewhere in the mass range 1900-2000 MeV, ALMEHED 72 found an effect at 2200 MeV, but with such a large width, 600 MeV, that a resonance interpretation is questionable.

S_{11}''' : In Fig. II.2 are shown the low spin $I = 1/2$ waves of AYED 74. They associate the broad effect at ~ 2300 MeV with a third S_{11} resonance:

$$M = 2283 \text{ MeV}, \Gamma_{\text{tot}} = 310 \text{ MeV}, x_{\text{el}} = 0.14 .$$

ALMEHED 72 find a similar effect at 2100 MeV.

P_{11}' : AYED 74 claim the "Roper" to be split:

$$M_{\text{low}} = 1413 \text{ MeV}, \Gamma_{\text{tot}} = 187 \text{ MeV}, x_{\text{el}} = 0.55 ;$$

$$M_{\text{high}} = 1532 \text{ MeV}, \Gamma_{\text{tot}} = 89 \text{ MeV}, x_{\text{el}} = 0.12 .$$

While such a splitting might be a natural explanation for the relatively low mass and broad width of the "Roper" as seen in production experiments, the evidence in Fig. II.2, especially the projections, seems only tentative. The new Saclay ηn data and partial-wave analysis of FELTESSE 75 are not inconsistent with the parameters of the high mass state. However, an equally acceptable fit to their data is obtained by introducing a new P_{13} resonance with $M = 1530$ MeV, $\Gamma = 79$ MeV, and $\sqrt{x_{\text{el}}} = 0.33$.

D_{13}'' : A somewhat larger effect in this wave at ~ 1700 MeV is identified by AYED 74 as a second D_{13} resonance:

$$M = 1710 \text{ MeV}, \Gamma_{\text{tot}} = 100 \text{ MeV}, x_{\text{el}} = 0.09 .$$

The projections reveal a shoulder in the imaginary part and a small dip in the real part. A similar effect exists in ALMEHED 72, but was not claimed as a resonance. Further evidence for this state is found by LANGBEIN 73, DEANS 75, and KNAESEL 75 ($K\bar{A}$ associated production), as well as recent pion

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photoproduction analyses and the LBL-SLAC $\pi\pi N$ analysis (see III and IV below).

D_{13}''' : Along with ALMEHED 72, AYED 74 find evidence for a D_{13} state at ~ 2000 MeV:

$$M = 2029 \text{ MeV}, \Gamma_{\text{tot}} = 116 \text{ MeV}, x_{\text{el}} = 0.10 .$$

The evidence in the projections seems less than for the effect at 1700 MeV. Evidence for coupling of this state to the $K\bar{\Sigma}$ channel is found by DEANS 75, but not by LANGBEIN 73.

D_{15}'' : The rapid decrease in the real part of this amplitude around 2100 MeV (see Fig. II.4) is associated with a second D_{15} resonance by AYED 74:

$$M = 2100 \text{ MeV}, \Gamma_{\text{tot}} = 220 \text{ MeV}, x_{\text{el}} = 0.08 .$$

These parameters are in reasonable agreement with those of ALMEHED 72.

F_{15}'' : The broad shoulder in the real part and dip in the imaginary part are identified with a second F_{15} resonance by AYED 74:

$$M = 1989 \text{ MeV}, \Gamma_{\text{tot}} = 179 \text{ MeV}, x_{\text{el}} = 0.08 .$$

ALMEHED 72 find an effect at ~ 2200 MeV, and LANGBEIN 73 and DEANS 75 find some evidence for coupling to the $K\bar{\Sigma}$ channel at ~ 2000 MeV.

Higher Waves. Besides confirming the already accepted $H_{3,11}$ and H_{19} , AYED 74 present evidence for two additional states in the 2200 MeV region:

$$G_{39}: M = 2174 \text{ MeV}, \Gamma_{\text{tot}} = 205 \text{ MeV}, x_{\text{el}} = 0.04 ;$$

$$G_{19}: M = 2133 \text{ MeV}, \Gamma_{\text{tot}} = 193 \text{ MeV}, x_{\text{el}} = 0.10 .$$

The G_{39} effect in Fig. II.5 is not too convincing. Notice that ALMEHED 72 claimed a P_{33} resonance in this same region; see $\Delta(2160)$ for further possibilities in this mass region. In contrast, the evidence for G_{19} , Fig. II.6, looks quite good; note, in particular, the behavior of the real part projection in the region 1900-2500 MeV.

References for Section II

1. R. L. Kelly, in New Directions in Hadron Spectroscopy (ANL-HEP-CP-75-58), eds. S. L. Kramer and E. L. Berger, 1975.
2. E. Pietarinen, Universität Karlsruhe Reports TKP 4/75 and TKP 15/75.

For other references see the Data Card Listings.

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In the isobar model, the amplitude for the reaction $\pi N \rightarrow N\pi\pi$ is written

$$T = \sum \left\{ \begin{array}{l} T_{\Delta\pi}^{\text{JILL}'}(E) BW_{\Delta}(E_{\pi\pi}) X_{\Delta\pi}^{\text{JILL}'} \\ + T_{\rho N}^{\text{JILL}'}(E) BW_{\rho}(E_{\pi\pi}) X_{\rho N}^{\text{JILL}'} \\ + T_{\epsilon N}^{\text{JILL}'}(E) BW_{\epsilon}(E_{\pi\pi}) X_{\epsilon N}^{\text{JILL}'} \end{array} \right\}.$$

In this expression the BW's denote either appropriate Breit-Wigner's or the corresponding two-body amplitude. The functions X contain all the angular information; these are well-defined functions depending only on which isobars are used in the model. The partial-wave amplitudes, $T_{\Delta\pi}^{\text{JILL}'}$, etc., may be indicated by

$$\Delta\pi, LL'_{2I2J} \quad \rho N, LL'_{2I2J} \quad \epsilon N, LL'_{2I2J}$$

where L is the incoming (πN) angular momentum, and L' is the outgoing angular momentum between the isobar [$\rho, \epsilon (= \pi\pi \text{ I} = 0, \text{ S wave}), \Delta$] and the remaining hadron (N or π); as usual I and J are the isospin and total spin ($\vec{J} = \vec{L} + \vec{S} = \vec{L}' + \vec{S}'$), respectively. Often the ρ has a subscript 1 or 3 which denotes twice the ρN total spin. The generalization of the model to include more two-body final-state interactions is obvious, but only the three included above, Δ , ρ , and ϵ , have actually been used in presently published analyses. It should be noted that one assumes the partial-wave amplitudes depend only on the c.m. energy E and not on the diparticle subenergies, $E_{\pi\pi}$ and $E_{\pi N}$.

The largest isobar analysis to date is that of the LBL-SLAC collaboration, LONGACRE 75 and Ref. 1, which analyzes 200K events of $\pi^- p \rightarrow \pi^- \pi^0 p$, $\pi^+ \pi^- n$ and $\pi^+ p \rightarrow \pi^+ \pi^0 p$, $\pi^+ \pi^+ n$ in the c.m. energy range 1300-2000 MeV. Details of this analysis, particularly the evolution of the preferred solution B, have been discussed in the previous edition of this review². Results from this solution are shown in Fig. III.1. In this edition we list resonance parameters, including pole positions, from LONGACRE 75 in the Data Card Listings. Breit-Wigner masses, widths, and resonance couplings were extracted in two ways by

LONGACRE 75, and we list the results of both methods, in addition to the pole positions. Method 1 used Breit-Wigner fitting to individual partial waves, and method 2 used a coupled-channel K-matrix fit to all waves with the same IJP. Among the most interesting results are further evidence for the existence of the $D_{13}(1700)$ and the $P_{33}(1690)$, further evidence for dominantly F-wave decay of the $F_{37}(1950)$ into $\Delta\pi$, and good agreement with quark-model predictions³.

Some information on isobar couplings is also available from earlier analysis. CHAVANON 74⁴ analyzed $\pi^+ p \rightarrow \pi^+ \pi^0 p$ in the range 1580-1970 MeV and obtained information on Δ -resonance decays to $\Delta\pi$ and $N\rho$. MEHTANI 72 analyzed the $\Delta\pi$ channel in the region 1820-2090 MeV and found the first evidence for strong F-wave decay of the $F_{37}(1950)$. The analysis of DIEM 70 covered the region 1550-1650 MeV.

In connection with all of these analyses it should be kept in mind that the isobar model is really a model and is not an exact representation of the $\pi N \rightarrow \pi\pi N$ amplitude. There are a number of ways in which the model could seriously fail to represent the physical amplitude. The problem that has received the most attention recently is the assumption that the isobar partial-wave amplitudes do not depend on the diparticle sub-energies, an assumption which is inconsistent with unitarity. For discussion of this problem see Refs. 2, 5, 6, and other references quoted therein.

References for Section III

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3. R. J. Cashmore et al., Nucl. Phys. B92, 37 (1975).
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For other references see the Data Card Listings.

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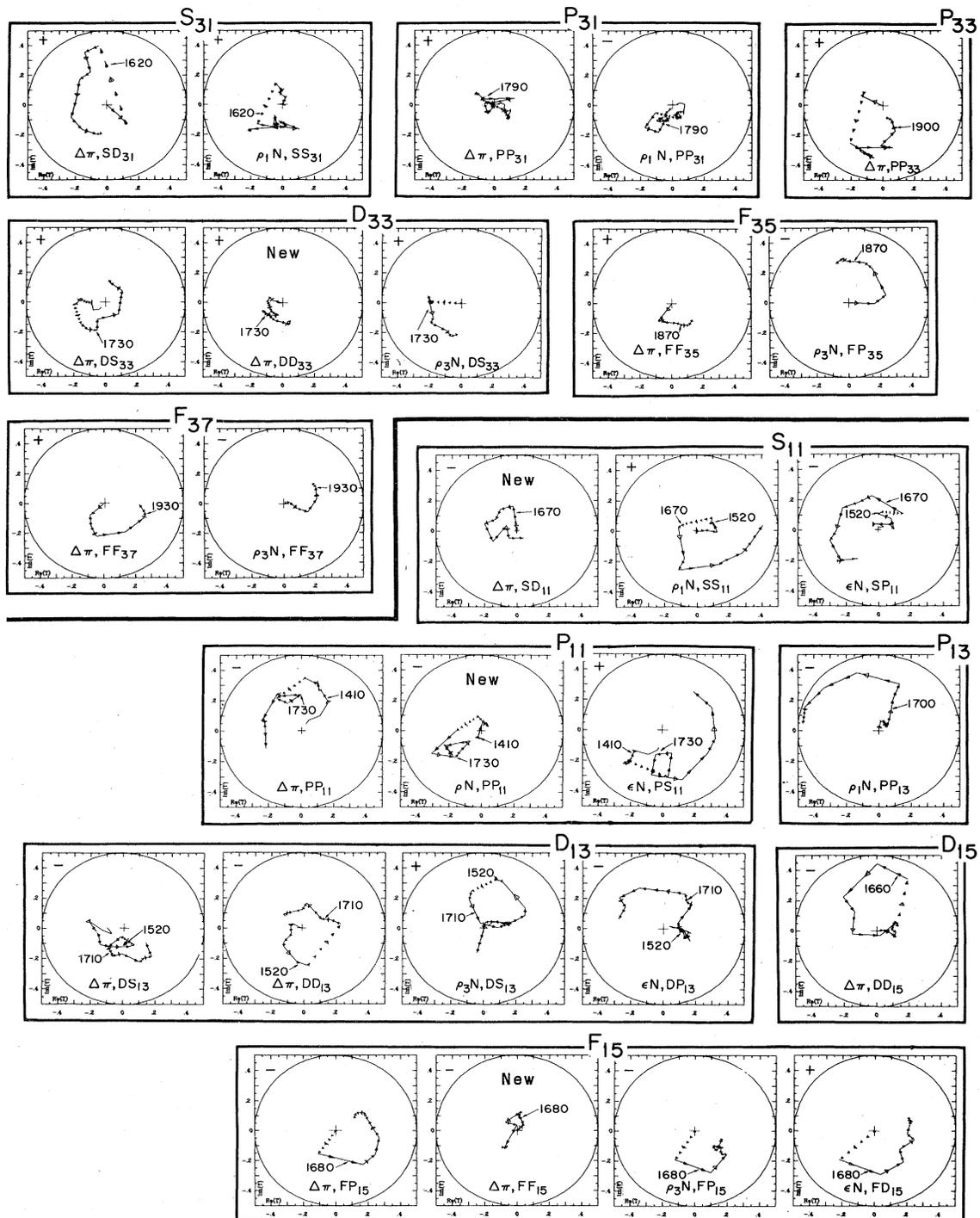


Fig. III.1. Solution B of the LBL-SLAC isobar analysis of $\pi N \rightarrow N\pi\pi$. See mini-review for notation. The four waves labeled "New" were not used in earlier solutions. The sign in the upper left-hand corner of each plot indicates how to transform from the sign conventions of this group to that of the "baryon-first" convention. (See main text, Sec. IV B.)

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IV. Photon Couplings

(R. L. Crawford and R. G. Moorhouse, Glasgow Univ.)

Photon couplings can be studied in formation reactions like

$$\gamma N \rightarrow N^* \rightarrow \pi N, \eta N, K\Lambda, \pi\Delta, \dots$$

A partial-wave analysis of these processes is the standard technique to determine the coupling strengths $g(N^* N \gamma)$. Up to now, almost all results are derived from the analysis of single-pion photoproduction. In the following we, therefore, define the conventions of pion photoproduction in which the results will be quoted.

The process $\gamma N \rightarrow N^* \rightarrow \pi N$ for a specific intermediate resonance can be symbolically described as

$$\langle \pi N | H_{\pi} | N^* \rangle \langle N^* | H_{\gamma} | \gamma N \rangle$$

The first term is measured in strong interactions, e.g., by partial-wave analysis of elastic πN scattering. A common feature of almost all analyses of pion photoproduction is a strong reliance on the knowledge of resonance parameters from πN partial-wave analyses. Pion photoproduction does not give a particularly accurate method of determining the resonance masses and widths, due partly to a lack of sufficiently precise data at many energies and partly because photoproduction is complicated by the fact that the photon has helicity states ± 1 and can react as an isoscalar and isovector particle. Consequently, several couplings for $N^* \rightarrow \gamma N$ have to be determined (2 for Δ and 4 for N).

Isospin Decomposition. Ignoring possible isotensor components, the reactions leading to the four possible final charge states are described by three isospin amplitudes. One set of these consists of the amplitudes A^{Δ} , A^P , and A^N which are respectively the amplitudes for the reaction to proceed by an $I = 3/2$ state, an $I = 1/2$ state with charge = +1, and an $I = 1/2$ state with charge = 0:

$$\begin{aligned} A(\gamma p \rightarrow \pi N) &= C_{\pi N}^{3/2} A^{\Delta} + C_{\pi N}^{1/2} A^P, \\ A(\gamma n \rightarrow \pi N) &= C_{\pi N}^{3/2} A^{\Delta} + C_{\pi N}^{1/2} A^N. \end{aligned} \quad (1)$$

The C-G coefficients, $C_{\pi N}^I$, for the coupling to a specific πN state are given explicitly by

$$\begin{aligned} A(\gamma p \rightarrow \pi^+ n) &= -\sqrt{1/3} A^{\Delta} - \sqrt{2/3} A^P, \\ A(\gamma p \rightarrow \pi^0 p) &= \sqrt{2/3} A^{\Delta} - \sqrt{1/3} A^P, \\ A(\gamma n \rightarrow \pi^- p) &= \sqrt{1/3} A^{\Delta} - \sqrt{2/3} A^N, \\ A(\gamma n \rightarrow \pi^0 n) &= \sqrt{2/3} A^{\Delta} + \sqrt{1/3} A^N. \end{aligned} \quad (2)$$

Walker's¹ amplitudes, A^{V3} , A^{V1} , and A^S for the production of an isospin eigenstate by the isovector and isoscalar parts of the electromagnetic current and the corresponding CGLN² amplitudes $A^{(3)}$, $A^{(1)}$, and $A^{(0)}$ are given by

$$\begin{aligned} A^{V3} &= A^{\Delta} = \sqrt{3} A^{(3)}, \\ A^{V1} &= 1/2(A^N - A^P) = \sqrt{1/3} A^{(1)}, \\ A^S &= 1/2(A^N + A^P) = \sqrt{2/3} A^{(0)}. \end{aligned} \quad (3)$$

Resonance Couplings. Since we are interested in intermediate resonances, we can approximate the energy dependence of the photoproduction partial waves by a Breit-Wigner form with added background. Using the helicity and parity eigenstates,³ $C_{\lambda}^{\ell\pm}(W)$, this gives

$$C_{\lambda}^{\ell\pm}(W) = \epsilon \left\{ \frac{\Gamma^{\lambda}(N^* \rightarrow \gamma N) \Gamma(N^* \rightarrow \pi N)}{kq} \right\}^{1/2} \frac{W}{W^2 - m_R^2 - iW\Gamma} + \text{background} \quad (4)$$

where ϵ is the sign of the amplitude and k, q are the c.m. momenta of the initial, final states. The resonance energy is m_R and Γ is its total width. In the following discussion of resonance couplings, we use the notation \tilde{A} for the imaginary part of the resonance coupling to an amplitude $A(W)$ evaluated at resonance ($W = m_R$). Thus

$$\tilde{C}_{\lambda}^{\ell\pm} = \epsilon \left\{ \frac{\Gamma^{\lambda} \Gamma_{\pi}}{kq |1|^2} \right\}^{1/2} \quad (5)$$

A dominant feature in pion photoproduction is the Born approximation which contains the nucleon pole in the s- and u-channels and the pion pole in the t-channel. It reproduces, for example, the experimentally observed forward peak in charged-pion photoproduction. In partial-wave analyses, the sign factor ϵ is well determined relative to the Born terms.

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Introducing helicity amplitudes A_{λ}^{jP} for the decay $N^*(j^P) \rightarrow (\gamma N)_{\lambda}$, where j^P labels the spin and parity of the N^* , we can calculate the radiative width⁴, $\Gamma_{\gamma}^{\lambda}$,

$$\Gamma_{\gamma}^{\lambda} = \frac{k^2}{\pi} \frac{m_N}{m_R} \frac{1}{2j+1} \left(\tilde{A}_{\lambda}^{jP} \right)^2, \quad (6)$$

where m_N is the nucleon mass. Introducing this expression into Eq. (5) we find

$$\tilde{C}_{\lambda}^{\ell\pm} = \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} \tilde{A}_{\lambda}^{jP}. \quad (7)$$

We quote the results of partial-wave analyses in terms of the amplitudes \tilde{A}_{λ}^{jP} in units of $\text{GeV}^{-1/2}$.

The total radiative width Γ_{γ} and the corresponding contribution $\sigma_T^{\ell\pm}$ of the partial waves $C_{\lambda}^{\ell\pm}$ to the total cross section are given by

$$\Gamma_{\gamma} = \sum_{\lambda=-3/2}^{3/2} \Gamma_{\gamma}^{\lambda} = \frac{k^2}{\pi} \frac{m_N}{m_R} \frac{2}{2j+1} \left\{ \left(\tilde{A}_{1/2}^{jP} \right)^2 + \left(\tilde{A}_{3/2}^{jP} \right)^2 \right\}, \quad (8)$$

$$\sigma_T^{jP} = 2 \left(C_{\pi N}^I \right)^2 \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \left\{ \left(\tilde{A}_{1/2}^{jP} \right)^2 + \left(\tilde{A}_{3/2}^{jP} \right)^2 \right\}. \quad (9)$$

The Hebb-Walker amplitudes¹, $\tilde{A}_{\ell\pm}^1$, $\tilde{B}_{\ell\pm}^1$, are related to the \tilde{A}_{λ}^{jP} by

$$\tilde{A}_{\ell\pm}^1 = \mp \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} C_{\pi N}^I \tilde{A}_{1/2}^{jP}, \quad (10)$$

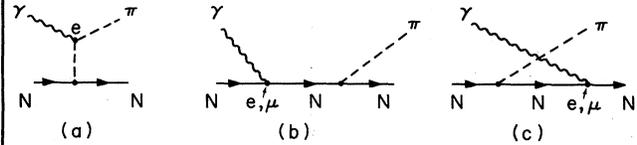
$$\tilde{B}_{\ell\pm}^1 = \pm \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} \times \left\{ \frac{16}{(2j-1)(2j+3)} \right\}^{1/2} C_{\pi N}^I \tilde{A}_{3/2}^{jP}. \quad (11)$$

We see from (2), (4), (5) and (7) that, apart from the factor $[16/(2j-1)(2j+3)]^{1/2}$, these are just the imaginary part of the resonance part of the amplitude (2) at resonance. Consequently, they are more directly related to experiment than the \tilde{A}_{λ}^{jP} , which incorporate an additional uncertainty in the partial width $\Gamma(N^* \rightarrow \pi N)$.

Methods of Partial-Wave Analysis. (a) Simple Isobar Model: In this method, the partial-wave amplitudes are formulated as in (4) (or similarly), with the resonance or resonances built partly

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Fig. IV.1. Feynman diagrams for the Born terms in pion photoproduction. The couplings of the photon to the pion and nucleon are indicated by an e to represent a coupling via charge and by a μ for a coupling via anomalous magnetic moment.

using knowledge of m_R , Γ , and $\Gamma(N^* \rightarrow \pi N)$ from πN elastic partial-wave analysis. The parametrization of the background is generally author-dependent, as is the detailed form of the resonance term, but the background is generally assumed to vary more slowly than the resonances. When the partial waves of Eq. (4) have been summed to form the angular-dependent helicity amplitudes, the Born amplitudes given by the Feynman diagrams of Fig. IV.1 must be added. [Alternatively, these Born amplitudes may previously have been projected into partial waves and included in the background terms of Eq. (4).]

As indicated in Fig. IV.1, the photon couples to the nucleon through both the electric charge, e , and the anomalous magnetic moment, μ . The minimum gauge invariant form, which includes the pion pole, is given by the pion pole plus the nucleon pole terms with the electric charge interaction only. Generally, only these terms, which we can refer to as the electric Born terms, are included explicitly because they reproduce the forward peak in charged-pion production. The magnetic terms can be assumed to be included in the real parts of the background and resonance terms in Eq. (4). However, this certainly raises problems of principle since some of the electric Born terms might also be subsumed in the background and resonances.

The best values of the resonance couplings, $\Gamma(N^* \rightarrow \pi N)$, and background parameters are found as those which give minimum χ^2 for all data within the energy range, in the usual manner for DPWA.

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(b) Fixed-t Dispersion Relations (FTDR): The difficulties of principle with the Born approximation in the simple isobar model are among the reasons which have led to the use of FTDR:

$$\text{Re}A_i^{(\pm,0)}(s,t) = B_i^{(\pm,0)}(s,t) + \frac{1}{\pi} P \int_{(m_N+m_\pi)^2}^{\infty} ds' \times \left\{ \frac{\text{Im}A_i^{(\pm,0)}(s',t)}{s'-s} + \xi_i \frac{\text{Im}A_i^{(\mp,0)}(s',t)}{s'-u} \right\} \quad (12)$$

for the invariant amplitudes $A_i^{(\pm,0)}(s,t)$ ($i=1, 2, 3, 4$), where $B_i^{(\pm,0)}(s,t)$ are the Born amplitudes⁵ of Fig. IV.1, $\xi_1 = \xi_2 = \xi_4 = 1$, $\xi_3 = -1$ and $\pm, 0$ denote $\gamma p \rightarrow \pi^+ n$, $\gamma n \rightarrow \pi^- p$, $\gamma p \rightarrow \pi^0 p$ respectively. $A_i^{(\pm,0)}(s,t)$ can be expressed in terms of the helicity amplitudes, and vice versa, by standard formulae⁵. $B_i^{(\pm,0)}(s,t)$ includes both the electric Born and anomalous magnetic moment terms, and thus no ambiguity arises as in the simple isobar model because all the real part (including the "resonant" real part) is manufactured by the unique prescription in Eq. (12).

The $\text{Im}A_i^{(\pm,0)}(s,t)$ are constructed using only the imaginary part of Eq. (4), or alternative formulae⁵. The $\text{Re}A_i^{(\pm,0)}(s,t)$ are obtained from Eq. (12) and minimization of χ^2 in fitting to experimental data to find the unknown parameters of Eq. (4) as is usual in DPWA.

A second advantage of FTDR over the simple isobar model is that the imaginary parts appear to be resonance-dominated^{1,5} (but not resonance-saturated), and this is not true of the real parts. The FTDR method only parametrizes the imaginary parts, and the smaller influence therein of background terms, relative to the size of the background in the real parts, is helpful. This advantage is diminished by the need, apparent from Eq. (12), to parametrize $\text{Im}A_i^{(\pm,0)}(s,t)$ outside the energy region of fitting for many analyses, thus effectively adding extra parameters to describe the resulting real background in the energy region being studied. This has been treated in two recent analyses^{6,7} by fitting data simultaneously in the resonance region and in the high energy (Regge) region.

Another, and in principle serious, disadvantage is that $\text{Im}A_i^{(\pm,0)}(s,t)$ in Eq. (12) is required for values of t outside the physical range corresponding to s' . Formally, $\text{Im}A_i^{(\pm,0)}(s',t)$ is given by the partial-wave expansion in $\cos\theta$ (θ is the production angle in the center-of-mass system), and this is the method used in practice to calculate the dispersion integrals. While the convergence of this expansion has not been proved for important parts of the unphysical region, Devenish, Lyth, and Rankin⁸ have surmised on the basis of the Mandelstam representation that the method is good up to $-t \leq 1$ (GeV/c)² in π^\pm production and to $-t \leq 1.5$ (GeV/c)² in π^0 production.

At low energies, the unitarity of the S-matrix imposes a phase condition on the partial-wave amplitudes known as Watson's theorem. It states that in the elastic region the complex phase of each $C_\lambda^{\ell\pm}(W)$ is equal to the scattering phase of the corresponding πN partial wave. This is important across the first resonance region in photoproduction since it imposes a strict phase condition on the two P_{33} partial waves and results in important non-resonant S-wave background in the imaginary parts contributing to the dispersion integrals. However, while the effects of Watson's theorem are significant in the first resonance region, the consequences in a FTDR calculation of ignoring it appear to be slight at higher energies.

The use of FTDR in pion photoproduction has a long history in the first resonance region⁹ where their role was envisaged as largely predictive or synthetic. The independently motivated extension by the Berkeley^{5,10,11,12}, Lancaster^{6,13}, and Glasgow^{7,14} based groups to a larger energy region has changed their role to one that is almost purely analytic.

(c) Energy-Independent Partial-Wave Analyses: Recent improvements in experimental data have allowed Berends and Donnachie¹⁵ to extend the energy-independent analysis of pion photoproduction to beyond the second resonance region. Previous analyses¹⁶ were restricted by the data then available to the first resonance region where Watson's theorem can be used to get a unique

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solution for the partial waves. In general, the energy-independent analyses are in good qualitative agreement with the energy-dependent analyses but there are some qualitative discrepancies. The significance of these is not yet clear.

Definitions of Resonance Parameters and

Errors: From analyses which use a form like that of Eq. (4) there are reported in the Data Card Listings and Table IV.1, photonic resonance couplings which are obtained through Eq. (4) from a Breit-Wigner partial width. The FTDR method fits only the imaginary part of Eq. (4) to experimental data. The FTDR method of the Berkeley group^{5,10,11,12} uses a K-matrix ansatz for the imaginary part and in these cases it is the corresponding K-matrix pole quantity that is reported. (Regretably, no authors give also partial widths corresponding to T-matrix residues.)

If no errors are assigned, the authors have given a unique result without quoting an error. The Berkeley^{5,10,11,12} and Glasgow⁷ analyses quote as an error the spread around a central value of a number of solutions. The Lancaster group^{6,13} estimate a "real error" for each parameter of a given solution that corresponds to the change in value required to increase "the best possible χ^2 " by 1%.

The variation between the central values of the various papers are generally about the same size as or greater than the quoted uncertainties in individual papers. This indicates that the systematic effects of different parametrizations and choices of experimental data are at least as important as statistical effects when estimating the uncertainty in the couplings. This is reflected in the errors given in Table IV.1.

Recent Partial-Wave Analyses: We describe only the most recent analyses by the various groups involved. A description of earlier analyses may be found in the previous edition of this Review.³

(a) Simple Isobar Model: The most important of these is still METCALF 74¹⁷ which uses the methods of Walker¹ to analyze data for $\gamma p \rightarrow \pi^+ n$,

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TABLE IV.1. Photon resonance couplings.

State	λ	Helicity Couplings (GeV) ^{-1/2} 10 ⁻³		Helicity Couplings (GeV) ^{-1/2} 10 ⁻³	
		Analyses Average ⁽¹⁾		Quarks ⁽²⁾	
		\tilde{A}_λ^p	\tilde{A}_λ^n	\tilde{A}_λ^p	\tilde{A}_λ^n
P ₁₁ ⁺ (1470)	1/2	-74±15	34±35	27	-18
D ₁₃ ⁺ (1520)	1/2	-10±15	-75±15	-34	-31
	3/2	171±15	-129±10	109	-109
S ₁₁ ⁺ (1535)	1/2	63±25	-49±35	156	-108
D ₁₅ ⁺ (1670)	1/2	16±10	-32±36	0	-38
	3/2	21±12	-62±40	0	-53
F ₁₅ ⁺ (1688)	1/2	-5±30	25±10	-10	30
	3/2	127±35	-16±20	60	0
S ₁₁ ⁰ (1700)	1/2	43±30	-37±40	0	30
D ₁₃ ⁰ (1700)	1/2	-20±45	18±55	0	-10
	3/2	19±45	18±80	0	-40
P ₁₁ ⁰ (1780)	1/2	18±40	18±50	-40	10
P ₁₃ ⁰ (1810)	1/2	-25±50	24±70	100	-30
	3/2	-31±50	-4±60	-30	0
		$\tilde{A}_\lambda^\Delta = \tilde{A}_\lambda^{\sqrt{3}}$	$\tilde{A}_\lambda^\Delta = \tilde{A}_\lambda^{\sqrt{3}}$		
P ₃₃ ⁺ (1232)	1/2	-139±5		-108	
	3/2	-256±5		-187	
S ₃₁ ⁺ (1650)	1/2	46±36		47	
D ₃₃ ⁺ (1670)	1/2	72±26		88	
	3/2	72±45		84	
P ₃₃ ⁰ (1690)	1/2	-2±40		23	
	3/2	-12±50		39	
F ₃₅ ⁺ (1890)	1/2	21±30		-20	
	3/2	-10±60		-90	
P ₃₁ ⁺ (1910)	1/2	-11±20		-30	
F ₃₇ ⁺ (1950)	1/2	-69±16		-50	
	3/2	-76±20		-70	

(1) Average of the couplings from MOORHOUSE 73, DEVENISH 73, MOORHOUSE 74, METCALF 74, KNIES 74, DEVENISH2 74, CRAWFORD 75, BARBOUR 76, and Ref. 12. The errors given are an estimate that takes into account the errors quoted in these analyses and the differences between the analyses.

(2) The naive, l -excitation, quark model is used in the 4-dimensional oscillator form of Feynman, Kislinger, and Ravndal. The non-relativistic quark model with recoil gives generally very similar results.

$\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ from the first through the fourth resonance regions. The partial waves are parametrized as in Eq. (4) with the background taken to be an independent number at each energy fitted. The electric Born terms are added explicitly.

Data Card Listings

For notation, see key at front of Listings.

Baryons N's and Δ's

Other recent isobar analyses have been on a much smaller scale. KRIVETS 74¹⁸ is an analysis of $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ across the first, second, and third resonance regions using a rather small data set consisting only of differential cross-section measurements. Several solutions are given with a large spread in the values of the couplings and those for the resonance couplings to γn are stated to be preliminary. HEMM1 73¹⁹ is a fit to forward $\gamma p \rightarrow \pi^0 p$ differential cross-section data across the second and third resonance regions and evaluates the $\lambda = 1/2$ couplings in this region. HEMM2 73²⁰ is an isobar analysis of $\gamma n \rightarrow \pi^0 n$ across the second resonance region. ROSSI 73²¹ and BENEVENTANO 73²² are analyses of data for $\gamma n \rightarrow \pi^- p$ across the second resonance region. Earlier analyses^{1,23} are not included in the Data Card Listings.

(b) Fixed-t Dispersion Relations:

MOORHOUSE 73¹⁰ and MOORHOUSE 74⁵ are FTDR analyses from threshold through the third resonance region ($1160 < W < 1780$ MeV) using data for $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$. The first uses the results of seven fits with a χ^2 per data point ranging from 9.7 with 52 variables to 5.7 with 74. The second uses the average of three solutions with a χ^2 per data point ranging from 4.0 (56 variables) to 3.0 (74 variables). KNIES 74¹¹ and Ref. 12 extend the energy range of the Berkeley analyses up to a center-of-mass total energy of 2 GeV, and thus include the energy region of the $\Delta(1950)F_{37}$. Several solutions are described with typical values of the χ^2 per data point of between 3.5 and 4.5. In all these analyses, the imaginary parts of the partial waves are parametrized using a K-matrix formalism.

DEVENISH 73¹³ analyze data from the first through the third resonance regions and parametrize the imaginary parts of the partial waves as Breit-Wigner forms without background except for the π^0 -production S wave. The dispersion integral is cut off at $W = 1.9$ GeV, but parametrized real background is added to allow for the integrals above 2 GeV. Three solutions are given with χ^2 per data point of 10.5, 3.0, and 4.7 for

37, 68, and 68 free parameters, respectively. DEVENISH2 74⁶ treats the problem of the high energy contributions to the dispersion integral by simultaneously fitting data in the resonance region and at high energies and evaluates the couplings for all resonances for $W \leq 2$ GeV. It uses data for $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ at all available energies and for $|t| < 1$ (GeV/c)². A χ^2 per data point of 4.8 is obtained with 48 resonance parameters and 50 high energy parameters.

CRAWFORD 75¹⁴ is an FTDR analysis from the first through the fourth resonance regions of $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$. Data are fitted for $W \leq 2$ GeV and $|t| \leq 1.5$ (GeV/c)². Three solutions are given with a best χ^2 per data point of 3.1 with 58 free parameters. The masses and widths of all resonances below 1.7 GeV are evaluated. Stable values are obtained that agree well with the values from elastic πN partial-wave analyses. BARBOUR 76⁷ is an extension of the previous analysis to include the fitting of high energy data for $W > 2.6$ GeV. One fit is presented with a χ^2 per data point of 2.6 for the resonance region data and 1.8 for the high energy data. No data are fitted for $2 \text{ GeV} < W < 2.6 \text{ GeV}$, thus omitting the region of the $H_{3,11}$.

(c) Energy-Independent Analyses: These have now been extended to the second resonance region by Berends and Donnachie¹⁵ as described above. Earlier analyses¹⁶ fitted only in the first resonance region. Resonance couplings have not been quoted in these analyses.

Electroproduction in the Resonance Region:

Although single-pion electroproduction can be considered in many ways to be similar to photoproduction²⁴, a partial-wave analysis of the process is considerably more difficult than in photoproduction. This is due to the less accurate and complete experimental data that are available and because the scalar part of the virtual photon in electroproduction contributes an extra (scalar) coupling for each resonance. The scalar amplitudes uncouple from the T-matrix in the photoproduction limit. Also, it is necessary, in principle, to make an analysis at each value of λ^2 , the (virtual photon mass)², at which there are

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

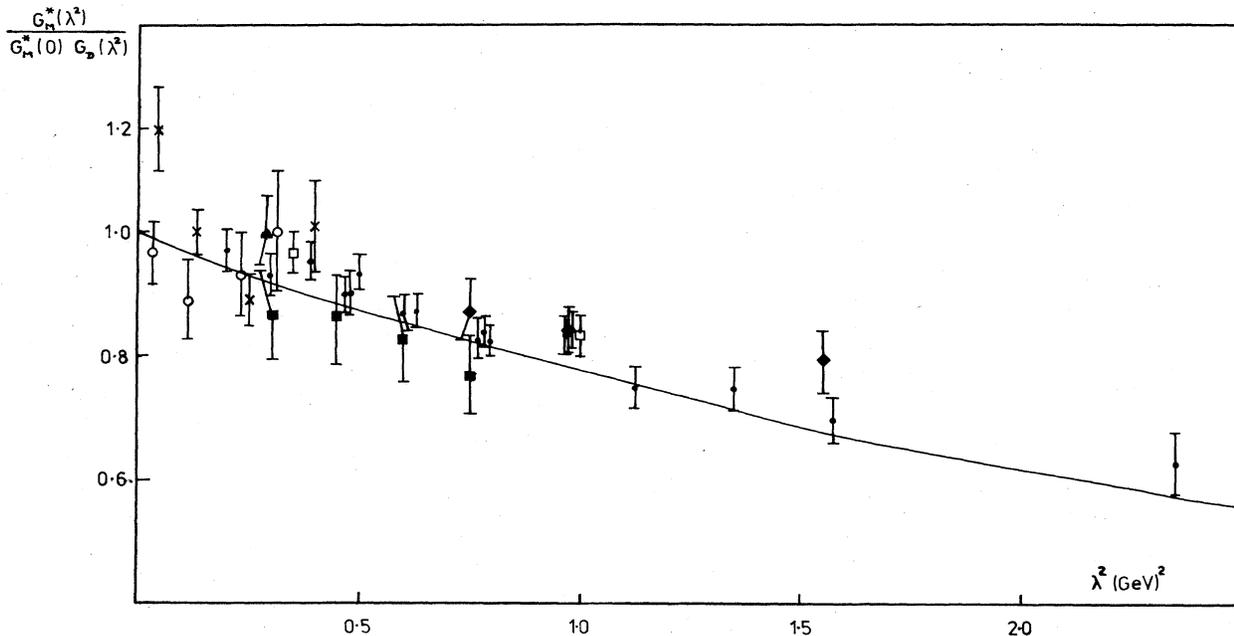


Fig. IV.2. The $\Delta(1232)$ magnetic moment form factor, $G_M^*(\lambda^2)$, normalized to $G_M^*(0) = 3$, compared to the nucleon dipole form factor, $G_D(\lambda^2) = (1 + \lambda^2/0.71)^{-2}$. The sources of the data points are \bullet ref. 25, \circ ref. 27, \times ref. 28, \blacksquare ref. 29, \square ref. 30, \blacklozenge ref. 31, \blacktriangle ref. 32, and \diamond ref. 33. The solid line is from Devenish and Lyth³⁴.

data. The information to be obtained from such an analysis is the size of the scalar couplings to the resonance and the form factors of the scalar transverse couplings, i.e., their λ^2 dependence. This information has been extracted in a number of ways, depending on the experimental data available and the energy range being studied.

(a) Total Cross Sections: In a single-arm experiment, in which only the final electron is detected, it is possible to measure only σ_T and σ_S , the total cross sections for the absorption of transverse and scalar photons. In the first resonance region, due to the dominance of the $M_{1+}^{(3)}$ multipole, σ_T can give the form factor for the magnetic coupling to the $P_{33}(1232)$. There have been many analyses of this type, both the form factors for γp and γn coupling to the $P_{33}(1232)$ having been measured. We quote as examples of these, the analyses of Bartel et al.²⁵ and of Köbberling et al.²⁶ Due to the unknown non-resonant scalar background and because the scalar coupling is relatively small, σ_S gives

only qualitative information about the scalar coupling to this resonance. Above the second resonance region, due to the large number of resonances, the total cross sections again give little more than qualitative information.

(b) Energy-Independent Analyses: Differential cross-section measurements have been made for π^0 and π^+ electroproduction in a number of coincidence experiments. Although the data are not good enough for a full partial-wave analysis, it is possible, due to the small number of partial waves involved, to use the dominance of the $M_{1+}^{(3)}$ multipole to measure $|M_{1+}|$, $\text{Re } E_{1+}^* M_{1+} / |M_{1+}|^2$, and $\text{Re } S_{1+}^* M_{1+} / |M_{1+}|^2$ for π^0 production across the first resonance region. This has been done by a number of groups²⁷⁻³². Because $M_{1+}^{(3)}$ is so large and is purely imaginary at the P_{33} mass, these terms can be taken to be respectively $\text{Im} M_{1+}^{(3)}$, $E_{1+}^{(3)} / M_{1+}^{(3)}$, and $S_{1+}^{(3)} / M_{1+}^{(3)}$ at the resonance.

(c) Energy-Dependent Analyses: Energy-dependent partial-wave analyses using FTDR have been made for pion electroproduction by Crawford³³

Data Card Listings

For notation, see key at front of Listings.

Baryons N's and Δ's

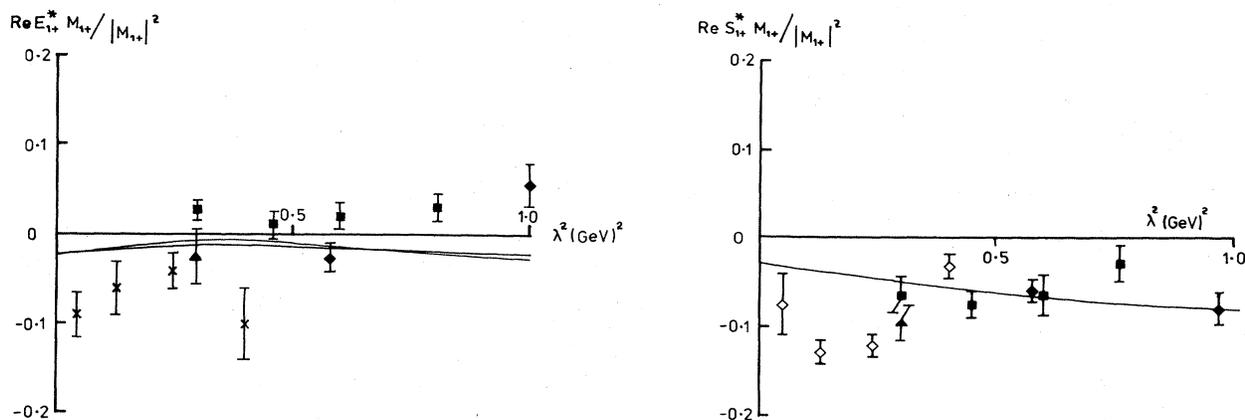


Fig. IV.3. The ratios $\text{Re}(E_{1+}^* M_{1+} / |M_{1+}|^2)$ and $\text{Re}(S_{1+}^* M_{1+} / |M_{1+}|^2)$ for π^0 electroproduction at the $\Delta(1232)$. The data points and curves are indicated as in Fig. IV.2. Crawford³³ and Devenish and Lyth³⁴ give the ratios $E_{1+}^{(3)} / M_{1+}^{(3)}$ and $S_{1+}^{(3)} / M_{1+}^{(3)}$.

and by Devenish and Lyth³⁴. The first fitted only in the first resonance region and over a relatively small range of λ^2 . The second is an extensive analysis from the first to the fourth resonance regions. Using the results of photo-production analyses to give a starting point, the form factors for the resonance couplings are given algebraic forms that reflect their analytic structure in the complex λ^2 plane. Data for π^0 and π^+ production and for a range of values of λ^2 are fitted simultaneously and the results of a number of fits are presented. The analysis evaluates the couplings for the $P_{33}(1232)$, $P_{11}(1470)$, $D_{13}(1520)$, $S_{11}(1535)$, $D_{33}(1670)$, $F_{15}(1688)$, $S_{11}(1700)$, $D_{13}(1700)$, $P_{13}(1780)$, $F_{37}(1950)$, and $D_{13}(2040)$ resonances. The results are discussed below.

The couplings and form factors for the $P_{33}(1232)$ are now established for $\lambda^2 < 1$ (GeV)² and there is agreement on all general features between the basic methods of analysis (Figs. IV.2 and IV.3). The λ^2 dependence of $M_{1+}^{(3)}$ is usually expressed in terms of the form factor $G_M^*(\lambda^2)$,²⁷

$$G_M^*(\lambda^2) = 2M \left(\frac{3}{2\alpha} \frac{|\vec{k}|}{\sin^2 \delta_{33}} \frac{\Gamma |M_{1+}^{(3)}|^2}{q^2} \right)^{1/2},$$

where \vec{k} and \vec{q} are respectively the center-of-mass three-momenta of the pion and the virtual photon.

The behavior of G_M^* for the charge = +1 state, normalized to $G_M^*(0) = 3$, is compared with the nucleon dipole form, $G_D(\lambda^2) = (1 + \lambda^2/0.71)^{-2}$, in Fig. IV.2. It is clearly seen that G_M^* falls off significantly faster with λ^2 than the nucleon dipole. Measurements of G_M^* for the charge = 0 state^{26,35} show that it has the same form as for the charge = +1 state. The measurements of the ratios of $E_{1+}^{(3)}$ and $S_{1+}^{(3)}$ to $M_{1+}^{(3)}$ are all consistent with these being small, although the $S_{1+}^{(3)}$ is clearly established as being non-zero with a value of about $-0.06 M_{1+}^{(3)}$. This is a small violation of the quark model selection rule that predicts that there should be no scalar excitation of the $P_{33}(1232)$ ³⁶.

In the second to the fourth resonance regions, Devenish and Lyth obtain definite results for the couplings to the charge = +1 states of the $P_{11}(1470)$, $D_{13}(1520)$, $S_{11}(1535)$, $F_{15}(1688)$, and $F_{37}(1950)$ resonances. These are expressed in terms of the multipole partial waves:

$P_{11}(1470)$ The form factor for the magnetic multipole rapidly vanishes and there is no evidence for a scalar coupling.

$D_{13}(1520)$ The electric multipole decreases rapidly with λ^2 in an approximate dipole form, but the electric dipole decreases significantly less quickly. There is therefore a rapid change in

Baryons

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the ratio of the transverse helicity partial waves. The scalar coupling is non-zero at $\lambda^2 = 0$ but decreases rapidly.

S₁₁ (1535) The electric multipole decreases slowly with λ^2 . The scalar coupling is comparable to the electric one at $\lambda^2 = 0$ but may decrease in an approximately dipole form.

F₁₅ (1688) The form factors for the transverse multipoles are similar to those for D₁₃ (1520). The scalar coupling decreases rapidly.

F₃₇ (1950) The magnetic multipole probably decreases more slowly than the dipole form. The electric multipole is small and its λ^2 behavior is not clearly determined. The scalar multipole is small and decreases rapidly.

The form factors for the S₁₁ (1535) can be checked in η electroproduction. This process proceeds mainly through S-wave excitation in the second resonance region and is thus dominated by the S₁₁ (1535). The results for the cross section of this process³⁷ again require that at least one of the form factors should fall off slowly with λ^2 .

The rapid change in the ratio of the transverse helicity amplitudes of the D₁₃ (1520) and the F₁₅ (1688), along with the almost constant value of the ratio for the P₃₃ (1232), gives qualitative agreement with quark model calculations^{4,38}. However, some quark models find that the form factor for the S₁₁ (1535) should decrease rapidly and that the one for the P₁₁ (1470) should stay up, and are therefore in conflict with the results above.

Information in this Edition: The Baryon Table contains the branching fractions, Γ_Y/Γ for 14 resonances.

The Data Card Listings contain the photon resonance couplings (\tilde{A}_λ^j)^P for p, n, and Δ from the analyses of Moorhouse and Oberlack¹⁰ (MOORHOUSE 73), Moorhouse, Oberlack, and Rosenfeld⁵ (MOORHOUSE 74), Knies, Moorhouse, and Oberlack¹¹ (KNIES 74), Devenish, Lyth, and Rankin^{13,6} (DEVENISH 73 and DEVENISH2 74),

Crawford¹⁴ (CRAWFORD 75), Barbour and Crawford⁷ (BARBOUR 76), Metcalf and Walker¹⁷ (METCALF 74), Hemmi et al.¹⁹ (HEMMI1 73), Hemmi et al.²⁰ (HEMMI2 73), Rossi et al.²¹ (ROSSI 73), and Beneventano et al.²² (BENEVENTANO 73).

The average of the listed results along with an estimate of the uncertainty in each from the spread of results is given in Table IV.1. We give results from a naive ($\bar{\Lambda}$ -excitation) quark model for comparison.

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For notation, see key at front of Listings.

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STATUS OF N* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --										OTHER CHANN.			
			TOTAL* CR.S.	PI	N	ETA	N	K	LAM	K	SIG	PI		DE		
N*(940)	P11	****														
N*(1470)	P11	****	****													EPS N
N*(1520)	D13	****	****	*												RHC N
N*(1535)	S11	****	****	***												RHO N
N*(1670)	D15	****	****	***												RHC N
N*(1688)	F15	****	****	*												RHC N
N*(1700)	S11	****	****	*												EPS N
N*(1700)	D13	**	**	*						*						EPS N
N*(1780)	P11	**	**	*						*						EPS N
N*(1810)	P13	**	**	*						*						RHO N
N*(1990)	F17	**	**	*						*						
N*(2000)	F15	*	*	*						*						
N*(2040)	D13	**	**	*						*						
N*(2100)	S11	*	*	*						*						
N*(2100)	D15	*	*	*						*						
N*(2190)	G17	***	***	***						*						
N*(2220)	H19	***	***	***						*						
N*(2650)		***	***	***						*						
N*(3030)		***	***	*						*						
N*(3245)		*	*	*						*						
N*(3690)		*	*	*						*						
N*(3755)		*	*	*						*						
DE(1232)	P33	****	****	****						F						****
DE(1650)	S31	****	****	****						O						RHO N
DE(1670)	D33	***	***	***						R						****
DE(1690)	P33	*	*	*						B						****
DE(1890)	F35	***	***	***						I						RHC N
DE(1900)	S31	*	*	*												****
DE(1910)	F31	***	***	***						D						RHC N
DE(1950)	F37	****	****	****						D						RHO N
DE(1560)	D35	**	**	**						E						****
DE(2160)	**	***	***	***						N						****
DE(2420)	F311	***	***	***						F						****
DE(2850)	***	***	***	***												****
DE(3230)	***	***	*	*						C						****

*** GOOD, CLEAR, AND UNMISTAKABLE.
** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
* NEEDS CONFIRMATION.
* WEAK.
* ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

S=0 I=1/2 NUCLEON STATES (N)

p

16 PROTON(938, J=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

n

17 NEUTRON(939, J=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

Baryons
N(1470)

Data Card Listings
For notation, see key at front of Listings.

N(1470)

61 N*1/2(1470, JP=1/2+) I=1/2
MASS AND WIDTH ARE BEST DETERMINED FROM PARTIAL WAVE ANALYSES. WE LIST PRODUCTION EXPERIMENTS SEPARATELY -- SEE BELOW.
AYED 74 CLAIM TWO P11 STATES IN THE 1500 MEV REGION. WE TENTATIVELY LIST BOTH HERE.

P11

Table with 4 columns: Mass (MEV), Author, Year, and Reference. Includes entries for Brandesen, Roper, Baryere, Donnach, Aved, Almed, and Crawford.

Table with 4 columns: Width (MEV), Author, Year, and Reference. Includes entries for Baryere, Donnach, Davies, Almed, Aved, Crawford, and Barbour.

Table with 4 columns: Real part of pole position (MEV), Author, Year, and Reference. Includes entries for Aved and Longacre.

Table with 4 columns: 2*Imag part of pole position (MEV), Author, Year, and Reference. Includes entries for Lee and Longacre.

Table with 4 columns: Absolute value of pole residue, Author, Year, and Reference. Includes entries for Lee.

Table with 4 columns: Phase of pole residue (radians), Author, Year, and Reference. Includes entry for Lee.

Table with 4 columns: Partial decay modes, Author, Year, and Reference. Lists various decay channels like N(1470) into pi N, eta N, etc.

Table with 4 columns: Branching ratios, Author, Year, and Reference. Lists ratios for various decay modes.

Table with 4 columns: Additional data/notes, Author, Year, and Reference. Includes entries for Aved, Davies, Makarov, Mickens, and Barbour.

Table with 4 columns: Decay mode, Author, Year, and Reference. Includes entries for R4, R5, R6, R7, R8, R9, R10, R11.

Table with 4 columns: Photon decay amplitudes, Author, Year, and Reference. Includes entries for A1, A2, A3, A4, A5, A6, A7, A8, A9, A10, A11, A12.

Table with 4 columns: References for N(1470), Author, Year, and Reference. Lists various scientific publications.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1470), N(1520)

Table listing authors and their affiliations for the 1470 MeV region. Includes names like BAREYRE, DALITZ, JOHNSON, etc.

THE FOLLOWING ARE THEORETICAL PAPERS CONCERNING THE N(1470) -- RESONANT STATE, IF ANY, THE HBC EXPERIMENTS SEE ENHANCEMENTS MAINLY IN THE P PI MASS PLOT. FOR ZERO CHARGE EXCHANGE, SUCH FINAL STATES ARE KNOWN TO HAVE LARGE DECK-TYPE BACKGROUND. THIS FACT COMPLICATES THE INTERPRETATION OF THIS BUMP AS A RESONANCE.

1470 MEV REGION - PRODUCTION EXPERIMENTS

91 N(1470, JP=) I=1/2 PRODUCTION EXPERIMENTS

THE BUMP SEEN IN PRODUCTION EXPERIMENTS AT LOW INVARIANT MASS MOST LIKELY CORRESPONDS TO THE P11 (SEE ABOVE) ENHANCEMENTS MAINLY IN THE P PI MASS PLOT. FOR ZERO CHARGE EXCHANGE, SUCH FINAL STATES ARE KNOWN TO HAVE LARGE DECK-TYPE BACKGROUND. THIS FACT COMPLICATES THE INTERPRETATION OF THIS BUMP AS A RESONANCE.

91 N(1470) MASS (MEV) (PROD. EXP.)

Main data table for the 1470 MeV region, listing mass (MEV), production method, and various parameters. Includes entries for APPROX, BELLETTIN, ANDERSON, etc.

91 N(1470) WIDTH (MEV) (PROD. EXP.)

Table listing width (MEV) for the 1470 MeV region, including authors like BELL, SHAPIRA, etc.

91 N(1470) PARTIAL DECAY MODES (PROD. EXP.)

Table listing partial decay modes for the 1470 MeV region, including decay channels like N(1470) INTO PI N, etc.

Table listing authors and their affiliations for the 1520 MeV region. Includes names like COCCONI, ADELMAN, ANKENBRANDT, etc.

REFERENCES FOR N(1470) (PROD. EXP.)

Table listing references for the 1470 MeV region, including authors like LILLETUN, ADELMAN, ANKENBRANDT, etc.

PAPERS NOT REFERRED TO IN CATA CARDS

Table listing references for the 1520 MeV region, including authors like SMITH, WOJCIK, COLTON, etc.

N(1520)

62 N(1520, JP=3/2-) I=1/2

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

62 N(1520) MASS (MEV)

Table listing mass (MEV) for the 1520 MeV region, including authors like BRANDSEN, ROPER, etc.

62 N(1520) WIDTH (MEV)

Table listing width (MEV) for the 1520 MeV region, including authors like BAREYRE, ODNAGH, etc.

62 N(1520) REAL PART OF POLE POSITION (MEV)

Table listing real part of pole position (MEV) for the 1520 MeV region, including authors like LANGGACRE, etc.

62 N(1520) 2*IMAG PART OF POLE POSITION (MEV)

Table listing 2*imag part of pole position (MEV) for the 1520 MeV region, including authors like LANGGACRE, etc.

Baryons
N(1520)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle ID, decay mode, and decay masses. Includes entries for N(1520) partial decay modes.

Table with columns for particle ID, branching ratios, and other parameters. Includes entries for N(1520) branching ratios.

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

ALMOST THE ENTIRE INELASTICITY IS IN N PI PI (CNLY N ETA COULD COMPLETE, AND IT DOESNT). THE N PI PI SEEMS TO BE MAINLY N*3/2(1232) PI, IN BTH R1 S AND D WAVES.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N*3/2(1232) PI/TOTAL.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N*3/2(1232) PI/(N PI PI).

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N EPSILON/TOTAL.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N ETA/TOTAL.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N RHO 1/TOTAL.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) from PI N TO K LAMBDA.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) from PI N TO N*3/2(1232) PI, S-WAVE.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) from PI N TO N*3/2(1232) PI, D-WAVE.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) from PI N TO N RHO, S=3/2, S-WAVE.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into GAM P, HELICITY=1/2 (GEV**1/2).

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into GAM P, HELICITY=3/2 (GEV**1/2).

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into GAM N, HELICITY=1/2 (GEV**1/2).

AVERAGE MEANINGLESS (SCALE FACTOR = 3.9)

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into GAM N, HELICITY=3/2 (GEV**1/2).

AVERAGE MEANINGLESS (SCALE FACTOR = 3.1)

REFERENCES FOR N(1520)

SEE A PREVIOUS EDITION (RMP 37, 633, 1965) FOR EARLIER REFERENCES.

BRANDSEN 65 PR 139 B1566 +GDCNELL, MOORHOUSE (DURHAM, RHEL) IJP
ROBERTS 67 PREPRINT R G ROBERTS (DURHAM)
THURNAUER 65 PRL 14 985 P G THURNAUER (ROCH)

KIRZ 66 PRIVATE COMM J KIRZ (LRL)
NUMBER EXTRACTED FROM CATA DISCUSSED IN KIRZ 63, Z 63.
NAMYSLOW 66 PR 157 1328 NAMYSLOSKI, RAZMI, ROBERTS (STAN, EDIN, LOIC)
OLSSON 66 PR 145 1309 M G OLSSON, G B YODH (MISC, UMD)

DAVIES 67 NC 52A 1112 A T DAVIES, R G MOORHOUSE (GLASGOW, RHEL)
ROBERTS 67 PREPRINT R G ROBERTS (DURHAM)
ROSENFELD 67 IRVING CONF A H ROSENFELD, P SOOING (LRL)

BAREYRE 68 PR 165 1731 P BAREYRE, C BRICMAN, G VILLET (SACLAY) IJP
DONNACHI 68 PL 268 161 B CARRERAS, A DONNACHIE (CERN) IJP
ALSO 68 VIENNA 139 DONNACHIE, RAPPOREUR-S TALK (GLAS)
MORGAN 68 PR 166 1731 R G KIRSOPP (EDIN)
(RHEL)

BOTKE 69 PR 180 1417 J C BOTKE (UCSB)
DEANS 69 PR 185 1797 S DEANS, J WOOTEN (UNIV S FLORIDA)

AYED 70 KIEV CONF. R AYED, P BAREYRE, G VILLET (SACL) IJP
CARRERAS 70 NP 168 35 B CARRERAS, A DONNACHIE (DARE, MCHS)
DAVIES 70 NP 821 359 A DAVIES (GLAS)
DIEM 70 KIEV CONF. + SMADJA, CHAVANON, DELER, DOLBEAU+ (SACL)
ALMEHED 72 NP 840 157 +LOVELACE (LUND, RUTG) IJP

DEVENISH 73 PL 478 53 DEVENISH, RANKIN, LYTH (LOUC+BOHN+LANC) IJP
HEMMLI 73 PL 438 79 +EMMI, INAGAKI+ (KYOTO) SAGA+KEK+TKYU) IJP
HEMMLI 73 NP 855 333 +INAGAKI, KIKUCHI, MAKI, MIYAKE+ (KYOTO, TOKYO) IJP
LEMOIGNE 73 PURDUE CONF. 93 +GRANET, MARTY, AYED, BAREYRE, BORGEAUD, + (SACL) IJP
MOORHOUS 73 PL 438 44 MOORHOUSE, OBERLACK (GLAS+BL) IJP
ROSSI 73 NC 134 59 +PIAZZA, SUSINNO, + (ROMA, FRAS, NAPL) IJP
ALSO 71 LNC 2 1183 CARBONARA, FIORE, + (NAPL, FRAS, PAVIA, ROMA) IJP

AYED 74 PRIVATE COMMCTN. AYED, BAREYRE (SACL) IJP
ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL) IJP
BENEVENT 74 NC 19A 529 BENEVENTANO, DANGELO, NOTARISTEFANI, + (RCMA) IJP
DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DES, NORDITA, LOUC)
DEVENISH 74 PL 528 227 DEVENISH, LYTH, RANKIN (DES, YLANC, BCNN) IJP
KNIES 74 PRD 9 2680 KNIES, MOORHOUSE, OBERLACK (LBL, GLAS) IJP
METCALF 74 NP 876 253 W J METCALF, R L WALKER (CTI) IJP
MOORHOUS 74 PRD 9 1 MOORHOUSE, OBERLACK, RCSENFELD (GLAS+LBL) IJP

CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS) IJP
FELTESSE 75 NP 893 242 +AYED, BAREYRE, BORGEAUD, DAVID, ERNEWIN+ (SACL) IJP
LONGACRE 75 PL 558 415 +ROSENFELD, LASINSKI, SMADJA, + (LBL, SLAC) IJP

BARBOUR 76 SBMTD. TO NP I. M. BARBOUR, R. L. CRAWFORD (GLAS) IJP
PAPERS NOT REFERRED TO IN CATA CARDS.

KIRZ 63 PR 130 2481 J KIRZ, J SCHWARTZ, R D TRIPP (LRL)
BAREYRE 65 PL 18 342 + BRICMAN, STIRLING, VILLET (SACLAY) IJP
CROUCH 65 DESY CONF II 21 + (BRONN, CE, HARVARD, MIT, PADOVA, WEIZMANN)
DERADD 65 ATHENS CONF 244 +KENNEY, LAMSA, + (NOTRE DAME, KENTUCKY)
MERLO 66 P ROY SOC 289 489 J P MERLO, G VALLACAS (SACLAY)
JOHNSON 67 UCL-17683 THESIS C H JOHNSON
DEANS 69 PRL 177 2623 S R DEANS (UNIV S FLORIDA)
DCNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
WALKER 69 PR 182 1729 R L WALKER (CAL TECH) IJP
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

Data Card Listings
For notation, see key at front of Listings.

Baryons
N(1535)

N(1535)

63 N*1/2(1535, JP=1/2-) I=1/2
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED. IT IS STRONGLY ASSOCIATED WITH THE ETA N CHANNEL.

S11

Table with columns for mass (MEV) and various parameters. Includes entries for HENDRY, MICHAEL, UCHIYAMA, BAREYRE, DONNACHI, DELCOURT, AYEY, DAVIES, ALMEHED, HICKS, CRAWFORD, LONGACRE, BARBOUR.

Table with columns for width (MEV) and various parameters. Includes entries for HENDRY, MICHAEL, UCHIYAMA, BAREYRE, DONNACHI, DELCOURT, AYEY, DAVIES, ALMEHED, HICKS, CRAWFORD, LONGACRE, BARBOUR.

Table with columns for real part of pole position (MEV) and various parameters. Includes entry for LONGACRE.

Table with columns for 2*imag part of pole position (MEV) and various parameters. Includes entry for LONGACRE.

Table with columns for partial decay modes and various parameters. Includes entries for INTO PI N, INTO N ETA, INTO N PI PI, INTO N EPSILON, INTO N RHO, INTO GAM P, INTO GAM N, INTO K LAMBDA.

Table with columns for branching ratios and various parameters. Includes entries for INTO (PI N)/TOTAL, DOMINANT INEL DECAY, INTO N ETA, INTO N RHO.

Table with columns for branching ratios and various parameters. Includes entries for INTO (N RHO)/TOTAL, PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING.

Table with columns for branching ratios and various parameters. Includes entry for INTO (N RHO)/TOTAL.

Table with columns for branching ratios and various parameters. Includes entry for INTO (N RHO)/TOTAL.

Table with columns for branching ratios and various parameters. Includes entry for INTO (N RHO)/TOTAL.

Table with columns for various parameters and values. Includes entries for INTO GAMMA PROTON TO ETA PROTON, FROM PI N INTO ETA N, FROM PI N TO K LAMBDA, FROM PI N TO N RHO, FROM PI N TO N EPSILON.

63 N*1/2(1535) PHOTON DECAY AMPLITUDE (GEV**1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for helicity and various parameters. Includes entries for INTO GAM P, HELICITY=1/2, INTO GAM N, HELICITY=1/2.

Table with columns for helicity and various parameters. Includes entries for INTO GAM N, HELICITY=1/2, CONVERTED TO OUR CONVENTIONS.

REFERENCES FOR N*1/2(1535)
HENRY 65 PL 18 171
MICHAEL 66 PL 21 93
UCHIYAMA 66 PR 149 1220
DAVIES 67 NC 52A 1112
BAREYRE 68 PR 165 1731
DONNACHI 68 PL 268 161
ALSO 68 VIENNA 139
ALSO 68 THESIS
DEANS 69 PR 185 1797
DELCOURT 69 PL 298 75
AYED 70 KIEV CONF
CARRERAS 70 NP 168 35
DAVIES 70 NP 821 359
DIEM 70 KIEV CONF.
ALMEHED 72 NP 840 157
DEANS 72 PN 3 217
DEVENISH 73 PL 478 53
HEMII 73 PL 438 79
HICKS 73 PR 7 2614
LEMOIGNE 73 PURDUE CONF. 93
MOORHOUSE 73 PL 438 44
ROSSI 73 NC 15A 59
ALSO 71 LNC 2 1183
AYED 74 PRIVATE COMM.
ALSO 73 AIX CONFERENCE
BENEVISE 74 NC 19A 529
DEVENISH 74 NP 881 330
DEVENISH 74 PL 528 227
KNIES 74 PR 9 2680
METCALF 74 NP 876 253
MOORHOUSE 74 PR 9 1
CRAWFORD 75 NP 897 125
FELTESSE 75 NP 893 242
LONGACRE 75 PL 558 415
BARBOUR 76 SBMTD. TO NP
BAREYRE 65 PL 18 342
BRANDESEN 65 PR 139 81566
JOHNSON 67 UCL-17683 THESIS
LOVELACE 67 HEIDELBERG C. 79
DONNACHI 69 NP 10 433
AYED 70 PL 318 598
A W HENDRY, R G MOORHOUSE (RHIL)
REVIEWS EARLY PHASE-SHEFT-ANALYSIS RESULTS AND PI- TO ETA N EXPERIMENTS. WE TAKE NUMBERS FROM THE SOLUTION USING BRANDSEN 65.
UCHIYAMA 66 PR 149 1220 F UCHIYAMA-CAMPBELL, R K LCCAN (OXF) (ILL)IJP
DAVIES 67 NC 52A 1112 A T DAVIES, R G MOORHOUSE (GLASGOW,RHEL)
BAREYRE 68 PR 165 1731 P BAREYRE, C BRICMAN, G VILLET (SACLAY)IJP
DONNACHI 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
ALSO 68 VIENNA 139 DONNACHIE RAPPORTEUR-S TALK (GLAS)
ALSO 68 THESIS R G KIRSOPP (EDIN)
DEANS 69 PR 185 1797 S DEANS, J WOCTEN (UNIV S FLORIDA)
DELCOURT 69 PL 298 75 DELCOURT,LEFRANCOIS,PEREZ-Y-JORBA,+ (ORS)
AYED 70 KIEV CONF R AYED,P BAREYRE, G VILLET (SACL)IJP
CARRERAS 70 NP 168 35 B CARRERAS, A DONNACHIE (DARE,MCHS)
DAVIES 70 NP 821 359 A DAVIES (GLAS)
DIEM 70 KIEV CONF. + SMADJA, CHAVANON, DELER, DOLBEAU+ (SACL)
ALMEHED 72 NP 840 157 +LOVELACE (LUND,RUTG)IJP
DEANS 72 PN 3 217 +JACOBS, LYONS, HICKS (U S FL TAMPA+CARN)
DEVENISH 73 PL 478 53 DEVENISH,RANKIN,LYTH (LOUC+BOON+LANC)IJP
HEMII 73 PL 438 79 FEMMI,INAGAKI+ (KYOTO+SAGA+KEK+TKY)IJP
HICKS 73 PR 7 2614 +DEANS, JACOBS, LYONS+ (CARN+ORNL+SOUTH FLA)IJP
LEMOIGNE 73 PURDUE CONF. 93 +GRANET, MARTY, AYED, BAREYRE, BORGEAUD,+ (SACL)IJP
MOORHOUSE 73 PL 438 44 MOORHOUSE, OBERLACK (GLAS+LBL)IJP
ROSSI 73 NC 15A 59 +PIAZZA,SUSTINNO,+ (ROMA,FRAS,NAPL,PAVIA)IJP
ALSO 71 LNC 2 1183 CARBONARA,FIORIO,+ (NAPL,FRAS,PAVIA,ROMA)IJP
AYED 74 PRIVATE COMM. AYED, BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)IJP
BENEVISE 74 NC 19A 529 BENEVENTANO,DANGELO,NOTARISTEFANI,+ (ROMA)IJP
DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
DEVENISH 74 PL 528 227 DEVENISH,LYTH,RANKIN (DESY,LANC,BOON)IJP
KNIES 74 PR 9 2680 KNIES,MOORHOUSE,OBERLACK (LBL+GLAS)IJP
METCALF 74 NP 876 253 W J METCALF,R L WALKER (CTI)IJP
MOORHOUSE 74 PR 9 1 MCOORHOUSE,OBERLACK,ROSENFELD (GLAS+LBL)IJP
CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
FELTESSE 75 NP 893 242 +AYED, BAREYRE, BORGEAUD, DAVID, ERNEINI+ (SACL)IJP
LONGACRE 75 PL 558 415 +ROSENFELD,LASINSKI,SMADJA (LBL,SACL)IJP
BARBOUR 76 SBMTD. TO NP I. M. BARBOUR, R. L. CRAWFORD (GLAS)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
BAREYRE 65 PL 18 342 + BRICMAN, STIRLING, VILLET (SACLAY)IJP
BRANDESEN 65 PR 139 81566 +DONNELL, MCOORHOUSE (OURHAM, RHEL)IJP
BASIS OF NUMBERS WE QUOTE FRGM HENDRY 65.
JOHNSON 67 UCL-17683 THESIS C H JOHNSON (LRL)
LOVELACE 67 HEIDELBERG C. 79 C LOVELACE (CERN)IJP
DONNACHI 69 NP 10 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

Baryons

N(1535), N(1670)

THE FOLLOWING ARTICLES DEAL WITH THE REACTIONS PI- P TO ETA N AND GAMMA P TO ETA P NEAR THRESHOLD. THE DATA AND THE THEORETICAL ARTICLES ARE USEFUL IN UNDERSTANDING THE BEHAVIOR OF THE S11 AMPLITUDE AS DETERMINED IN PI P PHASE-SHIFT ANALYSES. FURTHER REFERENCES MAY BE FOUND IN THEM.

Table listing experimental and theoretical references for N(1535) and N(1670) baryons, including authors, journals, and years.

1520 MEV REGION - PRODUCTION EXPERIMENTS

8 N*1/2(1520, JP=) I=1/2 PRODUCTION EXPERIMENTS

THIS INFORMATION REFERS TO EITHER THE D13 OR THE S11 STATE SEEN AT THIS MASS. FOR SPIN-PARITY ANALYSIS OF THIS MASS REGION, SEE JOHNSTAD 72.

8 N*1/2(1520) MASS (MEV) (PROD. EXP.)

Table of mass measurements for N*1/2(1520) baryon, listing mass values and associated references.

8 N*1/2(1520) WIDTH (MEV) (PROD. EXP.)

Table of width measurements for N*1/2(1520) baryon, listing width values and associated references.

8 N*1/2(1520) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for N*1/2(1520) baryon, listing decay channels and branching ratios.

8 N*1/2(1520) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for N*1/2(1520) baryon, listing ratios and associated references.

REFERENCES FOR N*1/2(1520) (PROD. EXP.)

Table of references for N*1/2(1520) baryon, listing authors and journal information.

Data Card Listings

For notation, see key at front of Listings.

Table listing data card entries for N(1670) baryon, including authors and journal information.

N(1670) 64 N*1/2(1670, JP=5/2-) I=1/2 D15 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

64 N*1/2(1670) MASS (MEV)

Table of mass measurements for N*1/2(1670) baryon, listing mass values and associated references.

64 N*1/2(1670) WIDTH (MEV)

Table of width measurements for N*1/2(1670) baryon, listing width values and associated references.

64 N*1/2(1670) REAL PART OF POLE POSITION (MEV)

Table of real part of pole position for N*1/2(1670) baryon, listing values and references.

64 N*1/2(1670) 2*IMAG PART OF POLE POSITION (MEV)

Table of imaginary part of pole position for N*1/2(1670) baryon, listing values and references.

64 N*1/2(1670) PARTIAL DECAY MODES

Table of partial decay modes for N*1/2(1670) baryon, listing decay channels and branching ratios.

64 N*1/2(1670) BRANCHING RATIOS

Table of branching ratios for N*1/2(1670) baryon, listing ratios and associated references.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1670), N(1688)

R8 N*1/2(1670) FROM PI N TO N*3/2(1232) PI, D-WAVE SORT(PI*P11) 11/75*
R8 L (-.45)OR -.50 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

SEE NOTE PRECEDING THE N*1/2(1688) INELASTIC DECAY MODE MEASUREMENTS.

64 N*1/2(1670) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for particle ID (A1, A2, A3, A4), helicity, and various decay amplitudes and phase shifts for N*1/2(1670) and N*1/2(1688).

Table listing references for N*1/2(1670) and N*1/2(1688) from various sources like BRANDSEN, TRIPP, BAREYRE, etc.

N(1688)

65 N*1/2(1688, JP=5/2+) I=1/2

F'15

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

65 N*1/2(1688) MASS (MEV)

Table listing mass measurements for N*1/2(1688) from various experiments like BRANDSEN, BAREYRE, etc.

65 N*1/2(1688) WIDTH (MEV)

Table listing width measurements for N*1/2(1688) from various experiments like W 1, W 3, etc.

65 N*1/2(1688) REAL PART OF POLE POSITION (MEV)

Table listing real part of pole position for N*1/2(1688) with RE (1688.) and LONGACRE 75 IPWA.

65 N*1/2(1688) 2*IMAG PART OF POLE POSITION (MEV)

Table listing imaginary part of pole position for N*1/2(1688) with IM (132.) and LONGACRE 75 IPWA.

65 N*1/2(1688) PARTIAL DECAY MODES

Table listing partial decay modes for N*1/2(1688) into various particles like PI N, N*1/2(1688) INTO N ETA, etc.

65 N*1/2(1688) BRANCHING RATIOS

Table listing branching ratios for N*1/2(1688) into various particles like R1 1, R1 3, etc.

MORE INFORMATION ON THE INELASTIC DECAY MODES OF THE 1690 MEV BUMP, AS SEEN IN PRODUCTION EXPERIMENTS, MAY BE FOUND BELOW

Table listing inelastic decay modes and branching ratios for N*1/2(1688) into various particles like R2, R3, R4, R5, etc.

Baryons

N(1688), N(1700)

Data Card Listings

For notation, see key at front of Listings.

R10 N*1/2(1688) FROM PI N TC N RHO,S=3/2,P-WAVE SQRT(P1*P14) 11/75*
R10 L (+.27)OR +.30 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

R11 N*1/2(1688) FROM PI N TO K SIGMA SQRT(P1*P15) 11/75*
R11 2 LESS THAN .001 DEANS 75 DPWA PI N TO K SIGMA 11/75*
R11 2 RANGE GIVEN IS FROM 3 OF 4 BEST SOLUTIONS, NOT PRESENT IN SLTN. 11/75*

65 N*1/2(1688) PHOTON DECAY AMPL(GEV**1/2)
FCR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*1/2(1688) INTO GAM P, HELICITY=1/2 (GEV**1/2)
A1 +.015 .023 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A1 (-.J03) HEMMIL 73 + FWD P10 PHTOPROD 2/74

A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

A2 N*1/2(1688) INTO GAM P, HELICITY=3/2 (GEV**1/2)
A2 +.146 .031 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A2 +.100 .012 MOORHUS 73 DPWA PI N PHOTO-PRCD 2/73

A2 AVERAGE MEANINGLESS (SCALE FACTOR = 3.0)

A3 N*1/2(1688) INTO GAM N, HELICITY=1/2 (GEV**1/2)
A3 +.035 .049 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A3 +.017 .014 MOORHUS 73 DPWA PI N PHOTO-PRCD 2/73

A3 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

A4 N*1/2(1688) INTO GAM N, HELICITY=3/2 (GEV**1/2)
A4 -.018 .039 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A4 -.005 .018 MOORHUS 73 DPWA PI N PHOTO-PRCD 2/73

A4 AVERAGE MEANINGLESS (SCALE FACTOR = 2.9)

***** REFERENCES FOR N*1/2(1688) *****

SEE A PREVIOUS EDITION (RMP 37, 633, 1965) FOR EARLIER REFERENCES.

BRANDSEN 65 PL 19 420 +ODDONNELL, MCOORHOUSE (DURHAM, RHEL) IJP
HEUSCH 66 PRL 17 1019 C A HEUSCH, C Y PRESCOTT, R F DASHEN (CIT)
TRIPP 67 NP B3 10 + LEITH, + (LRL,SLAC,CERN,HEID,SACLAY)

N(1700) 66 N*1/2(1700, JP=1/2-1) I=1/2 S11
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

66 N*1/2(1700) MASS (MEV)

M (1695.0) BRANDSEN 65 RVUE PHASE-SHIFT ANAL 9/66
M 1 (1700.0) MICHAEL 66 RVUE FITS BAREYRE S11 7/66
M 1 (1710.0) BAREYRE 68 RVUE PHASE-SHIFT ANAL 11/67

66 N*1/2(1700) WIDTH (MEV)

W 1 (240.0) MICHAEL 66 RVUE 7/66
W 1 (260.0) BAREYRE 68 RVUE 11/67
W 3 (300.0) DGNWACH1 68 RVUE 8/69

66 N*1/2(1700) REAL PART OF PCLE POSITION (MEV)

RE (1648.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

66 N*1/2(1700) 2*IMAG PART OF POLE POSITION (MEV)

IM (117.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

66 N*1/2(1700) PARTIAL DECAY MODES

P1 N*1/2(1700) INTO PI N DECAY MASSES 139+ 938
P2 N*1/2(1700) INTO N ETA 939+ 548
P3 N*1/2(1700) INTO LAMBDA K 1115+ 497

66 N*1/2(1700) BRANCHING RATIOS

R1 N*1/2(1700) INTO (PI N)/TOTAL (P1)
R1 (1.0) APPROX MICHAEL 66 RVUE 7/66
R1 3 (0.79) DGNWACH1 68 RVUE 8/69

R2 N*1/2(1700) FROM PI N TO K LAMBDA SQRT(P1*P3) 4/75*

R2 (+20) (.05) ORITO 69 RVUE 4/75*

R2 A (-21)OR .23 WAGNER 71 IPWA PI- P TO K LAMB 4/75*

R2 (-179) .033 DEVENISH 73 DPWA O FIXED T DISP REL 4/75*

R2 (.12) KNASEL 75 DPWA O PI- P TO KO LAM 11/75*

R3 N*1/2(1700) INTO (LAMBDA K)/TOTAL (P3)

R3 8 (0.028) APPROX. RUSH 68 MPWA T POLE + RESON. 8/69

R3 B PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

R4 N*1/2(1700) INTO (N ETA)/TOTAL (P2)

R4 B (0.013) BOTKE 69 MPWA T POLE + RESON. 10/69

R4 B (0.03) (0.02) DEANS 69 MPWA T POLE + RESON. 8/69

R4 C (0.19) OR 0.27 DEANS 70 MPWA T POLE + RESON. 5/70

R4 B PARAMETRIZATION USED COULD BE IN DANGER OF DOUBLE COUNTING

R4 C CARRERAS TO USES REGGE POLES + RESONANCES. VALUES SUSPICIOUSLY LARG

R5 N*1/2(1700) FROM GAMMA PROTON TO K LAMBDA SQRT(P3*P4) 9/73

R5 (0.002)OR LESS ORITO 69 CNTR K LAM PHOTOPROD 10/71

R5 (0.0072) SCHORSCH 70 DPWA K LAM PHOTOPROD. 10/71

R5 (.0060) DEANS 72 MPWA GAM P-K LM, SOL D 9/73

R6 N*1/2(1700) FROM GAMMA PRGTON TO ETA PROTON SQRT(P2*P4) 9/73

R6 2 (0.101) HICKS 73 MPWA GAM P-ETA P 9/73

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1700)

Table with columns for particle ID, quantum numbers, and decay parameters. Includes entries for N(1700) from various production methods like PI N TO K SIGMA and LANGBEIN.

66 N*1/2(1700) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table of decay amplitudes for N*1/2(1700) into various channels like GAM P, HELICITY=1/2 and GAM N, HELICITY=1/2. Includes average meaningless scale factors.

***** REFERENCES FOR N*1/2(1700) *****

Bibliography table listing references for N*1/2(1700) with author names, journal titles, and dates. Includes names like BRANDSEN, MICHAEL, BAREYRE, etc.

N(1700)

18 N*1/2(1700, JP=3/2-) I=1/2
AYED 74 AND LONGACRE 75 BOTH FIND EVIDENCE FOR THIS STATE. THERE IS ADDITIONAL EVIDENCE FROM PHOTOPRODUCTION AND ASSOCIATED PRODUCTION.

18 N*1/2(1700) MASS (MEV)

Table of mass measurements for N(1700) from various experiments like DONNACHIE, KIRSOPP, WAGNER, etc.

Table with columns for particle ID, quantum numbers, and decay parameters. Includes entries for N(1700) from LANGBEIN and LONGACRE.

18 N*1/2(1700) WIDTH (MEV)

Table of width measurements for N(1700) from experiments like DEANS, LANGBEIN, AYED, LONGACRE.

18 N*1/2(1700) REAL PART OF POLE POSITION (MEV)

Table of real part of pole position for N(1700) from LANGACRE.

18 N*1/2(1700) 2*IMAG PART OF POLE POSITION (MEV)

Table of imaginary part of pole position for N(1700) from LANGACRE.

18 N*1/2(1700) PARTIAL DECAY MODES

Table of partial decay modes for N(1700) into various channels like PI N, LAMBDA K, GAM P, etc.

18 N*1/2(1700) BRANCHING RATIOS

Table of branching ratios for N(1700) from experiments like DEANS, LANGBEIN, AYED, LONGACRE, etc.

18 N*1/2(1700) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table of decay amplitudes for N(1700) into various channels like GAM P, HELICITY=1/2 and GAM N, HELICITY=1/2.

Table of mass measurements for N(1700) from various experiments like DEANS, LANGBEIN, AYED, LONGACRE, etc.

Baryons
N(1700)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, and references. Includes entries like DONNACH2, KIRSOPP, WAGNER, DEANS, DEVENISH, LANGREIN, etc.

1700 MEV REGION - PRODUCTION EXPERIMENTS

20 N*1/2(1700, JP=) I=1/2 PRODUCTION EXPERIMENTS

PARTIAL WAVE ANALYSIS REQUIRES AT LEAST FOUR I=1/2 STAT IN THE 1670 TO 1780 REGION (D15, F15, S11, P11) AND AT LEAST ONE J=3/2 STATE (D33). OBVIOUSLY, DIFFERENT EXPERIMENTS ARE SEEING DIFFERENT STATES AND OFTEN IT IS NOT CLEAR WHAT ISOPIN STATE IS BEING OBSERVED. NO EFFORT WAS MADE TO SEPARATE THESE EXPERIMENTS ACCORDING TO JP, SINCE NONE OF THE REPORTED JP IS FIRMLY ESTABLISHED. WE LIST ALL THE INFORMATION HERE, BUT WE HAVE NOT USED IT IN THE BARYON TABLE.

20 N*1/2(1700) MASS (MEV) (PROD. EXP.)

Table listing mass measurements for N*1/2(1700) with columns for mass (MeV), error, and references. Includes entries like A-BORELLI, GALLCWAY, BARNES, etc.

20 N*1/2(1700) WIDTH (MEV) (PROD. EXP.)

Table listing width measurements for N*1/2(1700) with columns for width (MeV), error, and references. Includes entries like A-BORELLI, GALLCWAY, BARNES, etc.

Table with columns for particle name, mass, and references. Includes entries like OH, RONAT, ABE, DAVIDSON, LICHTMAN, BRAUNI, CAVALLI, MUSGRAVE, WEBB.

20 N*1/2(1700) PARTIAL DECAY MODES (PROD. EXP.)

Table listing partial decay modes for N*1/2(1700) with columns for mode, branching ratio, and references. Includes entries like INTO PI N, INTO N PI PI, INTO N*3/2(1232) PI, etc.

20 N*1/2(1700) BRANCHING RATIOS (PROD. EXP.)

Table listing branching ratios for N*1/2(1700) with columns for mode, ratio, and references. Includes entries like INTO (PI N)/(PI N*3/2(1232)), INTO (N PI)/(N PI PI), etc.

REFERENCES FOR N*1/2(1700) (PROD. EXP.)

Table listing references for N*1/2(1700) with columns for author, journal, and year. Includes entries like MADANSKY, ALEXANDER, ALLES-BORELLI, etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1700), N(1780)

Table listing particle data for N(1700) and N(1780) with columns for name, mass, width, and other properties.

N(1780)

14 N*1/2(1780, JP=1/2+) I=1/2 P11 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

Table for N(1780) MASS (MEV) with columns for mass, width, and other parameters.

Table for N(1780) WIDTH (MEV) with columns for width, mass, and other parameters.

Table for N(1780) REAL PART OF POLE POSITION (MEV) with columns for real part, width, and other parameters.

Table for N(1780) 2*IMAG PART OF POLE POSITION (MEV) with columns for imaginary part, width, and other parameters.

Table for N(1780) PARTIAL DECAY MODES with columns for decay mode, mass, width, and other parameters.

Table for N(1780) BRANCHING RATIOS with columns for branching ratio, mass, width, and other parameters.

Table listing particle data for N(1700) and N(1780) with columns for name, mass, width, and other properties.

Table for N(1780) PHOTON DECAY AMPL(1GEV**1/2) with columns for amplitude, mass, width, and other parameters.

Table listing references for N(1780) with columns for reference name, mass, width, and other parameters.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1780), N(1810), N(1990)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623 S R DEANS (UNIV S FLORIDA)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

N(1810)

15 N*1/2(1810, JP=3/2+) I=1/2

P13

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

15 N*1/2(1810) MASS (MEV)

Table with columns: M, X, W, A, M, M, M, M, M, M, M, M, L, L. Rows contain mass values and associated parameters like DGNACHI, LEA, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) WIDTH (MEV)

Table with columns: W, L. Rows contain width values and associated parameters like DGNACHI, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) REAL PART OF POLE POSITION (MEV)

Table with columns: RE, W, W, W, W, W, L. Rows contain real part values and associated parameters like LONGACRE.

15 N*1/2(1810) 2*IMAG PART OF POLE POSITION (MEV)

Table with columns: IM, W, W, W, W, W, L. Rows contain imaginary part values and associated parameters like LONGACRE.

15 N*1/2(1810) PARTIAL DECAY MODES

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10. Rows contain decay mode descriptions and associated parameters like DGNACHI, LEA, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) BRANCHING RATIOS

Table with columns: R1, R2, R3, R4, R5, R6, R7, R8. Rows contain branching ratio descriptions and associated parameters like DGNACHI, LEA, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns: A1, A2, A3, A4. Rows contain amplitude values and associated parameters like DGNACHI, LEA, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

REFERENCES FOR N*1/2(1810)

Table with columns: DGNACHI, RUSH, BOTKE, DEANS, LEA, AYED, CARRERAS, DAVIES, WAGNER, ALMEHED, DEANS, HICKS, AYED, DEVENISH, DEVENISH, METCALF, CRAWFORD, KNASEL, LONGACRE, BARBOUR. Rows contain references for various parameters.

N(1990)

17 N*1/2(1990, JP=7/2+) I=1/2

F17

THE MOST RECENT PI N PARTIAL WAVE ANALYSIS, AYED 74, FINDS EVIDENCE FOR THIS STATE. THERE IS ALSO SOME INDICATION IN ASSOCIATED PRODUCTION.

17 N*1/2(1990) MASS (MEV)

Table with columns: M, L. Rows contain mass values and associated parameters like DGNACHI, LEA, DAVIES, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1990), N(2000), N(2040)

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Includes entries for N(1990) and N(2000).

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Includes entries for N(2040).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists partial decay modes for N(1990).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists partial decay modes for N(2000).

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(1990).

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(2000).

Table with 4 columns: Particle ID, Photon Decay Amplitude, and other properties. Lists photon decay amplitudes for N(1990).

Table with 4 columns: Particle ID, References, and other properties. Lists references for N(2000).

Table with 4 columns: Particle ID, References, and other properties. Lists references for N(1990).

N(2040) THIS STATE IS NOW SEEN BY THE SACLAY GROUP, AYED 74.

Table with 4 columns: Particle ID, Mass (MeV), and other properties. Lists mass values for N(2040).

N(2000) THE MOST RECENT PI N PARTIAL WAVE ANALYSIS, AYED 74, FINDS EVIDENCE FOR THIS STATE. THERE IS ALSO SOME INDICATION IN ASSOCIATED PRODUCTION.

Table with 4 columns: Particle ID, Width (MeV), and other properties. Lists width values for N(2040).

Table with 4 columns: Particle ID, Mass (MeV), and other properties. Lists mass values for N(2000).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists partial decay modes for N(2040).

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(2040).

Baryons

N(2040), N(2100), N(2190)

Data Card Listings

For notation, see key at front of Listings.

16 N*1/2(2040) PHOTON DECAY AMPL(GEV**--1/2)
FER DEFINITION OF GAMMA-NUCLEON DECA AMPLITUDES, SEE MINI-
REVIEW PRECEDING THE BARYON LISTINGS.

REFERENCES FOR N*1/2(2040)
A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
DONNACHIE RAPPORTEUR.S TALK (GLAS)
R G KIRSOPP (EDIN)

N(2100) 04 N*1/2(2100, JP=1/2-) I=1/2 S11
NOW ALSO SEEN BY SACLAY, AYED 74.

04 N*1/2(2100) MASS (MEV)
M 7 (2070.) ROYCHOUD 71 DPWA 3/72
M 7 (2100.) ALMEHED 72 IPWA 2/72
M (2283.) AYED 74 IPWA 2/74

04 N*1/2(2100) WIDTH (MEV)
W 7 (200.) ALMEHED 72 IPWA 2/72
W (310.) AYED 74 IPWA 2/74

04 N*1/2(2100) PARTIAL DECAY MODES
P1 N*1/2(2100) INTO PI N 139+ 938

04 N*1/2(2100) BRANCHING RATIOS
R1 N*1/2(2100) INTO (PI N)/TOTAL (P1)
R1 7 (0.5) ALMEHED 72 IPWA 2/72
R1 (.14) AYED 74 IPWA 2/74

REFERENCES FOR N*1/2(2100)
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSEN (DURH)IJP
ALMEHED 72 NP 840 157 +LOVELACE (LUND,RUTG)IJP
AYED 74 PRIVATE COMMCTN. AVEC,BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP

N(2100) 05 N*1/2(2100, JP=5/2-) I=1/2 D15
NOW ALSO SEEN BY SACLAY, AYED 74.

05 N*1/2(2100) MASS (MEV)
M 7 (2100.) ALMEHED 72 IPWA 2/72
M (2100.) AYED 74 IPWA 2/74

05 N*1/2(2100) WIDTH (MEV)
W 7 (150.) ALMEHED 72 IPWA 2/72
W (220.) AYED 74 IPWA 2/74

05 N*1/2(2100) PARTIAL DECAY MODES
P1 N*1/2(2100) INTO PI N 139+ 938

05 N*1/2(2100) BRANCHING RATIOS
R1 N*1/2(2100) INTO (PI N)/TOTAL (P1)
R1 7 (0.2) ALMEHED 72 IPWA 2/72
R1 (.08) AYED 74 IPWA 2/74

N(2190) 71 N*1/2(2190, JP=7/2-) I=1/2 G17
ROYCHOUDHURY 71 FIND SOME INDICATION OF P11 AND FIT IN
THIS REGION.BRANSDEN 71 ALSO FIND P11,F15,AND G19 RESO-
NANT NEAR THIS MASS. AYED 74 FIND A G19.

71 N*1/2(2190) MASS (MEV)
M (2190.0) DIDDENS 63 CNTR P1+- P TOTAL
M (2210.0) HOHLER 64 RVUE DATA + DISP REL
M (2190.0) APPROX YOKOSAWA 66 CNTR P1- P DSIG + POL 7/66
M 3 (2265.0) DONNACH1 68 RVUE PHASE-SHIFT ANAL 6/68
M (2000.0) LEA 69 CNTR P1-P ELASTIC 8/69
M 2180. 25. ANDERSON 70 MMS - P1- P TO P1- MMS 2/71
M 6 (2158.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM
M (2260.0) HULL 70 MPWA SMALL ANGLE PI-P 1/71
M (2160.0) (50.0) AMALDI 71 CNTR P P AT 24 GEV 10/71
M (2160.) BRANSDEN 71 DPWA 3/72
M (2200.) ROYCHOUD 71 DPWA 3/72
M 7 (2225.) ALMEHED 72 IPWA 2/72
M (2190.0) OTT 72 MPWA 0 PI-P BKWD ELSTC 2/73
M 1 (2208.) HICKS 73 MPWA GAM P-ETA P 9/73
M 1 ONLY STATES FROM TABLE VII OF HICKS73 ARE INCLUDED IN LISTINGS. 9/73
M 1 M AND W ARE FROM SOLUTION C2,BR=SQRT(G1)/W WITH G FROM TABLE VII. 4/75*
M (2141.) (20.) ABE 74 + P+P->P+X, JCBN PK 2/74
M 2 (2133.) AYED 74 IPWA 2/74
M 2 THIS IS A G19 RESONANCE LISTED HERE UNTIL CONFIRMED. 2/74

71 N*1/2(2190) WIDTH (MEV)
W (200.0) DIDDENS 63 CNTR 7/66
W (200.0) HOHLER 64 RVUE 7/66
W (220.0) APPROX YOKOSAWA 66 CNTR 6/68
W 3 (298.0) DONNACH1 68 RVUE 2/71
W 275. 70. ANDERSON 70 MMS - PI- P TO P1- MMS 1/71
W 6 (325.0) AYED 70 IPWA 1/71
W (239.0) HULL 70 MPWA SMALL ANGLE PI-P 1/71
W 7 (150.) ALMEHED 72 IPWA 2/72
W 1 (193.) HICKS 73 MPWA GAM P-ETA P 9/73
W (243.) AYED 74 IPWA 2/74
W (193.) AYED 74 IPWA 2/74
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

71 N*1/2(2190) PARTIAL DECAY MODES
P1 N*1/2(2190) INTO PI N 139+ 938
P2 N*1/2(2190) INTO LAMBDA K 1115+1765
P3 N*1/2(2190) INTO N PI PI 938+ 139+ 139
P4 N*1/2(2190) INTO GAM P,HELICITY=3/2 0+ 938
P5 N*1/2(2190) INTO GAM P,HELICITY=1/2 0+ 938
P6 N*1/2(2190) INTO GAM N,HELICITY=3/2 0+ 939
P7 N*1/2(2190) INTO GAM N,HELICITY=1/2 0+ 939
P8 N*1/2(2190) INTO ETA N 548+ 938
P9 N*1/2(2190) INTO SIGMA K 493+1189

71 N*1/2(2190) BRANCHING RATIOS
R1 N*1/2(2190) INTO (PI N)/TOTAL (P1)
R1 (0.3) APPROX DIDDENS 63 CNTR 7/66
R1 (0.3) APPROX YOKOSAWA 66 CNTR 7/66
R1 3 (0.349) DONNACH1 68 RVUE 6/68
R1 6 (0.150) AYED 70 IPWA 1/71
R1 (0.09) HULL 70 MPWA SMALL ANGLE PI-P 1/71
R1 7 (0.35) ALMEHED 72 IPWA 2/72
R1 (.25) OTT 72 MPWA 0 PI-P BKWD ELSTC 2/73
R1 (1.611) AYED 74 IPWA 2/74
R1 (.095) AYED 74 IPWA 2/74

R2 N*1/2(2190) FROM GAMMA PROTON TO K LAMBDA SORT((P4+P5)*P2) 9/73
R2 (.0161) DEANS 72 MPWA GAM P-K LM,SOL D 9/73
R3 N*1/2(2190) FROM GAMMA PROTON TO ETA PROTON SQRT((P4+P5)*P8) 9/73
R3 (.0094) HICKS 73 MPWA GAM P-ETA P 9/73
R4 N*1/2(2190) FROM PI N TO K SIGMA SQRT(P1*P9) 11/75*
R4 4 (.014) TO .019 DEANS 75 DPWA PI N TO K SIGMA 11/75*
R4 4 RANGE GIVEN IS FRCH FOUR BEST SOLUTIONS. 11/75*

REFERENCES FOR N*1/2(2190)
DIDDENS 63 PRL 10 262 +JENKINS, KYCIA, RILEY (BNL) I
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
YOKOSAWA 66 PRL 16 714 +SUMA,FILL,ESTERLING, BOOTH (ANL,CHIC) JP
DONNACH1 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
ALSO 68 THESIS DONNACHIE RAPPORTEUR.S TALK (GLAS)
R G KIRSOPP (EDIN)
LEA 69 PL 298 584 LEA,GADES,WARD,COWAN,+ (RHEL,BRISTOL,DARE)
ANDERSON 70 PRL 25,699 +BLESER,BLIEDEN,COLLINS++ (BNL,CARN)
AYED 70 KIEV CONF R AYED,P BAREYRE, G VILLET (SACL)IJP
HULL 70 PR D2 1783 J HULL, R LEACOCK (ISU)

Data Card Listings

Baryons

For notation, see key at front of Listings. N(2190), N(2220), N(2650), N(3030), N₇(3245)

AMALDI 71 PL 348 435 +BIANCATELLI, BOSIO, + (I SANITA ROMA+GERN)
 BRANSDEN 71 NP 826 511 ,OGDEN (DURHI)JP
 ALSO 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANSDEN (DURHI)JP
 ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSDEN (DURHI)JP

ALME-ED 72 NP 840 157 +LOVELACE (LUND, RUTG)JP
 DEANS 72 PRD 6 1906 DEANS, JACOBS, LYONS, MCNTGOMERY (SOUTH FLA.)JP
 DTT 72 PL 428 133 +TRISCHUK, VAVRA, RICHARDS, + (MCGI, STLO, IOWA)JP
 ALSO 72 MCGILL THESIS J, VAVRA (MCGI) JP
 HICKS 73 PRD 7 2614 +DEANS, JACOBS, LYONS+ (CARA+ORNL+SOUTH FLA.)JP

ABE 74 PL 538 114 +ALSPECTOR, BCMBEROWITZ+ (RLTG, UPNJ, FSU)
 AYEY 74 PRIVATE COMMCTN. AYEY, BAREYRE (SACLII)JP
 ALSO 73 AIX CONFERENCE AYEY, BAREYRE (SACLII)JP
 DEANS 75 NP 896 90 +MITCHELL, MCNTGOMERY, + (SFLA, ALABAMA)JP

PAPERS NOT REFERRED TO IN DATA CARDS.

AYED 70 PL 318 598 +BAREYRE, VILLET (SACLAY)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN DATA CARDS.

BARGER 66 PRL 16 913 V BARGER, D CLINE (NISC) P
 CARROLL 66 PRL 16 288 +CORBETT, DAMERELL, MIDDLEMAS, + (RHEL, OXF)J-L
 CARROLL 66 PRL 17 1274 +CORBETT, DAMERELL, MIDDLEMAS, + (RHEL, OXF)J-L
 ERRATUM CHANGING THE RATHER WEAK DETERMINATION OF J-L TO +1 (2.) JP
 KORMANYO 66 PRL 16 709 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) JP
 BUSZA 67 NC 52A 331 +DAVIS, DUFF, HEYMANN, + (LOUC, WESTFIELD)

N(2220)

90 N*1/2(2220, JP=9/2+) I=1/2 **H₁₉**
 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

90 N*1/2(2220) MASS (MEV)

M	(2200.)	APPROX.	BUSZA	67 DSPK	LEG. POLYN. ANAL.	2/71
M	(2221.0)		AYED	70 IPWA		1/71
M	6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM		HULL	70 MPWA	SMALL ANGLE PI-P	1/71
M	(2245.0)		AYED	74 IPWA		2/74
M	(2249.)					

90 N*1/2(2220) WIDTH (MEV)

W	6	(258.0)	AYED	70 IPWA	SMALL ANGLE PI-P	1/71
W		(329.0)	HULL	70 MPWA		1/71
W		(347.)	AYED	74 IPWA		2/74

90 N*1/2(2220) PARTIAL DECAY MODES

P1	N*1/2(2220) INTO PI N	DECAY MASSES
P2	N*1/2(2220) INTO N PI	139+ 938
		939+ 548

90 N*1/2(2220) BRANCHING RATIOS

R1	6	N*1/2(2220) INTO (PI N)/TOTAL	(P1)
R1		(0.140)	AYED
R1		(0.15)	HULL
R1		(.204)	AYED

SMALL ANGLE PI-P

REFERENCES FOR N*1/2(2220)

BUSZA 67 NC 52A 331 +DAVIS, DUFF, HEYMANN, NIMMON + (LOUC+LOWC)
 AYEY 70 KIEV CONF. R AYEY, P BAREYRE, G VILLET (SACLII)JP
 HULL 70 PR D2 1783 J HULL, R LEACOCK (ISU)
 AYEY 74 PRIVATE COMMCTN. AYEY, BAREYRE (SACLII)JP
 ALSO 73 AIX CONFERENCE AYEY, BAREYRE (SACLII)JP

PAPERS NOT REFERRED TO IN DATA CARDS

AYED 70 PL 318 598 +BAREYRE, VILLET (SACLAY)

N(2650) BUMPS

72 N*1/2(2650, JP=) I=1/2 PRODUCTION EXPERIMENTS
 ROYCHOUDHURY 71 CLAIM F15(2400) AND G19(2400) TO BE POSSIBLE RESONANCES. BRANSDEN 71 FIND THE POSSIBLE RESONANT CANDIDATES S11(2520) AND H19(2590).

72 N*1/2(2650) MASS (MEV) (PROD. EXP.)

M	(2700.0)	ALVAREZ	64 CNTR	PI PHCTOPROD
M	(2660.0)	HOHLER	64 RVUE	DATA + DISP REL
M	(2600.0)	WAHLIG	64 DSPK	0 PI-P CH EX
M	(2633.0)	BARGER	66 FIT	TOTAL + CH EX
M	2649.0	CITRCN	66 CNTR	PI+- P TOTAL

11/67 7/66

72 N*1/2(2650) WIDTH (MEV) (PROD. EXP.)

W	(100.0)	ALVAREZ	64 CNTR	
W	(200.0)	HOHLER	64 RVUE	
W	(425.0)	BARGER	66 FIT	TOTAL + CH EX
W	360.0	CITRCN	66 CNTR	

7/66 11/67 7/66

72 N*1/2(2650) PARTIAL DECAY MODES (PROD. EXP.)

P1	N*1/2(2650) INTO PI N	DECAY MASSES
P2	N*1/2(2650) INTO LAMBDA K	139+ 938
P3	N*1/2(2650) INTO N PI PI	1115+ 497
		938+ 139+ 139

72 N*1/2(2650) BRANCHING RATIOS (PROD. EXP.)

R1 N*1/2(2650) INTO (PI N)/TOTAL (P1)
 R1 ONLY (J+1/2)*(PI N/TOTAL) MEASURED FOR THIS STATE
 R1 B (0.456) (0.018) BARGER 66 RVUE TOTAL + CH EXC. 11/67
 R1 B 0.436 0.028 CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
 R1 B (0.30) BARGER 67 RVUE USES KORMANYOS67 11/67
 R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
 R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
 R1 D (0.24) DIKMEN 67 RVUE USES KORMANYOS66 11/67
 R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
 R1 KORMANYOS 67 CNTR PI-P AT 180 DEG. 11/67

REFERENCES FOR N*1/2(2650) (PROD. EXP.)

ALVAREZ 64 PL 12 710 +BAR-YAM, KERN, LUCKEY, OSBORNE, + (MIT, CEA)
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)
 BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)
 KORMANYO 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P

PAPERS NOT REFERRED TO IN DATA CARDS.

BAECKE 67 NC 51A 761 J BAECKE, M YVERT (KARLSRUHE, ORSAY)J-L
 DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)
 WAHLIG 68 PR 168 1515 M A WAHLIG, I MANNELLI (MIT, PISA)

FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES CCMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

BRANSDEN 71 NP 826 511 ,OGDEN (DURHI)JP
 ALSO 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANSDEN (DURHI)JP
 ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSDEN (DURHI)JP

N(3030) BUMPS

73 N*1/2(3030, JP=) I=1/2 PRODUCTION EXPERIMENTS

73 N*1/2(3030) MASS (MEV) (PROD. EXP.)

M	(3080.0)	HOHLER	64 RVUE	DATA + DISP REL	7/66
M	(3030.0)	CITRON	66 CNTR	PI+- P TOTAL	7/66

73 N*1/2(3030) WIDTH (MEV) (PROD. EXP.)

W	(400.0)	CITRON	66 CNTR		7/66
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73 N*1/2(3030) PARTIAL DECAY MODES (PROD. EXP.)

P1	N*1/2(3030) INTO PI N	DECAY MASSES
P2	N*1/2(3030) INTO N PI PI	139+ 938
		938+ 139+ 139

73 N*1/2(3030) BRANCHING RATIOS (PROD. EXP.)

R1 N*1/2(3030) INTO (PI N)/TOTAL (P1)
 R1 ONLY (J+1/2)*(PI N/TOTAL) MEASURED FOR THIS STATE
 R1 B (0.088) (0.016) BARGER 66 RVUE TOTAL + CH EXC. 11/67
 R1 B (0.048) CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
 R1 B (0.12) BARGER 67 CNTR USES KORMANYOS66 11/67
 R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
 R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
 R1 D (0.016) DIKMEN 67 RVUE USES KORMANYOS67 11/67
 R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

REFERENCES FOR N*1/2(3030) (PROD. EXP.)

HGHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)
 CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
 BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
 DIKMEN 67 PRL 18 798 F N DIKMEN (MICH)

PAPERS NOT REFERRED TO IN DATA CARDS

KORMANYO 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P
 DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)

N₇(3245) BUMPS

74 N* /2(3245, JP=) PRODUCTION EXPERIMENTS

EXISTENCE NOT CONCLUSIVELY ESTABLISHED. I-SPIN NOT DETERMINED, BUT THE NARROW WIDTH PRECLUDES IDENTIFICATION WITH THE N*3/2(3230). OMITTED FROM TABLE.

74 N* /2(3245) MASS (MEV) (PROD. EXP.)

M	3245.0	10.0	KORMANYOS 67 CNTR	PI-P 180 DEG EL	6/68
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74 N* /2(3245) WIDTH (MEV) (PROD. EXP.)

W	(35.0)	OR LESS	KORMANYOS 67 CNTR		6/68
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Baryons

N₇(3245), N(3690), N₇(3755), Δ(1232)

Data Card Listings

For notation, see key at front of Listings.

74 N* /2(3245) PARTIAL DECAY MODES (PROD. EXP.)
P1 N* /2(3245) INTO PI N
DECAY MASSES
139+ 938

74 N* /2(3245) BRANCHING RATIOS (PROD. EXP.)
R1 N* /2(3245) INTO (PI N)/TOTAL (P1)
R1 J IS NOT KNOWN. FOLLOWING IS (J+1/2)*(PI N)/TOTAL (P1)
R1 (0.37) KORMANYOS 67 CNTR 6/68

REFERENCES FOR N* /2(3245) (PROD. EXP.)
KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P

N(3690) BUMPS

75 N*1/2(3690, JP=) I=1/2 PRODUCTION EXPERIMENTS
A BUMP SEEN IN THE INVARIANT MASS OF A VERY COMPLICATED STATE IN + SEVEN PIS, SO AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. NOT INCLUDED IN TABLE.

75 N*1/2(3690) MASS (MEV) (PROD. EXP.)
M 3690.0 10.0 BARTKE 67 HBC + PI+P 8 PRONGS 8/67

75 N*1/2(3690) WIDTH (MEV) (PROD. EXP.)
W 50.0 30.0 BARTKE 67 HBC + 8/67

75 N*1/2(3690) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*1/2(3690) INTO N + 7 PIS
DECAY MASSES

REFERENCES FOR N*1/2(3690) (PROD. EXP.)
BARTKE 67 PL 248 118 +CZYZEWSKI,DANYSZ,+ (CRACOW,ORSAY) I

N7(3755) BUMPS

76 N* /2(3755, JP=) PRODUCTION EXPERIMENTS
A SMALL PEAK IN THE (P P PBAR) INVARIANT MASS FROM 8.4 BEV/C PI+ P TO PI+ P P PBAR EVENTS. AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. OMITTED FROM TABLE.

76 N* /2(3755) MASS (MEV) (PROD. EXP.)
M 3755.0 8.0 EHRLICH 68 HBC + PI+ P P PBAR 6/68

76 N* /2(3755) WIDTH (MEV) (PROD. EXP.)
W 40.0 20.0 EHRLICH 68 HBC + 6/68

76 N* /2(3755) PARTIAL DECAY MODES (PROD. EXP.)
P1 N* /2(3755) INTO PI+ P P PBAR
DECAY MASSES
139+ 938+ 938+ 938

REFERENCES FOR N* /2(3755) (PROD. EXP.)
EHRLICH 68 PRL 20 686 R EHRLICH,R J PLANC,J B WHITTAKER (RUTGERS)

S=0 I=3/2 NUCLEON STATES (Δ)

Δ(1232)

33 N*3/2(1232, JP=3/2+) I=3/2 P33
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED. SEE CARTER 71 AND CARTER 73 FOR PI N CROSS-SECTION DATA IN THIS REGION.

33 N*3/2(1232) MASS (MEV)
M (1234.) ROPER 65 DPWA ++0 PHASE SHIFT AN. 2/72
M (1235.) ALMEHED 72 IPWA 2/74
M 3 (1243.3) (1241.7) CHENG 73 FIT CARTER 71 2/74
M 3 THE TWO ENTRIES ARE FROM TWO DIFFERENT PARAMETRIZATIONS OF THE 2/74
M 3 RESONANCE CONTRIBUTION TO THE P33 PHASE SHIFT. 2/74
M (1230.4) TSCHANG 73 FIT CARTER71 P33 1/74
M (1234.) AYED 74 IPWA 2/74
M++ 1236.0 0.55 OLSSON 65 RVUE ++ TOTAL-SIGMA DATA
M++ 2 1231.0 1.5 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
M++ 1 1231.1 +2 CARTER 73 IPWA ++ PI N 88-310 MEV 9/73
M++ 1 EXPERIMENTAL QUANTITY-SEE CARTER73 FOR COULOMB BARRIER CORRECTIONS 9/73
M++ 2 EXPERIMENTAL QUANTITY-SEE CARTER71 FOR COULOMB BARRIER CORRECTIONS 1/74
M++ AVERAGE MEANINGLESS (SCALE FACTOR = 8.4)

M+ (1231.8) BERENDS 75 IPWA + GA P TO PI NUC 4/75*
M+ 1230.6 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
M+ (1231.7) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*
MO 1236.45 0.65 OLSSON 65 RVUE 0
MO 1232.9 0.6 CARTER 71 MPWA 0 PI-P SIG. TOTAL 1/71
MO AVERAGE MEANINGLESS (SCALE FACTOR = 4.0)

33 N*3/2(1232) WIDTH (MEV)
W (120.) ROPER 65 DPWA ++0 PHASE SHIFT AN. 2/72
W (129.) ALMEHED 72 IPWA 2/74
W 3 (152.2) (145.8) CHENG 73 FIT CARTER 71 2/74
W (120.) TSCHANG 73 FIT CARTER71 P33 1/74
W (120.) AYED 74 IPWA 2/74
W++ 120.0 2.0 OLSSON 65 RVUE ++
W++ 2 111.1 1.8 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
W++ 1 111.5 .4 CARTER 73 IPWA ++ PI N 88-310 MEV 9/73
W++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

M+ 120.2 3.9 CRAWFORD 75 DPWA PI N PHOTO-PRGD 1/76*
M+ (117.4) BARBOUR 76 DPWA PI N PHOTO-PRGD 1/76*
MO 119.6 2.4 OLSSON 65 RVUE 0
MO 114.7 3.0 CARTER 71 MPWA 0 PI-P SIG TOT. 1/71
MO AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)

33 (N*0) - (N***) MASS DIFFERENCE (MEV)
D R (0.45) (0.85) OLSSON 65 RVUE 1/74
D 2 1.3 1.9 CARTER 71 MPWA ++ PI+P SIG. TOTAL 2/73
D 1 1.4 .4 CARTER 73 IPWA PI N 88-310 MEV 9/73
D R REDUNDANT WITH DATA IN MASS LISTING.
D AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 (N*0)-(N***) WIDTH DIFFERENCE (MEV) 9/73
WD 2 6.5 2.3 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
WD 1 10.3 1.2 CARTER 73 IPWA PI N 88-310 MEV 9/73
WD AVERAGE MEANINGLESS (SCALE FACTOR = 1.5)

33 N*3/2(1232) REAL PART OF PCLE POSITION(MEV)
REE M (1214.) MICHAEL 67 2/74
REE P (1211.1) BALL 72 2/73
REE P (1211.6) PDG 72 2/73
REE 3 (1210.7) (1210.7) CHENG 73 FIT DELTA 33 2/74
REE 3 (1214.5) 10. NOGOVA 73 FIT ALMEHED72 2/74
REE M (1213.1) SPEARMAN 74 FIT ZERO TRJCTRY 4/75*
REE M FIT INCLUDES OLSSON 65 PARAMETERS PLUS SCATTERING LENGTH PLUS 6
REE M PHASE SHIFT VALUES FOR TPI=120 TO 492 MEV
REE P ERROR EST. FROM FITS WITH SOMEWHAT VARYING ASSUMPTIONS
REE AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

R++ U 1211.5 .6 BALL 75 ++ FIT CARTER 73 11/75*
R++ U 1210.9 .8 LICHTENB 75 ++ FIT CARTER 73 11/75*
R++ C 1209.6 .5 VASAN 75 ++ FIT CARTER 73 1/76*
R++ C FROM FITS TO COULOMB CORRECTED CARTER 73 PHASE SHIFT. 1/76*
R++ U (1210.5) TO (1210.8) VASAN 75 ++ FIT CARTER 73 1/76*
R++ U FROM FITS TO UNCORRECTED CARTER 73 PHASE SHIFT. 1/76*
R++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)

REO U (1211.6) BALL 75 0 FIT CARTER 73 11/75*
REO U 1210.9 1.4 LICHTENB 75 0 FIT CARTER 73 11/75*
REO C 1210.75 .6 VASAN 75 ++ FIT CARTER 73 1/76*
REO U (1210.2) VASAN 75 ++ FIT CARTER 73 1/76*
REO AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 N*3/2(1232) IMAG PART OF PCLE POSITION(MEV)
IME M (52.) MICHAEL 67 2/74
IME (50.) BALL 72 2/73
IME P 49.5 1.8 PDG 72 FIT DELTA 33 2/73
IME 3 (50.7) (50.6) CHENG 73 FIT CARTER 71 2/74
IME (48.6) 5. NOGOVA 73 FIT ALMEHED72 2/74
IME (49.) SPEARMAN 74 FIT ZERO TRJCTRY 4/75*
IME AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

I++ U 50.1 .6 BALL 75 ++ FIT CARTER 73 11/75*
I++ U 49.6 .75 LICHTENB 75 ++ FIT CARTER 73 11/75*
I++ C 50.4 .5 VASAN 75 ++ FIT CARTER 73 1/76*
I++ U (49.9) TO (50.0) VASAN 75 ++ FIT CARTER 73 1/76*
I++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

IMO U (53.0) BALL 75 0 FIT CARTER 73 11/75*
IMO U 53.25 1.75 LICHTENB 75 0 FIT CARTER 73 11/75*
IMO C 52.8 .75 VASAN 75 ++ FIT CARTER 73 1/76*
IMO U (52.9) TO (53.1) VASAN 75 ++ FIT CARTER 73 1/76*
IMO AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 N*3/2(1232) ABSOLUTE VALUE OF PCLE RESIDUE (MEV)
ABS (53.) BALL 73 FIT DELTA 33 9/73
A++ C (52.4) TO (53.2) VASAN 75 ++ FIT CARTER 73 1/76*
A++ U (52.1) TO (52.4) VASAN 75 ++ FIT CARTER 73 1/76*
ABO C (54.8) TO (55.0) VASAN 75 ++ FIT CARTER 73 1/76*
ABO U (55.2) TO (55.3) VASAN 75 ++ FIT CARTER 73 1/76*

Baryons

$\Delta(1650)$, $\Delta(1670)$

Data Card Listings

For notation, see key at front of Listings.

$\Delta(1650)$

82 N*3/2(1650, JP=1/2-) I=3/2
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

S₃₁

Table with columns for mass (MEV), width (MEV), real part of pole position (MEV), imaginary part of pole position (MEV), and partial decay modes. Includes data for DEVLIN, BAREYRE, DONNACHI, and others.

Table with columns for width (MEV), real part of pole position (MEV), imaginary part of pole position (MEV), and partial decay modes. Includes data for BAREYRE, DCNNACHI, and others.

Table with columns for real part of pole position (MEV) and imaginary part of pole position (MEV). Includes data for LONGACRE.

Table with columns for imaginary part of pole position (MEV) and partial decay modes. Includes data for LONGACRE.

Table with columns for partial decay modes and decay masses. Includes data for various decay channels like INTO PI N, INTO GAM NUCLEON, etc.

Table with columns for branching ratios and partial decay modes. Includes data for various decay channels.

Table with columns for partial decay modes and decay masses. Includes data for various decay channels.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for branching ratios and partial decay modes. Includes data for various decay channels.

REFERENCES FOR N*3/2(1650)

Table listing references for N*3/2(1650) with columns for author, journal, and page number.

Table with columns for mass (MEV), width (MEV), real part of pole position (MEV), imaginary part of pole position (MEV), and partial decay modes. Includes data for CRAWFORD, ROSENFELD, BARBOUR, and others.

$\Delta(1670)$

10 N*3/2(1670, JP=3/2-) I=3/2
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

D₃₃

Table with columns for mass (MEV), width (MEV), real part of pole position (MEV), imaginary part of pole position (MEV), and partial decay modes. Includes data for DCNNACHI, BAREYRE, and others.

Table with columns for width (MEV), real part of pole position (MEV), imaginary part of pole position (MEV), and partial decay modes. Includes data for DCNNACHI, BAREYRE, and others.

Table with columns for real part of pole position (MEV) and imaginary part of pole position (MEV). Includes data for LONGACRE.

Table with columns for imaginary part of pole position (MEV) and partial decay modes. Includes data for LONGACRE.

Table with columns for partial decay modes and decay masses. Includes data for various decay channels like INTO PI N, INTO GAM NUCLEON, etc.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for branching ratios and partial decay modes. Includes data for various decay channels.

REFERENCES FOR N*3/2(1670)

Table listing references for N*3/2(1670) with columns for author, journal, and page number.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for branching ratios and partial decay modes. Includes data for various decay channels.

Data Card Listings
For notation, see key at front of Listings.

Baryons
Δ(1670), Δ(1690), Δ(1890)

A1 +.101 .011 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
A1 (+.120) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)
A2 N*3/2(1670) INTO GAM NUCLEON, FELICITY=3/2 (GEV**=1/2)
A2 .110 .039 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A2 +.022 .052 MOORHOUS 73 DPWA PI N PHOTO-PROD 2/73
A2 .072 .014 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*

Δ(1690) 19 N*3/2(1690, JP=3/2+) I=3/2 P33
RECENT ANALYSES INDICATE A POSSIBLE P33 RESONANCE SOMEWHERE IN THE 1600-2000 MEV REGION. SEE ALSO THE N*3/2(1690) LISTINGS.

19 N*3/2(1690) MASS (MEV)
M 3 (1690.) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
M 3 (1690.) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
M 3 WHERE MAX. ABSORPTION IS -DONNACH1, 2, KIRSOPP EYEBALL FIT CERN 1 10/69

19 N*3/2(1690) WIDTH (MEV)
W 3 (281.) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
W 3 (240.) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
W 6 (598.0) AYED 70 IPWA 1/71
W 7 (220.) ALMEHED 72 IPWA 2/72
W L (204.) AYED 74 IPWA 2/74
W L 1900. OR 1640. LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) REAL PART OF POLE POSITION (MEV)
RE (1609.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) 2*IMAG PART OF POLE POSITION (MEV)
IM (323.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) PARTIAL DECAY MODES
P1 N*3/2(1690) INTO PI N 139+ 938
P2 N*3/2(1690) INTO K SIGMA 493+1189
P3 N*3/2(1690) INTO N*3/2(1232) PI,P-WAVE 1232+ 139

19 N*3/2(1690) BRANCHING RATIOS
R1 N*3/2(1690) INTO (PI N)/TOTAL (P1)
R1 3 (.10) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
R1 3 (.08) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
R1 6 (0.135) AYED 70 IPWA 1/71
R1 7 (0.1) ALMEHED 72 IPWA 2/72
R1 1 (.194) AYED 74 IPWA 2/74
R2 N*3/2(1690) INTO (K SIGMA)/TOTAL (P2)
R2 1 (.0002)OR LESS FEUERBACH 70 RVUE PI P TO K+ SIG+ 7/70
R2 1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68
R2 1 MODEL USED MAY DOUBLE COUNT.

R3 N*3/2(1690) FROM PI N TO K SIGMA SQRT(PI*P2) 11/75*
R3 2 (.006)TO .042 DEANS 75 DPWA PI N TO K SIGMA 11/75*
R3 2 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS. 11/75*

19 N*3/2(1690) PHOTON DECAY AMPL(GEV**=1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*3/2(1690) INTO GAM NUCLEON, HELICITY=1/2 (GEV**=1/2)
A1 .016 .055 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A1 -.033 .037 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*
A1 .003 .015 KNIES 74 DPWA PI N PHOTO PRCD 2/74
A1 .0 .038 METCALF 74 DPWA PI N PHOTO-PROD 2/74
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
A2 N*3/2(1690) INTO GAM NUCLEON, HELICITY=3/2 (GEV**=1/2)
A2 .074 .064 DEVENISH 73 DPWA PI N PHOTO PRCD 2/74
A2 .308 .046 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*

REFERENCES FOR N*3/2(1690)
DONNACH2 68 VIENNA 139 DONNACHIE RAPPORTEUR'S TALK (GLAS)
KIRSOPP 68 THESIS R G KIRSOPP (EDIN)
AYED 70 KIEV CNF R AYED,P BAREYRE, G VILLET (SACL)IJP
FEUERBACH 70 NP 168 85 FEUERBACHER+HOLLADAY (VANDERBILT)
ALMEHED 72 NP 840 157 +LOVELACE (LUND,RUTG)IJP
DEVENISH 73 PL 478 53 MOORHOUS 73 PL 438 44 MOORHOUSE, OBERLACK (LOUC+BNN+LANC)IJP (GLAS+LBL)IJP
AYED 74 PRIVATE COMMCTN. AYED,BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP
DEVENIS2 74 PL 528 227 DEVENISH,LYTH,RANKIN (DESY,LANC,BONN)IJP
KNIES 74 PRD 9 2680 KNIES,MOORHOUSE,OBERLACK (LBL,GLAS)IJP
METCALF 74 NP 876 253 W J METCALF,R L WALKER (CIT)IJP
MOORHOUS 74 PRD 9 1 MOORHOUSE,OBERLACK,ROSENFELD (GLAS+LBL)IJP
R L CRAWFORD (GLAS)IJP
+MITCHELL,MONTGOMERY,+ (SFLA,ALABAMA)IJP
GAIOS MILLER (PURO)IJP
+ROSENFELD,LASINSKI,SMADJA+ (LBL,SLAC)IJP
I. M. BARBOUR,R. L. CRAWFORD (GLAS)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
DONNACHIE, R KIRSOPP (GLAS+EDIN)
BAREYRE,VILLET (SACLAY)
+CASHMERE (U. OXFORD)

Δ(1890) 11 N*3/2(1890, JP=5/2+) I=3/2 F35
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

11 N*3/2(1890) MASS (MEV)
M 3 (1913.0) DONNACH1 68 RVUE PHASE-SHIFT ANAL 8/69
M 6 (1837.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM DAVIDS 70 RVUE P-S ANAL SCL A 8/69
M 4 (1841.0) AYED 70 IPWA 2/72
M 7 (1875.) ALMEHED 72 IPWA 2/72
M 1 1890. TO 1986. MEHTANI 72 DPWA PI+P TO PI+PIOP 9/73
M (1890.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
M (1890.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M (1869.) AYED 74 IPWA 2/74
M L 1870. OR 1830. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
M L THE 2 SETS OF PARAMETERS ARE FROM METHODS 1 AND 2 OF LONGACRE 75. 11/75*

11 N*3/2(1890) WIDTH (MEV)
W 3 (350.0) DONNACH1 68 RVUE 8/69
W 6 (198.0) AYED 70 IPWA 1/71
W 4 (136.0) DAVIDS 70 RVUE SOL A 8/69
W 7 (250.) ALMEHED 72 IPWA 2/72
W 1 273. TO 322. MEHTANI 72 DPWA PI+P TO PI+PIOP 9/73
W (180.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
W (140.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (255.) AYED 74 IPWA 2/74
W L 255. OR 220. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
SEE NOTES ACCOMPANYING MASSES QUOTED AS FOR N*1/2(1910)

11 N*3/2(1890) REAL PART OF POLE POSITION (MEV)
RE (1813.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) 2*IMAG PART OF POLE POSITION (MEV)
IM (193.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) PARTIAL DECAY MODES
P1 N*3/2(1890) INTO PI N 139+ 938
P2 N*3/2(1890) INTO K SIGMA 938+ 139+ 139
P3 N*3/2(1890) INTO K SIGMA 493+1189
P4 N*3/2(1890) INTO N*3/2(1232) PI 1232+ 139
P5 N*3/2(1890) INTO GAM NUCLEON, HELICITY=1/2 0+ 938
P6 N*3/2(1890) INTO GAM NUCLEON, HELICITY=3/2 0+ 938
P7 N*3/2(1890) INTO N RHO 938+ 773
P8 N*3/2(1890) INTO N*3/2(1232) PI,P-WAVE 1232+ 139
P9 N*3/2(1890) INTO N*3/2(1232) PI,P-WAVE 1232+ 139
P10 N*3/2(1890) INTO N RHO,S=3/2,P-WAVE 938+ 773

Baryons

$\Delta(1890)$, $\Delta(1900)$, $\Delta(1910)$

11 N*3/2(1890) BRANCHING RATIOS

R1	N*3/2(1890) INTO (PI N)/TOTAL	(P1)	
R1	3 (0.16)	DCNNACH1 68 RVUE	8/69
R1	6 (0.147)	AYED 70 IPWA	1/71
R1	4 (0.20)	DAVIES 70 RVUE	8/69
R1	7 (0.19)	ALMEHD 72 IPWA	2/72
R1	1 (0.144)	AYED 74 IPWA	2/74
R2	N*3/2(1890) INTO (K SIGMA)/TOTAL	(P3)	
R2	1 (0.08) OR LESS	FEUERBACH 70 RVUE	PI P TO K+ SIG+ 7/70
R2	1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68		
R2	1 MODEL USED MAY DOUBLE COUNT.		
R3	N*3/2(1890) INTO (K SIGMA)*(PI N)/TOTAL**2	(P3*P1)	
R3	3 (0.016) OR LESS	KALMUS 70 DPWA	PI*P TO K+ SIG+ 1/71
R4	N*3/2(1890) FROM PI N TO K SIGMA	SQRT(P1*P3)	9/73
R4	1 (0.19 TO .23)	MEHTANI 72 DPWA	PI*P TO PI*P10P 9/73
R4	1 MOSTLY F WAVE DECAY		
R5	N*3/2(1890) FROM PI N TO K SIGMA	SQRT(P1*P3)	9/73
R5	1 (0.06)	LANGBEIN 73 IPWA	PI N-K SIG,SOL 1 9/73
R5	2 (0.21) TO .054	DEANS 75 DPWA	PI N TO K SIGMA 11/75*
R5	2 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS.		11/75*
R6	N*3/2(1890) FROM PI N TO N RHO,S=3/2,P-WAVE	SQRT(P1*P3)	11/75*
R6	1 (-.12) OR -.20	LONGACRE 75 IPWA	PI N TO 2PI N 11/75*
R7	N*3/2(1890) FROM PI N TO N RHO,S=3/2,P-WAVE	SQRT(P1*P3)	11/75*
R7	1 (-.28) OR -.33	LONGACRE 75 IPWA	PI N TO 2PI N 11/75*

11 N*3/2(1890) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1	N*3/2(1890) INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)			
A1	0.19	.027	DEVENIS2 74 DPWA	PI N PHOTO-PROD 4/75*
A1	.042	.016	KNIES 74 DPWA	PI N PHOTO-PROD 2/74
A1	+.047	.067	METCALF 74 DPWA	PI N PHOTO-PROD 2/74
A1	-.003	.009	CRAWFORD 75 DPWA	PI N PHOTO-PROD 1/76*
A1	(-.035)		BARBOUR 76 DPWA	PI N PHOTO-PROD 1/76*
A1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)			
A2	N*3/2(1890) INTO GAM NUCLEON, HELICITY=3/2 (GEV**1/2)			
A2	0.178	.020	DEVENIS2 74 DPWA	PI N PHOTO-PROD 4/75*
A2	-.022	.020	KNIES 74 DPWA	PI N PHOTO-PROD 2/74
A2	-.028	.066	METCALF 74 DPWA	PI N PHOTO-PROD 2/74
A2	-.021	.036	CRAWFORD 75 DPWA	PI N PHOTO-PROD 1/76*
A2	(-.013)		BARBOUR 76 DPWA	PI N PHOTO-PROD 1/76*
A2	AVERAGE MEANINGLESS (SCALE FACTOR = 2.2)			

REFERENCES FOR N*3/2(1890)

DCNNACH1 68 PL 268 161	A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN) IJP
ALSO 68 VIENNA 139	DCNNACHIE RAPPORTEUR'S TALK (GLAS)
ALSO 68 THESIS	R G KIRSOPP (EDIN)
AYED 70 KIEV CCFN	R AYED, P BAREYRE, G VILLET (SACL) IJP
DAVIES 70 NP 821 359	A DAVIES (GLAS)
FEUERBACH 70 NP 168 85	FEUERBACHER+HOLLADAY (VANDERBILT)
KALMUS 70 PR D2 1824	G KALMUS, G BORREANI, J LOUIE (LRL)
ALMEHD 72 NP 840 157	+LOVELACE (LUND, RUTG) IJP
MEHTANI 72 PRL 29 1634	+FUNG, KERNAN, SCHALK, + (UCR+LBL)
LANGBEIN 73 NP 853 251	LANGBEIN, WAGNER (MUNICH) IJP
AYED 74 PRIVATE COMMCTN.	AYED, BAREYRE (SACL) IJP
ALSO 73 AIX CONFERENCE	AYED, BAREYRE (SACL) IJP
DEVENIS2 74 PL 528 227	DEVENISH, LYTH, RANKIN (DESY, LANG, BONN) IJP
KNIES 74 PRD 9 2680	KNIES, MOORHOUSE, OBERLACK (LBL, GLAS) IJP
METCALF 74 NP 876 253	W J METCALF, R L WALKER (CIT) IJP
CRAWFORD 75 NP 897 125	R L CRAWFORD (GLAS) IJP
DEANS 75 NP 896 90	+MITCHELL, MCNTGOMERY, + (SFLA, ALABAMA) IJP
LONGACRE 75 PL 558 415	+ROSENFELD, LASINSKI, SMADJA + (LBL, SLAC) IJP
BARBOUR 76 SBMTD. TO NP	I. M. BARBOUR, R. L. CRAWFORD (GLAS) IJP
PAPERS NOT REFERRED TO IN DATA CARDS.	
AYED 70 PL 318 598	+BAREYRE+VILLET (SACL) IJP

$\Delta(1900)$ 30 N*3/2(1900, JP=1/2-) I=3/2 S_{31}^+ 9/73

THIS EFFECT IS SEEN IN TWO CHANNELS.

30 N*3/2(1900) MASS (MEV) 9/73

M	(1920.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
M	(1870.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
M	(2001.)	AYED 74 IPWA	2/74

30 N*3/2(1900) WIDTH (MEV) 9/73

W	(140.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
W	(160.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
W	(307.)	AYED 74 IPWA	2/74

SEE THE NOTES ACCOMPANYING MASSES QUOTED

30 N*3/2(1900) PARTIAL DECAY MODES 9/73

P1	N*3/2(1900) INTO PI N	139+ 938
P2	N*3/2(1900) INTO K SIGMA	493+1189

Data Card Listings

For notation, see key at front of Listings.

30 N*3/2(1900) BRANCHING RATIOS 9/73

R1	N*3/2(1900) FROM PI N TO K SIGMA	SQRT(P1*P2)	9/73
R1	3 (0.11)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
R1	6 (0.12)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
R1	1 (0.76)	DEANS 75 DPWA	PI N TO K SIGMA 11/75*
R1	1 VALUE GIVEN IS FROM SOLUTION 1, NOT PRESENT IN SOLUTIONS 2,3,4.		11/75*
R2	N*3/2(1900) INTO PI N/TOTAL	(P1)	
R2	1 (0.76)	AYED 74 IPWA	2/74

REFERENCES FOR N*3/2(1900)

LANGBEIN 73 NP 853 251	LANGBEIN, WAGNER (MUNICH) IJP
AYED 74 PRIVATE COMMCTN.	AYED, BAREYRE (SACL) IJP
ALSC 73 AIX CONFERENCE	AYED, BAREYRE (SACL) IJP
DEANS 75 NP 896 90	+MITCHELL, MCNTGOMERY, + (SFLA, ALABAMA) IJP

$\Delta(1910)$ 12 N*3/2(1910, JP=1/2+) I=3/2 P_{31}^-

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

12 N*3/2(1910) MASS (MEV)

M	3 (1934.0)	DCNNACH1 68 RVUE	PHASE-SHIFT ANAL 8/69
M	6 (1783.0)	AYED 70 IPWA	1/71
M	6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM		
M	4 (194.0)	DAVIES 70 RVUE	P-S ANAL SCL A 8/69
M	7 (1900.)	ALMEHD 72 IPWA	2/72
M	(1980.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
M	(1950.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
M	(1786.)	AYED 74 IPWA	2/74

12 N*3/2(1910) WIDTH (MEV)

W	3 (339.0)	DCNNACH1 68 RVUE	8/69
W	6 (308.0)	AYED 70 IPWA	1/71
W	4 (290.)	DAVIES 70 RVUE	SCL A 8/69
W	7 (200.)	ALMEHD 72 IPWA	493+1189 2/72
W	(190.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
W	(170.)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
W	(222.)	AYED 74 IPWA	2/74

SEE NOTES ACCOMPANYING MASSES QUOTED

12 N*3/2(1910) PARTIAL DECAY MODES

P1	N*3/2(1910) INTO PI N	139+ 938
P2	N*3/2(1910) INTO N PI	938+ 139+ 139
P3	N*3/2(1910) INTO K SIGMA	493+1189
P4	N*3/2(1910) INTO N*3/2(1232) PI	1232+ 139
P5	N*3/2(1910) INTO GAM NUCLEON, HELICITY=1/2	0+ 938
P6	N*3/2(1910) INTO N RHO	938+ 773

12 N*3/2(1910) BRANCHING RATIOS

R1	N*3/2(1910) INTO (PI N)/TOTAL	(P1)	
R1	3 (0.30)	DCNNACH1 68 RVUE	8/69
R1	6 (0.128)	AYED 70 IPWA	1/71
R1	4 (0.18)	DAVIES 70 RVUE	8/69
R1	7 (0.33)	ALMEHD 72 IPWA	2/72
R1	1 (0.158)	AYED 74 IPWA	2/74
R2	N*3/2(1910) INTO (K SIGMA)/TOTAL	(P3)	
R2	1 (0.08) OR LESS	FEUERBACH 70 RVUE	PI P TO K+ SIG+ 7/70
R2	1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68		
R2	1 MODEL USED MAY DOUBLE COUNT.		
R3	N*3/2(1910) FROM PI N TO K SIGMA	SQRT(P1*P3)	9/73
R3	1 (0.11)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 1 9/73
R3	2 (0.15)	LANGBEIN 73 IPWA	PI N-K SIG, SOL 2 9/73
R3	2 (0.82) TO .184	DEANS 75 DPWA	PI N TO K SIGMA 11/75*
R3	2 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS.		11/75*

12 N*3/2(1910) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1	N*3/2(1910) INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)			
A1	0.00	.025	DEVENIS2 74 DPWA	PI N PHOTO-PROD 4/75*
A1	0.10	.012	KNIES 74 DPWA	PI N PHOTO-PROD 2/74
A1	-.032	.065	METCALF 74 DPWA	PI N PHOTO-PROD 2/74
A1	-.009	.013	CRAWFORD 75 DPWA	PI N PHOTO-PROD 1/76*
A1	(-.031)		BARBOUR 76 DPWA	PI N PHOTO-PROD 1/76*
A1	AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)			

REFERENCES FOR N*3/2(1910)

DCNNACH1 68 PL 268 161	A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN) IJP
ALSO 68 VIENNA 139	DCNNACHIE RAPPORTEUR'S TALK (GLAS)
ALSO 68 THESIS	R G KIRSOPP (EDIN)
AYED 70 KIEV CCFN	R AYED, P BAREYRE, G VILLET (SACL) IJP
DAVIES 70 NP 821 359	A DAVIES (GLAS)
FEUERBACH 70 NP 168 85	FEUERBACHER+HOLLADAY (VANDERBILT)
ALMEHD 72 NP 840 157	+LOVELACE (LUND, RUTG) IJP
LANGBEIN 73 NP 853 251	LANGBEIN, WAGNER (MUNICH) IJP
AYED 74 PRIVATE COMMCTN.	AYED, BAREYRE (SACL) IJP
ALSO 73 AIX CONFERENCE	AYED, BAREYRE (SACL) IJP
DEVENIS2 74 PL 528 227	DEVENISH, LYTH, RANKIN (DESY, LANG, BONN) IJP
KNIES 74 PRD 9 2680	KNIES, MOORHOUSE, OBERLACK (LBL, GLAS) IJP
METCALF 74 NP 876 253	W J METCALF, R L WALKER (CIT) IJP

Data Card Listings

For notation, see key at front of Listings.

Baryons

Δ(1910), Δ(1950), Δ(1960)

CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
DEANS 75 NP 896 90 +MITCHELL,MCNTGCMERY,+ (SFLA,ALABAMA)IJP
BARBOUR 76 SBMTD. TO NP I. M. BARBOUR,R. L. CRAWFCRD (GLAS)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
CARYANN 65 PR 138 8433 CARAYANNOPOULOS,TAUTFEST,WILLMANN (PURD)
A PARTIAL WAVE ANALYSIS OF PI+P TO SIGMA+ K+ (PURD)
AYED 70 PL 318 958 +BAREYRE+VILLET (SACLAY)

Δ(1950)

83 N*3/2(1950, JP=7/2+) I=3/2 F37
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

83 N*3/2(1950) MASS (MEV)
M (1920.0) DUKE 65 CNTR PI-P EL + PCL 7/66
M (1950.0) APPROX YOKOSAWA 66 CNTR PI- P DSIG + PCL 7/66
M 1 (1975.0) BAREYRE 68 RVUE PHASE-SHIFT ANAL 11/67
M 3 (1946.0) WHERE CROSS SECTION IS GREATEST - EYEBALL FIT 6/68
M 6 (1931.0) DONNACHI 68 RVUE PHASE-SHIFT ANAL 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM
M 4 (1935.0) DAVIES 70 RVUE P-S ANAL SOL A 8/69
M (1950.0) (30.0) KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
M (1920.0) ROYCHOUD 71 DPWA 3/72
M 7 (1925.0) ALMEHED 72 IPWA 2/72
M 2 1920 TO 1951 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
M (1928.0) AYED 74 IPWA 2/74
M L 1930. OR 1925. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
M L THE 2 SETS OF PARAMETERS ARE FROM METHODS 1 AND 2 OF LONGACRE 75. 11/75*

83 N*3/2(1950) WIDTH (MEV)
W (170.0) DUKE 65 CNTR 7/66
W (200.0) APPROX YOKOSAWA 66 CNTR 7/66
W 1 (190.0) BAREYRE 68 RVUE 11/67
W 3 (221.0) DONNACHI 68 RVUE 6/68
W 6 (197.0) AYED 70 IPWA 1/71
W 4 (221.0) DAVIES 70 RVUE SOL A 8/69
W (300.0) (60.0) KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
W 7 (200.0) ALMEHED 72 IPWA 2/72
W 2 234 TO 269 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
W (237.0) AYED 74 IPWA 2/74
W L 235. OR 240. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

83 N*3/2(1950) REAL PART OF POLE POSITION (MEV) 11/75*
RE (1924.0) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

83 N*3/2(1950) 2*IMAG PART OF POLE POSITION (MEV) 11/75*
IM (258.0) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

83 N*3/2(1950) PARTIAL DECAY MODES
P1 N*3/2(1950) INTO PI N 139+ 938
P2 N*3/2(1950) INTO SIGMA K 1189+ 493
P3 N*3/2(1950) INTO N*3/2(1232) PI 1232+ 139
P4 N*3/2(1950) INTO Y*1(1365) K 1384+ 493
P5 N*3/2(1950) INTO N*3/2(1232) RHO 1232+ 773
P6 N*3/2(1950) INTO NEUTRON PI+ PI+ 939+ 139+ 139
P7 N*3/2(1950) INTO N*3/2(1232) PI PI (NOT RHO) 1232+ 938
P8 N*3/2(1950) INTO GAM NUCLEON, HELICITY=1/2 0+ 938
P9 N*3/2(1950) INTO GAM NUCLEON, HELICITY=3/2 0+ 938
P10 N*3/2(1950) INTO N RHO 938+ 773
P11 N*3/2(1950) INTO N*3/2(1232) PI, F-WAVE 1232+ 139
P12 N*3/2(1950) INTO N*3/2(1232) PI, H-WAVE 1232+ 139
P13 N*3/2(1950) INTO N RHO, S=3/2, F-WAVE 938+ 773

83 N*3/2(1950) BRANCHING RATIOS
R1 N*3/2(1950) INTO (PI N)/TOTAL (P1)
R1 (0.41) DUKE 65 CNTR VERY ENERGY DEP 7/66
R1 (0.4) APPROX YOKOSAWA 66 CNTR 7/66
R1 1 (0.57) BAREYRE 68 RVUE 11/67
R1 3 (0.386) DONNACHI 68 RVUE 6/68
R1 6 (0.496) AYED 70 IPWA 1/71
R1 4 (0.51) DAVIES 70 RVUE SOL A 8/69
R1 7 (0.4) ALMEHED 72 IPWA 2/72
R1 (0.405) AYED 74 IPWA 2/74
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

R2 N*3/2(1950) INTO (SIGMA K)+(PI N)/TOTAL**2 (P2*P1)
R2 (0.004) (0.008) BORREANI 68 HBC PI+P 1.35-1.68 10/69
R2 1 (0.004) (0.008) FEUERBACH 70 RVUE PI P TO K+ SIG+ 7/70
R2 1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68
R2 1 MODEL USED MAY DOUBLE COUNT.
R2 0.0081 0.0013 KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
R3 N*3/2(1950) FROM PI N TO N*3/2(1232) PI SQRT(PI*P3)
R3 2 .37 TO .48 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
R3 2 MOSTLY F WAVE DECAY 9/73
R4 N*3/2(1950) FROM PI N TO SIGMA K SQRT(PI*P2)
R4 (.04) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
R4 (.05) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
R4 5 (1.02) DEANS 75 IPWA PI N TO K SIGMA 11/75*
R4 5 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS.. 11/75*

R5 N*3/2(1950) FROM PI N TO N*3/2(1232) PI, F-WAVE SQRT(PI*P11) 11/75*
R5 L (-.2510R -.32 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

R6 N*3/2(1950) FROM PI N TO N RHO, S=3/2, F-WAVE SQRT(PI*P13) 11/75*
R6 L (-.1810R -.24 LONGACRE 75 IPWA PI N TO 2PI N 11/75*
MORE INFORMATION ON INELASTIC DECAY MODES OF BUMPS, SEEN IN PRODUCTION EXPERIMENTS AROUND 1950 MEV, MAY BE FOUND IN THE NEXT ENTRY

83 N*3/2(1950) PHOTON DECAY AMPL(GEV**1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*3/2(1950) INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)
A1 -.088 .025 DEVENIS2 74 DPWA PI N PHOTO-PRD 4/75*
A1 -.070 .012 KNIES 74 DPWA PI N PHOTO-PRD 2/74
A1 -.055 .029 METCALF 74 DPWA PI N PHOTO-PRD 2/74
A1 (-.080) MODRHOUS 74 DPWA PI N PHOTO-PRD 2/74
A2 -.038 .014 CRAWFORD 75 DPWA PI N PHOTO-PRD 1/76*
A2 (-.065) BARBOUR 76 DPWA PI N PHOTO-PRD 1/76*
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)
A2 N*3/2(1950) INTO GAM NUCLEON, HELICITY=3/2 (GEV**1/2)
A2 -.080 .021 DEVENIS2 74 DPWA PI N PHOTO-PRD 4/75*
A2 -.078 .010 KNIES 74 DPWA PI N PHOTO-PRD 2/74
A2 -.093 .024 METCALF 74 DPWA PI N PHOTO-PRD 2/74
A2 (-.180) MODRHOUS 74 DPWA PI N PHOTO-PRD 2/74
A2 -.038 .014 CRAWFORD 75 DPWA PI N PHOTO-PRD 1/76*
A2 (-.065) BARBOUR 76 DPWA PI N PHOTO-PRD 1/76*
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.5)

REFERENCES FOR N*3/2(1950)
DUKE 65 PRL 15 468 +JONES,KEMP,MURPHY,PRENTICE, + (RHEL,OXF)IJP
YOKOSAWA 66 PRL 16 714 +SUWA, FILL, ESTERLING, BOOTH (ANL,CHIC)IJP

BAREYRE 68 PR 165 1731 P BAREYRE, C BRICMAN, G VILLET (SACLAY)IJP
BORREANI 68 UCL 18350 BORREANI, KALMUS (LRL)
DONNACHI 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
ALSO 68 VIENNA 139 DONNACHIE, RAPPORTEUR'S TALK (GLAS)
ALSO 68 THESIS R G KIRSOPP (EDIN)

AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACL)IJP
DAVIES 70 NP 821 359 A DAVIES (GLAS)
FEUERBACH 70 NP 168 85 FEUERBACHER+HCLLDAY (VANDERBILT)
KALMUS 70 PR D2 1824 G KALMUS, G BORREANI, J LOUIE (LRL)
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSDEN (DURH)IJP

ALMEHED 72 NP 840 157 +LCVELACE (LUND,RUTG)IJP
MEHTANI 72 PRL 29 1634 +FUNG, KERNAN, SCHALK, + (UCR +LBL)
LANGBEIN 73 NP 853 251 LANGBEIN,WAGNER (MUNICH)IJP

AYED 74 PRIVATE COMMCTN. AYED,BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP
DEVENIS2 74 PL 528 227 DEVENIS,LYTH,RANKIN (DESY,LANC,BCN)IJP
KNIES 74 PRD 9 2680 KNIES,MODRHOUS,OVERLACK (LBL,GLAS)IJP
METCALF 74 NP 874 253 W J METCALF,R L WALKER (CIT)IJP
MODRHOUS 74 PRD 9 1 MODRHOUS,OVERLACK,ROSENFELD (GLAS+LBL)IJP

CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
DEANS 75 NP 896 90 +MITCHELL,MCNTGCMERY,+ (SFLA,ALABAMA)IJP
LONGACRE 75 PL 558 415 +ROSENFELD,LASINSKI,SMADJA+ (LBL,SLAC)IJP
BARBOUR 76 SBMTD. TO NP I. M. BARBOUR,R. L. CRAWFCRD (GLAS)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.

HCHLER 63 NP 48 470 G HOHLER, G EBEL (KARLSRUHE) I
LAYSON 63 NC 27 724 W H LAYSON (CERN) IJ
AUVIL 64 NC 33 473 P AUVIL, C LOVELACE (LOIC)IJP
HELLAND 64 PR 134 81062 +DEVLIN,HAGGE,LCNGO,MOYER,WOOD (LRL) IJ
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
HOLLADAY 65 PR 139 81348 W G HOLLADAY (VANDERBILT)
JOHNSON 67 UCL-17683 THESIS C H JOHNSON (LRL)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

Δ(1960)

13 N*3/2(1960, JP=5/2-) I=3/2 D35
THIS EFFECT IS SEEN IN TWO CHANNELS.

13 N*3/2(1960) MASS (MEV)
M 3 (1954.0) DONNACHI 68 RVUE PHASE-SHIFT ANAL 6/68
M 3 (1970.0) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
M X (1950.0) APPROX LEA 69 CNTR PI-P ELASTIC 8/69
M X SEE ALSO APLIN 70
M 3 WHERE MAX. ABSORPTION IS -DONNACHI, 2, KIRSOPP EYEBALL FIT CERN 1 10/69
M 7 (2200.0) ALMEHED 72 IPWA 2/72
M 1 (1560.0) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M 1 NOT SEEN IN SOLUTION 1 OF LANGBEIN73 9/73
M (1889.0) AYED 74 IPWA 2/74

13 N*3/2(1960) WIDTH (MEV)
W 3 (311.00) DONNACHI 68 RVUE 8/69
W 3 (400.0) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
W 7 (600.0) ALMEHED 72 IPWA 2/72
W 1 (150.0) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (1121.0) AYED 74 IPWA 2/74
SEE THE NOTES ACCOMPANYING MASSES QUOTED

Baryons

$\Delta(1960)$, $\Delta(2160)$

Data Card Listings

For notation, see key at front of Listings.

Table with columns for decay modes (P1, P2), branching ratios (R1-R7), and decay masses (139+938, 493+1189). Includes sub-sections for partial decay modes and branching ratios.

13 N*3/2(1960) PHOTON DECAY AMPL(GEV**1/2) FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for decay modes (A1, A2) and branching ratios (1/76*, 1/76*). Includes sub-sections for gamma nucleon decay amplitudes.

REFERENCES FOR N*3/2(1960)

Table of references for N*3/2(1960) including authors like DONNACHI, LEA, LANGBEIN, CRAWFORD, BARBOUR, DEANS, and APLIN.

1950 MEV REGION - PRODUCTION AND σ_{TOTAL} EXP'TS

70 N*3/2(1950, JP= 1 I=3/2 PRODUCTION EXPERIMENTS

70 N*3/2(1950) MASS (MEV) (PROD. EXP.)

Table of production experiments for N*3/2(1950) showing mass (MEV) and total cross-sections.

70 N*3/2(1950) WIDTH (MEV) (PROD. EXP.)

Table of production experiments for N*3/2(1950) showing width (MEV) and total cross-sections.

70 N*3/2(1950) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for N*3/2(1950) including decay masses and branching ratios.

70 N*3/2(1950) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for N*3/2(1950) including decay masses and branching ratios.

Table with columns for decay modes (R3, R4, R5, R6, R7, R8) and branching ratios (0.05, 0.013, 0.035, 0.015, 0.451, 0.551). Includes sub-sections for partial decay modes and branching ratios.

REFERENCES FOR N*3/2(1950) (PROD. EXP.)

Table of references for N*3/2(1950) including authors like COOL, BRISSON, DEVLIN, LEE, CHINOWSKY, GALLOWAY, COLTEN, COLLEY, BRAUNZ, CHUNG, and GAIDOS.

PAPERS NOT REFERRED TO IN DATA CARDS.

Table of references for N*3/2(1950) including authors like DEUTSCHMANN, AACH, BONN, BERL, CERN, CRAC, and EID.

$\Delta(2160)$

9 N*3/2(2160,) I=3/2

EARLY ANALYSES FOUND EVIDENCE FOR A RESONANCE NEAR THIS MASS IN THE P33 PARTIAL WAVE, AND UNLESS STATED OTHERWISE, ALL DATA CARDS BELOW APPLY TO THIS WAVE. IN ADDITION, ROYCHOUDHURY 71 FIND POSSIBLE EVIDENCE FOR P31, D33, AND D35 RESONANCES IN THIS MASS REGION. IN A SIMILAR ANALYSIS BRANDSEN 71 FOUND SOME EVIDENCE FOR S31, D33, AND D35 RESONANCES IN THIS REGION. VON SCHLIPPE 72 SUGGESTS A G39. THE MOST RECENT P1 N ANALYSIS FINDS A G39 RESONANCE IN THIS MASS REGION. THE EFFECT SEEN IN K SIG PROD. IS 100 MEV LOWER IN MASS. A PRONOUNCED SHARP DIP IS OBSERVED IN P1+ P BACKWARD SCATTERING AT 2200 MEV BY REY 74. DUAL INTERFERENCE MODEL ANALYSIS OF MA 75 FINDS SIGNAL FOR P33, P31, AND D35, BUT NOT FOR G39.

9 N*3/2(2160) MASS (MEV)

Table of mass (MEV) for N*3/2(2160) including authors like M 3, M 7, M 1, M 2, M 4, and M 4.

9 N*3/2(2160) WIDTH (MEV)

Table of width (MEV) for N*3/2(2160) including authors like W 3, W 7, W 1, and W 4.

9 N*3/2(2160) PARTIAL DECAY MODES

Table of partial decay modes for N*3/2(2160) including decay masses and branching ratios.

9 N*3/2(2160) BRANCHING RATIOS

Table of branching ratios for N*3/2(2160) including decay masses and branching ratios.

REFERENCES FOR N*3/2(2160)

Table of references for N*3/2(2160) including authors like KIRSOPP, ROYCHOUDHURY, ALMEHEID, LANGBEIN, and LANGBEIN.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Δ(2160), Δ(2420), Δ(2850)

AYED 74 PRIVATE COMMCTN. AYED, BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE (SACL)IJP
REY 74 PRD 32 908 REV, LENNOX, POIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 74 PRL 33 250 REV, LENNOX, PCIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 75 PRD 11 1777 LENNOX, POIRIER, REY, SANDER+ (NDAM+FNAL+ANL)IJP
DEANS 75 NP 896 90 +MITCHELL, MCNTGMEY,+ (SFLA, ALABAMA)IJP

BRANDEN 71 NP 826 511 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
ALSO 70 NP 816 461 +CCEN (DURH)IJP
VON SCHL 72 LNC 4 767 ROYCHOUDHURY, PERRIN, BRANDEN (DURH)IJP
BAKER 74 PRL 32 251 VON SCHLIPPE (LONC)IJP
MA 75 PRD 11 1852 BAKER, EARTLY, PRETAL, PRLSS,+ (FNAL, ANL, NDAM)IJP
MA, SHAW (UCSB, SLAC)IJP

PAPERS NOT REFERRED TO IN DATA CARDS.

Δ(2420)

84 N*3/2(2420, JP=11/2+) I=3/2 H3 11
BOTH ROYCHOUDHURY 71 AND BRANDEN 71 SEE A POSSIBLE
RESONANT F35 IN THIS MASS REGION. IN ADDITION BRANDEN
71 FIND A RESONANT P33 AT 2600 MEV.

84 N*3/2(2420) MASS (MEV) (PROD. EXP.)
M 6 (2312.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM
(2400.) BRANSDEN 71 CPWA 3/72
M (2400.) ROYCHOUD 71 DPWA 3/72
M (2440.) OTT 72 MPWA 0 PI-P BKWD ELSTC 2/73
M (2392.) AYED 74 IPWA 2/74
M 1 (2404.) (63.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

84 N*3/2(2420) WIDTH (MEV)
W 6 (347.0) AYED 70 IPWA 1/71
W (289.) (44.) OTT 72 MPWA 2/74
W 1 (484.) (75.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

84 N*3/2(2420) PARTIAL DECAY MODES
P1 N*3/2(2420) INTO PI N DECAY MASSES 139+ 938
P2 N*3/2(2420) INTO SIGMA K 1197+ 493

84 N*3/2(2420) BRANCHING RATIOS
R1 N*3/2(2420) INTO (PI N)/TOTAL (P1)
R1 6 (0.113) AYED 70 IPWA 1/71
R1 7 (4.) OTT 72 MPWA 0 PI-P BKWD ELSTC 2/73
R1 (1.09) AYED 74 IPWA 2/74
R1 1 (.157) (.070) (.035) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACL)IJP
BRANDEN 71 NP 826 511 +CCEN (DURH)IJP
ALSO 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANDEN (DURH)IJP
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANDEN (DURH)IJP
OTT 72 PL 428 133 +TRISCHUK, VAVRA, RICHARDS,+ (MCGI, STLO, IOWA)IJP
ALSO 72 MCGILL THESIS J. VAVRA (MCGI) IJP
AYED 74 PRIVATE COMMCTN. AYED, BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)IJP
REY 74 PRD 32 908 REV, LENNOX, POIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 74 PRL 33 250 REV, LENNOX, POIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 75 PRD 11 1777 LENNOX, POIRIER, REY, SANDER+ (NDAM+FNAL+ANL)IJP

BELLAMY 67 PRL 19 476 +BUCKLEY, DOBINSN,+ (WESTFIELD, LOUC) JP
AYED 70 PL 318 598 +BAREYRE+VILLET (SACL)IJP

PAPERS NOT REFERRED TO IN DATA CARDS.

2420 MEV REGION - PRODUCTION AND σTOTAL EXP'TS

69 N*3/2(2420, JP=) I=3/2 PRODUCTION EXPERIMENTS
69 N*3/2(2420) MASS (MEV) (PROD. EXP.)
M (2360.0) DIDDENS 63 CNTR PI+ P TOTAL
M (2520.0) (40.0) ALVAREZ 64 CNTR PI PHOTOPROD 7/66
M (2440.0) HOHLER 64 RVUE DATA P, DISP REL
M (2400.0) APPROX WAHLIG 64 OSPK 0 PI-P CH EX
M B (2452.0) BARGER 66 RVUE TOTAL + CH EX 11/67
M B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE
M B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
M 2423.0 10.0 CITRON 66 CNTR PI+ P TOTAL 7/66

69 N*3/2(2420) WIDTH (MEV) (PROD. EXP.)
W (200.0) DIDDENS 63 CNTR 7/66
W (245.0) HOHLER 64 RVUE 7/66
W B (275.0) BARGER 66 RVUE TOTAL + CH EX 11/67
W 310.0 20.0 CITRON 66 CNTR 7/66

69 N*3/2(2420) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*3/2(2420) INTO PI N DECAY MASSES 139+ 938
P2 N*3/2(2420) INTO SIGMA K 1197+ 493
P3 N*3/2(2420) INTO N*3/2(1232) PI 1232+ 139
P4 N*3/2(2420) INTO NEUTRON PI+ PI+ 939+ 139+ 139

69 N*3/2(2420) BRANCHING RATIOS (PRGD. EXP.)
R1 N*3/2(2420) INTO (PI N)/TOTAL (P1)
R1 ALVAREZ 64 PRL 12 710 APPROX DIDDENS 63 CNTR ASSUMING J=11/2 7/66
R1 0.113 0.0036 CITRON 66 CNTR ASSUMING J=11/2 7/66
R1 B (0.12) BARGER 67 FIT ASSUMING J=11/2 11/67
R1 D (0.163) DIKMAN 67 FIT ASSUMING J=11/2 11/67
R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (0.06) KORMANYOS 67 CNTR ASSUMING J=11/2 11/67
R2 N*3/2(2420) INTO (PI N)*(NEUTRON PI+ PI+)/TOTAL**2 (P1*P4)
R2 0.0195 0.0048 GALLOWAY 68 RVUE 6/68

REFERENCES FOR N*3/2(2420) (PROD. EXP.)
DIDDENS 63 PRL 10 262 +JENKINS, KYCIA, RILEY (BNL) I
ALVAREZ 64 PRL 12 710 +BAR-YAM, KERN, LUCKEY, CSBORNE, + (MIT, CEA)
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)
BARGER 66 PR 151 1123 V BARGER, M CLSSON (WISC)
CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMAN 67 PR 18 798 F N DIKMAN (MICH)
KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P
GALLOWAY 68 PL 268 334 K F GALLOWAY (INDIANA) I

PAPERS NOT REFERRED TO IN DATA CARDS.
BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE, ORSAY) J-L
DOBROWOL 67 PL 248 203 DOBROWOLSKI, GUSKOV, LIKHACHEV, + (DUBNA) P
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)
WAHLIG 68 PR 168 1515 M A WAHLIG, I MANNELLI (MIT, PISA)
FINAL VERSION OF DATA USED IN WAHLIG 64+. IN CONJUNCTION WITH
CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

Δ(2850) BUMPS

85 N*3/2(2850, JP=) I=3/2 PRODUCTION EXPERIMENTS
85 N*3/2(2850) MASS (MEV) (PROD. EXP.)
M (2870.0) HOHLER 64 RVUE DATA + DISP REL
M (2700.0) APPROX WAHLIG 64 OSPK 0 PI-P CH EX
M (2850.0) BARDADIN 66 HBC ++ N* TO P + 3 PIS 7/66
M 2850.0 12.0 CITRON 66 CNTR PI+ P TOTAL 7/66
M (2883.1) (26.) (28.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

85 N*3/2(2850) WIDTH (MEV) (PROD. EXP.)
W (150.0) BARDADIN 66 HBC ++ 7/66
W 400.0 40.0 CITRON 66 CNTR 7/66
W (380.) (141.) (201.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

85 N*3/2(2850) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*3/2(2850) INTO PI N DECAY MASSES 139+ 938
P2 N*3/2(2850) INTO P PI PI 938+ 139+ 139+ 139
P3 N*3/2(2850) INTO N PI PI 938+ 139+ 139

85 N*3/2(2850) BRANCHING RATIOS (PROD. EXP.)
R1 N*3/2(2850) INTO (PI N)/TOTAL (P1)
R1 ONLY (J+1/2) (PI N)/TOTAL MEASURED FOR THIS STATE
R1 B (0.224) (0.016) BARGER 66 RVUE TOTAL + CH EX. 11/67
R1 0.261 0.048 CITRON 66 CNTR TOTAL CROSS-SEC. 11/67
R1 B (0.40) BARGER 67 RVUE USES KORMANYOS66 11/67
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 C (0.49) DIKMAN 67 RVUE USES KORMANYOS67 11/67
R1 C USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 DOBROWOLSKI 67 CNTR PI+P AT 180 DEG
R1 (0.10) KORMANYOS 67 CNTR PI-P AT 180 DEG. 11/67
R1 D (0.06) OR LESS CL=-.95 HALDORSE 72 HBC PP 19 GEV/C 12/72
R1 D UPPER LIMIT CN ELASTICITY. ALSO FIND J=9/2 OR MORE.
R1 (.28) (.13) (.19) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

REFERENCES FOR N*3/2(2850) (PROD. EXP.)
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)
BARDADIN 66 PL 21 357 BARDADIN-OTWINSKA, CANYSZ, + (WARSAW)
BARGER 66 PR 151 1123 V BARGER, M CLSSON (WISC)
CITRON 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMAN 67 PR 18 798 F N DIKMAN (MICH)
DOBROWOL 67 PL 248 203 DOBROWOLSKI, GUSKOV, LIKHACHEV, + (DUBNA) P
KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P
HALDORSE 72 NC 10A 468 HALDORSEN, JACOBSEN (OSLO) IJ
REY 74 PR 32 908 REV, LENNOX, POIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 74 PRL 33 250 REV, LENNOX, POIRIER, PRETZL (NDAM+MPIM)IJP
ALSO 75 PRD 11 1777 LENNOX, POIRIER, REY, SANDER+ (NDAM+FNAL+ANL)IJP

Baryons

$\Delta(2850)$, $\Delta(3230)$, EXOTIC NUCLEONS, Z^* 's

Data Card Listings

For notation, see key at front of Listings.

PAPERS NOT REFERRED TO IN DATA CARDS.

BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE,ORSAYJJ-L
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)
WAHLIG 68 PR 168 1515 M A WAHLIG, I MANNELLI (MIT,PISA)
FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH
CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

$\Delta(3230)$
BUMPS

86 N*3/2(3230, JP=) I=3/2 PRODUCTION EXPERIMENTS

86 N*3/2(3230) MASS (MEV) (PRCD. EXP.)

M (3230.0) CITRDN 66 CNTR PI+ P TOTAL 7/66
(3296.) (79.) (78.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

86 N*3/2(3230) WIDTH (MEV) (PROD. EXP.)

W (440.0) CITRCA 66 CNTR 7/66
W (687.) (1043.) (323.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

86 N*3/2(3230) PARTIAL DECAY MODES (PROD. EXP.)

DECAY MASSES
P1 N*3/2(3230) INTO PI N 139+ 938
P2 N*3/2(3230) INTO N PI PI 938+ 139+ 139

86 N*3/2(3230) BRANCHING RATIOS

R1 N*3/2(3230) INTO (PI N)/TOTAL (PI)
R1 ONLY (J+1/2)* (PI N)/TOTAL MEASURED FOR THIS STATE
R1 B (0.03) (0.01) BARGER 66 RVUE TOTAL + CH EXC. 11/67
CITRDN 66 CNTR TOTAL CROSS. SEC. 11/67
R1 B (0.06) TO 0.1 CITRDN 66 CNTR USES KORMANYOS66 11/67
BARGER 67 CNTR USES KORMANYOS66 11/67
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 D (0.25) DIKMAN 67 RVUE USES KORMANYOS67 11/67
R1 C USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (.45) (.09) (.13) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

REFERENCES FOR N*3/2(3230) (PROD. EXP.)

BARGER 66 PR 151 1123 V BARGER, M CLSSON (WISC)
CITRDN 66 PR 144 1101 +GALBRAITH,KYCIA,LEONTIC,PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMAN 67 PRL 18 798 F N DIKMAN (MICH)
REY 74 PRL 32 908 REV,LENNOX,POIRIER,PRETZL (NDAM+MPI)IP
ALSO 74 PRL 33 250 REV,LENNOX,POIRIER,PRETZL (NCAM+MPI)IP
ALSC 75 PRD 11 1777 LENNCX,POIRIER,REV,SANDER* (NDAM+FNAL+ANL)IP

PAPERS NOT REFERRED TO IN DATA CARDS

KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)

EXOTIC NUCLEONS - 1640 MEV REGION

THIS IS NOT A COMPLETE LIST. WE WILL TABULATE EXOTICS FROM NOW ON

EX(1640, JP=) I=5/2

THIS IS NOT A COMPLETE LIST. WE TABULATE ONLY FROM 1970 ON.

IN A MISSING MASS EXPERIMENT, PI+ P TO PI- X+*,
BIRULEV 71 FIND NO EVIDENCE FOR EXOTIC (I=5/2) RESONANCES IN THE
MASS INTERVAL 1.2 TO 2.2 GEV.

EX(1640) MASS (MEV)

M A 29 1627. 12. PRICE 70 DBC -- K-D AT 4.91GEV/C 3/71
M A FOUR S. D. EFFECT

EX(1640) WIDTH (MEV)

W B 29 30. OR LESS CL=.90 PRICE 70 DBC -- PI-PI-N BUMP 3/71
W B CROSS SECTION 13.0+-3.9 MICROBARNS

EX(1640) CROSS SECTION LIMITS (MICROBARNS)

CS B 40. OR LESS BANNER 70 DSPK +++ PI+P,1.9 GEV/C 7/70
CS B I=5/2 LIMIT GIVEN ABOVE IS FOR MASS RANGE 1540-1750 MEV

REFERENCES FOR EX(1640)

BANNER 70 NP 815 205 +CHEZE,HAMEL,TEIGER,ZACCONA + (SACLAY)
PRICE 70 PL 338,595 +BERG,SALANT,WATERS,WEBSTER,WEINBERG (VAND)

PAPERS NOT REFERRED TO IN DATA CARDS

AMMANN 71 PL 348 533 +CARMONY,GARFINKEL,GLTAY,MILLER,YEN (PURD)
BIRULEV 71 SJNP 12 536 +VOVENKO,GUSKOV,DJBROVOLSKII,++ (JINR)
JOHNSON 71 PL 348 428 D JOHNSON (ANL)

Note on the S = +1 Baryon System

Models based on states of three quarks have been successful in the description of the spectrum and decay of the known baryon resonances. Three quarks cannot produce S = +1 baryon resonances (Z*s), and this has probably been the primary motivation for the great amount of experimental effort that has gone into S = +1 baryon physics below ~ 2 GeV/c during the last several years. In addition, the S = +1 system offers an opportunity to study single and double pion production in a small number of rather distinct quasi-two-body channels (K Δ , K*N, and K* Δ). Recent developments in the field are summarized below.

I = 1 System

The most recent K+p total cross-section data are from BNL (CARROLL 73; 0.41 to 1.06 GeV/c) and Arizona-Michigan (BOWEN 73; 0.57 to 1.16 GeV/c). Neither measurement exhibits the dip of approximately 1.5 mb at 0.7 GeV/c reported by BUGG 68 and BOWEN 70. Recent differential elastic cross-section measurements (including extensions and final results of previously reported results) have come from Glasgow-Bologna-Trieste (CAMERON 74; 0.13-0.76 GeV/c), Birmingham-RHEL (ADAMS 73; 0.43 to 0.94 GeV/c), Bristol-RHEL-Aarhus-Southampton (CHARLES 73; 0.7 to 1.9 GeV/c), Maryland-ANL (ABE 75; 0.86 to 2.12 GeV/c), Washington (ADAMS 75; 1.0-1.5 GeV/c at 180°), UCL-RHEL (BARBER 73; 1.37-2.26 GeV/c), Yale-BNL (PATTON 75; 1.7-3.0 GeV/c near 0° and 180°), and ANL (YUTA 74; 2.53, 2.76, and 3.20 GeV/c). Below 750 MeV/c (where the inelastic cross section is less than 0.1 mb), the total elastic cross sections from the first three above experiments also fail to exhibit the dip seen in earlier total cross-section experiments. Coulomb interference effects are observed by CAMERON 74 and ADAMS 73; both confirm the early conclusion of GOLDBABER 62 that the low energy S-wave interaction is repulsive. The most recent polarization measurements are those of Maryland-ANL-Northwestern-FNAL (BARNETT 73; 1.37-2.31 GeV/c) and Yale-BNL (PATTON 75; 1.7-3.0 GeV/c near 0° and 180°). Recent studies of pion production in K+p interactions are those of CERN-Saclay (BERTHON 73; KN π , 1.21-1.69 GeV/c), CIT-UCLA (LOKEN 72; KN π and KN $\pi\pi$, 1.37-2.17 GeV/c),

Data Card Listings

For notation, see key at front of Listings.

Baryons

Z*'s

France-Saclay-IC-Westfield (BRUNET 72 and LEWIS 73; $KN\pi$, $KN\pi\pi$, and $KN\pi\pi\pi$, 2.11-2.72 GeV/c), and ANL (MUSGRAVE 75; $KN\pi$, 2.53-3.20 GeV/c). The general conclusion is that most pion production is associated with $\Delta(1232)$ and/or $K^*(892)$ production, though other resonances are also produced at the higher energies. The $I=1$ total and partial cross sections are shown as dashed lines in Fig. 1.

Five partial-wave analyses of K^+p elastic scattering have been carried out since 1973. CAMERON 74 performed an elastic energy-independent analysis of data at momenta up to 870 MeV/c, including their new high-statistics cross-section data below 755 MeV/c. S and P waves were kept up to 400 MeV/c with D waves being added at higher momenta, and solutions at different energies were linked by the shortest-path method, yielding a unique overall solution. An energy-dependent fit, with an effective-range parametrization for the S wave and a zero-effective-range scattering-length parametrization for P and D waves, was also performed and agreed well with the energy-independent analysis. The χ^2 per data point (χ^2/ND) was well below 1.0 at most momenta in both fits. Threshold parameters were extracted, the dominant threshold effect being a repulsive S wave. The energy-dependent ADAMS 73 analysis included data at momenta up to 1.5 GeV/c, but was aimed primarily at determining the amplitudes below 1.0 GeV/c. Below about 0.5 GeV/c a zero-effective-range scattering-length parametrization was used for S and P waves, with D and F waves set to zero; at higher energies the η and δ parameters of the S, P, D, and F waves were taken to be polynomials in c.m. momentum. Waves with $l \geq 4$ were set to the values calculated from two- and three-pion exchange by ALCOCK 73. The best χ^2 values obtained were about 2250 for 1600 degrees of freedom, of which an estimated 300 comes from inconsistent data. Two solutions were found; they are essentially identical below 1 GeV/c but differ significantly at higher energies. MARTIN 75 have performed an energy-dependent analysis of K^+p scattering up to 1.5 GeV/c as part of a simultaneous fit of $I=1$ and $I=0$ KN scattering. Below the inelastic threshold the inverse amplitudes for waves with $J \leq 5/2$ were parametrized as rational functions of the c.m. momentum which obeyed elastic

unitarity and had correct elastic threshold behavior.

To allow for inelasticity a (fitted) negative term was added to the imaginary part of the inverse amplitudes above the inelastic threshold. The F_{17} , G_{17} , and G_{19} waves were also kept, and set to the values calculated by ALCOCK 73. Two similar solutions were found with $\chi^2/ND = 1.1$. ARNDT 74

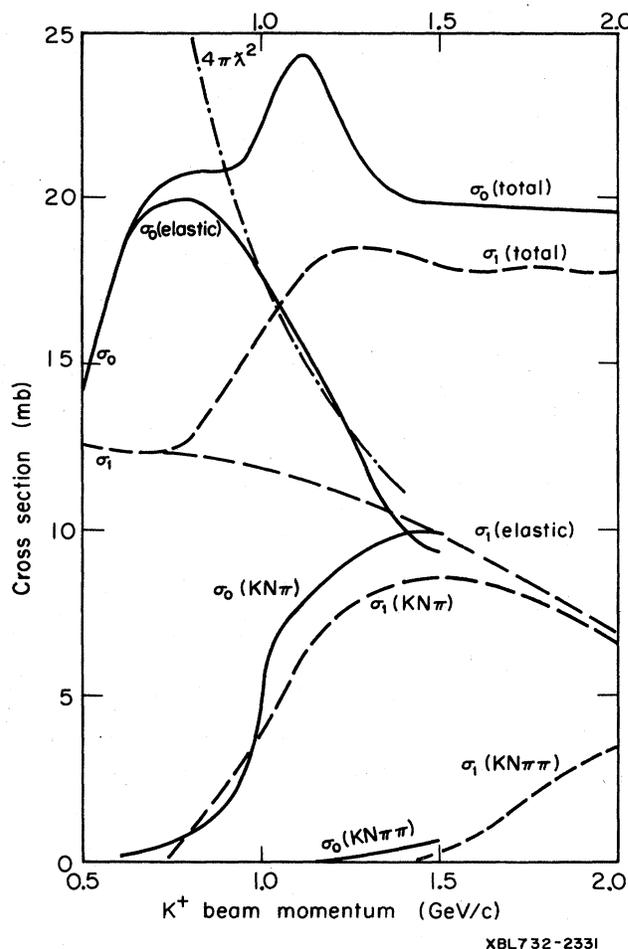


Fig. 1. KN total and partial cross sections. Subscripts indicate isospin; $I=1$ ($I=0$) cross sections are indicated by dashed (solid) curves. The total cross-section curves are from CARROLL 73, who use their recent data and that of BOWEN 73 as well as older data. The elastic $I=1$ curve is hand drawn through new and old elastic data. The inelastic $I=1$ curves are taken from LOKEN 72. The inelastic $I=0$ curves are taken from GIACOMELLI 72, and the elastic $I=0$ curve is obtained by subtracting these from the $I=0$ total cross-section curve of CARROLL 73.

Baryons

Z*'s

have analyzed data up to 2.0 GeV/c using an energy-dependent, two-channel K-matrix formalism in which the inelastic channel coupling to the S wave is taken to be K^*N , and that coupling to higher waves (with $\ell \leq 4$) is taken to be $K\Delta$. A unique solution

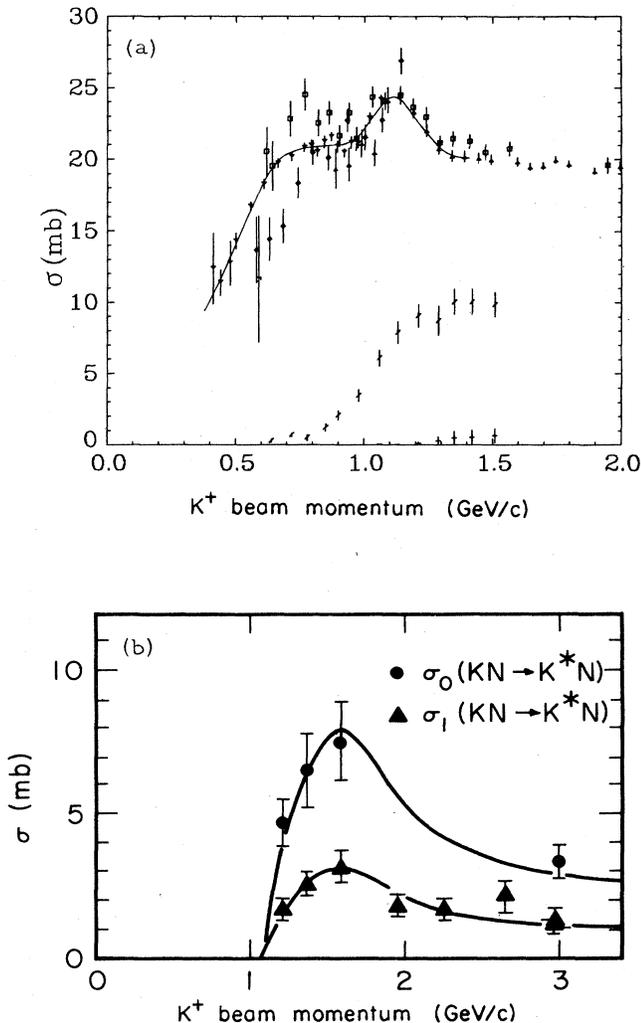


Fig. 2. (a) Unfolded $I=0$ cross sections as quoted by the various authors discussed in this mini-review:

- ◇ BOWEN 73 σ_T
- BUGG 68 σ_T (as unfolded by CARROLL 73)
- ▽ CARROLL 73 σ_T
- △ COOL 70 σ_T
- / GIACOMELLI 72 $\sigma(\pi KN)$
- GIACOMELLI 72 $\sigma(\pi\pi KN)$

(b) Energy dependence of the $I=0$ and $I=1$ cross sections for the reaction $KN \rightarrow K^*N$ (HIRATA 70).

Data Card Listings

For notation, see key at front of Listings.

with a χ^2 per degree of freedom (χ^2/DF) of 1.33 for 3822 data points was found, and it was possible to decrease this to 1.16 by dropping 70 data points with $\chi^2 > 8$ in the original fit. The parametrized P_{13} partial wave had a resonance pole at (1787 - 100i) MeV. The CUTKOSKY 76 analysis included data from 0.8 to 2.5 GeV/c. The ACE parametrization (see MILLER 72) was used along with contributions from P , ρ , A_2 , Λ , Σ , and low mass di-pion exchange, the latter being based on the results of ALCOCK 73. An iterative fitting procedure was followed in which energy-independent fitting was alternated with energy smoothing using a parametrization based on partial-wave dispersion relations which included $K\Delta$, K^*N , and $K^*\Delta$ threshold effects for all waves with $J \leq 7/2$. Typical χ^2/DF values were 1.0 to 1.1.

The question of the existence of a Z_1^* in the elastic channel cannot yet be considered settled. The most likely candidate has long been considered to be the P_{13} wave around 1900 MeV. ARNDT 74 found such an effect, but the three other analyses described above which reached high enough energy to observe the effect (ADAMS 73, MARTIN 75, and CUTKOSKY 76) failed to do so. See KELLY 75 for a discussion of the experimental implications of this discrepancy.

Results of the analyses of ADAMS 73 and CUTKOSKY 76 are shown in Figs. 3 and 4. Solution 2 of ADAMS 73 is plotted below 1.0 GeV/c with errors¹ on real and imaginary parts ranging from 3% to 50% for the various waves. The data set used for the results plotted included a preliminary version of the CAMERON 74 data, and the results are similar to those of CAMERON 74. This solution also agrees qualitatively with CUTKOSKY 76 at higher energies; solution 1 disagrees. The results² of CUTKOSKY 76 are plotted from 0.8 to 2.5 GeV/c. These results are quantitatively quite different from those of MARTIN 75 and ARNDT 74.

Two multi-channel partial-wave analyses of the $I=1$ KN , $K\Delta$, and K^*N channels at 1.21, 1.29, 1.38, and 1.69 GeV/c have been performed by LESQUOY 75. One analysis used a quasi-two-body approach, and the other the isobar model approach. It was not possible to draw any firm conclusions about the

Data Card Listings

For notation, see key at front of Listings.

Baryons
Z*'s

K*N channel or to discriminate between results of other analyses of the elastic channel alone. Results on the $K\Delta$ channel indicate four comparably large waves: SD_{11} , PP_{11} , PP_{13} , and either DD_{13} or DD_{15} depending on the method of analysis. This is in contrast to earlier analyses (e.g., GRIFFITH 72) which found this channel to be dominated by the PP_{13} wave over most of the 1.0 to 1.5 GeV/c region. BERTHON 73 find (like GRIFFITH 72) no indication in the $K\Delta$ channel of rapid phase variation signaling possible resonant effects.

I = 0 System

The experiments of CARROLL 73 and BOWEN 73 (see above) measured K^+d total cross sections in the same energy ranges as their K^+p measurements. These experiments agree rather well with each other and with older data except that the CARROLL 73 cross sections are systematically higher than those of BOWEN 70 (0.36 to 0.72 GeV/c). Unfolded I=0 cross sections, new and old, are shown in Fig. 2(a); the disagreement between CARROLL 73 and BOWEN 73 at low energies is primarily due to differences in the unfolding procedure. The I=0 cross section unfolded from the newer K^+p and K^+d data rises to a plateau in the 0.7 to 0.9 GeV/c region and has a hump of about 3 mb in the 0.9 to 1.3 GeV/c region. Older data had indicated a hump of comparable size in the plateau region. As shown in Figs. 1 and 2(b), the plateau is associated with a broad peak in the elastic cross section, while the hump is associated with a rising K^*N cross section.

In a partial-wave analysis of the I=0 elastic channel GIACOMELLI 74 fit data in the 0.38 to 1.51 GeV/c region, including the new K^+n differential cross-section data of the BGRT collaboration (GIACOMELLI 73; 0.64 to 1.51 GeV/c). Three types of partial-wave analysis (all with S, P, and D waves only) were used to analyze these data: an energy-dependent analysis (ED) using the parametrization of LEA 68 and I=1 partial waves from previous solutions of GIACOMELLI 70, ALBROW 71, and LOVELACE 71; an energy-independent analysis (EI-1) which included K^+p data and determined both the I=0 and I=1 partial waves; and an energy-independent analysis (EI-2) in which the I=1 partial waves were fixed at the values of ALBROW 71's

solution γ . In analyses EI-1 and EI-2 several variants of the LOVELACE 71 shortest-path technique were used to link solutions at different energies. The solutions found by these different methods fell into three main families denoted A, C, and D, the best χ^2/ND values being 1.4-1.9, 1.6-1.7, and 1.5-1.8, respectively. Examples of each class from an ED analysis using the I=1 partial waves of GIACOMELLI 70 are shown in Fig. 5. Some criteria for choice between these solutions are offered by comparison with particular features of the data. The single existing measurement of $K^+n \rightarrow K^0p$ polarization (RAY 69; 0.60 GeV/c) is consistent with classes C and D only. The I=0 total cross sections of CARROLL 73 and BOWEN 73 were not included in the fitting, but comparison is made with CARROLL 73 and the data are found to be consistent with class D only. LONDON 74 has calculated $K_{Lp} \rightarrow K_{Sp}$ amplitudes by combining the best (i.e., lowest χ^2) energy-dependent solution from each of GIACOMELLI 74's classes with existing I=1 KN and $\bar{K}N$ solutions and compared the results with $K_{Lp} \rightarrow K_{Sp}$ cross-section data. He finds that the class D solution is preferred. Similar calculations have been made more recently, and with much more regeneration data, by ALEXANDER 75 and CAMERON 75. The former favors class A, while the latter favors classes C and D but is unable to distinguish between them. No conclusive evidence for Z*'s emerges from the analysis of GIACOMELLI 74; the most likely candidate is the class-D P_{01} wave, but GIACOMELLI 74 point out that the resonance interpretation is questionable for some of the class-D solutions and for all of the class-C solutions. A Breit-Wigner plus quadratic background fit to the class-D P_{01} wave yields the following parameters: $M = 1740$ MeV, $\Gamma = 300$ MeV, $x = 0.85$.

A more recent effort is the energy-dependent simultaneous I=0 and I=1 analysis of MARTIN 75, which makes use of the new K^+n elastic and charge-exchange data of DAMERELL 75. The parametrization used is described above. Two qualitatively similar solutions are found. Both are quite different from all the GIACOMELLI 74 solutions, but the P_{01} wave most resembles that of the class-C solution (see Fig. 5). MARTIN 75 find no significant evidence for resonance-like energy dependence.

Baryons
Z*'s

Data Card Listings
For notation, see key at front of Listings.

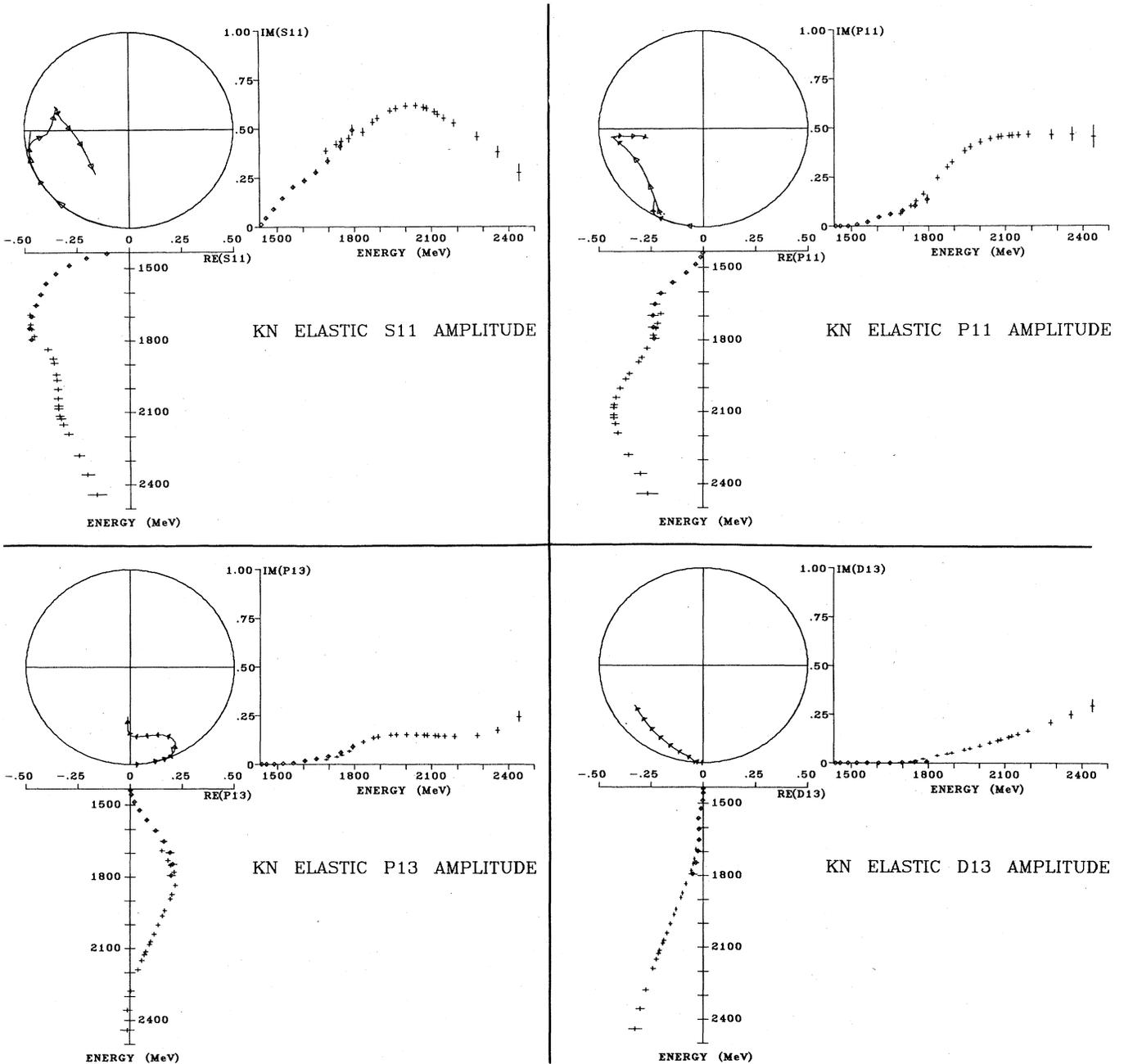


Fig. 3. Amplitudes for $I=1$ KN elastic scattering in the $J=1/2$ and $J=3/2$ waves from ADAMS 73 (ϕ) and CUTKOSKY 76 (\dagger). The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 100 MeV and a base-to-tip length of 20 MeV. All the energy axes run from elastic threshold to 2500 MeV.

Data Card Listings

For notation, see key at front of Listings.

Baryons

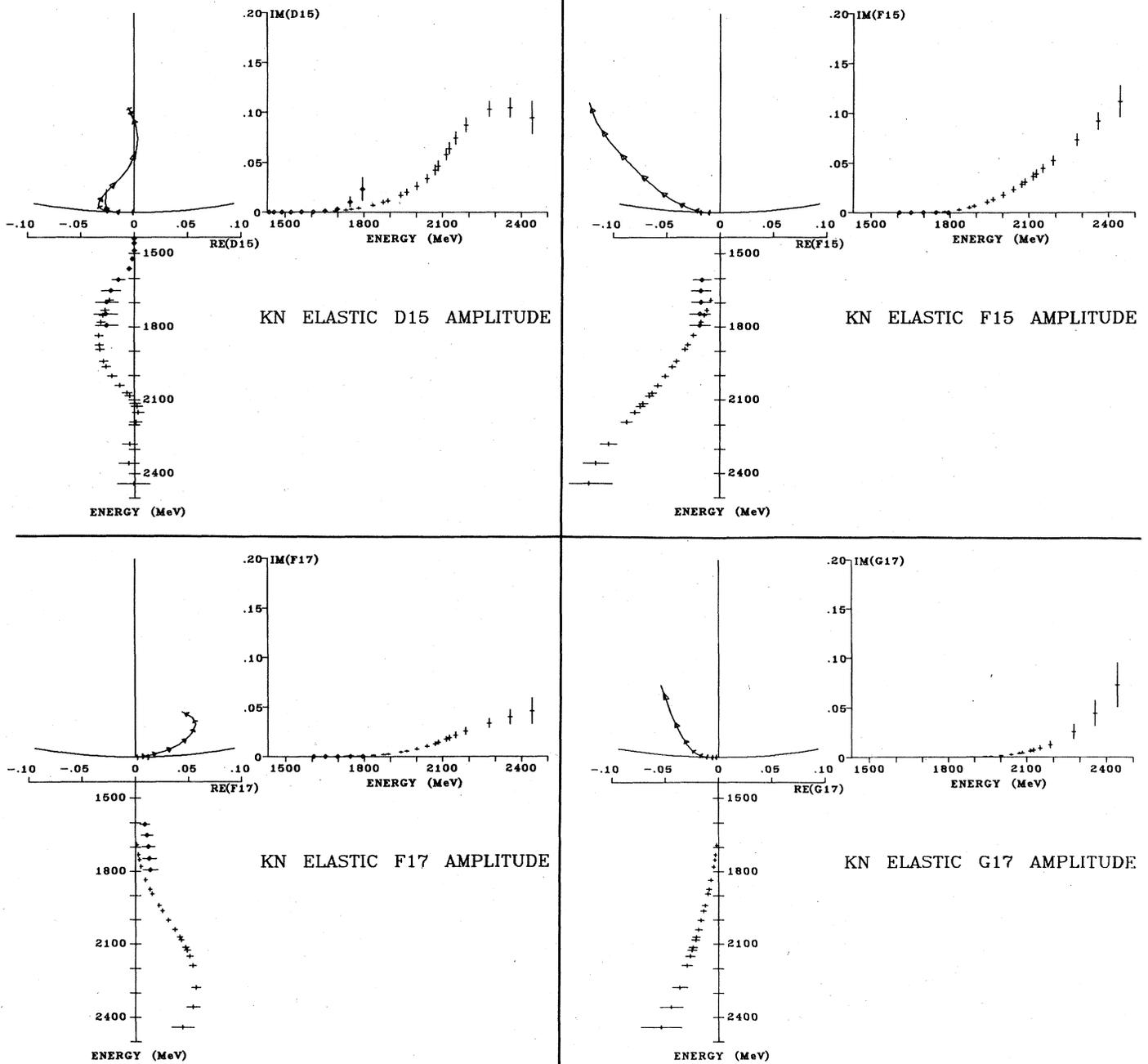
 Z^* 's

Fig. 4. Amplitudes for $I=1$ KN elastic scattering in the $J=5/2$ and $J=7/2$ waves from ADAMS 73 (ϕ) and CUTKOSKY 76 (\dagger). The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 100 MeV and a base-to-tip length of 20 MeV. All the energy axes run from elastic threshold to 2500 MeV.

Baryons
Z*'s

Data Card Listings
For notation, see key at front of Listings.

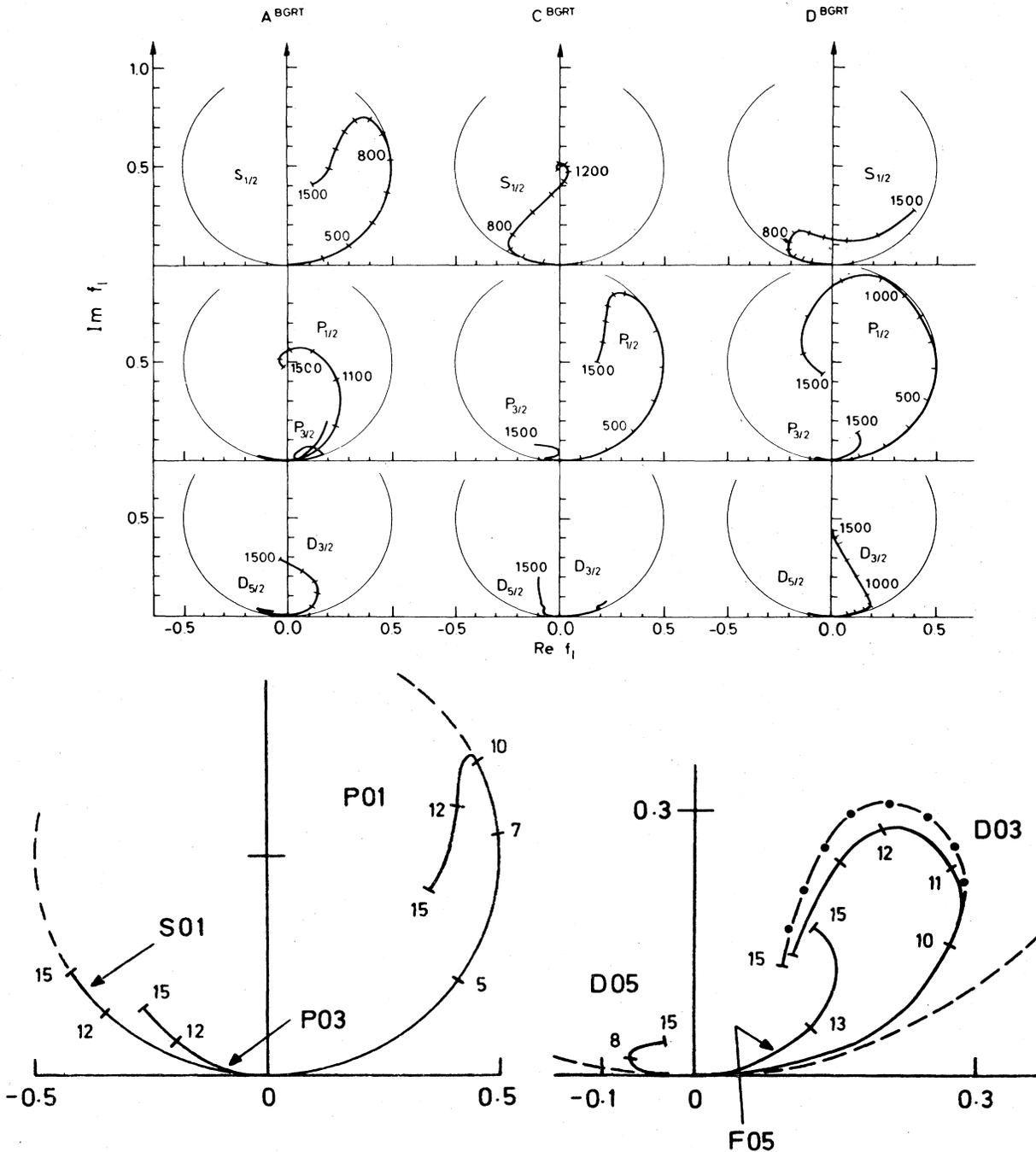


Fig. 5. Amplitudes for $I=0$ KN elastic scattering from GIACOMELLI 74 (upper plots) and MARTIN 75 (lower plots). The GIACOMELLI 74 results are from the energy-dependent analysis and exemplify the three main classes of solutions found in the complete analysis; the energy dependence of the amplitudes is displayed by tick marks at integral multiples of 100 MeV. The MARTIN 75 results are from the preferred solution 2, with the solution-1 D_{03} wave (which exhibited the largest difference between solutions 1 and 2 of any wave) shown as a broken line for comparison; the energy dependence of the amplitudes is displayed by tick marks at the indicated integral multiples of 100 MeV.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Z*^s, Z₀(1780), Z₀(1865), Z₁(1900)

Production Experiments

There are no new developments in production experiments, and the observation of ERNE 70 that present upper limits for Z* production are not small still holds. Cross-section limits for the production of broad Z*^s are comparable to the observed N* and Y* production cross sections.

References

- J. D. Dowell, private communication.
 - R. E. Cutkosky, private communication.
- See the Data Card Listings for other references.

S=1 I=0 EXOTIC STATES (Z₀)

Z₀(1780) 95 Z*0(1780, JP=1/2+) I=0 **P₀₁**

SEE THE MINI-REVIEW PRECEDING THIS LISTING.
WILSON 72 AND GIACOMELLI 74 FIND SOME SOLUTIONS WITH RESONANT-LIKE BEHAVIOR IN THE P₀₁ PARTIAL WAVE. THE EFFECT SEEN IN THE I=0 TOTAL CROSS SECTIONS, IF A RESONANCE, MUST HAVE SPIN=1/2, BECAUSE THE INELASTIC CROSS SECTION IS VERY SMALL AND THE TOTAL CROSS SECTION IS ABOUT 4*PI/K**2.

95 Z*0(1780) MASS (MEV)

M	1780.0	10.0	COOL	70 CNTR	K+P, D TOTAL	1/71
M	D	SEEN	DOWELL	70 CNTR	K+P, D TOTAL	7/70
M	D	SEE ALSO DISCUSSION OF LYNCH 70				7/70
M	W	(1800.)	WILSON	72 PWA	K+N P ₀₁ WAVE	3/72
M	W	ESTIMATE OF PARAMETERS FROM BW + QUADRATIC BACKGROUND FIT TO P ₀₁ .				3/72
M	1	(1750.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 1	9/73
M	1	(1825.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
M	1	FIT 1=FIT OF SINGLE L=1 BW+BACKGROUND TO I=0 TCS FROM 4-1.1 GEV/C				9/73
M	1	FIT 2=FIT OF L=1 AND L=2 BWS TO SAME DATA; SEE Z ₀ (1865) FOR L=2 PART				9/73
M	1	(1740.)	GIACOMEL	74 PWA	.38-1.51 GEV/C	10/74*

95 Z*0(1780) WIDTH (MEV)

W	W	(565.0)	COOL	70 CNTR	K+P, D TOTAL	1/71
W	W	(300.)	WILSON	72 PWA	K+N P ₀₁ WAVE	3/72
W	1	(600.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 1	9/73
W	1	(845.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
W	1	(300.)	GIACOMEL	74 PWA	.38-1.51 GEV/C	10/74*

95 Z*0(1780) PARTIAL DECAY MODES

P1	Z*0(1780) INTO K N	DECAY MASSES
		493+ 939
P2	Z*0(1865) INTO N K*(892)	938+ 892

95 Z*0(1780) BRANCHING RATIOS

R1	Z*0(1780) INTO (K N)/TOTAL	(P1)	
R1	(0.95)	COOL	70 CNTR K+P, D TOTAL 1/71
R1	(0.85)	WILSON	72 PWA K+N P ₀₁ WAVE 3/72
R1	(.75)	CARROLL	73 CNTR IF J=1/2; FIT 1 9/73
R1	(.91)	CARROLL	73 CNTR IF J=1/2; FIT 2 9/73
R1	(.85)	GIACOMEL	74 PWA .38-1.51 GEV/C 10/74*

REFERENCES FOR Z*0(1780)

COOL	70 DUKE CONF 47	R L COOL	(BNL)
ALSO	69 PL 308 564	ABRAMS, COOL, GIACOMELLI, KYCIA, LI +	(BNL)
ALSC	70 PR D1 1887	COOL, GIACOMELLI, KYCIA, LEONTIC, LI +	(BNL)
DOWELL	70 DUKE 53	J.D. DOWELL	(BIRM)
WILSON	72 NP 842 445	+GRIFFITHS, HIRATA +	(BGNA+GLAS+RCMA+TRST)
CARROLL	73 PL 458 531	+KYCIA, LI, MICHAEL, MOCKETT, RAHM +	(BNL)
GIACOMEL	74 NP 871 138	GIACOMELLI, +	(BGNA+GLAS+ROMA+TRST) IJP

PAPERS NOT REFERRED TO IN DATA CARDS

LYNCH	70 DUKE 9	G LYNCH (REVIEWER OF CR. SEC. DATA)	(LRL)
HIRATA	71 NP B30 157	+GOLDHABER, HALL, SEEGER, THRILLING, WOHL (LBL) IJP	
BCWEN	73 PR D7 22	+JENKINS, KALBACH, PETERSEN +	(ARIZ+MICH)
JOHNSON	74 PL 508 343	JOHNSON, VLASSOPULOS	(CFRN, DURH)
CAMERON	75 PALERM CGNF.	+CAPLIUPPI +	(BGNA+EDIN+GLAS+PISA+RHEL) IJP

EXPERIMENTS MAINLY ABOUT ELASTIC CHANNELS --
GOLDHABER 62 PRL 9 135 GOLDHABER, CHINDOSKY, GOLDHABER+ (LRL+UCLA) IJP
RAY 69 PR 183 1183 RAY, BURRIS, FISK, KRAEMER, HILL+ *CFRN+BNL
ARMITAGE 72 NAL PAPER 391 *ASTON, DUEROOTH, ELLISON, + (MCHS+DARE)
GIACOMEL 72 NP 842 437 GIACOMELLI + (BGNA+GLAS+RCMA+TRST)
GIACOMEL 73 NP 856 346 GIACOMELLI, + (BGNA+GLAS+RCMA+TRST)
ALSO 73 BGNA PPT. AE-73/4 GIACOMELLI, GRIFFITHS, + (BGNA+GLAS+ROMA+TRST)
LONDON 74 PRD 9 1565 LONDON (BNL) IJP
ALEXANDE 75 PL 588 484 ALEXANDER, BAR-NIR, BENARY+ (TELAF+EID) IJP
DAMERELL 75 NP 894 374 *POTCHKISS, WICKENS, BENTLEY+ (RHEL+BIRM)

EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS --
GIACOMEL 72 NP 837 577 GIACOMELLI + (BGNA+GLAS+ROMA+TRST)

Z₀(1865) 96 Z*0(1865, JP=3/2-) I=0 **D₀₃**
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K* N THRESHOLD. SEE HIRATA 68 AND 70. WILSON 72 AND GIACOMELLI 73 REPORT PARTIAL WAVE ANALYSES. AARN 73 CLAIMS A RESONANCE IN A MODEL DEPENDENT PWA. SEE ALSO Z*0(1780).

96 Z*0(1865) MASS (MEV)

M	(1860.0)	(15.0)	CARTER	67 THEO	DISPERSION REL.	8/67
M	(1868.0)	(10.0)	COOL	70 CNTR	K+P, D TOTAL	8/67
M	(1830.)		AARON	73 MPWA	I=0 KN .6-1.6G/C	9/73
M	1	(1840.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
M	1	FIT 2=FIT OF L=1 AND L=2 BWS TO I=0 TCS FROM 4-1.1 GEV/C.				9/73
M	1	SEE Z ₀ (1780) FOR FIT 1 AND L=1 PART OF FIT 2.				9/73

96 Z*0(1865) WIDTH (MEV)

W	(200.0)	(50.0)	CARTER	67 THEO		8/67
W	(160.0)	(30.0)	COOL	70 CNTR		8/67
W	(100.)		AARON	73 MPWA	I=0 KN .6-1.6G/C	9/73
W	1	(75.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73

96 Z*0(1865) PARTIAL DECAY MODES

P1	Z*0(1865) INTO K N	DECAY MASSES
		493+ 939
P2	Z*0(1865) INTO N K*(892)	938+ 892

96 Z*0(1865) BRANCHING RATIOS

R1	Z*0(1865) INTO (K N)/TOTAL	(P1)	
R1	(.155)	(.025)	CARTER 67 THEO IF J=3/2 9/73
R1	(.115)	(.025)	COOL 70 CNTR IF J=3/2 9/73
R1	1	(.085)	CARROLL 73 CNTR IF J=3/2; FIT 2 9/73
R2	Z*0(1865) INTO N K*(892)	(P2)	
R2	MAIN INELASTIC DECAY		HIRATA .68 HBC 11/68

REFERENCES FOR Z*0(1865)

CARTER	67 PRL 18 801	A A CARTER	(CAVENDISH)
HIRATA	68 PRL 21 1485	HIRATA, WOHL, GOLDHABER, TRILLING	(LRL)
COOL	70 PR D1 1887	COOL, GIACOMELLI, KYCIA, LEONTIC, LI +	(BNL)
ALSO	66 PRL 17 102	+GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, +	(BNL) I
ALSC	69 PL 308 564	ABRAMS, COOL, GIACOMELLI, KYCIA, LI +	(BNL)
AARON	73 PRD 7 1401	AARON, RICH, HOGAN, SRIVASTAVA	(LASL+NEAS) IJP
CARROLL	73 PL 458 531	+KYCIA, LI, MICHAEL, MOCKETT, RAHM +	(BNL) I

PAPERS NOT REFERRED TO IN DATA CARDS

HIRATA	70 DUKE 429	+GOLDHABER, SEEGER, TRILLING, WOHL	(LRL)
AARON	71 PRL 26 407	+AMADO+ILBAR	(NEAS, PENN, LASL) IJP
HIRATA-1	71 NP B33 445	+GOLDHABER, HALL, SEEGER, TRILLING, WOHL (LBL)	
GIACOMEL	72 NP 837 577	GIACOMELLI +	(BGNA+GLAS+ROMA+TRST)
WILSON	72 NP 842 445	+GRIFFITHS, HIRATA +	(BGNA+GLAS+ROMA+TRST)

S=1 I=1 EXOTIC STATES (Z₁)

Z₁(1900) 97 Z*1(1900, JP=3/2+) I=1 **P₁₃**
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K-DELTA THRESHOLD. SEE THE MINIREVIEW PRECEDING Z*0(1780)

97 Z*1(1900) MASS (MEV)

M	1	(1922.0)	AYED	70 IPWA	P13, SOL. I	6/70
M	1	(1899.0)	AYED	70 IPWA	P13, SOL. II	6/70
M	1	(2030.0)	AYED	70 IPWA	S11, SOL. III	6/70
M	1	THREE SCALNS IN ORDER OF DECREASING SIGNIFICANCE, THOUGH AYED 70				
M	1	GIVE PARAMETERS, THEY CONCLUDE RESONANT INTERPRETATION DOUBTFUL.				
M	2	(1830.)	BARNETT	70 IPWA	P13, SOLN III	9/73
M	2	RESONANCE SIGNAL BARELY ABOVE BACKGROUND DUE TO THE LARGE ERRORS				
M	2	IN THE AMPLITUDES RESULTING FROM THE ANALYSIS				
M		1900.0	COOL	70 CNTR	K+P TOTAL	1/71
M		(1880.)	10.0	ALSBRO	71 IPWA ++ SOL. GAMMA	10/71
M	K	(1890.)	KATO	71 IPWA	SOL II (FIT BW)	10/71
M	K	(2040.)	KATO	71 IPWA	SOL II (FIT BW)	10/71
M	K	KATO 71 ESTIMATE RESONANCE PARAMETERS -- UPDATED PHASE SHIFTS				3/72
M	K	PUBLISHED IN MILLER 72.				3/72

Baryons

Z₁(1900), Z₁(2150), Z₁(2500)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) WIDTH (MEV).

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) REAL PART OF POLE POSITION.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) IMAGINARY PART OF POLE POSITION.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) PARTIAL DECAY MODES.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) BRANCHING RATIOS.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) MAIN INELASTIC DECAY.

Table with columns for particle name, mass, width, and various parameters. Includes a section for REFERENCES FOR Z1(1900).

Table with columns for particle name, mass, width, and various parameters. Includes a section for TOTAL-CROSS-SECTION EXPERIMENTS.

Table with columns for particle name, mass, width, and various parameters. Includes a section for A K-MATRIX ANALYSIS OF SOME OF THE EARLY K+P DATA.

Table with columns for particle name, mass, width, and various parameters. Includes a section for THEORETICAL AND MODEL DEPENDENT ANALYSES.

Table with columns for particle name, mass, width, and various parameters. Includes a section for EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS.

Table with columns for particle name, mass, width, and various parameters. Includes a section for THE MAIN ELASTIC SCATTERING AND POLARIZATION EXPERIMENTS.

Table with columns for particle name, mass, width, and various parameters. Includes entries for BARNETT, CHARLES, CAMERON, YUTA, ABE, ADAMS, PATTON.

Table with columns for particle name, mass, width, and various parameters. Includes a section for PARTIAL-WAVE ANALYSES.

Table with columns for particle name, mass, width, and various parameters. Includes a section for EARLIER ANALYSES THAT DO NOT INCLUDE RECENT POLARIZATION DATA.

Table with columns for particle name, mass, width, and various parameters. Includes a section for LATEST REVIEW TALKS AND PAPERS.

Z1(2150) BUMPS 93 Z1(2150, JP= 1) I=1 A SMALL BUMP IN TOTAL CROSS SECTION AT PK=1.8 GEV/C

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2150) MASS (MEV).

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2150) WIDTH (MEV).

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2150) PARTIAL DECAY MODES.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2150) BRANCHING RATIOS.

Table with columns for particle name, mass, width, and various parameters. Includes a section for REFERENCES FOR Z1(2150).

Z1(2500) BUMPS 94 Z1(2500, JP= 1) I=1 A SMALL BUMP IN TOTAL CROSS SECTION AT PK=2.7 GEV/C

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2500) MASS (MEV).

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2500) WIDTH (MEV).

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2500) PARTIAL DECAY MODES.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(2500) BRANCHING RATIOS.

Table with columns for particle name, mass, width, and various parameters. Includes a section for REFERENCES FOR Z1(2500).

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's**Z₁ CROSS SECTION LIMITS**

SEE MINTREVIEW PRECEDING Z*0

CS	UNITS	MICROBARN	50.		BASSOMPIE 68 HBC	K+P TO Z** PI+	10/69
CS	A	LESS THAN	-2	+3	ANDERSON 69 ASPK +	PI-P TO K-Z**	10/69
CS	A	ABOVE LIMIT FOR	M=1.2 TO 1.4	GEV -	CL= 99 P.C.		
CS	B	LESS THAN	1.4	+1.9	ANDERSON 69 ASPK +	PI-P TO K-Z**	10/69
CS	B	ABOVE LIMIT FOR	M=1.5 TO 2.5	GEV			

REFERENCES FOR Z*1 CROSS SECTION LIMITS

BASSOMPI 68 PL 278 468
ANDERSON 69 PL 298 136

BASSOMPIERRE, + (CERN, BRUXELLES)
+BLESER, BLIEDEN, COLLINS, + (BNL, CARNEGIE)

PAPERS NOT REFERRED TO IN DATA CARDS

TYSON 67 PRL 19 255 +GREENBERG, HUGHES, LU, MINEHART, MCRI, (YALE)
MORI 68 PL 288 152 +GREENBERG, HUGHES, LU, ROTHBERG, + (YALE)
MORI 69 PR 185 1687 +GREENBERG, HUGHES, LU, MINEHART, + (YALE)
MORI 69 REPLACES TYSON 67 AND MORI 68.

Note on Λ 's and Σ 's

The number of confirmed Y^* states has not increased greatly in the last few years, although there has been a larger increase in the number of proposed, but unconfirmed, possible new resonances. Since our last edition (1974) we have made only one addition to the Y^* portion of the Baryon Table, the $\Lambda(1860)$, but have entered four new states into the Data Card Listings, and encountered several further indications that some of the states we list may really be more than one resonance. New high-statistics experiments are needed to clarify the situation.

Just as the recently discovered N^* 's are only weakly coupled to the $\pi N \rightarrow \pi N$ reaction, so also are the recently proposed Y^* 's only weakly coupled to the $\bar{K}N \rightarrow \bar{K}N$, $\bar{K}N \rightarrow \Lambda\pi$, and $\bar{K}N \rightarrow \Sigma\pi$ reactions. For this reason the newer Y^* 's are more difficult to uncover; in invariant mass distributions they usually appear as small peaks or make no appearance at all. Rather, when the 2-body reactions are partial-wave analyzed, some of the amplitudes are found to traverse resonance-like counterclockwise circles. The results of partial-wave analysis give J^P information, whereas a peak seen in an invariant mass distribution or a total cross section often cannot be analyzed for its quantum numbers. We will keep information coming from formation experiments and from production experiments separate whenever necessary.

Formation Experiments

Partial-wave analyses have been performed on $\bar{K}N$, $\Lambda\pi$, $\Sigma\pi$, EK , and $\Lambda\omega$, plus some quasi-two-body channels. Given the present accuracy of the data it is not possible to perform a completely energy-independent analysis, that is, solve for the partial-wave amplitudes at each energy in a model-independent way.

For the great majority of the analyses done so far, the energy range covered was rather limited, corresponding usually to a single bubble chamber experiment. A disturbing feature that appears when examining the partial waves obtained in a particular analysis is that they do not join smoothly with the partial waves given in analyses done for the same channel in a different energy range. One way to avoid this lack of continuity is to analyze the data in as large an energy range as possible. This has been done by six analyses which are described below and three of which are illustrated in Figs. 1-11. See the Data Card Listings for information on other analyses covering limited energy ranges.

Four of the analyses described below (RLIC 76, LEA 73, LANGBEIN 72, and KIM 71) have attempted a multi-channel approach using data on the three 2-body reactions $\bar{K}N$, $\Lambda\pi$, and $\Sigma\pi$, with a fictitious channel sometimes being introduced to account for the global effect of the remaining final states. The latter have large cross sections, so the present multi-channel analyses do not really impose any more stringent unitarity constraints than those already contained in the single-channel fits. However, there is an advantage in the determination of resonance parameters (masses, widths, and branching fractions) since they are fit simultaneously to data in all three channels. The other two analyses (BAILLON 75 and VAN HORN 75) deal only with the $\Lambda\pi$ channel which is, in principle, the easiest to handle as only one isospin has to be considered and both the differential cross section and baryon polarization are measured simultaneously in bubble chamber experiments. These two analyses cover the mass range 1540-2200 MeV and make use of the Legendre polynomial expansion of the angular distributions.

Baryons

Λ 's and Σ 's

a) The most recent analysis is the work done by the Rutherford Laboratory-Imperial College collaboration, RLIC 76. Here the $\bar{K}N$, $\Sigma\pi$, and $\Lambda\pi$ channels have been considered in the mass range 1480-2170 MeV. It is a conventional energy-dependent analysis, the presence of a resonance in a partial wave being detected by comparing the goodness of the fit when this wave is parametrized as a smooth background to the alternative fit when a Breit-Wigner is added to the background. The data used have been carefully selected in order to eliminate inconsistencies (usually the older and statistically less accurate points have been rejected). Angular distributions were directly used in the fit except when the quality of the data was such that no loss of information occurred by using Legendre coefficients (e.g., $K^-p \rightarrow \Sigma^0\pi^0$). Internal consistency is introduced by requiring that the mass and width of the resonances be the same for each of the three channels; a "weighted average" of the three values has been done when necessary and used as a fixed parameter in the final fit. Argand diagrams from the final solution are given in Figs. 1-11. Some suspected resonances are confirmed by this analysis, but many other reported "resonant effects" are not found and new possible resonances are proposed. The situation, in particular for the low partial waves, is still very confused even for masses as low as ~ 1600 MeV.

b) The most recent $\Lambda\pi$ analysis, BAILLON 75, is a semi-energy-independent analysis. The D_{15} and F_{17} partial waves were constrained to lie near the well known $D_{15}(1765)$ and $F_{17}(2030)$ resonances, with their generally accepted parameters, in the energy region where they are important. This constraint was found necessary in order to reduce the number of possible solutions and to fix the arbitrary phase. The latter is undetermined, as in all the inelastic channels, and if unconstrained would make it difficult to join amplitudes at neighboring energies. Above a mass of 1750 MeV only one solution is found. Below this mass, two possible solutions are proposed; one of them (solution 1) has a much greater number of "new resonant structures" and is somewhat preferred on the basis of the behavior of the S and P waves for very low masses

Data Card Listings

For notation, see key at front of Listings.

(around 1580 MeV). Using the method of Barrelet-zeroes^{1,2} it is found that only these 2 solutions are possible if resonances are required to exist in the D_{15} and F_{17} waves. To extract the resonance parameters, the amplitudes were fitted to a combination of Breit-Wigner's plus background. The parameters corresponding to both solutions are given in the data cards, and Argand diagrams are shown in Figs. 1-6.

c) VAN HORN 75 has done an energy-dependent fit of the $\Lambda\pi$ channel. In addition to the best fit, VAN HORN 75 also presents all the other ambiguous solutions that can be generated by the method of Barrelet-zeroes. Among the "ambiguous" solutions, seven of them are found to preserve the established resonances $D_{13}(1670)$, $D_{15}(1765)$, $F_{15}(1915)$, and $F_{17}(2030)$, but with couplings to the $\Lambda\pi$ channel which are sometimes very different from their generally accepted values (cf., BAILLON 75). Also, new resonant structures appear in all the waves, particularly in the lower spins. This analysis is instructive in so far as indicating that there could be an entire constellation of other possible resonances beyond those which appear at the primitive stage of the analysis. However, one should keep in mind that these possible resonances do not necessarily correspond to a simple or even plausible parametrization of the amplitudes. (The values we have listed on the data cards correspond to the average of the 20 best original fits of this analysis, all containing the established, plus a few probable, resonances).

d) The analysis of LEA 73 is a multi-channel energy-dependent partial-wave analysis with parametrized K-matrix elements. The momentum range covered is 440 to 1190 MeV/c, and 99 parameters were used in the fit. Established resonances were constrained to have parameters near their generally accepted values, taken from a previous edition of these tables (for this reason we do not list the results of LEA 73 for these resonances in the following data cards). New resonances are identified using poles of the K-matrix. The procedure used to get the resonance parameters is not described, the values of the fitted K-matrix elements are not given, and no explicit comparison of the results with alternative parametrizations is made.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Λ 's and Σ 's

e) LANGBEIN 72 performed single-energy fits at 40 momenta between 436 and 1226 MeV/c. This work is the nearest approximation to an energy-independent partial-wave analysis existing for the $S=-1$ system. The partial waves corresponding to established resonances were, however, constrained to have Breit-Wigner forms. Approximately 90 acceptable single-energy fits per energy were generated and were used in shortest-path searches. Several candidates for acceptable shortest paths were generated and a preferred path was chosen by rejecting those that failed to reproduce known resonance behavior. Resonances in this solution were identified by loops in Argand diagrams correlated with peaks in the ≥ 3 -body final-state cross section. Resonance parameters were then extracted by fitting Breit-Wigner's with both multiplicative and additive background. A disturbing characteristic of this work is that the amplitudes which were not specifically chosen to be resonant (sometimes even those, as in the case $D_{15}(1765) \rightarrow \Lambda\pi$) show an erratic behavior which cannot be justified simply by the fluctuation of the data.

f) KIM 71 fit data from threshold to 1226 MeV/c using the Ross and Shaw³ effective-range expansion of the inverse multi-channel K-matrix. The data in each of seven energy intervals bounded by 0, 534, 658, 806, 916, 1022, 1117, and 1226 MeV/c, were fit with a constant effective-range matrix. Only the $F_{15}(1915)$ was fixed to a Breit-Wigner form, all other waves included being parametrized by the K-matrix formalism. Resonances were identified by a method involving the appearance of loops in the Argand diagram, peaks in the speed plot, and poles of the K-matrix, the exact procedure not being reported.

None of the three older multi-channel fits above extend beyond a mass of 1.9 GeV (the highest-energy data from the only high-statistics experiment existing before 1971, ARMENTEROS 68). These analyses do not join smoothly with either older (LITCHFIELD 71) or more recent (HEMINGWAY 75) partial-wave analyses at higher energies. The three more recent analyses extend to 2.2 GeV, and for masses higher than 2.2 GeV, the only existing work is an energy-dependent analysis of the $\Lambda\pi$ and $\Sigma\pi$ channels for momenta between 1.93 and

2.5 GeV/c by the Saclay-College de France collaboration (BELLEFON1 75 and BELLEFON 76).

Production Experiments

This type of experiment is often difficult to analyze. Information on $I=0$ states is possible only when there is no $I=1$ state at a similar mass. The main controversies at the present time concern resonances in the 1600 to 1700 MeV region. See the mini-reviews on $\Sigma(1620)$ and $\Sigma(1670)$ in these Listings.

Figures

Argand plots of 15 $S=-1$ partial waves are shown in Figs. 1-11. The analyses shown were picked largely for illustrative purposes rather than on the basis of our judgment of their quality; as discussed above there are a number of analyses extant and no clear choice of the "best" ones is possible.

Errors on Masses and Widths

Often the quoted errors in partial-wave analyses are only statistical, and the values of masses and widths can change by more than these errors when a new parametrization is used. For this reason we report the values of M , Γ , and x_i obtained by different authors even if they analyze the same data. The spread of these masses and widths is certainly a better estimate of the uncertainties than the statistical errors. Sometimes the errors quoted are obtained by the inspection of various fits done with different hypotheses (see, for example, BERTHON 70, GALTIERI 70, VAN HORN 75, RLIC 76). For three states, $\Lambda(1520)$, $\Lambda(1815)$, and $\Sigma(1765)$, there are enough data available to perform an overall fit of the various x_i of the type discussed in the main text (Sec. VII B). In this case we are forced to use the errors, however small they may be, but we warn the reader that the final errors are not to be taken seriously.

In conclusion, in the Baryon Table we choose not to give errors on masses and total widths determined primarily by partial-wave analyses, but, whenever necessary, to show a range of values. As for the branching ratios, we use the errors when needed to perform an overall fit, but we caution the reader.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

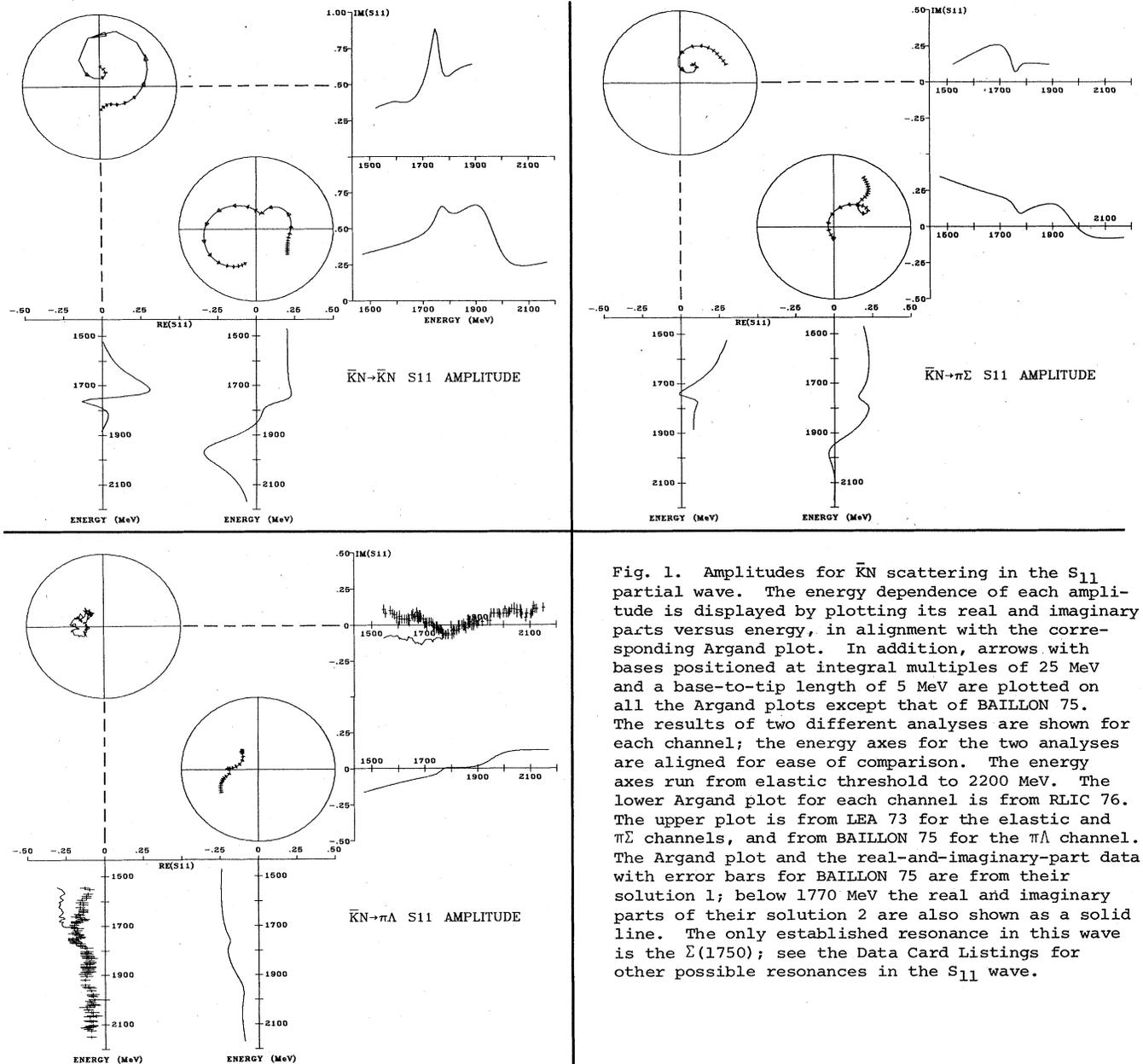


Fig. 1. Amplitudes for $\bar{K}N$ scattering in the S_{11} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. The only established resonance in this wave is the $\Sigma(1750)$; see the Data Card Listings for other possible resonances in the S_{11} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

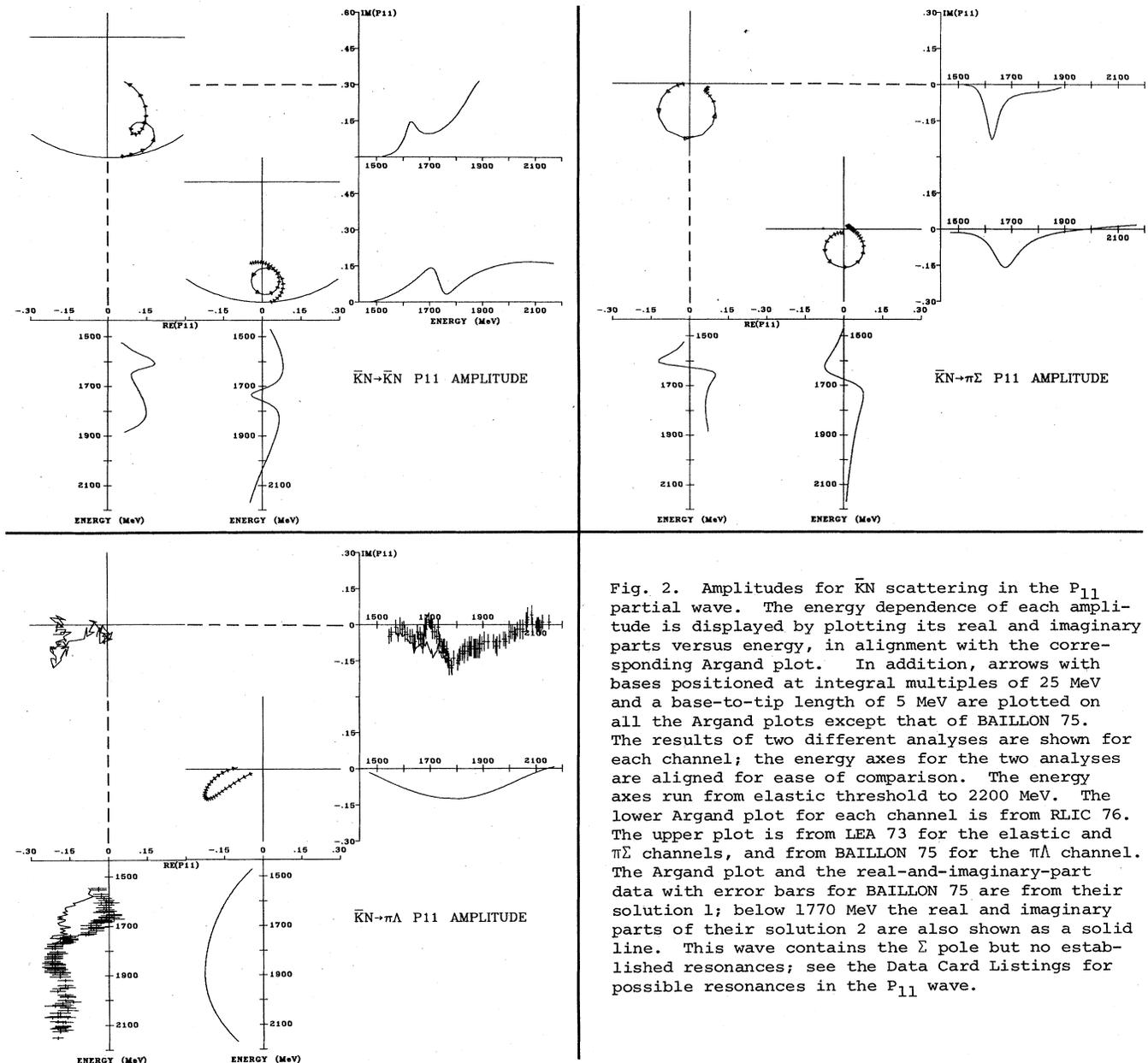
 Λ 's and Σ 's

Fig. 2. Amplitudes for $\bar{K}N$ scattering in the P_{11} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. This wave contains the Σ pole but no established resonances; see the Data Card Listings for possible resonances in the P_{11} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

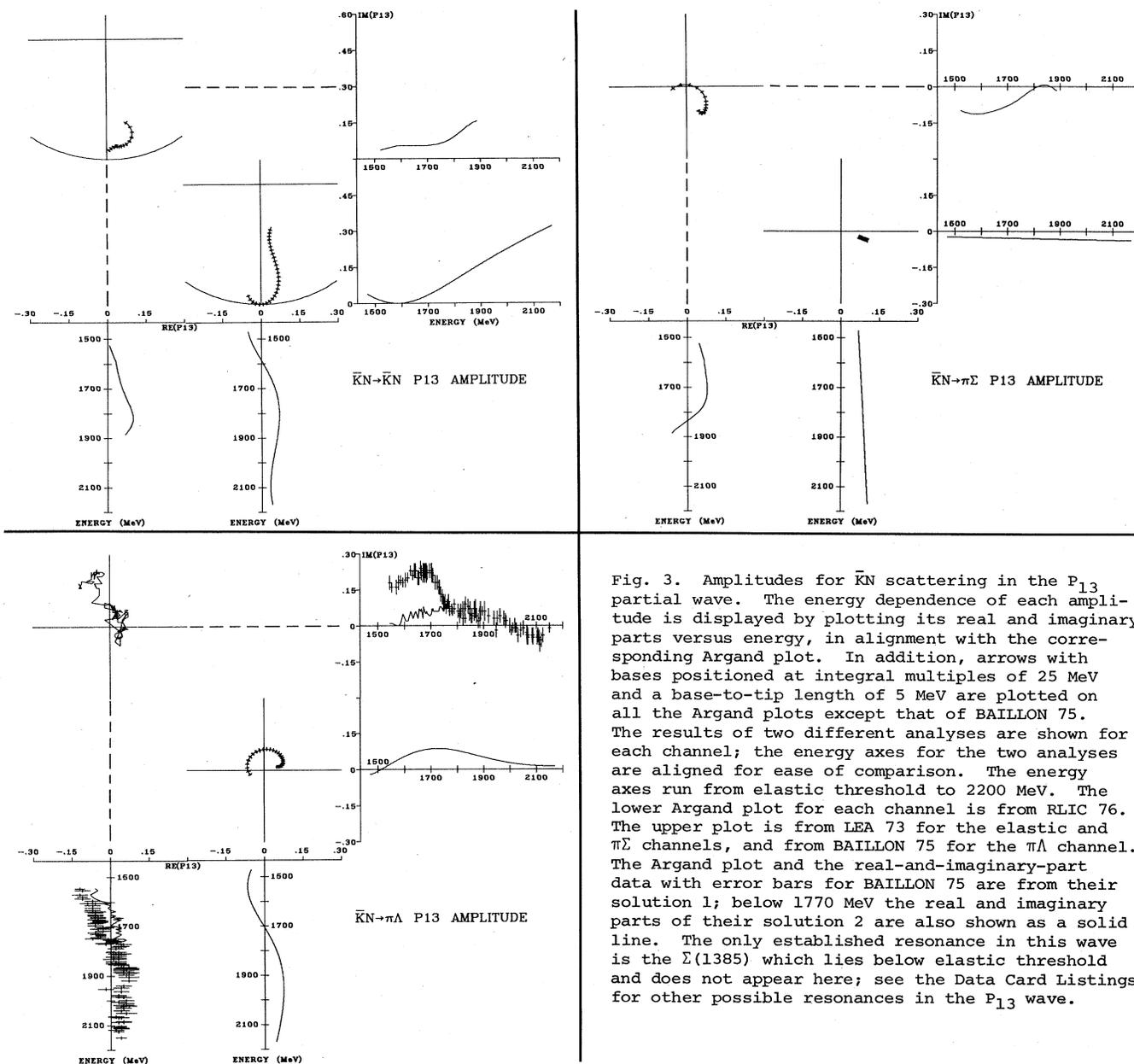


Fig. 3. Amplitudes for $\bar{K}N$ scattering in the P_{13} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. The only established resonance in this wave is the $\Sigma(1385)$ which lies below elastic threshold and does not appear here; see the Data Card Listings for other possible resonances in the P_{13} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

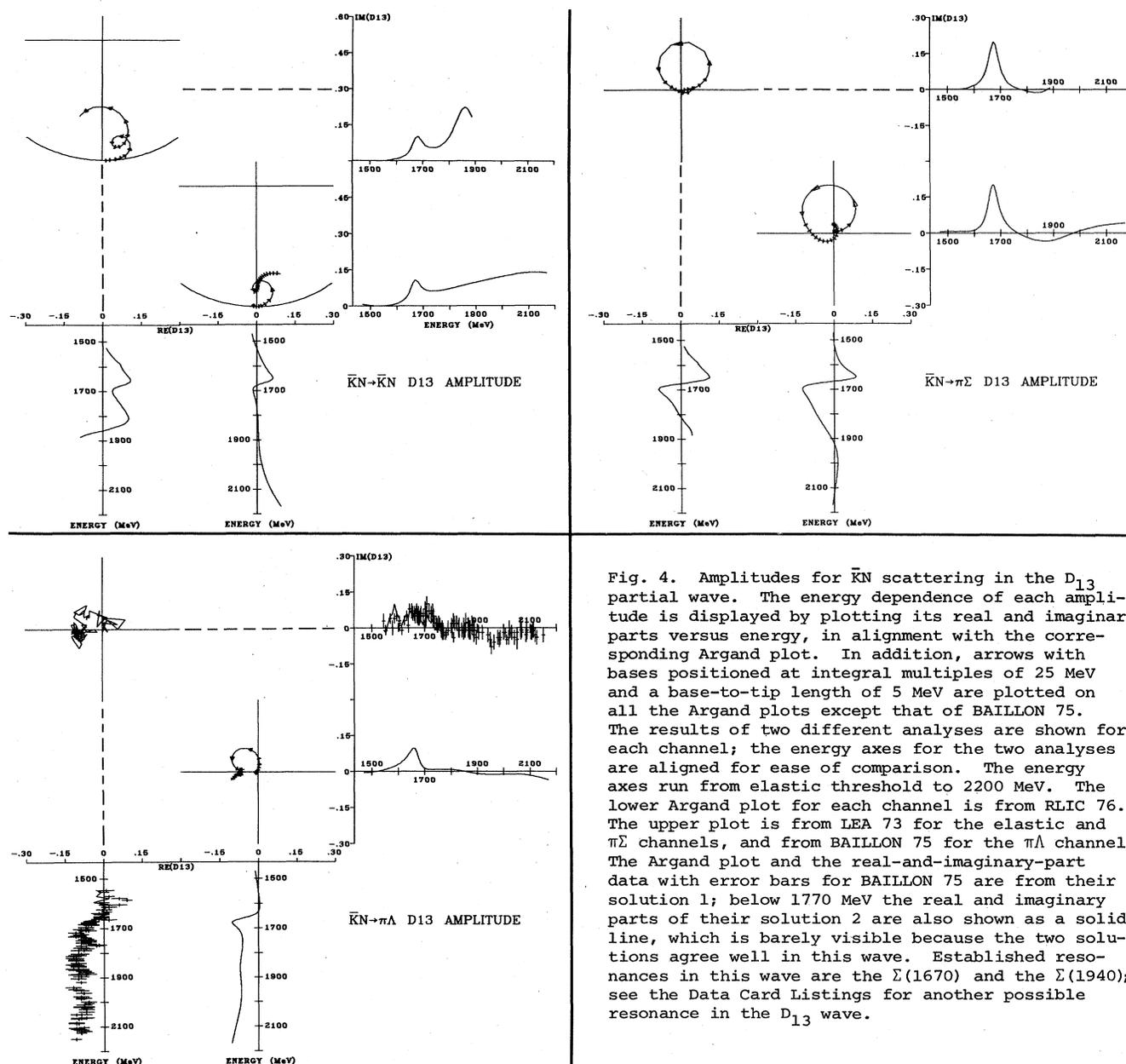
 Λ 's and Σ 's

Fig. 4. Amplitudes for $\bar{K}N$ scattering in the D_{13} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line, which is barely visible because the two solutions agree well in this wave. Established resonances in this wave are the $\Sigma(1670)$ and the $\Sigma(1940)$; see the Data Card Listings for another possible resonance in the D_{13} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

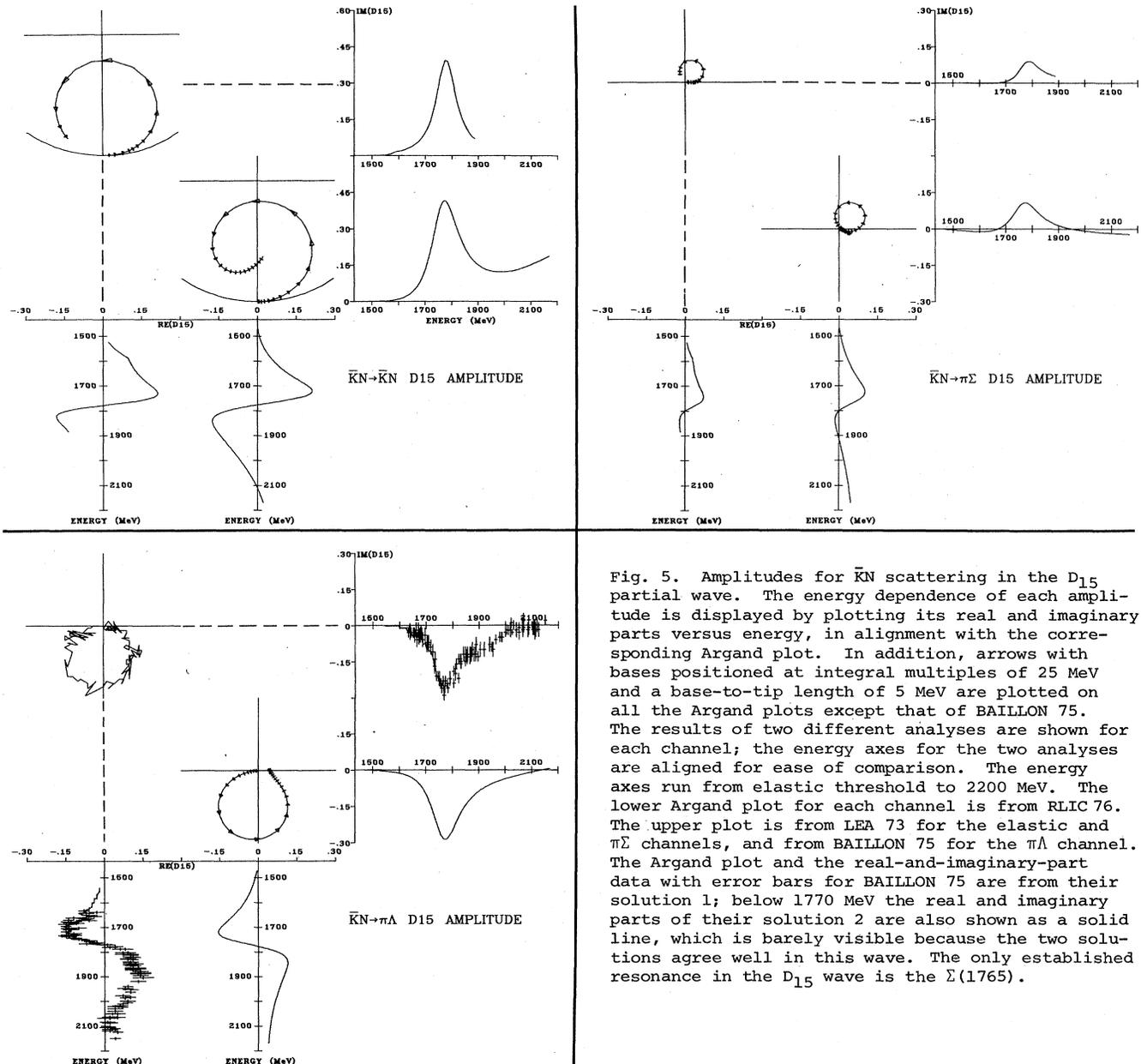


Fig. 5. Amplitudes for $\bar{K}N$ scattering in the D_{15} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line, which is barely visible because the two solutions agree well in this wave. The only established resonance in the D_{15} wave is the $\Sigma(1765)$.

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's

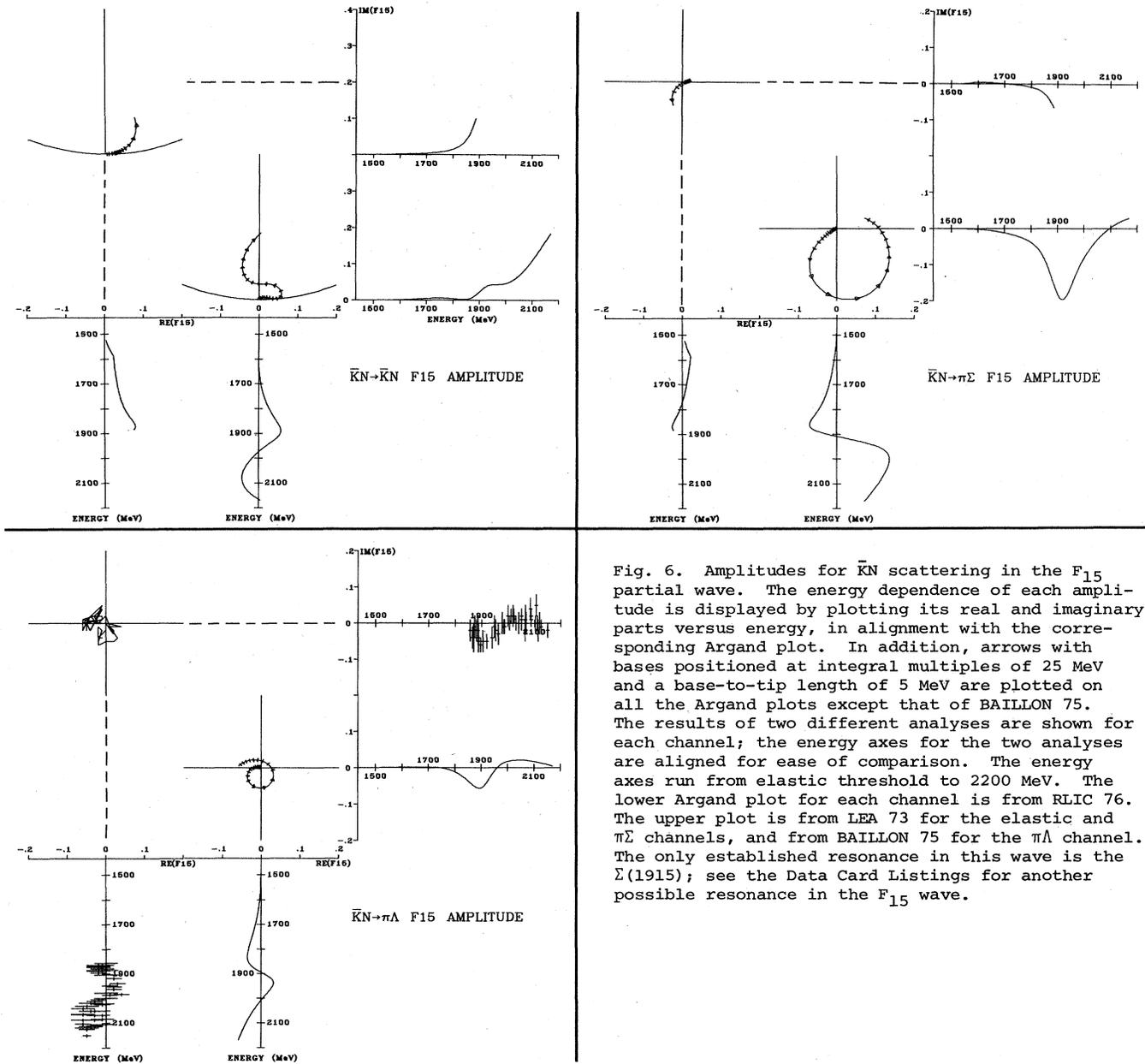
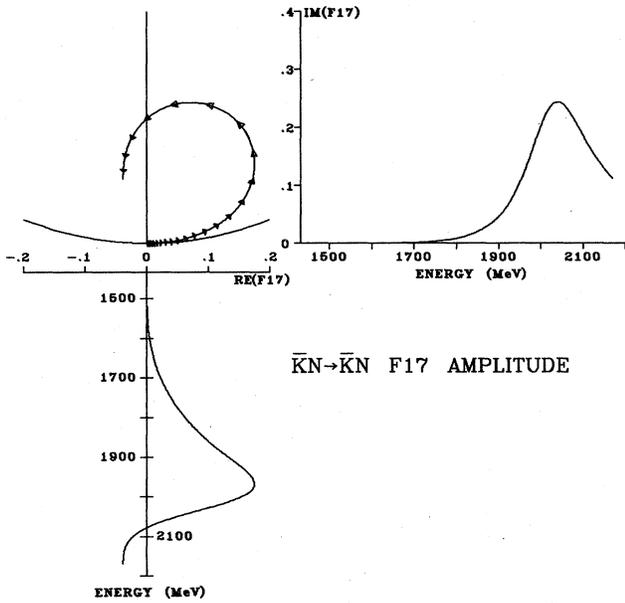


Fig. 6. Amplitudes for $\bar{K}N$ scattering in the F_{15} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The only established resonance in this wave is the $\Sigma(1915)$; see the Data Card Listings for another possible resonance in the F_{15} wave.

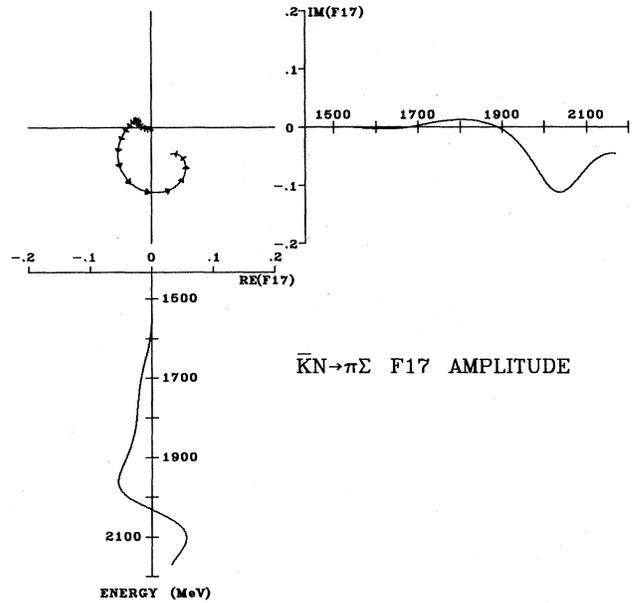
Baryons
 Λ 's and Σ 's

Data Card Listings

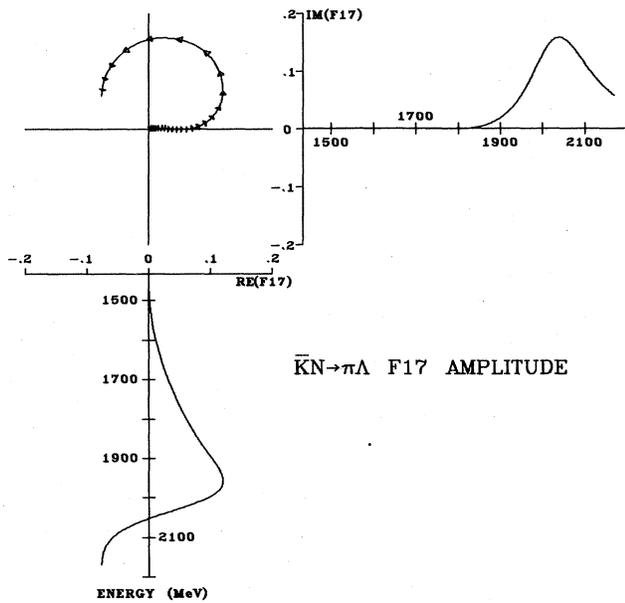
For notation, see key at front of Listings.



$\bar{K}N \rightarrow \bar{K}N$ F17 AMPLITUDE



$\bar{K}N \rightarrow \pi\Sigma$ F17 AMPLITUDE



$\bar{K}N \rightarrow \pi\Lambda$ F17 AMPLITUDE

Fig. 7. Amplitudes for $\bar{K}N$ scattering in the F_{17} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The energy axes run from elastic threshold to 2200 MeV. All the results shown are from RLIC 76. The only established resonance in the F_{17} wave is the $\Sigma(2030)$.

Data Card Listings

For notation, see key at front of Listings.

Baryons

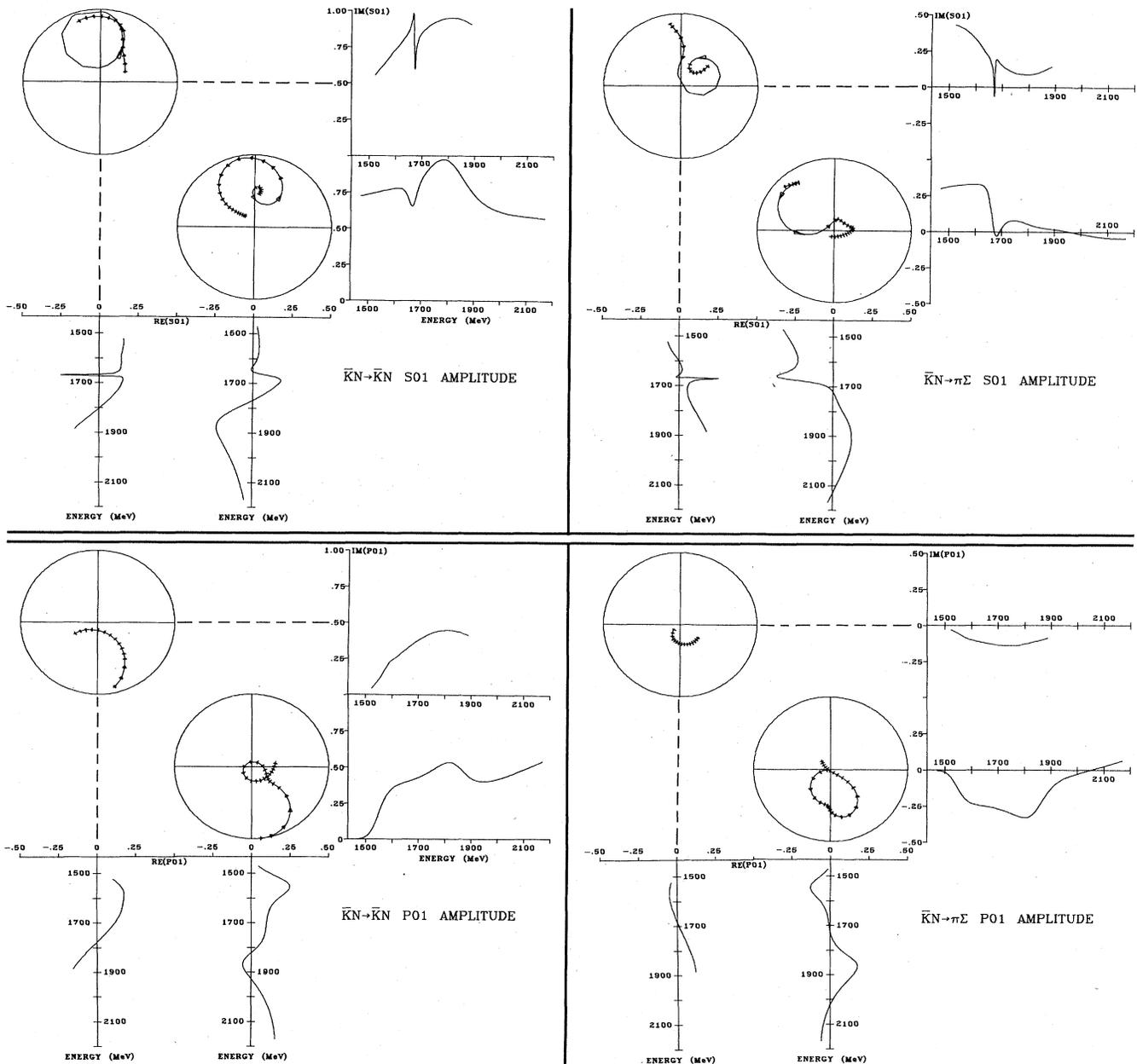
 Λ 's and Σ 's

Fig. 8. Amplitudes for $\bar{K}N$ scattering in the S_{01} and P_{01} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in the S_{01} wave are the $\Lambda(1405)$ (which lies below threshold and does not appear here), the $\Lambda(1670)$, and the $\Lambda(1870)$; the P_{01} wave contains the Λ pole but no established resonances. See the Data Card Listings for possible resonances in the P_{01} wave.

Baryons

Λ 's and Σ 's

Data Card Listings
For notation, see key at front of Listings.

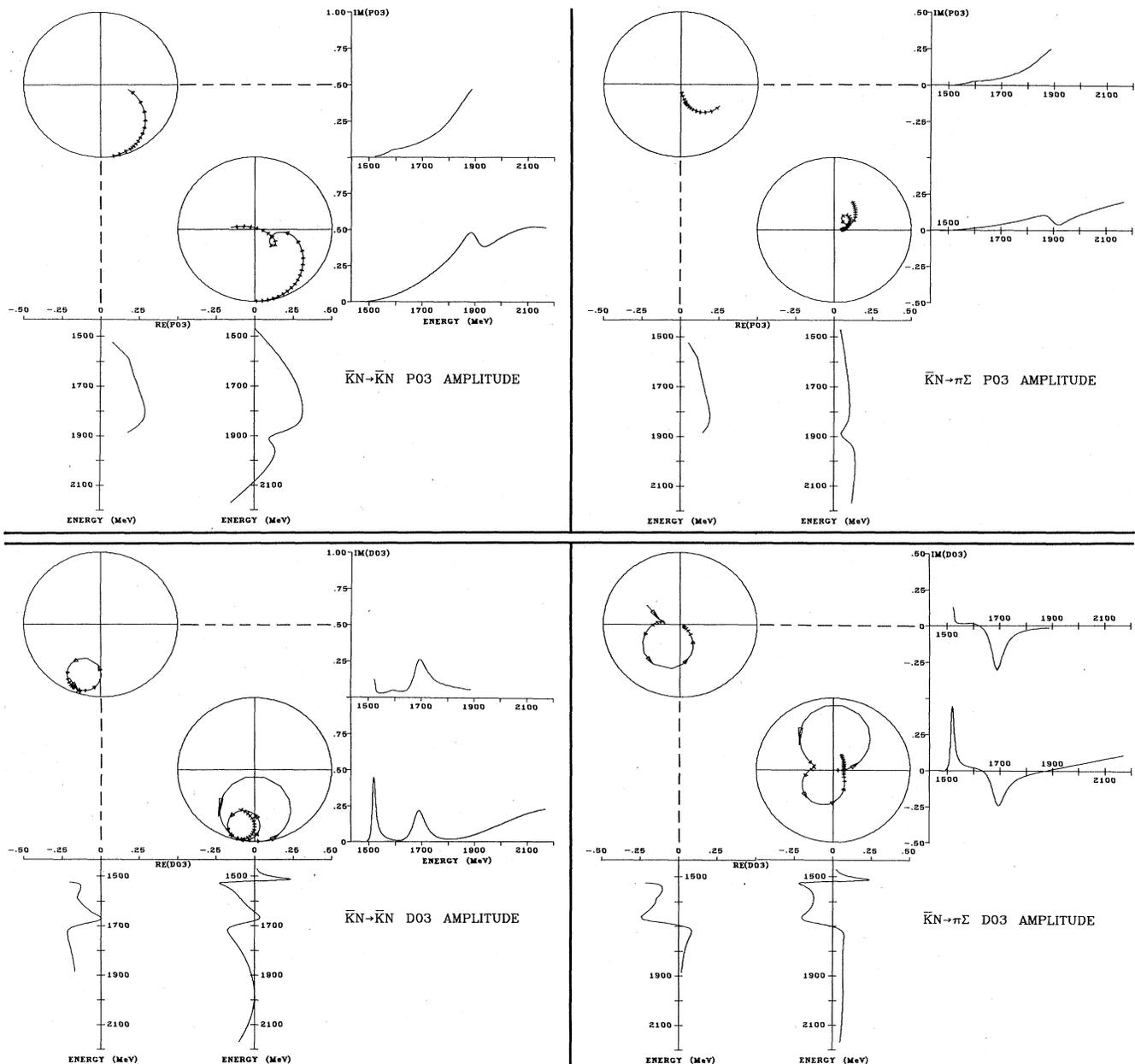


Fig. 9. Amplitudes for $\bar{K}N$ scattering in the P_{03} and D_{03} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in these waves are the $\Lambda(1860)$ in the P_{03} wave and the $\Lambda(1520)$ and $\Lambda(1690)$ in the D_{03} wave. See the Data Card Listings for another possible resonance in the D_{03} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

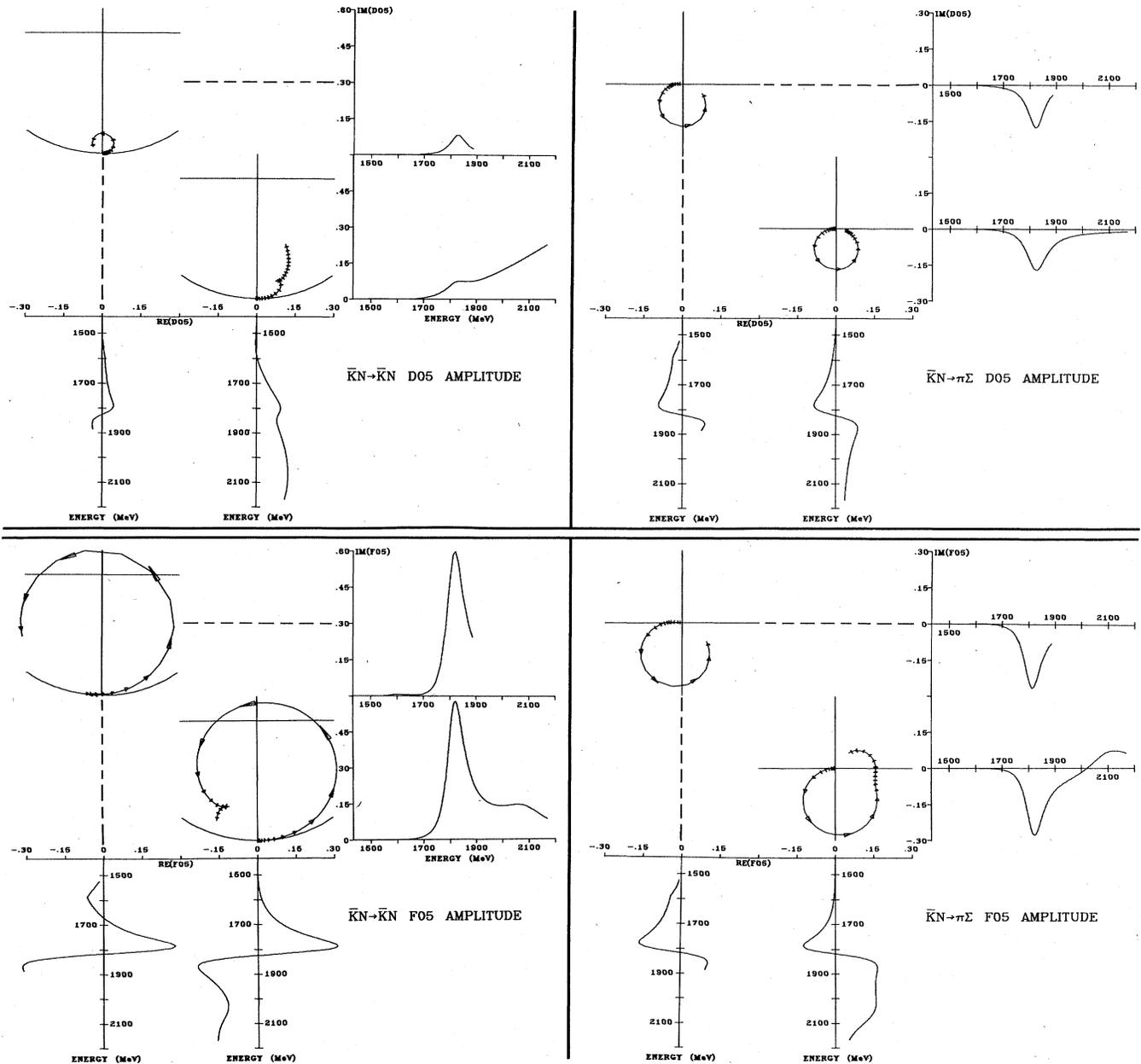
 Λ 's and Σ 's

Fig. 10. Amplitudes for $\bar{K}N$ scattering in the D₀₅ and F₀₅ partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in these waves are the $\Lambda(1830)$ in the D₀₅ wave and the $\Lambda(1815)$ in the F₀₅ wave. See the Data Card Listings for another possible resonance in the F₀₅ wave.

Baryons

 Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

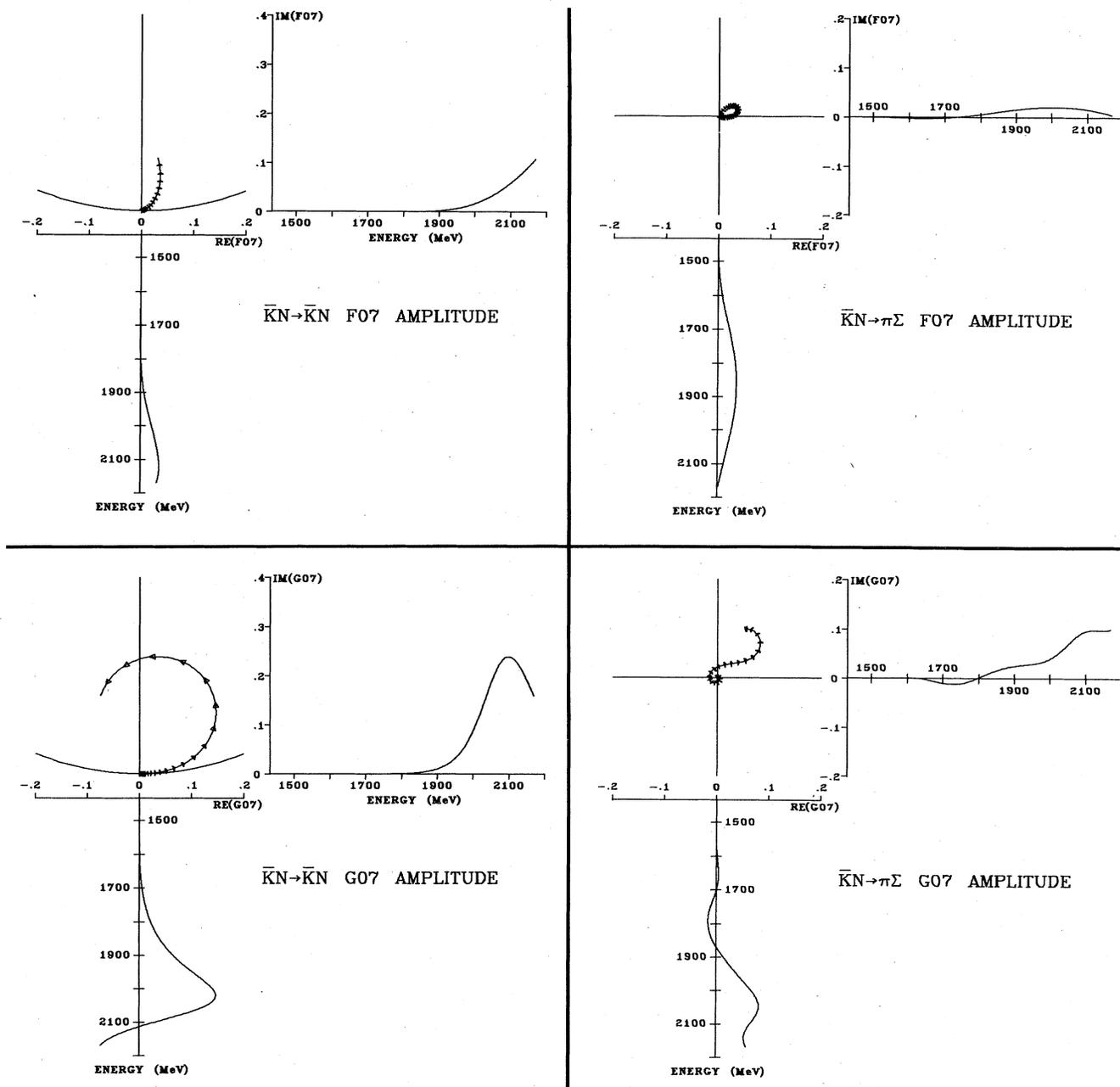


Fig. 11. Amplitudes for $\bar{K}N$ scattering in the F_{07} and G_{07} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The energy axes run from elastic threshold to 2200 MeV. All the results shown are from RLIC 76. The only established resonance in these waves is the $\Lambda(2100)$ in the G_{07} wave. See the Data Card Listings for a possible resonance in the F_{07} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Λ 's and Σ 's, Λ , $\Lambda(1330)$, $\Lambda(1405)$

Conclusions

Table I is an attempt to evaluate the status of the various Y^* 's. The evaluations are of course partly subjective. A blank indicates that there is no corresponding evidence at all. This may mean either that the relevant couplings are small or that the resonance does not really exist. The Baryon Table includes only the well-established resonances. It seems clear, however, that whereas any particular one of the questionable resonances may disappear with the next analysis, there are probably many new resonances underlying those already established.

References

1. E. Barrelet, Nuovo Cimento **8A**, 331 (1972).
2. P. Litchfield, Proceedings of the IInd Aix-en-Provence Conference (1973).
3. M. Ross and G. Shaw, Ann. Phys. (N.Y.) **13**, 147 (1961).

For other references see the Data Card Listings.

TABLE I. STATUS OF Y^* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --					OTHER CHANNELS
			TOTAL CR. SEC.	KBAR N	LAM PI	SIG PI		
LAM(1115) P01		****						WEAK TO N PI
LAM(1330)		DEAD						
LAM(1405) S01		****		****	F	****		LAM2PI, LAM GAM
LAM(1520) D03		****		****	D	****		
LAM(1600) P01		*			R	****		LAM ETA
LAM(1670) S01		****		****	B	****		LAM2PI, SIG2PI
LAM(1690) D03		****		****	I	****		
LAM(1800) P01		**						
LAM(1800) G09		*						
LAM(1815) F05		****		****	D	****		SIG(1385) PI
LAM(1830) D05		****		****	E	****		
LAM(1860) P03		***		**				
LAM(1870) S01		**		****	N	**		
LAM(2010)		**			F			LAM CMG
LAM(2020) F07		*			O			
LAM(2100) G07		****		****	R	****		XI K, LAM CMG
LAM(2110) F05		**		**	B	*		LAM CMG
LAM(2350)		****		****	I	****		
LAM(2350)		***		**	D	*		
SIG(1190) P11		****						WEAK TO N PI
SIG(1385) P13		****				****		
SIG(1440) PE		DEAD						
SIG(1480) PE		*		*		*		
SIG(1580) D13		**	**	**		**		
SIG(1620) S11		**	**	**		**		
SIG(1660) P11		**	**	**		**		
SIG(1670) D13		****	**	****		****		SEVERAL OTHERS
SIG(1670) PE		**	**	**		**		SEVERAL OTHERS
SIG(1690) PE		**	*	**		**		LAM 2-PI
SIG(1750) S11		****	****	****		****		SIG ETA
SIG(1765) D15		****	****	****		****		SEVERAL OTHERS
SIG(1770) P11		*	*	*		*		
SIG(1840) P13		*	*	*		*		
SIG(1880) P11		**	**	**		**		
SIG(1915) F15		****	***	****		****		
SIG(1940) D13		****	***	****		****		
SIG(2000) S11		*	*	*		*		
SIG(2030) F17		****	****	****		****		XI K
SIG(2070) F15		*	*	*		*		
SIG(2080) P13		**	**	**		**		
SIG(2100) G17		*	*	*		*		
SIG(2250)		****	****	*		*		
SIG(2455)		***	***	*		*		
SIG(2620)		***	***	*		*		
SIG(3000)		**	**	*		*		

**** GOOD, CLEAR, AND UNMISTAKABLE.
 *** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
 ** NEEDS CONFIRMATION.
 * WEAK.
 * ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

S=-1 I=0 HYPERON STATES (Λ)

Λ 18 LAMBDA(1115, JP=1/2+) I=0
 SEE STABLE PARTICLE DATA CARD LISTINGS

$\Lambda(1330)$ BUMPS 87 $Y^*(1330, JP=)$ I=0 PRODUCTION EXPERIMENTS
 SEE THE MINI-REVUE AT THE START OF THE Y^* LISTINGS.
 A PEAK WAS SEEN NEAR 1330 MEV IN THE LAMBDA GAMMA SPECTRUM IN THREE PI-PROPANE EXPERIMENTS (YUNG-CHANG 64, BUBELEV 67, AND BOZOKI 68). ALL MORE RECENT RESULTS INDICATE THAT THERE IS NO RESONANCE NEAR THIS MASS VALUE.

REFERENCES FOR $Y^*(1330)$ (PROD. EXP.)
 Y-CHANG 64 DUBNA CONF I 615 YUNG-CHANG, IN, KLADNITSKAYA, + (DUBNA)
 BUBELEV 67 PL 248 246 +CHADRAA, CHUVILO, + (JINR, BUCHAREST, CERN)
 DAHL 67 PR 163 1377 CAHL, HARDY, FESS, KIRZ, MILLER (LRL)
 BOZOKI 68 PL 288 360 +FENYVES, GEMESY, + (BUDAPEST, DUBNA)
 TAN 69 PRL 23 101 T H TAN (SLAC)
 MAYEUR 70 PL 338 441 +VAN BINST, WILQUET++ (BRUX, CERN, TUFT)
 COLAS 75 NP 891 253 COLAS, FARWELL, FERRER, SIX (DRSA)

$\Lambda(1405)$ 37 $Y^*(1405, JP=1/2-)$ I=0 PRODUCTION EXPERIMENTS **S₀₁**
 THIS RESONANCE CAN BE IDENTIFIED WITH THE VIRTUAL BOUND STATE IN THE KBAR-N SYSTEM FOUND IN THE ANALYSIS OF LOW ENERGY K-P INTERACTION. WE LIST SUCH EXPERIMENTS SEPARATELY BELOW. WE USE ONLY PRODUCTION EXPERIMENTS FOR AVERAGING OF MASSES AND WIDTHS.

37 $Y^*(1405)$ MASS (MEV) (PROD. EXP.)

M	(1405.0)	ALSTON 61 HBC	K-P 1.15 BEV/C
M	(1410.0)	ALEXANDER 62 HBC	PI-P 2.1 BEV/C
M	(1405.0)	ALSTON 62 HBC	K-P 1.2-5 BEV/C
M	(1382.0)	ENGLER 65 HBC	PI-P, PI+D 1.6B
M	1400.0	MUSGRAVE 65 HBC	PBAR P 3-4 BEV/C
M	67 1400.0	BIRMINGHAM 66 HBC	K-P 3.5
M	120 1405.0	GALTIERI 68 HBC	K-D 2.1-2.7BEV/C
M	AVG 1402.4	3.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
M	STUDENT 1402.4	3.9	AVERAGE USING STUDENT 10(H/1.11) -- SEE TEXT

37 $Y^*(1405)$ WIDTH (MEV) (PROD. EXP.)

W	(20.0)	ALSTON 61 HBC	7/66
W	35.0	ALEXANDER 62 HBC	
W	(50.0)	ALSTON 62 HBC	
W	(89.0)	ENGLER 65 HBC	7/66
W	60.0	MUSGRAVE 65 HBC	7/66
W	67 50.0	BIRMINGHAM 66 HBC	9/67
W	120 35.0	GALTIERI 68 HBC	6/68
W	AVG 38.1	3.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W	STUDENT 37.9	4.3	AVERAGE USING STUDENT 10(H/1.11) -- SEE TEXT

37 $Y^*(1405)$ PARTIAL DECAY MODES (PROD. EXP.)

PI	$Y^*(1405)$ INTO SIGMA PI	DECAY MASSES
		1197+139
REFERENCES FOR $Y^*(1405)$ (PROD. EXP.)		
ALSTON 61 PRL 6 698	+ALVAREZ, EBERHARD, GODO, GRAZIANO, + (LFL) I	
ALEXANDE 62 PRL 8 447	ALEXANDER, KALBFLEISCH, MILLER, SMITH (LRL) I	
ALSTON 62 CERN CCNF 311	+ALVAREZ, FERRO-LUZZI, ROSENFELD, + (LRL) I	
ENGLER 65 PRL 15 224	+FISK, KRAEMER, MELTZER, WESTGARD, + (CERN, BNL) IJ	
MUSGRAVE 65 NC 35 735	+PETMEZAS, + (BIRM, CERN, EPOL, LOIC, SACLAY)	
BIRMINGHAM 66 PR 152 1148	BIRMINGHAM, GLASGOW, LOIC, CXFORD, RUTH-ERFORD	
GALTIERI 68 PRL 21 573	BARBARO-GALTIERI, CHADWICK + (LRL, SLAC)	

Baryons

$\Lambda(1405)$, $\Lambda(1520)$

Data Card Listings

For notation, see key at front of Listings.

1405 MEV REGION: EXTRAPOLATIONS BELOW THRESHOLD

24 Y*(1405, JP=1/2-) I=0
EXTRAPOLATION BELOW THRESHOLD

SEE NOTE IN Y*(1405) PRODUCTION EXPERIMENTS -THE DIFFICULTIES IN EXTRAPOLATING FROM THE PHYSICAL REGION TO THE RESONANCE LOCATION ARE DISCUSSED BY DALITZ 67.

THE QUESTION ON WHETHER Y*(1405) IS A KBAR-N BOUND STATE OR A CDD POLE (DALITZ 70, RAJASEKARAN 72 HAS BEEN INVESTIGATED BY CLINE 71, MARTIN 71, GALTIERI 72, AND COBSON 72. THE LAST TWO PAPERS CONCLUDE THAT THE DATA CANNOT TELL THE DIFFERENCE.

24 Y*(1405) MASS (MEV)

Table with columns for mass values and associated references (KIM, SAKITT, KITTEL, etc.)

24 Y*(1405) WIDTH (MEV)

Table with columns for width values and associated references (KIM, SAKITT, KITTEL, etc.)

REFERENCES FOR Y*(1405) (FROM EXTRAPOLATIONS)

Table listing references for Y*(1405) with author names and journal abbreviations.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards, including authors like ABRAMS, DONALD, KADYK, etc.

$\Lambda(1520)$

D03

38 Y*(1520, JP=3/2-) I=0

PRODUCTION AND FORMATION EXPERIMENTS AGREE QUITE WELL WITH EACH OTHER, THEREFORE, THEY HAVE NOT BEEN SEPARATED FOR THIS PARTICLE.

THE DECAY MODE LAMBDA PI PI IS LARGELY DUE TO Y*(1385) PI. ONLY THE VALUES OF Y*(1385) PI/(LAMBDA PI PI) GIVEN BY MAST 72 AND CORDEN 75 ARE BASED ON REAL 3-BODY PARTIAL WAVE ANALYSES (THE OLDER RESULTS BEING OBTAINED USING CRUDE METHODS). THE DISCREPANCY BETWEEN THE 2 RESULTS IS ESSENTIALLY DUE TO THE DIFFERENT HYPOTHESIS MADE CONCERNING THE SHAPE OF THE EPSILON MESON.

38 Y*(1520) MASS (MEV)

Table with columns for mass values and associated references (GALTIERI, WATSON, ALMEIDA, etc.)

38 Y*(1520) WIDTH (MEV)

Table with columns for width values and associated references (WATSON, MUSGRAVE, BIRMINGHA, etc.)

38 Y*(1520) PARTIAL DECAY MODES

Table listing partial decay modes for Y*(1520) with branching fractions and references.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_i² + δP_j²), while the off-diagonal elements are the normalized correlation coefficients (δP_iδP_j)/(δP_iδP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table showing fitted partial decay mode branching fractions and correlation coefficients.

38 Y*(1520) BRANCHING RATIOS

Large table listing branching ratios for Y*(1520) into various channels (SIGMA PI, LAMBDA PI, etc.) with associated references and error factors.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1520)$, $\Lambda(1600)$, $\Lambda(1670)$

R9 Y*0(1520) TO Y*1(1385) PI TO LM PI PI/LM PI PI (P8)/(P3) 9/73
R9 MORE THAN 0.10 CLINE 69 DBC K-D TO 2PI LAM N 9/69
R9 B 0.35 0.10 BURKHARDT 71 HBC LAM. 3PI PRD. 3/71
R9 C (1.05) CHAN 72 IPWA K-P TO LAM 2PI 2/73
R9 M 0.82 0.10 MAST 73 IPWA K-P TO 2PI LAM 12/72
R9 .58 .22 CORDEN 75 CBC K-0 1.4-1.8GV/C 4/75*

R10 Y*0(1520) INTO (Y*1(1385) PI)/TOTAL (P7)
R10 0.041 0.005 CHAN 72 HBC K-P TO LAM 2PI 3/71

R11 Y*0(1520) INTO (LAMBDA PI PI)/TOTAL (P3)
R11 0.10 (0.02) COLLEY 71 DBC K-N 1.5 GEV PRD 10/71
R11 .11 .01 MAST 73 IPWA K-P TO 2PI LAM 9/73
R11 1 BASED ON ASSUMED ELASTICITY OF .467/-02 9/73
R11 3 .091 .006 CORDEN 75 DBC K-0 1.4-1.8GV/C 4/75*

R11 1 SUPERSEDES COLLEY 71 1/76*
R11 AVG 0.0960 0.0084 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
R11 STUDENT 0.0956 0.0062 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R11 FIT 0.0950 0.0043 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R12 Y*0(1520) INTO (LAMBDA EPSILON)/TOTAL (P9) 4/75*
R12 .20 .08 CORDEN 75 DBC K-D 1.4-1.8GV/C 4/75*

REFERENCES FOR Y*0(1520)
GALTIERI 63 PL 6 296 A BARBARO-GALTIERI, A HUSSAIN, R D TRIPP (LRL)
WATSON 63 PR 131 2248 M B WATSON, M FERRO-LUZZI, R D TRIPP (LRL) IJP
ALMEIDA 64 PL 9 204 S P ALMEIDA, G R LYNCH (CERN)
ARMENTER 65 PL 19 338 ARMENTEROS, F-LUZZI, * (CERN, HEID, SACLAY)
MUSGRAVE 65 NC 35 735 *PETMEZAS, * (BIRM, CERN, EPOL, LOIC, SACLAY)

BIRMINGHAM 66 PR 152 1148 BIRMINGHAM, GLASGOW, I. C., OXFORD, RUTHERFORD
DAHL 67 PR 163 1377 DAHL, HADY, HESS, KIRZ, MILLER (LRL)
DAUBER 67 PL 248 525 *MALAMUD, SCHLEIN, SLATER, STOKR (UCLA)
UHLIG 67 PR 155 1448 *CHARLTON, CORDON, GLASSER, YODH, * (UMD, NRL)
MAST 68 PR 21 1715 MAST, ALSTON, BANGERTER, GALTIERI * (LRL)
SCHEUER 68 NP 88 903 SABRE COLLAB. (SACL+AMST+BGNA+REHO+EPOL)

BURKHARDT 69 NP 814 106 *FILTHUTH+KLUGE+... (HEID+EFI+CERN+SACLAY)
CLINE 69 LNC 2 407 *LAUMANN+MAPP (WISC)
GALTIERI 69 LUND 352 BARBARO-GALTIERI, BANGERTER, MAST, TRIPP (LFL)
ALSO 70 DUKE 95 R D TRIPP (LRL)
BURKHARDT 71 NP 827 64 *FILTHUTH, KLUGE, OBERLACK+... (HEID+CERN+SACLAY)
COLLEY 71 NP 831 61 *COX, EASTWOOD, FRY+... (BIRM+EDIN+GLAS+LOIC)
KIM 71 PRL 27 356 J K KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP

CHAN 72 PRL 28 256 *BUT--SHAFFER, HERTZBACH, KOFLER+... (MASA, YALE)
MAST 73 PR D7 5 *ALSTON-GARNJOST, BANGERTER, +... (LBL) IJP
MAST2 73 PR D7 3212 *BANGERTER, ALSTON-GARNJOST, + (LBL) IJP
BERTHON 74 NC 21A 146 *BERTHON, TRISTRAM, + (CDFE+RHEL+SACL+STRB) U
CORDEN 75 NP 884 306 CORDEN, COX, CARTNELL, KEVON, ONEALE, * (BIRM)

MAST 76 LBL-4294 (PRPNT) MAST, ALSTON-GARNJOST, BANGERTER+ (LBL)
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP
PAPERS NOT REFERRED TO IN DATA CARDS
BERLEY 70 PR D1 1996 *YAMIN, KOFLER, MANN, MEISNER+ (BNL, MASA, YALE) IJP
GOLOWICH 74 PR D1 3861 EUGEN GOLOWICH (SLAC)

REFERENCES FOR Y*0(1600)
GALTIERI 63 PL 6 296 A BARBARO-GALTIERI, A HUSSAIN, R D TRIPP (LRL)
WATSON 63 PR 131 2248 M B WATSON, M FERRO-LUZZI, R D TRIPP (LRL) IJP
ALMEIDA 64 PL 9 204 S P ALMEIDA, G R LYNCH (CERN)
ARMENTER 65 PL 19 338 ARMENTEROS, F-LUZZI, * (CERN, HEID, SACLAY)
MUSGRAVE 65 NC 35 735 *PETMEZAS, * (BIRM, CERN, EPOL, LOIC, SACLAY)

BIRMINGHAM 66 PR 152 1148 BIRMINGHAM, GLASGOW, I. C., OXFORD, RUTHERFORD
DAHL 67 PR 163 1377 DAHL, HADY, HESS, KIRZ, MILLER (LRL)
DAUBER 67 PL 248 525 *MALAMUD, SCHLEIN, SLATER, STOKR (UCLA)
UHLIG 67 PR 155 1448 *CHARLTON, CORDON, GLASSER, YODH, * (UMD, NRL)
MAST 68 PR 21 1715 MAST, ALSTON, BANGERTER, GALTIERI * (LRL)
SCHEUER 68 NP 88 903 SABRE COLLAB. (SACL+AMST+BGNA+REHO+EPOL)

BURKHARDT 69 NP 814 106 *FILTHUTH+KLUGE+... (HEID+EFI+CERN+SACLAY)
CLINE 69 LNC 2 407 *LAUMANN+MAPP (WISC)
GALTIERI 69 LUND 352 BARBARO-GALTIERI, BANGERTER, MAST, TRIPP (LFL)
ALSO 70 DUKE 95 R D TRIPP (LRL)
BURKHARDT 71 NP 827 64 *FILTHUTH, KLUGE, OBERLACK+... (HEID+CERN+SACLAY)
COLLEY 71 NP 831 61 *COX, EASTWOOD, FRY+... (BIRM+EDIN+GLAS+LOIC)
KIM 71 PRL 27 356 J K KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP

REFERENCES FOR Y*0(1600)
KIM 71 PRL 27 356 J K KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP
LANGBEIN 72 NP 847 477 *WAGNER (MPIM) IJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP

REFERENCES FOR Y*0(1670)
GALTIERI 63 PL 6 296 A BARBARO-GALTIERI, A HUSSAIN, R D TRIPP (LRL)
WATSON 63 PR 131 2248 M B WATSON, M FERRO-LUZZI, R D TRIPP (LRL) IJP
ALMEIDA 64 PL 9 204 S P ALMEIDA, G R LYNCH (CERN)
ARMENTER 65 PL 19 338 ARMENTEROS, F-LUZZI, * (CERN, HEID, SACLAY)
MUSGRAVE 65 NC 35 735 *PETMEZAS, * (BIRM, CERN, EPOL, LOIC, SACLAY)

40 Y*0(1670) MASS (MEV)
M M (1666.0) OR (1675.0) BERLEY 65 HBC 0 K-P TO LAM ETA 7/66
M M THE FIRST VALUE ASSUMES THE BRANCHING RATIO INTO LAMBDA ETA IS SMALL, THE SECOND THAT IT IS LARGE. BECAUSE THE RESONANCE IS NEAR THE LAMBDA ETA THRESHOLD, THE BRANCHING RATIO AFFECTS THE MOMENTUM DEPENDENCE OF THE TOTAL WIDTH, AND THUS ALSO THE RESONANCE PARAMETERS OBTAINED BY FITTING TO THE DATA.
M N (1663.0) (3.0) ARMENT-1 68 HBC 0 ELASTIC, CH EXCH 11/68
M N (1678.0) (2.0) ARMENT-2 68 HBC 0 K-P TO SIGMA PI 11/68
M A 1674.0 (5.0) ARMENT-3 69 HBC 0 MULTICHANNEL 9/69
M N 1662.0 (3.0) ARMENT-4 69 HBC 0 ELAST, CH, EXC. ED 9/69
M N 1680.0 (1.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 9/69
M M 1674.0 (5.0) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
M M 1683.0 (5.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
M M 1670.0 (5.0) KIM 71 DPWA K-MATRIX ANAL. 3/71
M M 1640.0 (40.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M M 1700.0 (1.0) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
M M 1672.0 (1.0) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
M M 1665.0 (5.0) PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*
M M 1670.0 (5.0) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

40 Y*0(1670) WIDTH (MEV)
W M (22.0) OR (15.0) BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
W N (26.0) (8.0) ARMENT-1 68 HBC 0 SEE NOTE N ABOVE 11/68
W N (26.0) (5.0) ARMENT-2 68 HBC 0 9/69
W A 23.0 (3.0) ARMENT-3 69 HBC 0 9/69
W N 38.0 (15.0) ARMENT-4 69 HBC 0 ELAST, CH, EXC. ED 9/69
W A 35.0 (5.0) ARMENT-4 69 HBC 0 K-P TO SIG PI, ED 6/70
W N 31.0 (5.0) BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
W W 25.0 (5.0) GALTIERI 70 HBC 0 SIG PI, ED PWA 7/70
W W 35.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
W W 45.0 (20.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W W 65.0 (20.0) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
W W 19.0 (2.0) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
W W 19.0 (5.0) PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*
W W 45.0 (10.0) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

40 Y*0(1670) PARTIAL DECAY MODES
P1 Y*0(1670) INTO KBAR N 497+ 939
P2 Y*0(1670) INTO LAMBDA ETA 115+ 548
P3 Y*0(1670) INTO SIGMA PI 1189+ 139
P4 Y*0(1670) INTO SIGMA(1385) PI 139+1384

40 Y*0(1670) BRANCHING RATIOS
R1 Y*0(1670) INTO (KBAR N)/TOTAL (P1)
R1 P (0.14) (0.04) ARMENT-1 68 HBC 0 OLD DATA 11/68
R1 0.17 ARMENT-3 69 HBC 0 9/69
R1 P 0.14 (0.04) ARMENT-4 69 HBC 0 NEW DATA 9/69
R1 A (0.39) (0.05) CENFORTO 71 HBC 0 K-P, ELAST, CEX 6/70
R1 0.28 KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 0.35 (0.06) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 .36 (0.03) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
R1 .20 (0.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

40 Y*0(1670) FROM KBAR N TO LAMBDA ETA
R2 Y*0(1670) FROM KBAR N TO LAMBDA ETA SORT(P1*P2)
R2 M (0.20) OR 0.23 BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
R2 (0.26) ARMENT-3 69 HBC 0 9/69
R2 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 +.20 (.05) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
SEE THE NOTES ACCOMPANYING MASSES QUOTED
R3 Y*0(1670) FROM KBAR N TO SIGMA PI SORT(P1*P3)
R3 1 (-0.25) (0.06) ARMENT-2 68 HBC 0 OLD DATA 9/69
R3 1 (-0.27) ARMENT-3 69 HBC 0 9/69
R3 1 (-0.30) (0.03) ARMENT-4 69 HBC 0 NEW DATA 9/69
R3 1 PUBLISHED SIGN CHANGED TO AGREE WITH LUND 1969 CONVENTION (SEE TEXT) 10/74*

40 Y*0(1670) FROM KBAR N TO SIGMA(1385) PI SORT(P1*P4)
R4 Y*0(1670) FROM KBAR N TO SIGMA(1385) PI SORT(P1*P4)
R4 -18 .05 PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*

$\Lambda(1600)$ 101 Y*0(1600, JP=1/2+) I=0 P_{01}^1 1/76*

101 Y*0(1600) MASS (MEV) 1/76*
M 1 (1570.) KIM 71 DPWA K-MATRIX ANAL. 1/76*
M 1 POSSIBLE EFFECT IN SIGMA PI AND KBAR N CHANNELS.
M 1620.0 10.0 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
M 1573. 25. RLIC 76 DPWA KBAR N MULTICHNL 1/76*

101 Y*0(1600) WIDTH (MEV) 1/76*
W 1 (50.) KIM 71 DPWA K-MATRIX ANAL. 1/76*
W 60.0 10.0 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
W 147. 50. RLIC 76 DPWA KBAR N MULTICHNL 1/76*

101 Y*0(1600) PARTIAL DECAY MODES 1/76*
P1 Y*0(1600) INTO KBAR N 497+ 939
P2 Y*0(1600) INTO SIGMA PI 1197+ 139

101 Y*0(1600) BRANCHING RATIOS 1/76*
R1 Y*0(1600) INTO (KBAR N)/TOTAL (P1) 1/76*
R1 0.25 0.15 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
R1 .24 .04 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

101 Y*0(1600) FROM KBAR N INTO SIGMA PI SORT(P1*P2) 1/76*
R2 Y*0(1600) FROM KBAR N INTO SIGMA PI SORT(P1*P2) 1/76*
R2 0.28 0.09 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
R2 -.16 -.04 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

Baryons

$\Lambda(1670)$, $\Lambda(1690)$, $\Lambda(1800)$

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR $\Lambda(1670)$

BERLEY 65 PRL 15 641 +CCNNOLLY, HART, RAUM, STCNEHILL, + (BNL)IJP
ARMENT-1 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-2 68 NP 88 223 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-3 69 LUND PAPER 229 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
VALUES ARE QUOTED IN LEVI SETTI 69.

BIRMINGHAM 66 PR 152 1148 (BIRMINGHAM, GLASGOW, LOIC, OXFORD, RUTHERFD)
LEVISETTI 69 LUND 339 R LEVI SETTI (RAPPORTEUR) (CHICAGO)

$\Lambda(1690)$

55 $\Lambda(1690)$, JP=3/2- I=0 D03
SEE THE MINI-REVIEW AT THE START OF THE Λ LISTINGS.
THIS RESONANCE IS WELL ESTABLISHED.

55 $\Lambda(1690)$ MASS (MEV)

Table with columns for mass (M), width (W), and various experimental references (ARMENT-1, BARTLEY, etc.) and theoretical notes.

55 $\Lambda(1690)$ WIDTH (MEV)

Table with columns for width (W) and various experimental references (ARMENT-1, BARTLEY, etc.) and theoretical notes.

55 $\Lambda(1690)$ PARTIAL DECAY MODES

Table listing partial decay modes (P1, P2, P3, P4, P5, P6) and their corresponding decay masses.

55 $\Lambda(1690)$ BRANCHING RATIOS

THE SUM OF ALL THE QUOTED BRANCHING RATIOS IS MORE THAN 1.0. THE TWO-BODY RATIOS ARE FROM PARTIAL WAVE ANALYSES, AND THUS PROBABLY ARE MORE RELIABLE THAN THE THREE-BODY RATIOS, WHICH ARE DETERMINED FROM BUMPS IN CROSS SECTIONS.

Table with columns for branching ratios (R1, R1 M, R1 N, R1, R1, R1, R1, R1, R1, R1 N) and various experimental references (ARMENT-1, BUGG, etc.) and theoretical notes.

Table with columns for mass (M), width (W), and various experimental references (R2, R2 1, R2 1, R2 1, R2, R2, R2, R2, R2, R2, R2).

Table with columns for mass (M), width (W), and various experimental references (R3, R3, R3).

Table with columns for mass (M), width (W), and various experimental references (R4, R4).

Table with columns for mass (M), width (W), and various experimental references (R5, R5).

Table with columns for mass (M), width (W), and various experimental references (R6, R6).

REFERENCES FOR $\Lambda(1690)$

ARMENT-1 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-2 68 NP 88 216 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY) I
BARTLEY 68 PRL 21 1111 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
BUGG 68 PR 168 1466 ACHU,DOMO, GREENE, + (TIJFUS,US, BRANDEIS) I
ALSO 67 PRL 18 62 +GILMORE, KNIGHT, + (BIRM,CAVE,RHEL) I
CONFORTO 68 NP 88 265 DAVIES, DOWELL, + (BIRM,CAVE,RHEL) I
+PARMSEN, LASINSKI, + (CHICAGO,HEIDEL)IJP

PAPERS NOT REFERRED TO IN DATA CARDS
PREVOST 71 AMSTERDAM CONF + CHS COLLABORATION (CERN+HEID+SACL)

$\Lambda(1800)$

77 $\Lambda(1800)$, JP=1/2+ I=0 P01
SEE THE MINI-REVIEW AT THE START OF THE Λ LISTINGS.

THE EVIDENCE FOR THIS STATE IS SOMEWHAT CONFUSED. IT WAS FIRST SUGGESTED IN A PARTIAL WAVE ANALYSIS OF KBAR N DATA BY THE BEHAVIOUR OF THE P01 AMPLITUDE WHEN IT WAS PARAMETRIZED AS A TWO-STRAIGHT-LINE BACKGROUND (ARMENTEROS 68).

77 $\Lambda(1800)$ MASS (MEV)

Table with columns for mass (M), width (W), and various experimental references (ARMENTERO, BAILEY, etc.) and theoretical notes.

77 $\Lambda(1800)$ WIDTH (MEV)

Table with columns for width (W) and various experimental references (ARMENTERO, BAILEY, etc.) and theoretical notes.

77 $\Lambda(1800)$ PARTIAL DECAY MODES

Table listing partial decay modes (P1, P2, P3) and their corresponding decay masses.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1800)$, $\Lambda(1815)$

77 Y*(1800) BRANCHING RATIOS

Table with columns for particle name, mass, width, and branching ratios. Includes entries for R1, R2, R3, and R4 with various sub-labels like (P1), (P2), and (P3).

REFERENCES FOR Y*(1800)
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
DAVID SAAL BAILEY (LRL LIVERMORE) IJP
ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJP
A BARBARO-GALTIERI (LRL) IJP
J. K. KIM (HARV) IJP
J. K. KIM (HARV) IJP
*WAGNER (MPI) IJP

$\Lambda(1800)$

G₀₉

THIS NARROW RESONANCE SEEMS NECESSARY TO FIT A PEAK IN THE Λ COEFFICIENT OF THE K-P ANGULAR DISTRIBUTION. IT IS NOT REQUIRED IN ANY OTHER CHANNEL.

Table with columns for mass (MEV) and width (MEV). Values: 1808, 5, 27, 5.

Table with columns for partial decay modes. Values: 102, 5, 102.

Table with columns for partial decay modes and decay masses. Values: 497+939, 1197+139.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

REFERENCES FOR Y*(1800)
GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP

$\Lambda(1815)$

F₀₅

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THIS STATE IS WELL ESTABLISHED. MOST OF THE QUOTED ERRORS ARE STATISTICAL ONLY. THE SYSTEMATIC ERRORS DUE TO THE PARTICULAR PARAMETRIZATION USED IN THE P.W.A. ARE NOT INCLUDED. FOR THIS REASON WE DO NOT CALCULATE WEIGHTED AVERAGES FOR MASS AND WIDTH.

Table with columns for mass (MEV) and width (MEV). Values: 1813.0, 1816.0, 1817.0, 1819.0, 1825.0, 1819.0, 1830.0, 1820.0, 1818.0, 1810.0, 1823.0, 1819.0, 1822.0.

39 Y*(1815) WIDTH (MEV)

Table with columns for particle name, mass, width, and branching ratios. Includes entries for W, N, and R with various sub-labels like (P1), (P2), and (P3).

39 Y*(1815) PARTIAL DECAY MODES

Table with columns for partial decay modes and decay masses. Values: 497+939, 1189+139, 139+1384, 1192+139+139, 548+1115.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i , where $\delta P_i = \sqrt{(\delta P_i)^2 + (\delta P_j)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$. For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns for P1, P2, P3, P4, P5. Values: 0.6005+-0.0214, -0.154, 0.1155+-0.0080, -0.1541, 0.0795, 0.1390+-0.0261, -0.4139, 0.0216, -0.7638, 0.1292+-0.0308, -0.0642, 0.0331, 0.0099, -0.2505, 0.0153+-0.0086.

39 Y*(1815) BRANCHING RATIOS

ERRORS QUOTED BY EXPERIMENTERS DO NOT INCLUDE UNCERTAINTY DUE TO PARAMETRIZATION USED IN THE P.W.A. THEY SHOULD BE INCREASED.

Table with columns for branching ratios. Values: 0.62, 0.02, 0.65, 0.02, 0.58, 0.02, 0.63, 0.01, 0.52, 0.01, 0.47, 0.02, 0.57, 0.02, 0.601, 0.023, 0.614, 0.010, 0.601, 0.021.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

Table with columns for branching ratios. Values: 102, 5, 102.

REFERENCES FOR Y*(1815)
BIRGE 65 ATHENS CONF 296 +ELY, KALMUS, KERNAN, LOUIE, SAHOURIA, + (LRL) IJP
ARMENT-1 67 PL 248 198 ARMENTEROS, F LUZZI, + (CERN, HEIDEL, SACLAY) IJP
ARMENT-2 67 ZEIT PHYS 202 486 ARMENTEROS, F LUZZI, + (CERN, HEIDEL, SACLAY) IJP
BELL 67 PRL 19 936 R B BELL (LRL) IJP
ARMENT-3 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
ARMENT-4 68 NP 88 216 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHEL+BIRM+CAVE) I

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1815)$, $\Lambda(1830)$, $\Lambda(1860)$

BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
BRICMAN 70 PL 338 511 +FERRO-LUZZI, LAGNAUX (CERN)
COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJJP

CHAMBERL 62 PR 125 1696 CHAMBERLAIN, CROWE, KEEFE, KERTH, + (LRL) I
GALTIERI 63 PL 6 296 A BARBARO-GALTIERI, A HUSSAIN, RD TRIPP (LRL) IJJP
SODICKSON 64 PR 133 8757 SODICKSON, MANNELLI, FRISCH, WAHLIG (MIT(BNL)) J
HOLLEY 65 UCLR-16274 THESIS W R HOLLEY (LRL) J
BIRMINGHAM 66 PR 152 1148 BIRMINGHAM, GLASGOW, I.C., OXFORD, RUTHERFORD
+GIACOMELLI, KYCIA, LEONTIC, LUNDBY + (BNL) I
+HARMSEN, LEVI-SETTI, PREDAZZI + (EFI, LANL)
ARMENTEROS, FERRO-LUZZI + (CERN, HEID, SACLAY) IJJP
+HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJJP
LASINSKI 68 PR 163 1792 LASINSKI, LEVI SETTI, PREDAZZI (CHICAGO) JP
PREVOST 71 AMSTERDAM CONF + CHS COLLABORATION (CERN+HEID+SACL)

1830 56 Y*(1830, JP=5/2-) I=0 D05
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THE BEST EVIDENCE FOR THIS RESONANCE COMES FROM THE SIGMA PI CHANNEL. IT IS WELL ESTABLISHED.

1860 60 Y*(1860, JP=3/2+) I=0 P03
THE JP=3/2+ ASSIGNMENT IS CONSISTENT WITH ALL AVAILABLE DATA (INCLUDING POLARIZATION) AND RECENT PARTIAL WAVE ANALYSES. THE DOMINANT INELASTIC MODES REMAIN UNKNOWN. SEE ALSO Y*(2010) MINI-REVIEW.
60 Y*(1860) MASS (MEV)
M A F07 1864.0 2.0 ARMENTERO 68 DPWA 0 ELASTIC, CH EXCH 11/68
M N 1870.0 5.0 BUGG 68 CNTR 0 K-P TOTAL 7/68
M A F07 1871.0 6.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
M 1870.0 6.0 BRICMAN 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
M N 1883.0 10.0 CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
M 1 (1710.) 10.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1850.0 20.0 LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 2 (1868.) 20.0 LEA 73 DPWA MULTICHNL K-MTRX 9/73
M 1894. 10.0 HEMINGWA 75 DPWA 0 K-P TO KBAR N 11/75*
M 3 (1900.) 5.0 NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*
M 1900. 5.0 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

1830 56 Y*(1830) MASS (MEV)
M N 1827.0 (3.0) ARMENTERO 67 HBC 0 K-P TO SIGMA PI 8/67
M N 1837.0 (11.0) BELL 67 HBC 0 K-P TO SIGMA PI 11/67
M N 1807.0 (10.0) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
M 1840.0 (15.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
M N 1831.0 (5.0) CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
M 1830. KIM 71 DPWA K-MATRIX ANAL. 3/71
M K (1720.) KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1832.0 (5.0) KANE 72 DPWA 0 K-P TO PI SIG 10/71
M 1810.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 1825. (10.) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
M K POSSIBLE EFFECT MAINLY IN SIGMA PI. NOT CLEAR IF UNCORRELATED
M WITH THE 1830 EFFECT
M N ERROR STATIST. ONLY- NO ERROR DUE TO PARTICULAR P.W. ANAL. INCLUDED 1/71

1830 56 Y*(1830) WIDTH (MEV)
W 75.0 (9.0) ARMENTERO 67 HBC 0 K-P TO SIGMA PI 8/67
W 74.0 (18.0) BELL 67 HBC 0 K-P TO SIGMA PI 8/67
W 123.0 (32.0) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
W 150.0 (30.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
W 104.0 (35.0) CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
W 80. KIM 71 DPWA K-MATRIX ANAL. 3/71
W K (20.) KIM 71 DPWA K-MATRIX ANAL. 3/71
W 88.0 (10.0) KANE 72 DPWA 0 K-P TO PI SIG 10/71
W 60.0 (20.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 94. (10.) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
SEE THE NOTES ACCOMPANYING MASSES QUOTED

1830 56 Y*(1830) PARTIAL DECAY MODES
P1 Y*(1830) INTO KBAR N 497+ 939
P2 Y*(1830) INTO SIGMA PI 1189+ 139
P3 Y*(1830) INTO Y*(1385) PI D-WAVE 139+1384
P4 Y*(1830) INTO ETA LAMBDA 548+1115

1830 56 Y*(1830) BRANCHING RATIOS
R1 Y*(1830) INTO (KBAR N)/TOTAL (P1)
R1 0.09 (0.01) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
R1 0.03 (0.02) BRICMAN 70 DPWA SIGTOT, ELAS, CHEX 1/71
R1 0.05 (0.02) CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
R1 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 0.10 (0.03) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 .04 (0.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

1830 56 Y*(1830) BRANCHING RATIOS (CONT)
R2 Y*(1830) FROM KBAR N INTO SIGMA PI SQRTP(P1*P2)
R2 A (-0.15) (0.02) ARMENTERO 67 DPWA 0 K-P TO SIGMA PI 10/74*
R2 A PUBLISHED SIGN CHANGED TO AGREE WITH LUND 1969 CONVENTION (SEE TEXT) 10/74*
R2 0.19 (0.01) BELL 67 DPWA 0 K-P TO SIGMA PI 11/67
R2 -0.16 (0.03) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
R2 0.15 (0.01) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 -0.138 (0.0318) KANE 72 DPWA 0 K-P TO PI SIG 10/71
R2 0.27 (0.07) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R2 -0.17 (0.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

1830 56 Y*(1830) FROM KBAR N TO ETA LAMBDA SQRTP(P1*P4)
R3 -0.044 .020 RADER 73 MPWA 9/73
1830 56 Y*(1830) FROM KBAR N TO Y*(1385) PI D-WAVE SQRTP(P1*P3)
R4 +.13 .03 PREVOST 74 DPWA 0 K-N TO S(1385) IPI 10/74*

REFERENCES FOR Y*(1830)
ARMENTEROS, F.-LUZZI, + (CERN, HEIDEL, SACLAY) IJJP
R B BELL (LRL) IJJP
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJJP
+HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJJP
IS SUPERSEDED BY CONFORTO 71.
BRICMAN 70 PL 338 511 +FERRO-LUZZI, LAGNAUX (CERN)
GALTIERI 70 DUKE CONF 173 +BARBARO-GALTIERI (LRL) IJJP

CCNFORTO 71 NP 834 41 +LEVI SETTI, LASINSKI, OBERLACK + (EFI+HEID) IJJP
KIM 71 PRL 27 356 J. K. KIM (HARV) IJJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJJP
KANE 72 PR D5 1583 D. F. KANE (LBL) IJJP
LANGBEIN 72 NP 847 577 +WAGNER (MPIM) IJJP
RADER 73 NC 16A 178 +BARLOUTAUD, + (SACL+HEID+CERN+RHEL+CDEF)

1860 60 Y*(1860, JP=3/2+) I=0 P03
AVAILABLE DATA (INCLUDING POLARIZATION) AND RECENT PARTIAL WAVE ANALYSES. THE DOMINANT INELASTIC MODES REMAIN UNKNOWN. SEE ALSO Y*(2010) MINI-REVIEW.

60 Y*(1860) MASS (MEV)
M A F07 1864.0 2.0 ARMENTERO 68 DPWA 0 ELASTIC, CH EXCH 11/68
M N 1870.0 5.0 BUGG 68 CNTR 0 K-P TOTAL 7/68
M A F07 1871.0 6.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
M 1870.0 6.0 BRICMAN 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
M N 1883.0 10.0 CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
M 1 (1710.) 10.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1850.0 20.0 LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 2 (1868.) 20.0 LEA 73 DPWA MULTICHNL K-MTRX 9/73
M 1894. 10.0 HEMINGWA 75 DPWA 0 K-P TO KBAR N 11/75*
M 3 (1900.) 5.0 NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*
M 1900. 5.0 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

60 Y*(1860) WIDTH (MEV)
W A F07 39.0 7.0 ARMENTERO 68 DPWA 0 ELASTIC, CH EXCH 11/68
W N 40.0 10.0 BUGG 68 CNTR 0 K-P TOTAL 7/68
W A F07 24.0 15.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
W N 37.0 10.0 BRICMAN 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
W 80.0 20.0 CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
W 1 (20.) KIM 71 DPWA K-MATRIX ANAL. 3/71
W 125.0 20.0 LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 2 (323.8) 10.0 LANGBEIN 73 DPWA MULTICHNL K-MTRX 9/73
W 107. 10.0 HEMINGWA 75 DPWA 0 K-P TO KBAR N 11/75*
W 3 (100.) 10.0 NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*
W 72. 10.0 RLIC 76 DPWA KBAR N MULTICHNL 1/76*
SEE THE NOTES ACCOMPANYING MASSES QUOTED
AVERAGE MEANINGLESS (SCALE FACTOR = 2.5)

60 Y*(1860) PARTIAL DECAY MODES
P1 Y*(1860) INTO KBAR N 497+ 939
P2 Y*(1860) INTO SIGMA PI 1189+ 139
P3 Y*(1860) INTO LAMBDA OMEGA 1115+ 783

60 Y*(1860) BRANCHING RATIOS
R1 Y*(1860) INTO (KBAR N)/TOTAL (P1)
R1 A F07 0.12 0.02 ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
R1 (J+1/2)P1=0.40 BUGG 68 CNTR 0 7/68
R1 A F07 0.07 0.02 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
R1 0.14 0.02 BRICMAN 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
R1 N 0.25 0.03 CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
R1 0.37 0.05 LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 2 (.32) LEA 73 DPWA MULTICHNL K-MTRX 9/73
R1 .24 .04 HEMINGWA 75 DPWA 0 K-P TO KBAR N 11/75*
R1 .18 .02 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

60 Y*(1860) BRANCHING RATIOS (CONT)
R2 Y*(1860) INTO SIGMA PI (P2)
R2 P PROBABLY SEEN GALTIERI 68 CBC 0 K-N TO SIG PI PI 11/68
R2 (0.03) OR LESS LANGBEIN 72 IPWA MULTICHANNEL 12/72
R2 P POSSIBLY THIS BUMP SEEN AT 1840+-10 MEV WITH A WIDTH OF 35+-10 MEV
R2 IS THE Y*(1830), WHICH DECAYS STRONGLY TO SIGMA PI. HOWEVER THE
R2 NARROW WIDTH HERE ARGUES FOR ITS BEING THE Y*(1860).
R3 Y*(1860) FROM KBAR N TO SIGMA PI SQRTP(P1*P2) 9/73
R3 2 (-.15) .03 LEA 73 DPWA MULTICHNL K-MTRX 9/73
R3 -.09 .03 LANGBEIN 72 IPWA KBAR N MULTICHNL 1/76*

60 Y*(1860) FROM KBAR N TO LAMBDA OMEGA SQRTP(P1*P3) 1/76*
R4 3 (-0.32) NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*

REFERENCES FOR Y*(1860)
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJJP
+GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I
BARBARO-GALTIERI, MATISON, + (LRL, SLAC)
BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
BRICMAN 70 PL 338 511 +FERRO-LUZZI, LAGNAUX (CERN)
CONFORTO 71 NP 834 41 +LEVI SETTI, LASINSKI, OBERLACK + (EFI+HEID) IJJP
KIM 71 PRL 27 356 J. K. KIM (HARV) IJJP
ALSO DUKE 161 J. K. KIM (HARV) IJJP
LANGBEIN 72 NP 847 477 +WAGNER (MPIM) IJJP

Baryons

$\Lambda(2100)$, $\Lambda(2110)$

Data Card Listings

For notation, see key at front of Listings.

41 $\Lambda(2100)$ WIDTH (MEV)

W	(145.0)	WCHL	66 HBC		7/66	
W	(80.0)	BURGUN	68 DPWA	0 K-P TO XI K	10/69	
W	140.0 (15.0)	BERTHONI	70 DPWA	0 K-P TO SIGMA PI	10/70	
W	60.0 (25.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA PI	7/70	
W	(170.) TO 300.	LITCHFIE	71 DPWA	K-P TO KBAR N	10/71	
W	B LARGER VALUE CORRESPONDS TO PURE B.W. LOWER VALUE TO B.W. + BCKGRD					
W	140.0 (50.0) (30.0)	LITCHFIE	71 DPWA	K-P TO SIG PI	10/71	
W	1	208. TO 229.	BRANDSTE	72 DPWA	0 K-P TO LAM. OMG.	1/74
W	144.0 (26.0)	KANE	72 DPWA	0 K-P TO PI SIG	10/71	
W	241. (30.)	HEMINGWA	75 DPWA	0 K-P TO KBAR N	11/75*	
W	2	244. OR 302.	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.	11/75*
W	250. (30.1)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*	

SEE THE NOTES ACCOMPANYING MASSES QUOTED

41 $\Lambda(2100)$ PARTIAL DECAY MODES

P1	$\Lambda(2100)$ INTO KBAR N	497* 939	DECAY MASSES
P2	$\Lambda(2100)$ INTO SIGMA PI	1197* 139	
P3	$\Lambda(2100)$ INTO XI K	1321* 497	
P4	$\Lambda(2100)$ INTO LAMBDA OMEGA	1115* 783	
P5	$\Lambda(2100)$ INTO ETA LAMBDA	548*1115	

41 $\Lambda(2100)$ BRANCHING RATIOS

R1	$\Lambda(2100)$ INTO (KBAR N)/TOTAL	(P1)				
R1	(0.25)	WCHL	66 HBC		7/66	
R1	D (0.33)	DAUM	68 CNTR	K-P ELA,POL,SIGT	7/70	
R1	0.30 (0.03)	LITCHFIE	71 DPWA	K-P TO KBAR N	10/71	
R1	1	31 (0.03)	HEMINGWA	75 DPWA	0 K-P TO KBAR N	11/75*
R1	1	30 (0.03)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*
R1	D DAUM 68 ASSUMES (J=1/2)*X VALUE SEEN IN TOTAL CRCS SECTION.					

$\Lambda(2100)$ FROM KBAR N INTO SIGMA PI

R2	$\Lambda(2100)$ FROM KBAR N INTO SIGMA PI	SQRT(P1*P2)			
R2	L (+0.16) (0.02)	BERTHONI	70 DPWA	0 K-P TO SIGMA PI	10/70
R2	+0.06 (0.03)	GALTIERI	70 DPWA	0 K-P TO SIGMA PI	7/70
R2	L 0.16 (0.03)	LITCHFIE	71 DPWA	K-P TO SIG PI	10/71
R2	+0.096 (0.037)	KANE	72 DPWA	0 K-P TO PI SIG	10/71
R2	+1.2 (0.04)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*

$\Lambda(2100)$ FROM KBAR N TO XI K

R3	$\Lambda(2100)$ FROM KBAR N TO XI K	SQRT(P1*P3)			
R3	(0.05)	TRIPP	67 RVUE	0 K-P TO XI K	8/67
R3	B (0.09) (0.01)	BURGUN	68 DPWA	0 K-P TO XI K	10/69
R3	(0.003)	MULLER	69 DPWA	0	7/70
R3	0.035 0.018	LITCHFIE	71 DPWA	K-P TO XI K	3/72
R3	B BURGUN 68 UPDATED BY LITCHFIE 71, WHO TAKES SOLUTION C OF BURGUN				3/72

$\Lambda(2100)$ FROM KBAR N INTO LAMBDA OMEGA

R4	$\Lambda(2100)$ FROM KBAR N INTO LAMBDA OMEGA	SQRT(P1*P4)			
R4	1 (0.5) 0.11	BRANDSTE	72 DPWA	0 K-P TO LAM. OMG.	1/74
R4	2 (1.22)OR .154	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.	11/75*

$\Lambda(2100)$ FROM KBAR N TO ETA LAMBDA

R5	$\Lambda(2100)$ FROM KBAR N TO ETA LAMBDA	SQRT(P1*P5)			
R5	-0.50 0.20	RADER	73 MPWA		9/73

SEE THE NOTES ACCOMPANYING MASSES QUOTED

 REFERENCES FOR $\Lambda(2100)$

WCHL 66 PRL 17 107 C G WOHL, F T SOLMITZ, M L STEVENSON (LRL)JJP
 TRIPP 67 NP 83 10 + LRL,SLC,CERN,HEIDEL,SACLAY
 BURGUN 68 NP 88 447 +MEYER,PAULI, + (SACLAY,COLFRANCE,RHEL)
 DAUM 68 NP 87 19 +ERNE, LAGNAUX, SENS, STEUER, UDD (CERN)JP
 CONFIRMS THE SPIN-PARITY ASSIGNMENT.
 MULLER 69 THESIS,UCRL 19372 R A MULLER (LRL)

BERTHONI 70 NP 824 417 +VRANA, BUTTERNORTH, + (CCEF, RHEL, SACLAY)JJP
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LFL)JJP
 LITCHFIE 71 NP 830 125 LITCHFIE,....+LESQUOY,+. (RHEL+CDEF+SACL)JJP

BRANDSTE 72 NP 839 13 BRANSTETTER,....+TALLINI (RHEL,CDEF,SACL) JJP
 KANE 72 PR D5 1583 D F KANE (LBL)JJP
 RADER 73 NC 16A 178 +BARLOUTAUD, + (SACL+HEID+CERN+RHEL+CDEF)

HEMINGWA 75 NP 891 12 +HEMINGWAY,EADES,HARMSEN+ (CERN,HEID,MPIM)JJP
 NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)JJP
 RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCIPHERSON,ROSS+ (RHEL+LOIC)JJP

$\Lambda(2110)$

F⁰⁵

BERTHONI 70 FIND EITHER F05 OR D05 POSSIBLE IN THE SIG PI CHANNEL, WITH F05 SLIGHTLY PREFERRED. IN THE KBAR N CHANNEL, LITCHFIE 71 (SAME GROUP) FIND ONLY D05. AS USUAL, THE STATISTICS ARE MUCH BETTER IN THE ELASTIC CHANNEL. ALTHOUGH KANE 72 FINDS AN F05 EFFECT, THE UNUSUALLY BROAD WIDTH MAY INVALIDATE A RESONANT INTERPRETATION. HOWEVER RLIC 76 AND BELLEFCN 76 ALSO FIND AN F05. THE WEIGHT OF THE EVIDENCE IS THUS IN FAVOR OF F05. SEE ALSO THE $\Lambda(2100)$ MINI-REVIEW.

35 $\Lambda(2110)$ MASS (MEV)

M	(2110.)	(10.)	BERTHONI	70 DPWA	- K-P TO SIG PI	1/71
M	D05 2140.	40.	LITCHFIE	71 DPWA	K-P TO KBAR N	10/71
M	A (2141.0)	(6.0)	KANE	72 DPWA	0 K-P TO PI SIG	10/71
M	1 (2103.)	(20.)	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.	1/76*
M	(2140.)	(20.)	BELLEFCN	76 DPWA	0 K-P TO SIG PI	1/76*
M	(2100.)	(50.)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*
M	A RESONANCE OUTSIDE RANGE OF DATA.					
M	1 FOUND IN ONE OF TWO BEST SOLUTIONS.					1/76*

35 $\Lambda(2110)$ WIDTH (MEV)

W	(185.)	(30.)	BERTHONI	70 DPWA	- K-P TO SIG PI	1/71
W	D05 120.	40.	LITCHFIE	71 DPWA	K-P TO KBAR N	10/71
W	A (504.0)	(10.0)	KANE	72 DPWA	0 K-P TO PI SIG	10/71
W	1 (391.)	(20.)	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.	1/76*
W	(140.)	(20.)	BELLEFCN	76 DPWA	0 K-P TO SIG PI	1/76*
W	(200.)	(50.)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*

35 $\Lambda(2110)$ PARTIAL DECAY MODES

P1	$\Lambda(2110)$ INTO KBAR N	497* 939	DECAY MASSES
P2	$\Lambda(2110)$ INTO SIGMA PI	1197* 139	
P3	$\Lambda(2110)$ INTO LAMBDA OMEGA	1115* 783	

35 $\Lambda(2110)$ BRANCHING RATIOS

R1	$\Lambda(2110)$ FROM KBAR N TO SIGMA PI	SQRT(P1*P2)			
R1	A (+0.156) (0.013)	BERTHONI	70 DPWA	K-P TO SIG PI	1/71
R1	(+0.14) (0.01)	KANE	72 DPWA	0 K-P TO PI SIG	10/71
R1	(+0.10) (0.03)	BELLEFCN	76 DPWA	0 K-P TO SIG PI	1/76*
R1		RLIC	76 DPWA	KBAR N MULTICHNL	1/76*

$\Lambda(2110)$ INTO (KBAR N)/TOTAL

R2	$\Lambda(2110)$ INTO (KBAR N)/TOTAL	(P1)			
R2	D05 0.14 0.04	LITCHFIE	71 DPWA	K-P TO KBAR N	10/71
R2	(0.7) (0.03)	RLIC	76 DPWA	KBAR N MULTICHNL	1/76*

$\Lambda(2110)$ FROM KBAR N INTO LAMBDA OMEGA

R3	$\Lambda(2110)$ FROM KBAR N INTO LAMBDA OMEGA	SQRT(P1*P3)			
R3	1 (0.112)	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.	1/76*

 REFERENCES FOR $\Lambda(2110)$
 BERTHONI 70 NP 824 417 +VRANA,BUTTERNORTH,+. (CDEF,RHEL,SACLAY)JJP
 LITCHFIE 71 NP 830 125 LITCHFIE,....+LESQUOY,+. (RHEL+CDEF+SACL)JJP
 KANE 72 PR D5 1583 D F KANE (LBL)JJP
 NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)JJP
 BELLEFCN 76 SBMTD. TO NP DE BELLEFCN,BILLOIR,BERTHON+ (CDEF+SACL)JJP
 RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCIPHERSON,ROSS+ (RHEL+LOIC)JJP

2100 MEV REGION - PRODUCTION AND σ_{TOTAL} EXP'TS

25 $\Lambda(2100)$, JP= 1 I=0 PRODUCTION EXPERIMENTS
 SEE THE MINI-REVIEW AT THE START OF THE Λ LISTINGS.

SEE THE NOTE TO THE G07 $\Lambda(2100)$, WHICH PRECEDES THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE $\Lambda(2100)$, BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

25 $\Lambda(2100)$ MASS (MEV) (PROD. EXP.)

M	(2097.0)	(6.0)	BOCK	65 HBC	PBAR P 5.7 BEV/C	7/66
M	2100.0	(7.0)	BUGG	68 CNTR	K-P, D TOTAL	6/68
M	2121.0	(5.0)	BRICMAN	70 CNTR	0 TOTAL AND CH EX	6/70
M	2107.0	(10.0)	COOL	70 CNTR	K-P, D TOTAL	10/70
M	(2135.0)	(20.0)	LU	70 CNTR	0 GAMMA P TO K+ Λ	1/71

25 $\Lambda(2100)$ WIDTH (MEV) (PROD. EXP.)

W	(24.0)	(14.0)	(24.0)	BOCK	65 HBC	INTO KBAR N (P1)	7/66
W	140.0	(15.0)		BUGG	68 CNTR	0 TOTAL AND CH EX	6/68
W	147.0	(15.0)		BRICMAN	70 CNTR	0 TOTAL AND CH EX	6/70
W	185.0			COOL	70 CNTR	0 TOTAL	10/70
W	(40.0)			LU	70 CNTR	0 GAMMA P TO K+ Λ	1/71

25 $\Lambda(2100)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$\Lambda(2100)$ INTO KBAR N	497* 939	DECAY MASSES
P2	$\Lambda(2100)$ INTO KBAR N PI	497* 939* 139	
P3	$\Lambda(2100)$ INTO LAMBDA ETA	1115* 548	
P4	$\Lambda(2100)$ INTO LAMBDA OMEGA	1115* 783	

25 $\Lambda(2100)$ BRANCHING RATIOS (PROD. EXP.)

R1	$\Lambda(2100)$ INTO (KBAR N)/TOTAL	(P1)			
R1	THESE VALUES OF ELASTICITIES ASSUME J=7/2 --				
R1	0.305	BUGG	68 CNTR	0 TOTAL AND CH EX	6/68
R1	0.24 (0.02)	BRICMAN	70 CNTR	0 TOTAL AND CH EX	6/70
R1	0.4	COOL	70 CNTR	K-P, D TOTAL	10/70

$\Lambda(2100)$ INTO KBAR N PI

R2	$\Lambda(2100)$ INTO KBAR N PI	(P2)			
R2	SEEN	BOCK	65 HBC		

$\Lambda(2100)$ FROM KBAR N INTO LAMBDA ETA

R3	$\Lambda(2100)$ FROM KBAR N INTO LAMBDA ETA	SQRT(P1*P3)			
R3	(0.09) OR LESS	FLATTE 2	67 HBC	0 K-P TO LAM ETA	6/68

$\Lambda(2100)$ INTO (LAMBDA OMEGA)/TOTAL

R4	$\Lambda(2100)$ INTO (LAMBDA OMEGA)/TOTAL	(P4)			
R4	(0.1) OR LESS	FLATTE 1	67 HBC	0 K-P TO LAM OMEGA	8/67

 REFERENCES FOR $\Lambda(2100)$ (PROD. EXP.)
 BOCK 65 PL 17 166 +COOPER,FRENCH,KINSON, + (CERN,SACLAY)
 FLATTE 1 67 PR 155 1517 S M FLATTE (LRL)
 FLATTE 2 67 PR 163 1441 S M FLATTE, C G WOHL (LRL)
 BUGG 68 PR 168 1466 +GILMORE,KNIGHT, + (RHEL,BIRM,CAVE) I

BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN,CAEN,SACLAY)
 COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 LU 70 PR D2 1846 +GREENBERG, HUGHES, MINEHART, MORI, + (YALE)

PAPERS NOT REFERRED TO IN DATA CARDS
 COOL 66 PRL 16 1228 +GIACOMELLI,KYCIA,LEGNTIC,LI,LUNDBY, + (BNL) I
 SUPERSEDED BY COOL 70.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(2350)$, $\Lambda(2585)$, Σ^+ , Σ^- , Σ^0 , $\Sigma(1385)$

$\Lambda(2350)$ BUMPS

42 Y*(2350, JP=) I=0 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
DAUM 68 FAVORS JP=7/2- OR 9/2+. BRICHAN 70 FAVORS 9/2+.
LASINSKI 71 SUGGESTS THREE STATES IN THIS REGION
USING A PGMERK + RESONANCES MODEL. THERE IS NOW ALSO
ONE FORMATION EXPERIMENT WHICH WE INCLUDE HERE, DE BELLEFCN 76, WHICH
FINDS 9/2- FRGM A DPWA OF KBAR N TO SIGMA P1.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Bugg, Brichan, Cool, and Bellefcn data for Lambda(2350).

Table with 4 columns: W, mass (MeV), width (MeV), and production experiments. Rows include Bugg, Brichan, Cool, and Bellefcn data for Lambda(2350).

Table with 2 columns: P1, P2 and decay masses. Rows show decay into Kbar N and Sigma P1.

Table with 4 columns: R1, R2, branching ratios, and production experiments. Rows show branching ratios into Kbar N and Sigma P1.

REFERENCES FOR Y*(2350) (PROD. EXP.)
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I
DAUM 68 NP 87 19 +ERNE, LAGNAUX, SENS, STEUER, UDO (CERN) J P
BRICHAN 70 PL 318 152 +FERRO LUZZI, PERRERAU, + (CERN, CAEN, SACLAY) I
COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
LU 70 PR D2 1846 +GREENBERG, HUGES, MINEHART, MORI, + (YALE)

$\Lambda(2585)$ BUMPS

7 Y*(2585, JP=) I=0 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Abrams and Lu data for Lambda(2585).

Table with 4 columns: W, mass (MeV), width (MeV), and production experiments. Rows include Abrams and Lu data for Lambda(2585).

Table with 2 columns: P1, P2 and decay masses. Rows show decay into Kbar N.

Table with 4 columns: R1, R2, branching ratios, and production experiments. Rows show branching ratios into Kbar N and Sigma P1.

REFERENCES FOR Y*(2585) (PROD. EXP.)
COOL 66 PRL 16 1228 +GIACOMELLI, KYCIA, LEONTIC, LUNDBY + (BNL) I
ABRAMS 70 PR D1 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
BRICHAN 70 PL 318 152 +FERRO LUZZI, PERRERAU, + (CERN, CAEN, SACLAY) I
LU 70 PR D2 1846 +GREENBERG, HUGES, MINEHART, MORI, + (YALE)

S=-1 I=1 HYPERON STATES (Σ)

Σ^+ 19 SIGMA+(1189, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

Σ^- 20 SIGMA-(1198, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

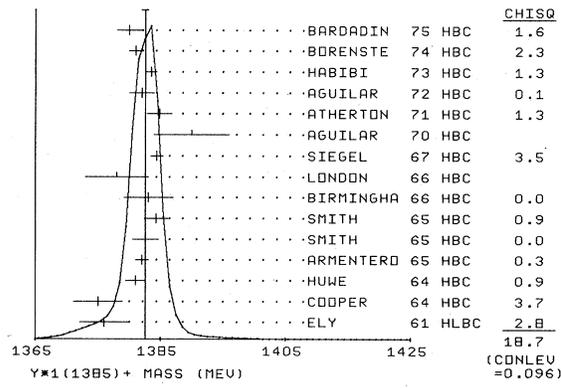
Σ^0 21 SIGMA0(1193, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

$\Sigma(1385)$ 43 Y*(1385, JP=3/2+) I=1
FOR DISCUSSION OF INCONSISTENCY OF ERRORS AND OUR
MODIFICATIONS, SEE NUTE CN K*(892)

FOR THE TABLES WE USE ONLY THE UNSTARRED DATA, WHICH
ATTEMPTS TO OBTAIN THE SEPARATE CHARGE-STATE MASSES AN
WIDTHS. SEE HOWEVER THE IDEOGRAMS INSERTED IN LISTING
THESE INDICATE SERVICUS SYSTEMATICS, PERHAPS ARISING FROM INTERFERENCE E
FFECTS THAT CHANGE WITH PRODUCTION MECHANISM AND BEAM MOMENTUM.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Alston, Borge, Martin, Colley, Baltay, Musgrave, Ammann, Atherton, Curtis, Thmas, Borenste, and others for Sigma(1385).

WEIGHTED AVERAGE = 1382.53 ± 0.50
ERROR SCALED BY 1.2



Baryons

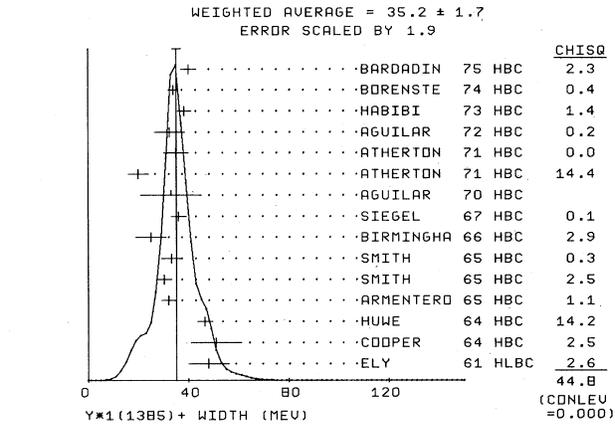
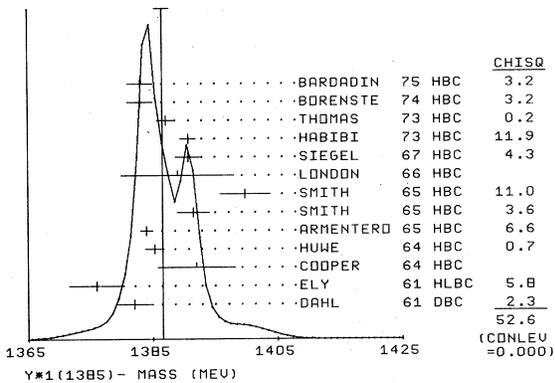
$\Sigma(1385)$

Data Card Listings

For notation, see key at front of Listings.

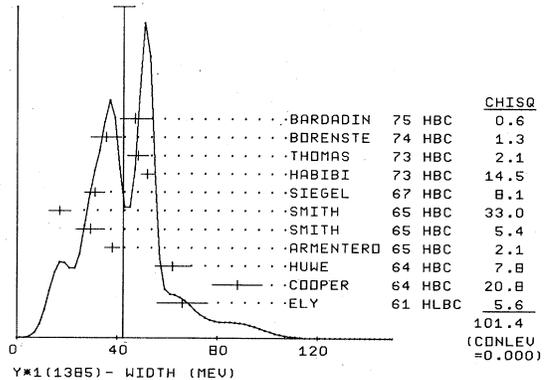
M-	93	1382.0	3.0	DAHL	61 DBC	-	K-C 0.45 BEV/C		
M-	E 224	1376.0	4.4	ELY	61 HLBC	-			
M-	E	200	1392.0	6.2	COOPER	64 HBC	-	ERROR OF 3.0 ENLARGED BY US, BECAUSE < STATIST. ERR. 10/69	
M-	1086	1385.3	1.5	HUWE	64 HBC	-			
M-	1380	1384.0	1.0	ARMENTERO	65 HBC	-			
M-	S 120	1391.5	2.6	SMITH	65 HBC	-	K-P 1.8 BEV/C	9/66	
M-	S 58	1399.8	4.0	SMITH	65 HBC	-	K-P 1.95 BEV/C	9/66	
M-	S	1389.0	9.0	LCNDON	66 HBC	-	K-P AT 2.1 GEV/C	10/69	
M-	370	1390.7	2.0	SIEGEL	67 HBC	-			
M-	1900	1390.7	1.2	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73	
M-	722	1387.1	1.3	THOMAS	73 HBC	-	PI-P TO PI-K+LM	9/73	
M-	3060	(1389.1)	(1.1)	BERTHON	74 HBC	0	QUASI 2 BODY CS	10/74*	
M-	2303	1385.1	2.1	BORENSTE	74 HBC	+OK-P TO(1385)+PIS		10/74*	
M-	1383	1383	2.1	BARDADIN	75 HBC	-	K-P 14.3 GEV/C	1/76*	
M-	AVG	1386.6	1.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.3)					
M-	STUDENT	1386.15	0.80	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)					

WEIGHTED AVERAGE = 1386.6 ± 1.2
ERROR SCALED BY 2.3



W-	(40.0)		DAHL	61 DBC	-			
W-	66.0	10.0	ELY	61 HLBC	-			
W-	88.0	10.0	COOPER	64 HBC	-			
W-	82.0	7.0	HUWE	64 HBC	-			
W-	38.0	3.0	ARMENTERO	65 HBC	-			
W-	29.2	5.7	SMITH	65 HBC	-	K-P 1.80 BEV/C	9/66	
W-	17.1	4.4	SMITH	65 HBC	-	K-P 1.95 BEV/C	9/66	
W-	370	31.0	SIEGEL	67 HBC	-	K-P AT 2.1 GEV/C	10/69	
W-	1900	51.9	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73	
W-	722	48.2	THOMAS	73 HBC	-	PI-P TO PI-K+LM	9/73	
W- 1	2303	35.5	BORENSTE	74 HBC	+OK-P TO(1385)+PIS		10/74*	
W-		47.1	BARDADIN	75 HBC	-	K-P 14.3 GEV/C	1/76*	
W-	AVG	42.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.2)					
W-	STUDENT	42.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)					

WEIGHTED AVERAGE = 42.4 ± 4.2
ERROR SCALED BY 3.2



D R	(0.0)	(4.2)	ELY	61 HLBC	+K-P 1.11 BEV/C	8/66	
D R	(17.1)	(7.1)	COOPER	64 HBC	-	10/69	
D R	(4.3)	(2.2)	HUWE	64 HBC	+K-P 1.22 BEV/C	8/66	
D R	(2.0)	(1.5)	ARMENTERO	65 HBC	+K-P 1.9-1.2 BEV/C	8/66	
D R	(7.2)	(2.1)	SMITH	65 HBC	+K-P 1.8 BEV/C	9/66	
D R	(17.2)	(2.0)	SMITH	65 HBC	+K-P 1.95 BEV/C	9/66	
D R	(11.0)	(9.0)	LONDON	66 HBC	+K-P 2.24 BEV/C	8/66	
D R	9.0	6.0	LCNDON	66 HBC	+ LAMBDA 3 PI EVTS	7/66	
D R	(6.3)	(2.0)	SIEGEL	67 HBC	K-P AT 2.1 GEV/C	10/69	
D R	(7.2)	(1.4)	HABIBI	73 HBC	K-P TO 2PI LAM	9/73	
D R	REDUNDANT WITH DATA IN MASS LISTING.						

43 Y*1(1385) WIDTH (MEV)

W	(64.0)		ALSTEN	60 HBC	+			
W	(40.0)		BERGE	61 HBC	+			
W	(20.0)	OR LESS	MARTIN	61 HBC	+			
W	(80.0)	(10.0)	COLLEY	62 HLBC	-			
W	(30.0)	(9.0)	CURTIS	63 OSPK	0			
W	(26.0)	(5.0)	BALTAY	65 HBC	+		7/66	
W	(38.0)	(9.0)	MUSGRAVE	65 HBC	+0		7/66	
W	61.	10.	AMMANN	73 DBC	K-N 4.5 GEV/C	9/73		
W	89.	23.	ATHERTON	75 HBC	PBAR P 5.7 GEV/C	1/76*		
W	AVG	65.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)					
W	STUDENT	65.2	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					

W0	330	39.3	6.1	THOMAS	73 HBC	PI-P TO PI0K0LM	9/73
W0	3100 CONSISTENT WITH W0=W+*W-			BORENSTE	74 HBC	+OK-P TO(1385)+PIS	10/74*

W+	48.0	8.0	ELY	61 HLBC	+			
W+	51.0	10.0	COOPER	64 HBC	+			
W+	46.5	3.0	HUWE	64 HBC	+			
W+	32.0	3.0	ARMENTERO	65 HBC	+			
W+	30.3	3.1	SMITH	65 HBC	+	K-P 1.8 BEV/C	9/66	
W+	33.1	3.8	SMITH	65 HBC	+	K-P 1.95 BEV/C	9/66	
W+	40	25.0	BIRMINGHA	66 HBC	+	K-P 3.5	9/67	
W+	1260	36.0	SIEGEL	67 HBC	+	K-P AT 2.1 GEV/C	10/69	
W+	33.	12.0	AGUILAR	70 HBC	+	K-P 4 GEV/SIG.PI	5/70	
W+ T	40	20.	ATHERTON	71 HBC	+	LAM PI+ + C.C.	10/71	
W+ R	FIT B.W. + PHASE SPACE BCKGRD							
W+ R	40	35.	5.	ATHERTON	71 HBC	LAM PI+ + C.C.	10/71	
W+	32.5	6.0	AGUILAR	72 HBC	+	K-P TO LAM+PIS	10/74*	
W+	2300	38.3	HABIBI	73 HBC	+	K-P TO 2PI LAM	9/73	
W+ 1	6846	34.	BORENSTE	74 HBC	+	+OK-P TO(1385)+PIS	10/74*	
W+ 1			BARDADIN	75 HBC	+	K-P 14.3 GEV/C	1/76*	
W+ 1	RESULTS FROM LAM PI+ PI- AND LAM PI+ PI- PIO COMBINED BY US.							
W+ R	FIT B.W. AND NO BCKGRD							
W+ R	40	35.2	1.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)				
W+ R	AVG	35.2	1.2	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)				
W+ R	STUDENT	35.0						

43 Y*1(1385) REAL PART OF POLE POSITION				4/75*	
RE+	1379.0	1.0	LICHTENB 74	+ EXTRAP HABIBI73	4/75*
RE-	1383.0	1.0	LICHTENB 74	- EXTRAP HABIBI73	4/75*

43 Y*1(1385) IMAGINARY PART OF POLE POSITION				4/75*	
IM+	17.5	1.5	LICHTENB 74	+ EXTRAP HABIBI73	4/75*
IM-	22.5	1.5	LICHTENB 74	- EXTRAP HABIBI73	4/75*

43 Y*1(1385) PARTIAL DECAY MODES				
P1	Y*1(1385) INTO LAMBDA PI		DECAY MASSES	
P2	Y*1(1385) INTO SIGMA PI		1115+ 139	
P3	Y*1(1385) INTO LAMBDA GAMMA		1115+ 0	
P4	Y*1(1385) INTO KBAR N		493+ 938	
P5	Y*1(1385) INTO SIGMA GAMMA		1197+ 0	

Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(1385), Σ(1440), Σ(1480), Σ(1580)

43 Y*(1385) BRANCHING RATIOS

Table with columns for particle name, branching ratio, and reference. Includes entries for Y*(1385) INTO (SIGMA PI)/(LAMBDA PI) and Y*(1385) INTO LAMBDA GAMMA.

WEIGHTED AVERAGE = 0.140 ± 0.015
ERROR SCALED BY 1.0

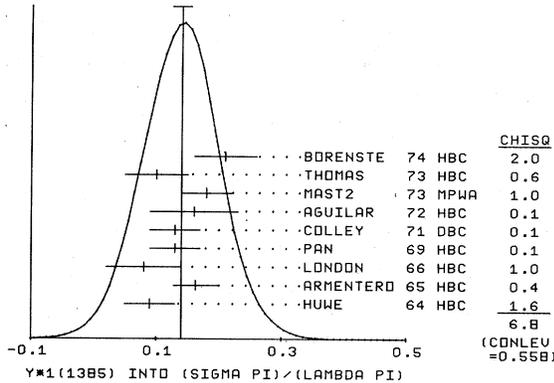


Table with columns for particle name, branching ratio, and reference. Includes entries for Y*(1385) FROM KBAR N TO LAMBDA PI and Y*(1385) INTO LAMBDA GAMMA/(LAMBDA PI).

REFERENCES FOR Y*(1385)

Large table of references for Y*(1385) from various experiments and institutions, including ALSTON, BASTIEN, BERGHE, etc.

Σ(1440) BUMPS

80 Y*(1440, JP= 1/2-) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

CLINE 68 FIND A NARROW PEAK AT 1440 MEV (JUST ABOVE THE KBAR N THRESHOLD) IN THE LAMBDA PI INVARIANT MASS FOR K- D TO LAMBDA PI- P EVENTS. THEY DISCUSS ALTERNATE INTERPRETATIONS - THAT IT IS A RESONANCE OR A KINEMATIC EFFECT. IN CLINE 68 THE K- BEAM MOMENTUM IS 0.4 GEV/C. IN A STUDY OF THE SAME REACTION WITH A MOMENTUM OF 1.1 GEV/C, ALEXANDER 69 FIND NO PEAK. IN ADDITION, THEY ARE ABLE TO EXPLAIN THE RESULTS OF BOTH EXPERIMENTS WITHOUT INVOKING A NEW RESONANCE. A REANALYSIS OF THE CLINE 68 DATA MADE BY BUNNEL 70 SHOW AGREEMENT OF THE DATA WITH THE ALEXANDER 69 INTERPRETATION.

REFERENCES FOR Y*(1440) (PROD. EXP.)

Table of references for Y*(1440) production experiments, including CLINE 68 PRL 21 1372 and ALEXANDER 69 PRL 22 483.

Σ(1480) BUMPS

23 Y*(1480, JP= 1/2-) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

PEAKS ARE SEEN IN LAMBDA PI AND SIGMA PI SPECTRA IN THE REACTION PI+P TO K+ PI Y AT 1.7 GEV/C. ALSO THE Y POLARIZATION OSCILLATES IN THE SAME REGION.

SEE MILLER 70 FOR A DISCUSSION OF THIS STATE. HE SUGGESTS A POSSIBLE ALTERNATE EXPLANATION IN TERMS OF A REFLECTION OF N(1216) DECAY TO LAMBDA K. HOWEVER, SUCH AN EXPLANATION FOR THE K+ SIGMA+ PI0 CHANNEL SEEMS UNLIKELY (SEE PAN 70). IN TERMS OF KN(1670) DECAY INTO SIGMA K, IN ADDITION SUCH REFLECTIONS WOULD ALSO HAVE TO ACCOUNT FOR THE OSCILLATION OF THE Y POLARIZATION IN THE 1480 MASS REGION. HANSON 71, WITH FEWER DATA THAN PAN 70, CAN NEITHER CONFIRM NOR DENY THE EXISTENCE OF THIS STATE. MAST 75 SEES NO STRUCTURE IN THIS MASS REGION IN K- P TO LAMBDA PI0.

23 Y*(1480) MASS (MEV) (PROD. EXP.)

Table of Y*(1480) mass measurements from various experiments, including PAN 70 HBC and CLINE 73 MPWA.

23 Y*(1480) WIDTH (MEV) (PROD. EXP.)

Table of Y*(1480) width measurements from various experiments, including PAN 70 HBC and CLINE 73 MPWA.

23 Y*(1480) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for Y*(1480), including decays into Kbar N, Lambda PI, and Sigma PI.

23 Y*(1480) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for Y*(1480) decays, including data from PAN 70 HBC and CLINE 73 MPWA.

REFERENCES FOR Y*(1480) (PROD. EXP.)

Table of references for Y*(1480) production experiments, including PAN 70 PR D2, 49 and CLINE 73 LNC 6 205.

Σ(1580)

00 Y*(1580, JP=3/2-) I=1 D13 4/75*

OBSERVED IN K- N I=1 TOTAL CS AT BNL (LI 73, CARROLL 73), AND IN PWA OF K- P -> LAMBDA PI FOR CM ENERGIES=1560-1600 MEV BY LITCHFIELD 74. LITCHFIELD 74 FINDS JP=3/2-.

00 Y*(1580) MASS (MEV)

Table of Y*(1580) mass measurements, including LITCHFIELD 74 DPWA 0 K- P TO LAM PI.

00 Y*(1580) WIDTH (MEV)

Table of Y*(1580) width measurements, including LITCHFIELD 74 DPWA 0 K- P TO LAM PI.

Baryons

$\Sigma(1580)$, $\Sigma(1620)$

Data Card Listings

For notation, see key at front of Listings.

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00 Y*1(1580) PARTIAL DECAY MODES
-----
P1 Y*1(1580) INTO KBAR N          497+ 939
P2 Y*1(1580) INTO LAMBDA PI      1115+ 139
P3 Y*1(1580) INTO SIGMA PI       1197+ 139
-----
00 Y*1(1580) BRANCHING RATIOS
-----
R1 Y*1(1580) INTO KBAR N/TOTAL (P1) 4/75*
R1 L +.03 .01 LITCHFIELD 74 DPWA KBAR N MULTICHNL 4/75*
R1 L MAIN EFFECT OBSERVED BY LITCHFIELD 74 IS IN PI LAMBDA FINAL STATE, 4/75*
R1 L KBAR N AND SIGMA PI COUPLINGS ALSO ESTIMATED FROM MULTICHANNEL FIT 4/75*
R1 L INCLUDING TOTAL CROSS SECTION DATA (LI 73). 4/75*
R2 Y*1(1580) FROM KBAR N TO LAMBDA PI SQRT(P1*P2) 4/75*
R2 L +.10 .02 LITCHFIELD 74 DPWA O K- P TO LAM PI 4/75*
R3 Y*1(1580) FROM KBAR N TO SIGMA PI SQRT(P1*P3) 4/75*
R3 L +.03 .04 LITCHFIELD 74 DPWA KBAR N MULTICHNL 4/75*
*****
REFERENCES FOR Y*1(1580)
LITCHFIE 74 PL 51E 509 LITCHFIELD (CERN)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
CARROLL 73 APS BRKLY MTG 208 CARRCLL,CHIANG,KYCIA,LI,MAZUR,MICHAEL+(BNL)I
LI 73 PURDUE CONF. 285 LI (BNL)I
*****

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Note on $\Sigma(1620)$

This state was first suggested by the BNL-CCNY collaboration (CRENNELL 68) who presented evidence for it in the reaction $K^- n \rightarrow \Sigma(1620)^+ \pi^- \pi^-$ with $\Sigma(1620)^+$ decaying into $\Lambda \pi^+$. Since then there have been conflicting reports about this state.

Total Cross-Section Experiment

A measurement of the $K^- p$ and $K^- d$ total cross sections in the 0.4 to 1.1 GeV/c range has been reported by the BNL group (CARROLL 73, LI 73). A clear bump about 40 MeV wide and 3-4 mb high [corresponding to $(J+1/2)x \sim 0.1$] is seen at a c.m. energy of 1590 MeV in the $I=1 K^- N$ cross section.

Formation Experiments

There is evidence from several partial-wave analyses for one or two fairly narrow states within ~ 50 MeV of the effect seen in production; see the entries for $\Sigma(1580, 3/2^-)$, $\Sigma(1620, 1/2^-)$, and $\Sigma(1660, 1/2^+)$. Note however that the various analyses do not agree on the widths and branching ratios of these states.

Production Experiments

A good review of the production experiments has been given by MILLER 70. Here the evidence is only in the $\Lambda \pi$ channel. The BNL-CCNY collaboration, with increased data, CRENNELL 69, still claim the effect in the $\Lambda \pi$ channel (no evidence seen in $\bar{K} N$ or $\bar{K} N \pi$). SABRE 70 studied the same reaction at 3.0

GeV/c with comparable statistics and do not see any evidence for it in the $\Lambda \pi$ channel; on the contrary, they believe it to be a spurious peak resulting from misidentified Σ^0 from the production of $\Sigma(1670)$ decaying into $\Sigma^0 \pi^+$. CRENNELL 69 give counter arguments to show that this is not the case in their data and the controversy goes on. AMMANN 70 studied the same reaction at 4.5 GeV/c and report a state at 1640 MeV, again decaying only into $\Lambda \pi$ (no evidence seen in $\Sigma \pi$ or $\bar{K} N$ channels). Upper limits on production cross sections for a 25 GeV/c Σ^- beam are reported by HUNGERBUHLER 74.

In conclusion, for understanding of the $\Sigma(1620)$ we probably have to wait for more data and for a more complete understanding of the entire mass region from 1600 to 1700 MeV. The closeness of the $\Sigma(1620)$ mass to 1670 MeV is suggestive that this effect may be related to what goes on in that region (see discussion below).

$\Sigma(1620)$ S11

32 Y*1(1620, JP=1/2- I=1 THE S11 STATE AT 1647 MEV REPORTED BY VANHORN 75 IS INTERMEDIATE IN MASS BETWEEN THE SIGMA(1620) AND SIGMA(1750). WE TENTATIVELY LIST IT UNDER SIGMA(1750).

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-----
32 Y*1(1620) MASS (MEV)
M (1620.) KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1630.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
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-----
32 Y*1(1620) WIDTH (MEV)
W (40.1) KIM 71 DPWA K-MATRIX ANAL. 3/71
W 65.0 (20.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
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-----
32 Y*1(1620) PARTIAL DECAY MODES
P1 Y*1(1620) INTO KBAR N          497+ 939
P2 Y*1(1620) INTO SIGMA PI      1197+ 139
P3 Y*1(1620) INTO LAMBDA PI     1115+ 134
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-----
32 Y*1(1620) BRANCHING RATIOS
R1 Y*1(1620) INTO KBAR N (P1)
R1 (0.05) KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 A 0.05 OR LESS WONG 71 DPWA K-PP-LAM+PI 10/71
R1 0.22 (0.02) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 A K-MATRIX FIT(NEGLECTS 3-BODY CHANNELS) REQUIRES NO RESONANCE 10/71
R2 Y*1(1620) FROM KEAR N TO SIGMA PI SQRT(P1*P2)
R2 (0.08) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 0.40 (0.06) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R3 Y*1(1620) FROM KBAR N TO LAMBDA PI SQRT(P1*P3)
R3 (0.15) KIM 71 DPWA K-MATRIX ANAL. 3/71
R3 NOT SEEN BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
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*****
REFERENCES FOR Y*1(1620)
KIM 71 PRL 27 356 J K KIM (HARV)IJP
ALSO 70 DUKE 161 J. K. KIM (HARV)IJP
WONG 71 NC 2A 353 N S WONG (YALE)IJP
LANGBEIN 72 NP 847 477 +WAGNER (MPI)IJP
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN,RHEL)IJP
PAPERS NOT REFERRED TO IN DATA CARDS
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL)IJP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL)IJP
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Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Sigma(1620)$, $\Sigma(1660)$, $\Sigma(1670)$

1620 MEV REGION - PRODUCTION EXPERIMENTS

78 $\Sigma(1620)$, JP=) I=1 PRODUCTION EXPERIMENTS
 SEE THE MINI-REVUE AT THE START OF THE Σ^* LISTINGS.
 THIS RESONANCE NEEDS CONFIRMATION. THE RESULTS OF CRENNELL 69 AT 3.9 GEV/C ARE NOT CONFIRMED BY THE SABRE COLLABORATION AT 3.0 GEV/C (SABRE 70). HOWEVER IN AN EXPERIMENT AT 4.5 GEV/C, AMMANN 70 SEE A PEAK AT 1642 MEV WHICH ON THE BASIS OF BRANCHING RATIOS THEY DO NOT ASSOCIATE WITH THE $\Sigma^*(1670)$. SEE MILLER 70 FOR A REVIEW OF THESE CONFLICTS.

78 $\Sigma(1620)$ MASS (MEV) (PROD. EXP.)

M	N	(1616.0)	(8.0)	CRENNELL 68 DBC	+ K-D 3.9 GEV/C	11/68
M	N	EVENTS OF CRENNELL 68	ARE IN THE LARGER SAMPLE OF CRENNELL 69.			
M		20 1618.0	3.0	BLUMENFEL 69 HBC	+ KO LONG + PROTON	9/69
M		1619.0	8.0	CRENNELL 69 DBC	+ K-N TO LAM 3 PI	9/69
M		1642.0	12.0	AMMANN 70 DBC	K-N 4.5 GEV/C	9/73
M	AVG	1619.4	3.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)		
M	STUDENT	1619.1	3.0	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		

78 $\Sigma(1620)$ WIDTH (MEV) (PROD. EXP.)

M	N	(66.0)	(16.0)	CRENNELL 68 DBC	+ SEE NOTE N ABOVE	11/68
M		20 30.0	10.0	BLUMENFEL 69 HBC		9/69
M		72.0	22.0	CRENNELL 69 DBC		9/69
M		55.0	24.0	AMMANN 70 DBC	K-N 4.5 GEV/C	9/73
M	AVG	41.3	12.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)		
M	STUDENT	40.7	10.4	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT		

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED

78 $\Sigma(1620)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$\Sigma(1620)$ INTO KBAR N	497+ 939
P2	$\Sigma(1620)$ INTO LAMBDA PI	1115+ 139
P3	$\Sigma(1620)$ INTO $\Sigma(1385)$ PI	1384+ 139
P4	$\Sigma(1620)$ INTO LAMBDA PI P1	1115+ 139+ 139
P5	$\Sigma(1620)$ INTO SIGMA PI	1197+ 139
P6	$\Sigma(1620)$ INTO $\Sigma(1405)$ PI	1405+ 139

78 $\Sigma(1620)$ BRANCHING RATIOS (PROD. EXP.)

R1	$\Sigma(1620)$ INTO (LAMBDA PI P1)/(LAMBDA PI)	(P4)/(P3)	
R1	14 (2.5) APPROX	BLUMENFEL 69 HBC +	
R2	$\Sigma(1620)$ INTO (KBAR N)/(LAMBDA PI)	(P1)/(P2)	
R2	(0.0) (0.1)	CRENNELL 68 DBC +	
R2	0.4	AMMANN 70 DBC	K-P 4.5 GEV/C 6/70
R3	$\Sigma(1620)$ INTO LAMBDA PI LARGE	CRENNELL 68 DBC +	11/68
R4	$\Sigma(1620)$ INTO ($\Sigma(1385)$ PI)/(LAMBDA PI)	(P3)/(P2)	
R4	(0.2) (0.1)	CRENNELL 68 DBC +	11/68
R4	(0.3) OR LESS CL=.95	AMMANN 70 DBC	K-P 4.5 GEV/C 6/70
R5	$\Sigma(1620)$ INTO (SIGMA PI)/(LAMBDA PI)	(P5)/(P2)	
R5	(1.1)(95 PC UPPER LIMIT)	AMMANN 70 DBC	K-N 4.5 GEV/C 9/73
R6	$\Sigma(1620)$ INTO ($\Sigma(1405)$ PI)/(LAMBDA PI)	(P6)/(P2)	
R6	0.7 0.4	AMMANN 70 DBC	K-P 4.5 GEV/C 6/70

REFERENCES FOR $\Sigma(1620)$ (PROD. EXP.)

CRENNELL 68 PRL 21 648 +DELANEY, FLAMINGO, KARSHCN, + (BNL,CUNY) I
 BLUMENFEL 69 PL 298 58 R LEVI SETTI (RAPPORTEUR) EFINS (BNL) I
 CRENNELL 69 LUND PAPER 183 +KARSHON, LAI, ONEIL, SCARR, + (BNL,CUNY) I
 RESULTS ARE QUOTED IN LEVI SETTI 69.

AMMANN 70 PRL 24 327 + GARFINKEL, CARMONY, GUTAY,+ (PURDUE,IND)
 ALSO 73 PRD 7 1345 AMMANN,CARMONY,GARFINKEL, (PURDUE,IUPUI)

PAPERS NOT REFERRED TO IN DATA CARDS

ARMENTERO 68 NP 88 183 ARMENTEROS,BAILLON + (CERN+HEID+SACL)
 LEVISETTI 69 UCL 19361 R LEVI SETTI (RAPPORTEUR) EFINS (HARV) IUP
 TRIPP 69 UCL 19361 R D TRIPP (RAPPORTEUR) EFINS (HARV) IUP
 ARMENTERO 70 DUKE 123 ARMENTEROS,BAILLON + (CERN+HEID+SACL)
 MILLER 70 DUKE 229 D H MILLER (REVIEW TALK) (PURDUE)
 SABRE 70 NP 816 201 SABRE COLLAB. (SACL,AMST,BGNA,REHC,EPOL)
 HUNGERBU 74 PRD 10 2051 HUNGERBUHLER,MAJKA,+ (YALE,FNAL,BNL,PITT)

$\Sigma(1660)$

P_{11}^*

79 $\Sigma(1660)$, JP=1/2+ I=1
 SEE THE MINI-REVUE AT THE START OF THE Σ^* LISTINGS.
 THE PARTIAL-WAVE ANALYSIS OF K- N TO SIGMA PI BY ARMENTEROS 70 SUGGEST SUCH A RESONANCE. NOW FOUND ALSO IN SOME, BUT NOT ALL, MORE RECENT ANALYSES.

79 $\Sigma(1660)$ MASS (MEV)

M		1500. -- 1600.	ARMENTERO 70 HDBC	-O K-N TO SIGMA PI	6/70
M		(1670.)	KIM 71 DPWA	K-MATRIX ANAL.	3/71
M	2	(1621.)	LEA 73 DPWA	MULTICHNL K-MTRX	9/73
M	2	ONLY UNCONSTRAINED STATES FROM	TABLE 1 OF LEA73 ARE IN LISTINGS.		9/73
M		1658. (4.)	HART 73 DPWA	EL+CX, 7--8GEV/C	2/74
M	1	(1660.) (30.)	BAILLON 75 IPWA	KBAR N TO LAM PI	11/75*
M	1	FROM SOLUTION 1 OF BAILLON 75,	NOT PRESENT IN SOLUTION 2		11/75*
M	3	(1671.) (2.)	PONTE 75 DPWA	O K- P TO LAM PI	1/76*
M	3	FROM SOLUTION 2 OF PONTE 75,	NOT PRESENT IN SOLUTION 1.		1/76*
M		1668. (25.)	VANHORN 75 DPWA	O K- P TO LAM PIO	11/75*
M		1676. (15.)	RLIC 76 DPWA	KBAR N MULTICHNL	1/76*

79 $\Sigma(1660)$ WIDTH (MEV)

W		(50.0)	ARMENTERO 70 HDBC	-O K-N TO SIGMA PI	6/70
W		(50.1)	KIM 71 DPWA	K-MATRIX ANAL.	3/71
W	2	(51.8)	LEA 73 DPWA	MULTICHNL K-MTRX	9/73
W		40. (10.)	HART 73 DPWA	EL+CX, 7--8GEV/C	2/74
W	1	(80.) (40.)	BAILLON 75 IPWA	KBAR N TO LAM PI	11/75*
W	3	(81.) (10.)	PONTE 75 DPWA	O K- P TO LAM PI	1/76*
W		230. (165.) (60.)	VANHORN 75 DPWA	O K- P TO LAM PIO	11/75*
W		120. (20.)	RLIC 76 DPWA	KBAR N MULTICHNL	1/76*

79 $\Sigma(1660)$ PARTIAL DECAY MODES

P1	$\Sigma(1660)$ INTO KBAR N	497+ 939
P2	$\Sigma(1660)$ INTO SIGMA PI	1197+ 139
P3	$\Sigma(1660)$ INTO LAMBDA PI	1115+ 139

79 $\Sigma(1660)$ BRANCHING RATIOS

R1	$\Sigma(1660)$ FROM KBAR N TO SIGMA PI	SQRT(P1*P2)	
R1	(+0.2)	ARMENTERO 70 HDBC	-O K-N TO SIGMA PI 6/70
R1	(0.24)	KIM 71 DPWA	K-MATRIX ANAL. 3/71
R1	(-.21)	LEA 73 DPWA	MULTICHNL K-MTRX 9/73
R1	-.16 (0.03)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*
R2	$\Sigma(1660)$ INTO KBAR N	(P1)	
R2	(0.14)	KIM 71 DPWA	K-MATRIX ANAL. 3/71
R2	(.10)	LEA 73 DPWA	MULTICHNL K-MTRX 9/73
R2	.11 (0.02)	HART 73 DPWA	EL+CX, 7--8GEV/C 2/74
R2	LESS THAN .04	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*
R3	$\Sigma(1660)$ FROM KBAR N TO LAMBDA PI	SQRT(P1*P3)	
R3	(0.0)	KIM 71 DPWA	K-MATRIX ANAL. 2/73
R3	(+.07)	LEA 73 DPWA	MULTICHNL K-MTRX 9/73
R3	(-.04) (0.02)	BAILLON 75 IPWA	KBAR N TO LAM PI 11/75*
R3	(+.16) (0.01)	PONTE 75 DPWA	O K- P TO LAM PI 1/76*
R3	.12 (0.12) (0.04)	VANHORN 75 DPWA	O K- P TO LAM PIO 11/75*
R3	LESS THAN .04	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

 REFERENCES FOR $\Sigma(1660)$
 ARMENTER 70 DUKE 123 ARMENTEROS, BAILLON, + (CERN,HEIDEL IJJP
 KIM 71 PRL 27 356 J K KIM (HARV) IJJP
 ALSO 70 DUKE 161 J. K. KIM (HARV) IJJP
 HART 73 PURDUE CGNF. 311 +RICE,BACASTON,FUNG,+ (TENN+UCR+MASA+BUFF) IJJP
 LEA 73 NP 856 77 +MARTIN,MOORHOUSE+ (RHEL+LOUC+GLAS+AARHUS) IJJP
 BAILLON 75 NP 894 39 P. BAILLON,P. J. LITCHFIELD (CERN,RHEL IJJP
 PONTE 75 PRD 12 2597 +HERTZBACH,BUTTON-SHAFER+ (MASA+TENN+UCR) IJJP
 VANHORN 75 NP 887 145 A. J. VAN HORN (LBL) IJJP
 LEA 73 NP 887 157 A. J. VAN HORN (LBL) IJJP
 RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCPHERSON,ROSS+ (RHEL+LOUC) IJJP

Note on $\Sigma(1670)$

Production Experiments

The measured $\Sigma\pi/\Sigma\pi\pi$ branching ratio for produced $\Sigma(1670)$'s is strongly dependent on momentum transfer. This was first discovered by EBERHARD 69 who suggested the existence of two Σ_1^* 's with the same mass and quantum numbers; one object with a large $\Sigma\pi\pi$ [mainly $\Lambda(1405)\pi$] decay mode produced peripherally, and another one with a large $\Sigma\pi$ decay mode produced at larger angles. This observation has been confirmed by AGUILAR-BENITEZ 70, APSELL 74, ESTES 74, and BLOKZIJL 75. When determined, the most likely quantum numbers are $3/2^-$ [for both $\Sigma\pi$ and $\Lambda(1405)\pi$]. There is also the possibility of a third Σ_1^* state, referred to as $\Sigma(1690)$ in the Data Card Listings, with a large $\Lambda\pi/\Sigma\pi$ branching ratio and somewhat larger mass. The large branching ratio is the main justification for this hypothesis and needs confirmation. These problems have been reviewed by EBERHARD 73 and MILLER 70.

Baryons
Sigma(1670)

Data Card Listings
For notation, see key at front of Listings.

Formation Experiments

Two states are also observed near this mass in formation. One of these, the Sigma(1670, 3/2-), has the same quantum numbers as those observed in production and a large Sigma pi/Sigma pi pi branching ratio. It may well correspond to the produced Sigma(1670) seen at larger angles. The other, the Sigma(1660, 1/2+) has different quantum numbers from those seen in production, and its Sigma pi/Sigma pi pi branching ratio is unknown. Thus its relation to the produced Sigma(1670) remains obscure.

[See also the mini-review on Sigma(1620)].

Sigma(1670) 44 Y*(1670, JP=3/2-) I=1 D13

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE NOTE ABOVE

WELL ESTABLISHED RESONANCE. IT HAS BEEN SEEN IN BOTH FORMATION AND PRODUCTION EXPERIMENTS. HOWEVER THE BRANCHING RATIOS OBTAINED BY THESE TWO METHODS SHOW LARGE INCONSISTENCIES.

SEE LISTING OF PRODUCTION EXPERIMENTS BELOW

AS FOR THE QUANTUM NUMBERS, THE ANALYSES OF LAMBDA PI CHANNEL (IN FORMATION EXP.) AS WELL AS THE SIGMA PI CHANNEL AGREE ON JP=3/2-.

Table with 4 columns: Mass (MEV), Width (MEV), and various production channels (e.g., BERLEY 64 HBC, ARMENTER 68 HBC).

Table with 4 columns: Mass (MEV), Width (MEV), and various production channels (e.g., BERLEY 64 HBC, ARMENTER 68 HBC).

44 Y*(1670) PARTIAL DECAY MODES

Table with 2 columns: Decay mode (e.g., Y*(1670) INTO KEAR N) and Decay masses.

44 Y*(1670) BRANCHING RATIOS

Table with 4 columns: Branching ratio, Production channel, and other parameters (e.g., ARMENTER 68 HBC).

Table with 4 columns: Production channel, Mass (MEV), Width (MEV), and other parameters (e.g., Y*(1670) INTO SIGMA PI).

REFERENCES FOR Y*(1670)

Table with 4 columns: Reference name, Mass (MEV), Width (MEV), and other parameters (e.g., BERLEY 64 DUBNA CONF I 565).

Sigma(1670) BUMPS

51 Y*(1670, JP=) I=1 PRODUCTION EXPERIMENTS

SEE NOTE PRECEDING Y*(1670) PROBABLY THERE ARE TWO STATES AT SAME MASS WITH SAME QUANTUM NUMBERS...

51 Y*(1670) MASS (MEV) (PROD. EXP.)

Table with 4 columns: Mass (MEV), Width (MEV), and other parameters (e.g., ALEXANDER 62 HBC).

Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(1670), Σ(1690)

M 1200 1688+-2. OR 1683+-5. BERTHON 74 HBC 0 QUASI 2 BODY CS 4/75*
M 3 1665. 3. BLOKZIJL 75 HBC K-P AT 4.2 GEV 1/76*
M 3 PARAMETERS DETERMINED FROM (SIGMA PI+) PRODUCTION. 1/76*

51 Y*(1670) MASS (MEV) (PROD. EXP.)
W (45.0) ALEXANDER 62 HBC -0
W 40.0 10.0 ALVAREZ 63 HBC +
W (30.0) (15.0) BUGG 68 CNTR 11/66

51 Y*(1670) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*(1670) INTO KBAR N 497+ 939
P2 Y*(1670) INTO LAMBDA PI 1115+ 139
P3 Y*(1670) INTO SIGMA PI 1197+ 139

51 Y*(1670) BRANCHING RATIOS (PROD. EXP.)
R1 Y*(1670) INTO (KBAR N)/(SIGMA PI) (P1)/(P3)
R1 0 (0.19) OR LESS ALVAREZ 63 HBC + K-P 1.15 BEV/C
R1 (0.51)+ .25 OR MORE SMITH 63 HBC -0

R2 Y*(1670) INTO (LAMB.PI)/(SIG PI) (P2)/(P3)
R2 130 (1.20) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R2 (1.2) SMITH 63 HBC -0
R2 0.15 0.07 MUSEL 64 HBC +

R3 Y*(1670) INTO (LAMB. PI PI)/(SIG PI) (P4)/(P3)
R3 90 (0.56) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R3 (0.17) SMITH 63 HBC -0

R4 Y*(1670) INTO (SIGMA PI PI)/(SIG PI) (P5)/(P3)
R4 180 (0.56) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R4 LARGEST AT SMALL ANGLES ESTES 74 HBC 0 K-P 2.1+2.6GEV/C 11/75*

R5 Y*(1670) INTO (Y*(1405) PI)/(SIG PI) (P7)/(P3)
R5 50 3. 1.6 LONDON 66 HBC + K-P 2.25 BEV/C 7/66
R5 P 17 (0.58) (0.20) PRIMER 68 HBC + K-P 4.6-5. GEV/C 7/68

R6 Y*(1670) INTO (SIGMA PI)/(SIGMA PI PI) (P3)/(P5)
R6 .4 OR LESS BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67
R6 0.30 0.15 LONDON 66 HBC + K-P 2.25 BEV/C 7/66

R7 Y*(1670) INTO (Y*(1405) PI)/(SIGMA PI PI) (P7)/(P5)
R7 0.90 0.10 0.16 EBERHARD 65 HBC + K-P 2.45 BEV/C 7/66
R7 B 1.00 -0.02 APSELL 74 HBC K-P 2.87 GEV/C 4/75*

R8 Y*(1670) INTO (Y*(1405) PI)/(Y*(1385) PI) (P7)/(P6)
R8 (0.8) OR LESS EBERHARD 65 HBC + K-P 2.45 BEV/C 7/66

R9 Y*(1670) INTO (LAMBDA PI PI)/(SIGMA PI PI) (P4)/(P5)
R9 0.35 0.2 BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67

R10 Y*(1670) INTO (LAMBDA PI)/(SIGMA PI PI) (P2)/(P5)
R10 (1.2) OR LESS BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67

R11 Y*(1670) INTO (LAMBDA PI)/(LAMBDA PI + SIG PI) (P2)/(P2+P3)
R11 (0.6) OR LESS AGUILAR 70 HBC 5/70

R12 Y*(1670) INTO (Y*(1385) PI)/(SIGMA PI) (P6)/(P3)
R12 (.251)+/--.05) OR LESS BLOKZIJL 75 HBC K-P AT 4.2 GEV 1/76*

51 Y*(1670) QUANTUM NUMBER DETERMINATION (PROD. EXP.)
Q1 JP=3/2+ LEVEQUE 65 HBC INTO Y*(1405)+PI 11/68
Q3 JP=3/2- EBERHARD 67 HBC + INTO Y*(1405) PI 11/68
Q4 400 JP=3/2- BUTTCH-SH 68 HBC + INTO SIGZERO+PI 11/68

REFERENCES FOR Y*(1670) (PROD. EXP.)
ALEXANDE 62 CERN CONF 320 ALEXANDER, JACOBS, KALBFLEISCH, MILLER, + (LRL) I
ALVAREZ 63 PRL 10 184 +ALSTON, FERRO-LUZZI, HUWE, + (LRL) I
SMITH 63 ATHENS CONF 67 G A SMITH (LRL)

APSELL 74 PRD 10 1419 APSELL, FCRO, GUREVITCH, LBRAN, UME, SYRA, TUFT II
BERTHON 74 NC 21A 146 BERTHON, TRISTRAM, + (CONF+RHEL+SACL+STRB) L
ESTES 74 LBL-3827 (THESIS) R. D. ESTES (LBL)

BLOKZIJL 75 ANL-HEP-CR-75-58 +DEGROOT, HODGLAND+ (AMST+CERN+NIJH+OXF)
ALSO 75 PALERMO CONF. +DEGROOT, HODGLAND+ (AMST+CERN+NIJH+OXF)

LEVEQUE 65 PL 18 69 + (SACLAY, EPOL, GLASGOW, LCLC, CFX, RHEL) JP
LEE 66 PRL 17 45 Y Y LEE, D D REEDER, R W HARTUNG (WISC) JP
EBERHARD 67 PR 163 1446 +PRIPSTEIN, SHIVELY, KRUSE, SWANSON (LRL, ILL) JP
MILLER 70 HDQ 229 D H MILLER (REVIEW TALK) (PURDUE)
EBERHARD 73 PURDUE CONF. 247 EBERHARD (LBL) JP
HUNGERBU 74 PRD 10 2051 HUNGERBUHLER, MAJKA, + (YALE, FNAL, BNL, PITT)

Σ(1690) BUMPS
58 Y*(1690, JP= 1) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.
SEE NOTE PRECEDING Y*(1670) LISTINGS, SEEN IN PRC. EXPERIMENTS GNLV, MAIN DECAY MODE IS LAMBDA PI.

58 Y*(1690) MASS (MEV) (PROD. EXP.)
M 30(1715.0) (12.0) COLLEY 67 HBC + K-P 6 GEV/C 8/67
M P 60(1694.0) (24.0) PRIMER 68 HBC + K-P 4.6-5 GEV/C 7/68
M N (1700.0) (6.0) SIMS 68 HBC - K-N TO LAMB PI 11/68
M 46(1682.0) (2.0) BLUMENFEL 69 HBC + KO LONG + PROTON 9/69
M (1700.0) (20.0) MOTT 69 HBC + K-P 5.5 GEV/C 9/69

58 Y*(1690) WIDTH (MEV) (PROD. EXP.)
W 30 (100.0) (35.0) COLLEY 67 HBC + 8/67
W 60 (105.0) (35.0) PRIMER 68 HBC + 7/68
W N 66 (62.0) (14.0) SIMS 68 HBC - SEE NOTE N ABOVE 11/68
W 46 (25.0) (10.0) BLUMENFEL 69 HBC + 9/69
W (130.0) (25.0) MOTT 69 HBC + 9/69

58 Y*(1690) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*(1690) INTO KBAR N 497+ 939
P2 Y*(1690) INTO LAMBDA PI 1115+ 139
P3 Y*(1690) INTO SIGMA PI 1197+ 139
P4 Y*(1690) INTO Y*(1385) PI 1384+ 139
P5 Y*(1690) INTO LAMBDA PI PI (INCLUDING P4) 1115+ 139+ 139

58 Y*(1690) BRANCHING RATIOS (PROD. EXP.)
R1 Y*(1690) INTO (KBAR N)/(LAMBDA PI) (P1)/(P2)
R1 18 0.4 0.25 COLLEY 67 HBC + 6/30 EVENTS 8/67
R1 (0.2) OR LESS MOTT 69 HBC + 9/69

R2 Y*(1690) INTO (SIGMA PI)/(LAMBDA PI) (P3)/(P2)
R2 0.3 0.3 COLLEY 67 HBC + 4/30 EVENTS 8/67
R2 (0.4) OR LESS CL=-.90 MOTT 69 HBC + 9/69

R3 Y*(1690) INTO (Y*(1385) PI)/(LAMBDA PI) (P4)/(P2)
R3 (0.5) OR LESS MOTT 69 HBC + 9/69

R4 Y*(1690) INTO (LAMBDA PI PI)/(LAMBDA PI) (P5)/(P2)
R4 0.5 0.25 COLLEY 67 HBC + 15/30 EVENTS 8/67
R4 2.0 0.6 BLUMENFEL 69 HBC + 31/15 EVENTS 9/69

R4 AVG 0.72 0.53 AVERAGE (ERROR INCLUDES SCALE FACTOR CF 2.5)
R4 STUDENT 0.67 0.28 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

R5 Y*(1690) INTO (Y*(1385) PI)/(LAMBDA PI PI) (P4)/(P5)
R5 SMALL COLLEY 67 HBC + K-N TO L2PI 8/67
R5 LARGE SIMS 68 HBC - 11/68

Baryons

$\Sigma(1690)$, $\Sigma(1750)$, $\Sigma(1765)$

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR $\Sigma(1690)$ (PROD. EXP.)

COLLEY 67 PL 248 489 (BIRM, GLAS, LOIC, MUNICH, OXFORD, RHEL) I
DERRICK 67 PRL 18 266 (ARGONNE, NORTHWEST) I
REPLACED BY HOIT 65.
PRIMER 68 PRL 20 610 (SYRACUSE, BNL) I
SIMS 68 PRL 21 1413 (FSU, TUFTS, BRANDEIS) I
BLUMENFE 69 PL 298 58 E J BLUMENFELD, G R KALBFLEISCH (BNL) I
HOIT 69 PR 177 1966 *AMMAR, DAVIS, KROPAC, +INCRTHWEST, ARGONNE) I

PAPERS NOT REFERRED TO IN DATA CARDS

AGUILAR 70 PRL 25 58 AGUILAR-BENITEZ, BARNES, BASSAND+ (BNL+SYRA)

$\Sigma(1750)$

S11

SEE THE MINI-REVIEW AT THE START OF THE Σ LISTINGS.

THIS STATE CORRESPONDS TO THE SIGMA ETA THRESHOLD BUMP BUT ITS INTERPRETATION IN TERMS OF A RESONANCE IS NOT CONCLUSIVE (JONES 74) -- MORE DATA ARE NEEDED. BY ANALOGY WITH THE SIMILAR N ETA AND LAMBDA ETA THRESHOLD EFFECTS, WHICH ARE ALMOST CERTAINLY RESONANCES, IT SEEMS VERY LIKELY THAT THIS TOO IS A RESONANCE.

THERE IS ALSO EVIDENCE FOR THIS STATE IN MANY PARTIAL-WAVE ANALYSES, MOSTLY IN THE LAMBDA PI CHANNEL. THIS EVIDENCE SHOULD BE CONSIDERED WITH CAUTION AS THE MASSES AND/OR BRANCHING RATIOS OF THE REPORTED STATES ARE OFTEN INCONSISTENT.

57 $\Sigma(1750)$ MASS (MEV)

Table with columns for mass values and references. Includes entries for CLINE, MEYER, ARMENTERO, CONFORTO, KIM, LANGBEIN, BAXTER, CHU, JONES, PREVOST, BAILLON, VANHORN, and RLIC.

57 $\Sigma(1750)$ WIDTH (MEV)

Table with columns for width values and references. Includes entries for MEYER, ARMENTERO, CONFORTO, KIM, LANGBEIN, BAXTER, CHU, JONES, PREVOST, BAILLON, VANHORN, and RLIC.

57 $\Sigma(1750)$ PARTIAL DECAY MODES

Table with columns for decay modes and references. Includes entries for INTO KBAR N, INTO SIGMA ETA, INTO LAMBDA PI, INTO SIGMA PI, and INTO SIGMA(1385) PI.

57 $\Sigma(1750)$ BRANCHING RATIOS

Table with columns for branching ratios and references. Includes entries for INTO (KBAR N)/TOTAL, FROM KBAR N INTO SIGMA ETA, FROM KBAR N INTO LAMBDA PI, FROM KBAR N TO SIGMA PI, and FROM KBAR N TO SIGMA(1385) PI.

REFERENCES FOR $\Sigma(1750)$

CLINE, OLSSON (WISCONSIN) IIP
MEYER (RAPPOURTEUR) (SACLAY) IIP
ARMENTERO, BAILLON, + (CERN, HEIDEL) IIP
*LEVI SETTI, LASINSKI, OBERLACK+ (EFI+HEID) IIP
J. K. KIM (HARV) IIP
J. K. KIM (HARV) IIP
*WAGNER (NPI) IIP

BAXTER, BUCKINGHAM, CORBETT, CUNN, + (OXFORD) IIP
CHU, BARTLEY, + (SUNY PLATTSBURGH+TUFTS+BRAN) IIP
JONES (U. CHICAGO) IIP
DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
PREVOST, BARLOUTAUD, + (SACL+ CERN+HEID)

BAILLON 75 NP B94 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IIP
VANHORN 75 NP B87 145 A. J. VAN HORN (LBL) IIP
ALSO 75 NP B87 157 A. J. VAN HORN (LBL) IIP

RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IIP

PAPERS NOT REFERRED TO IN DATA CARDS

FERRO-LU 66 BERKELEY CONF 183 M FERRO LUZZI (RAPPOURTEUR) (CERN)
ARMENTERO 68 NP 88 183 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IIP
ARMENTER 69 LUND CNFN PAPER ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IIP
HARRISON 70 FSU-HEP TO 3 1 W.C. HARRISON (THEISS) (FSU)

$\Sigma(1765)$

D15

SEE THE MINI-REVIEW AT THE START OF THE Σ LISTINGS.

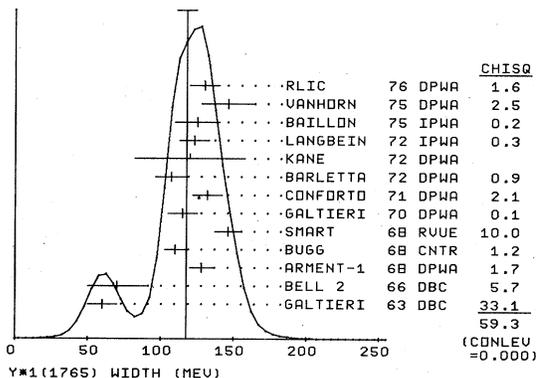
45 $\Sigma(1765)$ MASS (MEV)

Table with columns for mass values and references. Includes entries for GALTIERI, ARMENTERO, BELL, ARMENT-1, BUGG, SMART, COOL, GALTIERI, CONFORTO, KIM, BARLETTA, KANE, LANGBEIN, BAILLON, VANHORN, and RLIC.

45 $\Sigma(1765)$ WIDTH (MEV)

Table with columns for width values and references. Includes entries for GALTIERI, BELL, ARMENT-1, BUGG, SMART, COOL, GALTIERI, CONFORTO, KIM, BARLETTA, KANE, LANGBEIN, BAILLON, VANHORN, and RLIC.

WEIGHTED AVERAGE = 117.5 ± 6.8
ERROR SCALED BY 2.3



Data Card Listings
For notation, see key at front of Listings.

Baryons
Sigma(1765), Sigma(1770)

45 Y*(1765) PARTIAL DECAY MODES
Table with columns for decay mode (P1-P7), branching fraction, and decay masses.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i + 6P_i, where 6P_i = sqrt(6P_i * 6P_i), while the off-diagonal elements are the normalized correlation coefficients (6P_i * 6P_j) / (6P_i * 6P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns P 1 through P 6 and rows of numerical values representing branching fractions and correlations.

45 Y*(1765) BRANCHING RATIOS

ERRORS QUOTED BY EXPERIMENTERS DO NOT INCLUDE UNCERTAINTY DUE TO PARAMETRIZATION USED IN THE P.W.A. THEY SHOULD BE INCREASED.

Table listing branching ratios for various decay channels (e.g., GALTIERI, UHLIG, ARMENT-1) with associated errors and references.

WEIGHTED AVERAGE = 0.409 +/- 0.018
ERROR SCALED BY 2.8

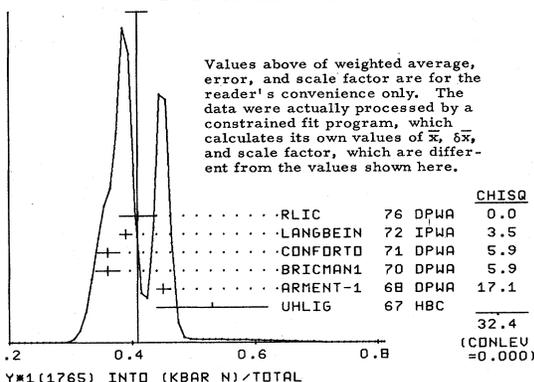


Table with columns for decay mode (R2-R3), branching fraction, and decay masses for Y*(1765) FROM KBAR N INTO LAMBDA PI.

Table with columns for decay mode (R3-R4), branching fraction, and decay masses for Y*(1765) FROM KBAR N INTO Y*(1520) PI.

Table with columns for decay mode (R4-R5), branching fraction, and decay masses for Y*(1765) FROM KBAR N INTO Y*(1385) PI D-WAVE.

Table with columns for decay mode (R5-R8), branching fraction, and decay masses for Y*(1765) FROM KBAR N INTO SIGMA PI.

REFERENCES FOR Y*(1765)
List of references for various decay channels, including GALTIERI, ARMENTEROS, UHLIG, etc.

Sigma(1770) 100 Y*(1770, JP=1/2+) I=1 P11

100 Y*(1770) MASS (MEV)
Table with columns for mass (M), branching fraction, and decay masses for Sigma(1770).

100 Y*(1770) WIDTH (MEV)
Table with columns for width (W), branching fraction, and decay masses for Sigma(1770).

100 Y*(1770) PARTIAL DECAY MODES
Table with columns for decay mode (P1-P3), branching fraction, and decay masses for Sigma(1770).

100 Y*(1770) BRANCHING RATIOS
Table with columns for decay mode (R1-R2), branching fraction, and decay masses for Sigma(1770).

Baryons

$\Sigma(1770)$, $\Sigma(1840)$, $\Sigma(1880)$, $\Sigma(1915)$

R2 Y*1(1770) INTO (KBAR N)/TOTAL 1/76*
R2 .14 .04 RLIC 76 DPWA (P1) KBAR N MULTICHNL 1/76*

REFERENCES FOR Y*1(1770)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IJJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJJP

$\Sigma(1840)$

01 Y*1(1840, JP=3/2+) I=1 P13

SEE THE MINI-REVIEWS PRECEDING THE Y*0'S.
FOR THE TIME BEING, WE LIST ALL RESONANCE CLAIMS IN THE
P13 WAVE IN THE 1700-1900 MEV MASS REGION TOGETHER UNDER THIS HEADING.

01 Y*1(1840) MASS (MEV)
M 1840.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 1 (1720.) (30.) BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

01 Y*1(1840) WIDTH (MEV)
W 120.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 1 (120.) (30.) BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

01 Y*1(1840) PARTIAL DECAY MODES
P1 Y*1(1840) INTO KBAR N 497+ 939
P2 Y*1(1840) INTO SIGMA PI 1197+ 139
P3 Y*1(1840) INTO LAMBDA PI 1115+ 134

01 Y*1(1840) BRANCHING RATIOS
R1 Y*1(1840) INTO (KBAR N)/TOTAL (P1) 12/72
R1 0.37 (0.13) LANGBEIN 72 IPWA MULTICHANNEL
R2 Y*1(1840) FROM KBAR N INTO SIGMA PI SQR(P1*P2) 12/72
R2 0.15 (0.04) LANGBEIN 72 IPWA MULTICHANNEL

REFERENCES FOR Y*1(1840)
LANGBEIN 72 NP 847 477 *WAGNER (MPIM) IJJP
DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IJJP
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL) IJJP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL) IJJP

$\Sigma(1880)$

67 Y*1(1880, JP=1/2+) I=1 P11

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL
PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE
IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN
BARYON TABLE.

67 Y*1(1880) MASS (MEV)
M 1882.0 40.0 SMART 68 DPWA -0 K- N TO LAM PI 7/68
M (1850.0) BAILEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
M ABOUT 1850.0 ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70

67 Y*1(1880) WIDTH (MEV)
W 222.0 150.0 SMART 68 DPWA -0 K- N TO LAM PI 7/68
W BAILEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
W ABOUT 30.0 ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70

Data Card Listings

For notation, see key at front of Listings.

67 Y*1(1880) PARTIAL DECAY MODES

P1 Y*1(1880) INTO KBAR N 497+ 935
P2 Y*1(1880) INTO SIGMA PI 1115+ 134
P3 Y*1(1880) INTO SIGMA PI 1197+ 139

67 Y*1(1880) BRANCHING RATIOS

R1 Y*1(1880) INTO (KBAR N)/TOTAL (P1)
R1 (0.22) BAILEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
R1 (0.20) ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70
R1 2 (.51) LEA 73 DPWA MULTICHNL K-MTRX 9/73

R2 Y*1(1880) FROM KBAR N INTO LAMBDA PI SQR(P1*P2)
R2 -0.11 0.03 SMART 68 DPWA -0 K- N TO LAM PI 7/68
R2 -0.09 0.04 GALTIERI 70 DPWA -0 K- N TO LAM PI 7/70
R2 -0.14 0.03 LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70

R3 Y*1(1880) FROM KBAR N INTO SIGMA PI SQR(P1*P3) 9/73
R3 (-108) KANE 72 DPWA K-P TO SIGMA PI 9/73
R3 2 NOT SEEN LEA 73 DPWA MULTICHNL K-MTRX 9/73

REFERENCES FOR Y*1(1880)
SMART 68 PR 169 1330 W M SMART (LRL) IJJP
BAILEY 69 THESIS UCRL-50617 DAVID SAAL BAILEY (LRL LIVERMORE) IJJP
ARMENTERO 70 DUKE CONF 123 ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJJP
GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJJP
LITCHFIELD 70 NP 822 269 P J LITCHFIELD (RUTHERFORD) IJJP
KANE 72 PR 05 1583 D F KANE (LBL)

$\Sigma(1915)$

46 Y*1(1915, JP=5/2+) I=1 F15

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THIS RESONANCE WAS FIRST SEEN IN THE TOTAL-CROSS-SEC-
TION MEASUREMENTS OF COOL 66. IN THIS ENTRY, HOWEVER,
WE LIST ONLY THE RESULTS FROM PARTIAL-WAVE ANALYSES.
SEE THE NEXT ENTRY FOR THE PARAMETERS OF PEAKS SEEN AROUND 1900-1950
MEV IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. WE MAKE THIS
SEPARATION BECAUSE ONLY THE PARTIAL-WAVE ANALYSES ISOLATE THE F15 WAVE
(OR AT LEAST ATTEMPT TO -- THE SIGNAL IS WEAK). THIS MASS REGION IS
COMPLICATED AND POORLY UNDERSTOOD AND THE PEAKS MAY CONTAIN MORE THAN
JUST THE Y*1(1915). SEE ALSO THE NOTE TO THE NEXT ENTRY.

46 Y*1(1915) MASS (MEV)

M 1902.0 11.0 SMART 68 DPWA -0 K- N TO LAMBDA PI 7/68
M 1910.0 20.0 BERTHON 70 DPWA 0 K-P TO LAMBDA PI 7/70
M 1900.0 15.0 BERTHON 70 DPWA 0 K-P TO SIGMA PI 10/70
M N 1936.0 (3.0) BRICMANI 70 DPWA SIGTOT, ELAS, CHEX 1/71
M 1903.0 10.0 COX 70 DPWA - K- N TO LAMBDA PI 6/70
M 1905.0 30.0 GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
M 1895.0 10.0 LITCHFIELD 70 DPWA -0 K- N TO LAMBDA PI 6/70
M B (1985.0) (21.0) ISLAM 71 DPWA KN--PI-SIG 12/72
M B DISCREPANCY DUE POSSIBLY TO INSUFFICIENT STATISTICS
M 1910. 15. LITCHFIELD 71 DPWA K-P TO KBAR N 10/71
M 1925.0 8.0 KANE 72 DPWA 0 K-P TO PI SIG 10/71
M 1920. 30. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

46 Y*1(1915) WIDTH (MEV)

W A (50.0) (20.0) ARMENTERO 67 DPWA 0 ELASTIC, CH EXCH 11/67
W 52.0 25.0 SMART 68 DPWA -0 K- N TO LAMBDA PI 7/68
W 60.0 20.0 BERTHON 70 DPWA 0 K-P TO LAMBDA PI 7/70
W 75.0 20.0 BERTHON 70 DPWA 0 K-P TO SIGMA PI 10/70
W 135.0 12.0 BRICMANI 70 DPWA SIGTOT, ELAS, CHEX 1/71
W 77.0 27.0 COX 70 DPWA - K- N TO LAMBDA PI 6/70
W 70.0 20.0 GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
W 70.0 15.0 LITCHFIELD 70 DPWA -0 K- N TO LAMBDA PI 6/70
W B (159.0) (80.0) ISLAM 71 DPWA KN--PI-SIG 12/72
W 70. 15. LITCHFIELD 71 DPWA K-P TO KBAR N 10/71
W 146.0 22.0 KANE 72 DPWA 0 K-P TO PI SIG 10/71
W 70. 20. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

46 Y*1(1915) PARTIAL DECAY MODES

P1 Y*1(1915) INTO KBAR N 497+ 939
P2 Y*1(1915) INTO LAMBDA PI 1115+ 139
P3 Y*1(1915) INTO SIGMA PI 1197+ 139

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Sigma(1915)$, $\Sigma(1940)$

46 Y*(1915) BRANCHING RATIOS

Table with columns for particle ID, mass, and branching ratios for various decay channels like K-bar N, Lambda PI, Sigma PI, etc.

REFERENCES FOR Y*(1915)

Table listing references for Y*(1915) including authors like Armenter, Berton, Brichman, etc., and their respective publications.

1915 MEV REGION - PRODUCTION AND σ_{TOTAL} EXPTS

29 Y*(1915, J_p=) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS. SEE THE NOTES TO THE Y*(1915) AND Y*(1940), WHICH IMMEDIATELY PRECEDE AND FOLLOW THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS SEEN IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE ALMOST CERTAINLY ASSOCIATED WITH THE P15 Y*(1915) SEEN IN PARTIAL-WAVE ANALYSES. THE INVARIANT-MASS PEAKS SEEM MORE LIKELY TO BE ASSOCIATED WITH THE NOT-COMpletely-ESTABLISHED WITH THE D13 Y*(1940).

29 Y*(1915) MASS (MEV) (PRCD. EXP.)

Table showing mass measurements for Y*(1915) from various experiments, including cross-section peaks and invariant-mass distributions.

29 Y*(1915) WIDTH (MEV) (PRCD. EXP.)

Table showing width measurements for Y*(1915) from various experiments, including cross-section peaks and invariant-mass distributions.

29 Y*(1915) PARTIAL DECAY MODES (PRCD. EXP.)

Table listing partial decay modes for Y*(1915) with associated decay masses.

29 Y*(1915) BRANCHING RATIOS (PROD. EXP.)

Table with columns for particle ID, mass, and branching ratios for various decay channels like K-bar N, Lambda PI, Sigma PI, etc.

REFERENCES FOR Y*(1915) (PROD. EXP.)

Table listing references for Y*(1915) including authors like Bock, Cool, Bugg, etc., and their respective publications.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards, including authors like Primer, Aguilars, etc.

$\Sigma(1940)$

98 Y*(1940, J_p=3/2-) I=1

D₁₃

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES IN THIS REGION. THIS EFFECT IS PERHAPS ASSOCIATED WITH THE BUMPS SEEN IN PRODUCTION EXPERIMENTS NEAR THIS MASS. SEE THE PRECEDING ENTRY.

98 Y*(1940) MASS (MEV)

Table showing mass measurements for Y*(1940) from various experiments, including cross-section peaks and invariant-mass distributions.

98 Y*(1940) WIDTH (MEV)

Table showing width measurements for Y*(1940) from various experiments, including cross-section peaks and invariant-mass distributions.

98 Y*(1940) PARTIAL DECAY MODES

Table listing partial decay modes for Y*(1940) with associated decay masses.

98 Y*(1940) BRANCHING RATIOS

Table with columns for particle ID, mass, and branching ratios for various decay channels like K-bar N, Lambda PI, Sigma PI, etc.

Baryons

$\Sigma(1940)$, $\Sigma(2000)$, $\Sigma(2030)$

Data Card Listings

For notation, see key at front of Listings.

R3 2 Y*1(1940) INTO KBAR N (P1) 9/73
 R3 2 (-21) LEA 73 DPWA MULTICHLN K-MTRX 9/73
 R3 LESS THAN .04 RLIC 76 DPWA KBAR N MULTICHLN 1/76*
 R3 NO SIGNAL FOR THIS STATE WITH X LARGER THAN ABOUT .03 IN THE 11/75*
 R3 ANALYSIS OF HEMINGWAY 75. 11/75*

R4 Y*1(1940) FROM KBAR N TO Y*0(1520) FI P-WAVE SQRT(P1*P4)
 R4 1 -.11 .04 LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 R4 1 ASSUMES LAMBDA(1520) ELASTICITY=.45,SIGN RLTV. TC SIG(2030) DECAY. 10/74*

R5 Y*1(1940) FROM KBAR N TO Y*0(1520) PI F-WAVE SQRT(P1*P5)
 R5 1 -.08 .04 LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 R5 1 ASSUMES LAMBDA(1520) ELASTICITY=.45,SIGN RLTV. TC SIG(2030) DECAY. 10/74*

R6 Y*1(1940) FROM KBAR N TO KBAR DELTA(1232) S-WAVE SQRT(P1*P6)
 R6 3 -.16 .05 LITCHF13 74 DPWA O K-P TO KBAR DEL 10/74*
 R6 3 SIGN RELATIVE TO SIGMA(2030) DECAY 10/74*

R7 Y*1(1940) FROM KBAR N TO KBAR DELTA(1232) D-WAVE SQRT(P1*P7)
 R7 3 -.14 .05 LITCHF13 74 DPWA O K-P TO KBAR DEL 10/74*
 R7 3 SIGN RELATIVE TO SIGMA(2030) DECAY 10/74*

 REFERENCES FOR Y*1(1940)
 GALTIERI 70 DUKE CCNF 173 A BARBARO-GALTIERI (LRL)IJP
 LITCHFIE 70 NP B22 269 P J LITCHFIELD (RUTHERFORD)IJP
 KANE 72 PR D5 1583 D F KANE (LBL)IJP
 LEA 73 NP B56 77 *MARTIN,MOORHOUSE* (RHEL+LOUC+GLAS+AARHUS)IJP

DEVENISH 74 NP B81 330 DEVENISH,FROGGATT,MARTIN(DESY,NORDITA,LOUC)
 LITCHF12 74 NP B74 19 LITCHFIELD,HEMINGWAY,BAILLON,+ (CERN+HEID)IJP
 LITCHF13 74 NP B74 39 LITCHFIELD,HEMINGWAY,BAILLON,+ (CERN+HEID)IJP

BAILLON 75 NP B94 39 P. BAILLON,P. J. LITCHFIELD (CERN,RHEL)IJP
 VANHORN 75 NP B87 145 A. J. VAN HORN (LBL)IJP
 ALSO 75 NP B87 157 A. J. VAN HORN (LBL)IJP

RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMLS,MCPHERSON,ROSS+ (RHEL+LOIC)IJP
 PAPERS NOT REFERRED TO IN DATA CARDS.

HEMINGWAY 75 NP B91 12 *HEMINGWAY,EADES,HARMSEN+ (CERN,HEID,MPIM)IJP

$\Sigma(2000)$ 02 Y*1(2000, JP=1/2-) I=1 S_{11}^{++}

02 Y*1(2000) MASS (MEV)
 M 2004. 40. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 M 1955. 15. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

02 Y*1(2000) WIDTH (MEV)
 W 116. 40. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 W 170. 40. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

02 Y*1(2000) PARTIAL DECAY MODES
 P1 Y*1(2000) INTO KBAR N 497+ 939
 P2 Y*1(2000) INTO LAMBDA PI 1115+ 134
 P3 Y*1(2000) INTO SIGMA PI 1197+ 139

02 Y*1(2000) BRANCHING RATIOS
 R1 Y*1(2000) FROM KBAR N INTO LAMBDA PI SQRT(P1*P2)
 R1 NOT SEEN BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
 R1 +.07 .02 VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 R1 .08 .03 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

R2 Y*1(2000) INTO (KBAR N)/TOTAL (P1) 1/76*
 R2 .44 .05 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

R3 Y*1(2000) FROM KBAR N INTO SIGMA PI SQRT(P1*P3) 1/76*
 R3 +.20 .04 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

 REFERENCES FOR Y*1(2000)
 BAILLON 75 NP B94 39 P. BAILLON,P. J. LITCHFIELD (CERN,RHEL)IJP
 VANHORN 75 NP B87 145 A. J. VAN HORN (LBL)IJP
 ALSO 75 NP B87 157 A. J. VAN HORN (LBL)IJP

RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMLS,MCPHERSON,ROSS+ (RHEL+LOIC)IJP

$\Sigma(2030)$ 47 Y*1(2030, JP=7/2+) I=1 F_{17}

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
 THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES. PARAMETERS OF PEAKS SEEN IN CROSS-SECTIONS AND INVARIANT-MASS DISTRIBUTIONS AROUND 2030 MEV ARE GIVEN IN THE NEXT ENTRY. EVENTUALLY THE PARTIAL-WAVE ANALYSES SHOULD GIVE THE BEST RESULTS, AS THEY ISOLATE THE FIT WAVE. THIS SUPERIORITY IS, HOWEVER, PROBABLY NOT YET ATTAINED, AND WE RELY ON BOTH ENTRIES FOR PARAMETERS GIVEN IN THE MAIN BARYON TABLE.

47 Y*1(2030) MASS (MEV)
 M (2030.0) (20.0) WOHL 66 HBC O K-P TO LAM P10 7/66
 M 2032.0 6.0 SMART 68 DPWA - K-N TO LAMBDA PI 6/68
 M 2030.0 10.0 BERTHON 70 DPWA O K-P TO LAMBDA PI 7/70
 M 2035.0 10.0 BERTHON1 70 DPWA O K-P TO SIGMA PI 10/70
 M 2027.0 6.0 COX 70 DPWA - K-N TO LAMBDA PI 6/70
 M 2010.0 15.0 GALTIERI 70 DPWA O K-P TO LAMBDA PI 7/70
 M 2000.0 20.0 GALTIERI 70 DPWA O K-P TO SIGMA PI 7/70
 M 2022.0 4.0 LITCHFIEL 70 DPWA -O K-N TO LAMBDA PI 6/70
 M 2025. 15. LITCHFIE 71 DPWA K-P TO KBAR N 10/71
 M 2034.0 14.0 KANE 72 DPWA O K-P TO PI SIG 10/71
 M 2025. 10.0 LITCHF11 74 DPWA O K-P TO L(1815)PI 10/74*
 M 2035. 10. LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 M 2020. 30. LITCHF13 74 DPWA O K-P TO KBAR DEL 10/74*
 M 2035. 15. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
 M 2038. 10. HEMINGWA 75 DPWA O K-P TO KBAR N 11/75*
 M 2042. 11. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 M 2040. 5. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

47 Y*1(2030) WIDTH (MEV)
 W (170.0) WOHL 66 HBC O 7/66
 W 160.0 16.0 SMART 68 DPWA - K-N TO LAMBDA PI 6/68
 W 165.0 30.0 15.0 BERTHON 70 DPWA O K-P TO SIGMA PI 7/70
 W 150.0 20.0 BERTHON1 70 DPWA O K-P TO LAMBDA PI 10/70
 W 158.0 16.0 COX 70 DPWA - K-N TO LAMBDA PI 6/70
 W 115.0 15.0 GALTIERI 70 DPWA O K-P TO LAMBDA PI 7/70
 W 100.0 40.0 GALTIERI 70 DPWA O K-P TO SIGMA PI 7/70
 W 170.0 15.0 LITCHFIEL 70 DPWA -O K-N TO LAMBDA PI 6/70
 W 200.0 30.0 LITCHFIE 71 DPWA K-P TO KBAR N 10/71
 W 118.0 12.0 KANE 72 DPWA O K-P TO PI SIG 10/71
 W 70. TO 125. LITCHF11 74 DPWA O K-P TO L(1815)PI 10/74*
 W 160. 20. LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 W 200. 30. LITCHF13 74 DPWA O K-P TO KBAR DEL 10/74*
 W 180. 20. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
 W 172. 15. HEMINGWA 75 DPWA O K-P TO KBAR N 11/75*
 W 178. 13. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 W 190. 10. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.7)

47 Y*1(2030) PARTIAL DECAY MODES
 P1 Y*1(2030) INTO KBAR N 497+ 939
 P2 Y*1(2030) INTO LAMBDA PI 1115+ 134
 P3 Y*1(2030) INTO SIGMA PI 1197+ 139
 P4 Y*1(2030) INTO XI K 1321+ 497
 P5 Y*1(2030) INTO Y*0(1815) PI P-WAVE 134+1820
 P6 Y*1(2030) INTO Y*0(1520) PI D-WAVE 134+1518
 P7 Y*1(2030) INTO Y*0(1520) PI G-WAVE 134+1518
 P8 Y*1(2030) INTO KBAR DELTA(1232) F-WAVE 493+1232
 P9 Y*1(2030) INTO KBAR DELTA(1232) H-WAVE 493+1232

47 Y*1(2030) BRANCHING RATIOS
 R1 Y*1(2030) INTO (KBAR N)/TOTAL (P1)
 R1 (0.25) WOHL 66 HBC O K-P CH EX 7/66
 R1 D (0.11) DAUM 68 CNTR K-P ELA,POL,SIG 7/70
 R1 0.17 0.04 CAMPBELL 71 DBC - K- NEUTRON ELAST 1/71
 R1 0.18 0.02 LITCHFIE 71 DPWA K-P TO KBAR N 10/71*
 R1 .18 .03 HEMINGWA 75 DPWA O K-P TO KBAR N 11/75*
 R1 .24 .02 RLIC 76 DPWA KBAR N MULTICHLN 1/76*
 R1 D DAUM 68 ASSUMES (J+1/2)*P1 VALUE SEEN IN TOTAL CROSS SECTION.

R2 Y*1(2030) FROM KBAR N INTO LAMBDA PI SQRT(P1*P2)
 R2 (0.20) WOHL 66 HBC O K-P TO LAMBDA PI 7/66
 R2 +0.21 0.01 SMART 68 DPWA - K-N TO LAMBDA PI 6/68
 R2 +0.2 0.02 BERTHON 70 DPWA O K-P TO LAMBDA PI 7/70
 R2 +0.19 0.01 COX 70 DPWA - K-N TO LAMBDA PI 6/70
 R2 +0.16 0.03 GALTIERI 70 DPWA O K-P TO LAMBDA PI 7/70
 R2 +0.20 0.008 LITCHFIE 70 DPWA -O K-N TO LAMBDA PI 6/70
 R2 +.195 .053 DEVENISH 74 O FIXED T DISP REL 4/75*
 R2 +.18 .02 BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
 R2 +.20 .01 VANHORN 75 DPWA O K-P TO LAM P10 11/75*
 R2 +.18 .02 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

R3 Y*1(2030) FROM KBAR N INTO SIGMA PI SQRT(P1*P3)
 R3 L (-0.09) (0.02) BERTHON1 70 DPWA O K-P TO SIGMA PI 10/70
 R3 -0.052 0.010 GALTIERI 70 DPWA O K-P TO SIGMA PI 7/70
 R3 -0.10 0.03 LITCHFIE 71 DPWA K-P TO SIG PI 3/72
 R3 L LITCHFIELD 71 IS AN UPDATE OF BERTHON1 70
 R3 -0.086 0.014 KANE 72 DPWA O K-P TO PI SIG 10/71
 R3 -.15 .03 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

AVERAGE MEANINGLESS (SCALE FACTOR = 2.1)

R4 Y*1(2030) FROM KBAR N INTO XI K SQRT(P1*P4)
 R4 (0.05) OR LESS TRIPP 67 RVUE O K-P TO XI K 8/67
 R4 (0.05) OR LESS BURGUN 68 DPWA O K-P TO XI K 10/69
 R4 (0.023) MULLER 69 DPWA O 7/70

R5 Y*1(2030) FROM KBAR N TO Y*0(1815) PI P-WAVE SQRT(P1*P5)
 R5 1 .18 .04 LITCHF11 74 DPWA O K-P TO L(1815)PI 10/74*
 R5 1 ASSUMES LAMBDA(1815) ELASTICITY=.6 10/74*
 R5 .14 .02 CORDEN2 75 DBC - KBAR PI- NUCLEON 11/75*
 R5 .5

R5 AVG 0.148 0.018 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
 R5 STUDENT 0.148 0.020 AVERAGE USING STUDENT(10)/(1.11) -- SEE TEXT

R6 Y*1(2030) FROM KBAR N TO Y*0(1520) PI D-WAVE SQRT(P1*P6)
 R6 2 -.14 .03 LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 R6 2 ASSUMES LAMBDA(1520) ELASTICITY=.45 10/74*
 R6 3 (.10) (.03) CORDEN2 75 DBC - KBAR PI- NUCLEON 11/75*
 R6 3 UPPER LIMIT 11/75*

R7 Y*1(2030) FROM KBAR N TO Y*0(1520) PI G-WAVE SQRT(P1*P7)
 R7 2 .02 .02 LITCHF12 74 DPWA O K-P TO L(1520)PI 10/74*
 R7 2 ASSUMES LAMBDA(1520) ELASTICITY=.45 10/74*

Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(2030), Σ(2070), Σ(2080)

R8 Y*1(2030) FRCH KEAR N TO KBAR DELTA(1232) F-WAVE SQRT(P1*P8) 10/74*
R8 3 .16 .03 LITCHF13 74 DPWA 0 K-P TO KBAR DEL 11/75*
R8 3 (-17) (.03) CORDEN2 75 DBC - KBAR PI-NUCLEON 11/75*

R9 Y*1(2030) FROM KEAR N TO KBAR DELTA(1232) H-WAVE SQRT(P1*P9) 10/74*
R9 .00 .02 LITCHF13 74 DPWA 0 K-P TO KBAR DEL 10/74*

REFERENCES FOR Y*1(2030)
WOHL 66 PRL 17 107 C G WOHL, F T SOLMITZ, M L STEVENSON (LRL)IJP
TRIPP 67 NP 83 10 + LEITH, + (LRL,SLAC,CERN,HEIDEL,SACLAY)
BURGUN 68 NP 88 447 +MEYER,PAULI,TALLINI + (SACL+CDEF+RHEL)
DAUM 68 NP 87 19 +ERNE,LAGNAUX,SENS,STEUER,UDO (CERN)IJP

CONFIRMS THE SPIN-PARITY ASSIGNMENT.
SMART 68 PR 169 1336 W M SMART (LRL)IJP
MULLER 69 THESIS,UCLR 19372 R A MULLER (LRL)IJP

+RANGAN, VRANA, + (COL FRANCE, RHEL, SACLAY)IJP
+VRANA, BUTTERWORTH, + (CDEF, RHEL, SACLAY)IJP
+ISLAM, COLLEY, + (BIRM,EDIN,GLAS,LOIC)IJP
A BARBARO-GALTIERI (LRL)IJP
P J LITCHFIELD (RUTHERFORD)IJP

+MORTON, NEGUS, GOYAL, MILLER (GLAS, LOIC)IJP
LITCHFIELD, ...+LESQUOY, ... (RHEL+CDEF+SACL)IJP
D F KANE (LBL)IJP

DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
LITCHF11 74 NP 874 12 COX, DARTNELL, KENYON, ONEALE, SUMOROK + (BIRM)IJP
LITCHF12 74 NP 874 19 HEMINGWAY, EADES, HARMSEN + (CERN, HEID, MPT)IJP
LITCHF13 74 NP 874 39 A. J. VAN HORN (LBL)IJP
A. J. VAN HORN (LBL)IJP

RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS + (RHEL+LOIC)IJP

2030 MEV REGION - PRODUCTION AND σTOTAL EXP'TS

28 Y*1(2030, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE THE NOTE TO THE F17 Y*1(2030), WHICH PRECEDES THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE Y*1(2030), BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

28 Y*1(2030) MASS (MEV) (PROD. EXP.)
M (2022.0) (20.0) BLANPIED 65 CNTR 0 GAMMA P TO K+ Y* 6/68
M 2020.0 7.0 BUGG 68 CNTR K-P, D TOTAL 6/70
M 2049.0 4.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
M 2025.0 10.0 COOL 70 CNTR K-P, D TOTAL 10/70
M (2025.0) (20.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71

28 Y*1(2030) WIDTH (MEV) (PROD. EXP.)
W (120.0) (20.0) BLANPIED 65 CNTR 0 6/68
W 130.0 10.0 BUGG 68 CNTR 6/70
W 126.0 11.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
W 165.0 0.0 COOL 70 CNTR K-P, D TOTAL 10/70
W (80.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71

28 Y*1(2030) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*1(2030) INTO KBAR N 497+ 929
P2 Y*1(2030) INTO KBAR N PI 497+ 935+ 139

28 Y*1(2030) BRANCHING RATIOS (PROD. EXP.)
R1 Y*1(2030) INTO (KEAR N)/TOTAL
R1 THESE VALUES OF ELASTICITIES ASSUME J=7/2 --- (P1)
R1 0.131 BUGG 68 CNTR 6/68
R1 0.27 (0.02) BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
R1 0.12 COOL 70 CNTR K-P, D TOTAL 10/70

REFERENCES FOR Y*1(2030) (PROD. EXP.)
BLANPIED 65 PRL 14 741 +GREENBERG, HUGHES, KITCHING, LU, + (YALE, CEA) I
COOL 66 PRL 16 1228 +GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, + (BNL) I
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I

BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
COOL 70 PR D1 1867 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
LU 70 PR D2 1846 +GREENBERG, HUGHES, WINEHART, MORI, + (YALE)

Σ(2070)

F15

THIS STATE HAS BEEN SUGGESTED BY ONLY ONE PARTIAL WAVE ANALYSIS ACROSS THIS REGION. IT NEEDS CONFIRMATION THE RESONANCE PROPOSED BY KANE IS TOO BROAD TO BE USED AS EVIDENCE.

34 Y*1(2070) MASS (MEV)
M (2070.) (10.) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
M (2057.0) KANE 72 DPWA K-P TO SIGMA PI 1/73

34 Y*1(2070) WIDTH (MEV)
W (140.) (20.) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
W (906.0) KANE 72 DPWA K-P TO SIGMA PI 1/73

34 Y*1(2070) PARTIAL DECAY MODES
P1 Y*1(2070) INTO KBAR N 497+ 939
P2 Y*1(2070) INTO SIGMA PI 1197+ 139

34 Y*1(2070) BRANCHING RATIOS
R1 Y*1(2070) FROM KEAR N TO SIGMA SQRT(P1*P2)
R1 (+.12) (-.02) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
R1 (+0.104) KANE 72 DPWA K-P TO SIGMA PI 1/73

REFERENCES FOR Y*1(2070)
BERTHON1 70 NP 824 417 +VRANA, BUTTERWORTH, + (CDEF, RHEL, SACLAY)IJP
KANE 72 PR D5 1583 D F KANE (LBL)IJP

Σ(2080)

P13

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

88 Y*1(2080) MASS (MEV)
M (2082.0) (4.0) COX 70 DPWA - K- N TO LAM PI 6/70
M (2070.0) (30.0) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
M 1 2120. 40. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
M 1 FROM SOLUTION 1 OF BAILLON 75. 1/76*
M 2 2140. 40. BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
M 2 FROM SOLUTION 2 OF BAILLON 75. 1/76*
M 2140. 30. BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*

88 Y*1(2080) WIDTH (MEV)
W (87.0) (20.0) COX 70 DPWA - K- N TO LAM PI 6/70
W (250.0) (40.0) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
W 1 240. 50. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
W 2 200. 50. BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
W 180. 20. BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*

88 Y*1(2080) PARTIAL DECAY MODES
P1 Y*1(2080) INTO KBAR N 497+ 939
P2 Y*1(2080) INTO LAMBDA PI 1115+ 139

88 Y*1(2080) BRANCHING RATIOS
R1 Y*1(2080) FROM KEAR N TO LAMBDA PI SQRT(P1*P2)
R1 (-0.16) (0.03) COX 70 DPWA - K- N TO LAM PI 6/70
R1 (-0.09) (0.03) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
R1 1 -.13 .04 BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
R1 2 -.13 .04 BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
R1 .19 .03 BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*

REFERENCES FOR Y*1(2080)
COX 70 NP 819 61 +ISLAM, COLLEY, + (BIRM, EDIN, GLAS, LOIC)IJP
LITCHF12 70 NP 822 269 P J LITCHFIELD (RUTHERFORD)IJP
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL)IJP
BELLEFOL 75 NP 890 1 DE BELLEFOL, BERTHON, BRUNET + (CDEF, SACL)IJP

Baryons

$\Sigma(2100)$, $\Sigma(2250)$, $\Sigma(2455)$

Data Card Listings

For notation, see key at front of Listings.

$\Sigma(2100)$

G₁₇

26 Y*1(2100, JP=7/2-) I=1
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

Table with 2 columns: Mass (MEV) and branching ratios for Y*1(2100). Rows include M(2060.0), M(2120.0), W(70.0), W(135.0).

Table with 2 columns: Width (MEV) and branching ratios for Y*1(2100). Rows include W(70.0), W(135.0).

Table with 2 columns: Partial Decay Modes and Decay Masses for Y*1(2100). Rows include P1, P2, P3.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2100). Rows include R1, R2.

REFERENCES FOR Y*1(2100)
GALTIERI 70 DUKE CNF 173 A BARBARO-GALTIERI (LRL)JJP

$\Sigma(2250)$ BUMPS

48 Y*1(2250, JP=) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THE PARTIAL-WAVE ANALYSIS RESULTS ARE TOO WEAK TO WARRANT SEPARATING THEM FROM THE PRODUCTION AND CROSS-SECTION EXPERIMENTS.

LASINSKI 71 IN KBAR N, USING A POMERON-RESONANCES MODEL, AND BELLEFON1 75*76 IN DPWA'S OF LAMBDA PI AND SIGMA PI SUGGEST THE PRESENCE OF TWO RESONANCES AROUND THIS MASS VALUE.

Table with 2 columns: Mass (MEV) (PROD. EXP.) and branching ratios for Y*1(2250). Rows include M(2245.0), M(2299.0), M(2250.0), M(2280.0), M(2237.0), M(2255.0), M(2250.0), M(2260.0), M(2215.0), M(2300.0), M(2251.0), M(2270.0), M(2210.0).

Table with 2 columns: Width (MEV) (PROD. EXP.) and branching ratios for Y*1(2250). Rows include W(150.0), W(21.0), W(230.0), W(100.0), W(164.0), W(170.0), W(125.0), W(100.0), W(60.0), W(130.0), W(192.0), W(70.0), W(60.0).

Table with 2 columns: Width (MEV) (PROD. EXP.) and branching ratios for Y*1(2250). Rows include W(102.5), W(96.0).

Table with 2 columns: Partial Decay Modes (PROD. EXP.) and Decay Masses for Y*1(2250). Rows include P1, P2, P3, P4, P5.

48 Y*1(2250) BRANCHING RATIOS (PROD. EXP.)

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R1, R2, R3, R4.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R3, R4.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R4, R5.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R6, R7.

REFERENCES FOR Y*1(2250) (PROD. EXP.)
BLANPIED 65 PRL 14 741 *GREENBERG, HUGHES, KITCHING, + (YALE) (CEA) J
BOCK 65 PL 17 166 *COOPER, FRENCH, KINSON, + (CERN, SACLAY) J
BUGG 68 PR 168 1466 *GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I
BARNES 69 PRL 22 479 *FLAMINIO, MONTANET, SAMIOS + (BNL+SYR) A

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R3, R4.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R3, R4.

Table with 2 columns: Branching Ratios and branching ratios for Y*1(2250). Rows include R3, R4.

$\Sigma(2455)$ BUMPS

53 Y*1(2455, JP=) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THERE IS ALSO SOME SLIGHT EVIDENCE FOR Y* STATES IN THIS MASS REGION FROM THE REACTION GAMMA + P TO K+ + MISSING MASS -- SEE GREENBERG 68.

Table with 2 columns: Mass (MEV) (PROD. EXP.) and branching ratios for Y*1(2455). Rows include M(2455.0), M(2455.0), M(2455.0), M(2455.0).

Table with 2 columns: Width (MEV) (PROD. EXP.) and branching ratios for Y*1(2455). Rows include W(100.0), W(140.0).

Table with 2 columns: Partial Decay Modes (PROD. EXP.) and Decay Masses for Y*1(2455). Rows include P1.

Table with 2 columns: Branching Ratios (PROD. EXP.) and branching ratios for Y*1(2455). Rows include R1, R2, R3, R4, R5, R6.

Table with 2 columns: Branching Ratios (PROD. EXP.) and branching ratios for Y*1(2455). Rows include R1, R2, R3, R4.

Data Card Listings

Baryons

For notation, see key at front of Listings. $\Sigma(2620), \Sigma(3000), \text{EXOTIC HYPERONS}, \Xi^{\prime}s, \Xi^{-}, \Xi^0$

**$\Sigma(2620)$
BUMPS**

54 Y*1(2620, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

54 Y*1(2620) MASS (MEV) (PROD. EXP.)

M	2620.0	15.0	ABRAMS	70 CNTR	K-P, D TOTAL	10/70
M	2542.	22.	DIBIANCA	75 DBC	XI K PI	1/76*
M	AVERAGE MEANINGLESS (SCALE FACTOR = 2.9)					

54 Y*1(2620) WIDTH (MEV) (PROD. EXP.)

W	(175.0)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70
W	221.	81.	DIBIANCA	75 DBC	XI K PI	1/76*

54 Y*1(2620) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y*1(2620) INTO KBAR N	DECAY MASSES
		497+ 939

54 Y*1(2620) BRANCHING RATIOS (PROD. EXP.)

R1	Y*1(2620) INTO (KBAR N)/TOTAL	(P1)	
R1	J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*P1.		
R1	(0.32)	ABRAMS	70 CNTR K-P, D TOTAL
R1	0.36	BRICMAN	70 CNTR 0 TOTAL AND CH EX

REFERENCES FOR Y*1(2620) (PROD. EXP.)

ABRAMS 67 PRL 19 678 +CGOL, GIACOMELLI, KYCIA, LECNTIC, LI, + (BNL)
SUPERSEDED BY ABRAMS 70.
ABRAMS 70 PR 10 1917 +CGOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
BRICMAN 70 PL 31B 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
DIBIANCA 75 NP 89B 137 DIBIANCA, ENDORFR (CERN)

**$\Sigma(3000)$
BUMPS**

59 Y*1(3000, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
ENHANCEMENT IN LAMBDA PI AND KBAR N INVARIANT MASS SPECTRA AND IN MISSING MASS OF NEUTRALS RECOLLING AGAINST K0. EVIDENCE NOT CONCLUSIVE. OMITTED FROM TABLE.

59 Y*1(3000) MASS (MEV) (PROD. EXP.)

M	(3000.0)	EHRlich	66 HBC	0 P1-P	7.91 BEV/C	9/66
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59 Y*1(3000) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y*1(3000) INTO KBAR N	DECAY MASSES
P2	Y*1(3000) INTO LAMBDA PI	497+ 939 1115+ 139

REFERENCES FOR Y*1(3000) (PROD. EXP.)

EHRlich 66 PR 152 1194 R EHRlich, W SELOVE, H YUTA (PENN(BNL)) I

EXOTIC HYPERON CROSS SECTION LIMITS

THIS IS NOT A COMPLETE LIST. WE TABULATE ONLY FROM 1970 ON.

CS	UNITS MICROBARN			
CS	G (20.) OR LESS	GALTIERI 68 DBC	K-N TO SG-PI-PI0	7/70
CS	G ABOVE LIMIT FOR MASS < 2.15 GEV AND GAMMA < 60 MEV- (2.1 GEV/C K-)			7/70
CS	A (40.) OR LESS	GALTIERI 68 DBC	K-N TO SG-PI-PI0	7/70
CS	A ABOVE LIMIT FOR MASS < 2.3 GEV AND GAMMA < 120 MEV- (2.7 GEV/C K-)			7/70

REFERENCES FOR EXOTIC HYPERONS

GALTIERI 68 PRL 21 573 A. BARBARO-GALTIERI, CHADWICK + (LRL+SLAC)

Note on Ξ Resonances

The Ξ resonance situation has long been unsettled. This is because 1) they can only be produced as part of a final state, $K^+p \rightarrow \Xi^* + \text{others}$, and 2) they are so produced with very small cross sections ($< 50 \mu\text{b}$). Thus the numbers of events available are small, and the analysis is more complicated than if direct formation were possible. Only the $\Xi(1530)$ is really well established. There are apparently at least two Ξ states in the 1800-2000 MeV region, and there are indications of several more above 2000 MeV, but the situation is unclear. Two large bubble chamber experiments reported at the 1975 ANL Symposium on New Directions in Hadron Spectroscopy (BLOKZIJL 75 and MORRIS 75) may be able to clarify the situation. In the meantime, we are forced to group together rather disparate observations and await new results. Figures in the listings point out disagreements among various experiments. The table following this note gives our evaluation of the status of the Ξ resonances, based on the data available at this time.

STATUS OF XI* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --				OTHER CHANNELS
			XI PI	LAM K	SIG K	XI* PI	
XI(1320)	P11	****					
XI(1530)	P13	****					WEAK TO LAM PI
XI(1630)		**	**				
XI(1820)		***	***	**	**	**	
XI(1940)		***	***	**	**	**	
XI(2030)		**	**	**	**	**	3-BODY DECAYS
XI(2250)		*	*	*	*	*	3-BODY DECAYS
XI(2500)		**	**	**	**	**	3-BODY DECAYS

**** GOOD, CLEAR, AND UNMISTAKABLE.
*** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
** NEEDS CONFIRMATION.
* WEAK.

S=-2 I=1/2 HYPERON STATES (Ξ)

Ξ^-	22 XI-(1321, JP=1/2) I=1/2 SEE STABLE PARTICLE DATA CARD LISTINGS
Ξ^0	23 XI(1314, JP=1/2) I=1/2 SEE STABLE PARTICLE DATA CARD LISTINGS

Baryons

$\Xi(1530)$, $\Xi(1630)$

$\Xi(1530)$

49 $\Xi^{*1/2}(1530)$, JP=3/2+ I=1/2 P13
THIS IS THE ONLY REALLY WELL-ESTABLISHED Ξ^{*} . THE QUANTUM NUMBERS 3/2+ ARE FAVORED BY THE DATA.

WE DO NOT USE DETERMINATIONS OF THE MASS AND THE WIDTH OF THIS STATE UNLESS THEY ARE ACCOMPANIED BY SOME DISCUSSION OF SYSTEMATICS AND RESOLUTION.

49 $\Xi^{*1/2}(1530)$ MASS (MEV)

Table with columns for mass (M), error (M-), fit (M-FIT), and various decay channels (PJERRU, BADIEN, LONDON, BALTAY, KIRSCH, ROSS2, BERTHON, BELLEF02) with their respective quantum numbers and branching ratios.

49 ($\Xi^{*1/2}$) - (Ξ) MASS DIFFERENCE (MEV)

Table showing mass differences for various decay channels (PJERRU, LONDON, MERRILL, BALTAY, KIRSCH) and average values.

49 $\Xi^{*1/2}(1530)$ WIDTH (MEV)

Table showing widths for various decay channels (SCHLEIN, BERGE, LONDON, BADIEN, BALTAY, BORENSTEIN, KIRSCH, ROSS2, BELLEF02) and average values.

WEIGHTED AVERAGE = 9.14 ± 0.4B
ERROR SCALED BY 1.0

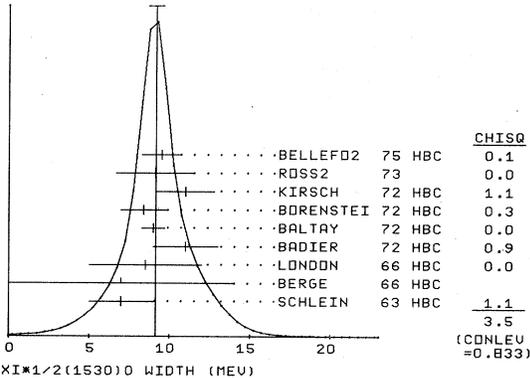


Table showing the weighted average and error scaled by 1.0 for the width of Xi*1/2(1530).

49 $\Xi^{*1/2}(1530)$ REAL PART OF POLE POSITION

Table showing the real part of the pole position for Xi*1/2(1530) with values for Lichtenb 74 and extrap habib173.

Data Card Listings

For notation, see key at front of Listings.

49 $\Xi^{*1/2}(1530)$ IMAGINARY PART OF POLE POSITION

Table showing the imaginary part of the pole position for Xi*1/2(1530) with values for Lichtenb 74 and extrap habib173.

49 $\Xi^{*1/2}(1530)$ PARTIAL DECAY MODES

Table showing partial decay modes for Xi*1/2(1530) into Xi pi and Xi gamma.

49 $\Xi^{*1/2}(1530)$ BRANCHING RATIOS (MEV)

Table showing branching ratios for Xi*1/2(1530) into Xi gamma and Xi pi.

REFERENCES FOR $\Xi^{*1/2}(1530)$

- List of references for Xi*1/2(1530) including PJERRU, SCHLEIN, BERGE, LONDON, MERRILL, BADIEN, BERTHON, ROSS2, KIRSCH, BELLEF02, SHAFER, HABIBI, HUNGERBU, BRIEFEL, BORENSTEIN, LICHTEBN, BALLEF02, KALBFLEI, etc.

PAPERS NOT REFERRED TO IN DATA CARDS
SHAFER 66 PR 142 883 BLTTCN-SHAHER, LINDSEY, MURRAY, SMITH (LRL) JP
A SPIN-PARITY DETERMINATION.
HABIBI 73 NEVIS 199(THESIS) HABIBI (COLU)
HUNGERBU 74 PRD 10 2051 HUNGERBUHLER, MAJKA, (YALE, FNAL, BNL, PITT)
BRIEFEL 75 PRD 12 1859 +GUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)

$\Xi(1630)$

21 $\Xi^{*1/2}(1630)$, JP= 1 I=1/2

THIS EFFECT NEEDS CONFIRMATION.
BARTSCH 69 SEE A SMALL, BROAD ENHANCEMENT NEAR 1650 MEV - IT IS NOT CLEAR THAT IT IS THE SAME PHENOMENON AS BMST 70, WHO FIND CS=3.6±1.6 MICROBARN AT 2.87 GEV/C INCIDENT K- MOMENTUM.
BORENSTEIN 72 SEE NO EFFECT IN THIS REGION. THEY FIND CS<2 MICROBARN AT 2.18.
ROSS 72 ARGUE THAT THE EFFECT THEY SEE IS NOT THE SAME AS THAT SEEN BY BMST 70, AND FIND CS=2±1 MICROBARN AT 3.3 GEV/C.
BELLEFON 75B FIND A CS OF AROUND 10 MICROBARN NEAR 2 GEV/C, BUT LESS THAN 3 MICROBARN AROUND 2.3 GEV/C.

21 $\Xi^{*1/2}(1630)$ MASS (MEV)

Table showing mass values for Xi*1/2(1630) from various experiments (M, M-, M-FIT, M-STUDENT).

21 $\Xi^{*1/2}(1630)$ WIDTH (MEV)

Table showing width values for Xi*1/2(1630) from various experiments (W, W-, W-FIT, W-STUDENT).

21 $\Xi^{*1/2}(1630)$ PARTIAL DECAY MODES

Table showing partial decay modes for Xi*1/2(1630) into Xi pi.

REFERENCES FOR $\Xi^{*1/2}(1630)$

- List of references for Xi*1/2(1630) including APSELL, BARTSCH, KALBFLEI, BORENSTEIN, SCHMIDT, HUNGERBU, BRIEFEL, etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Xi(1820)$, $\Xi(1940)$

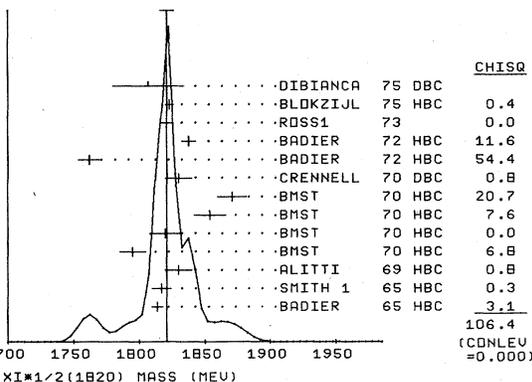
$\Xi(1820)$

50 XI*1/2(1820, JP=) I=1/2

AS THE ACCOMPANYING IDEOGRAMS ILLUSTRATE, THE SITUATION IS CONFUSED, UNTIL SOME FUTURE CLARIFICATION, WE LIST UNDER XI(1820) EVERYTHING REPORTED IN THE MASS RANGE 1750-1875 MEV. WHEN BRANCHING RATIOS ARE REPORTED, WE QUOTE THEM, BUT ONLY THE MOST QUALITATIVE CONCLUSIONS ARE JUSTIFIED.

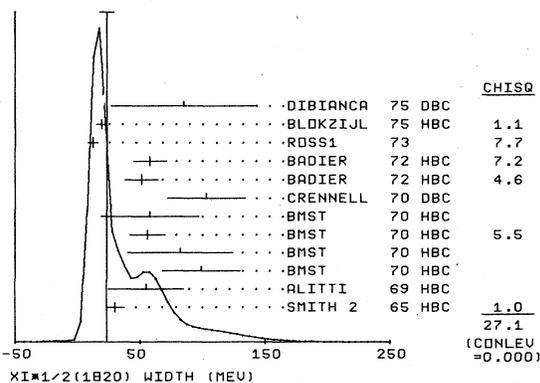
50 XI*1/2(1820) MASS (MEV)

M	(1770.0)		HALSTEINS	63 FBC	-0 K-FR 3.5 GEV/C	
M	30 1814.0	4.0	BADIER	65 HBC	0 LAMBDA KBAR	
M	29 1817.0	7.0	SMITH 1	65 HBC	-0 LAMBDA KBAR	
M	40 1830.0	10.0	ALITTI	69 HBC	- LAM, SIG KBAR	9/69
M	65 1795.0	10.0	BMST	70 HBC	0 XI-PI* (2.9 K-P)	7/70
M	55 1820.0	12.0	BMST	70 HBC	XI(1530) PI	7/70
M	35 1854.0	12.0	BMST	70 HBC	- SIGMA-KO BAR	7/70
M	65 1871.0	11.0	BMST	70 HBC	0 LAMBDA KO BAR	7/70
M	25 1830.0	10.0	CRENNELL	70 DBC	-0 3.6, 3.9 GEV/C	10/70
M	28 1762.0	8.0	BADIER	72 HBC	XI PI, XI2PI, K Y	10/71
M	38 1838.0	5.0	BADIER	72 HBC	XI PI, XI2PI, K Y	10/71
M	8		BADIER	72	ADDS ALL CHANNELS AND DIVIDES PEAK IN LOWER AND HIGHER M	
M	1 30 1821.0	5.0	ROSSI	73	LAMBDA K-KBAR	2/74
M	1				LESS SIGNIFICANT ENHANCEMENT (M=1825, W=100) SEEN IN XI(1530)+PI	2/74
M	100 1823.0	3.0	BLOKZIYL	75 HBC	K- P AT 4.2 GEV	1/76*
M	1807.0	27.0	DIBIANCA	75 DBC	XI 2PI, XI* PI	1/76*
M					AVERAGE MEANINGLESS (SCALE FACTOR = 3.1)	
M					(SEE IDEOGRAM BELOW)	



50 XI*1/2(1820) WIDTH (MEV)

W	(80.0)	OR LESS	HALSTEINS	63 FBC	-0 K-FR 3.5 GEV/C	
W	(12.0)	(4.0)	BADIER	65 HBC	0 LAMBDA KBAR	
W	30.0	7.0	SMITH 2	65 HBC	-0	
W	55.0	40.0	ALITTI	69 HBC	- LAM, SIG KBAR	9/69
W	65 99.0	31.0	BMST	70 HBC	XI-PI* (2.9 K-P)	7/70
W	55 82.0	42.0	BMST	70 HBC	XI(1530) PI	7/70
W	25 56.0	14.0	BMST	70 HBC	- SIGMA-KO BAR	7/70
W	65 58.0	39.0	BMST	70 HBC	0 LAMBDA KO BAR	7/70
W	103.0	38.0	CRENNELL	70 DBC	-0 3.6, 3.9 GEV/C	10/70
W	51.0	13.0	BADIER	72 HBC	LOWER MASS	10/71
W	58.0	13.0	BADIER	72 HBC	HIGHER MASS	10/71
W	12.0		ROSSI	73	LAMBDA K-KBAR	2/74
W	100 19.0	4.0	BLOKZIYL	75 HBC	K- P AT 4.2 GEV	1/76*
W	85.0	58.0	DIBIANCA	75 DBC	XI 2PI, XI* PI	1/76*
W					SEE THE NOTES ACCOMPANYING THE MASSES QUOTED	
W					AVERAGE MEANINGLESS (SCALE FACTOR = 2.3)	
W					(SEE IDEOGRAM BELOW)	



50 XI*1/2(1820) PARTIAL DECAY MODES

P1	XI*1/2(1820)	INTO LAMBDA KBAR	1115+ 497
P2	XI*1/2(1820)	INTO XI PI	1321+ 139
P3	XI*1/2(1820)	INTO SIGMA KBAR	1197+ 457
P4	XI*1/2(1820)	INTO XI*1/2(1530) PI	1533+ 139
P5	XI*1/2(1820)	INTO XI PI PI (EXCLUDING P4)	1321+ 139+ 139

50 XI*1/2(1820) BRANCHING RATIOS

R1	XI*1/2(1820)	INTO (LAMBDA KBAR)/TOTAL	(P1)	
R1	0.3	0.15	ALITTI	69 HBC - 9/69
R2	XI*1/2(1820)	INTO (XI PI)/TOTAL	(P2)	
R2	0.1	0.1	ALITTI	69 HBC - 9/69
R3	XI*1/2(1820)	INTO (SIGMA KBAR)/TOTAL	(P3)	
R3	(0.02) OR LESS	0.15	TRIPP	67 RVUE 8/67
R3	0.3	0.15	ALITTI	69 HBC - 9/69
R4	XI*1/2(1820)	INTO (XI*1/2(1530) PI)/TOTAL	(P4)	
R4	0.3	0.15	ALITTI	69 HBC - 9/69
R4	(0.25) OR LESS	0.15	DAUBER	69 HBC - K-P 2.7 BEV/C 9/69
R5	XI*1/2(1820)	INTO (XI PI)/(LAMBDA KBAR)	(P2)/(P1)	
R5	0.20	0.20	BADIER	65 HBC 7/66
R5	(.42)+/-.02) OR LESS	0.20	BLOKZIYL	75 HBC K- P AT 4.2 GEV 1/76*
R6	XI*1/2(1820)	INTO (XI*1/2(1530) PI)/(LAMBDA KBAR)	(P4)/(P1)	
R6	0.26	0.13	SMITH 1	65 HBC
R6	.77	.20	BLOKZIYL	75 HBC K- P AT 4.2 GEV 1/76*
R6				
R6	AVG	0.41	0.23	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.1)
R6	STUDENT	0.39	0.14	AVERAGE USING STUDENT10(M/1.11) -- SEE TEXT
R7	XI*1/2(1820)	INTO (XI PI PI)/(LAMBDA KBAR)	(P5)/(P1)	
R7	(0.1) OR MORE		SMITH 1	65 HBC
R8	XI*1/2(1820)	INTO (XI PI)/(XI*1/2(1530) PI)	(P2)/(P4)	
R8	1.5	0.6	APSELL	70 HBC 0 6/70
R9	XI*1/2(1820)	INTO (XI PI)/(XI*1/2(1530) PI)	(P5)/(P4)	
R9	0.3	0.5	APSELL	70 HBC 0 6/70
R10	XI*1/2(1820)	INTO (SIGMA KBAR)/(LAMBDA KBAR)	(P3)/(P1)	
R10	.23	.06	BLOKZIYL	75 HBC K- P AT 4.2 GEV 1/76*

REFERENCES FOR XI*1/2(1820)

HALSTEIN 63 SIENA CONF 173 HALSTEINSLID, (BERGEN, CEPN, EPOL, RHEL, LOUC) I
 BADIER 65 PL 16 171 + DEMOULIN, GOLDBERG, + (EPOL, SACLAY, AMST) I
 SMITH 1 65 PRL 16 25 + LINDSEY, BUTTCH-SHAFFER, MURRAY (LRL) JP
 SMITH 2 65 ATHENS CONF 251 G A SMITH, J S LINDSEY (LRL)
 TRIPP 67 NP 83 10 + LEITH, + (LRL, SLAC, CERN, HEIDEL, SACLAY)
 LSES DATA OF SMITH 1.

ALITTI 69 PRL 22 79 +BARNES, FLAMINIO, METZGER, + (BNL, SYRACUSE) I
 DAUBER 69 PR 179 1262 +BERGE, HUBBARD, MERRILL, MULLER (LRL)
 APSELL 70 PRL 24 777 + (BRANDEIS, MARYLAND, SYRACUSE, TUFTS) I
 70 DUKE 317 BRANDEIS+MARYLAND+SYRACUSE+TUFTS COLLABOR.
 CRENNELL 70 PR 10 847 +KARSHON, LAI, ONEALL, SCARR, SCHUMANN (BNL)
 BADIER 72 NP 837 429 +BARRELET, CHARLTON, VIDEAU (EPOL)
 ROSSI 73 PURDUE CONF. 345 ROSS, LLOYD, PADOJICIC (OXFORD)

BLOKZIYL 75 ANL-HEP-CP-75-58 +DEGROOT, HODGLAND+ (AMST+CERN+NIJ+OXF)
 ALSO 75 PALERMO CONF. +DEGROOT, HODGLAND+ (AMST+CERN+NIJ+OXF)
 DIBIANCA 75 NP 898 137 DIBIANCA, ENDORFER (CERN)

PAPERS NOT REFERRED TO IN DATA CARDS

MERRILL 68 PR 167 1202 D W MERRILL, J BUTTON-SHAFFER (LRL)
 WEAK EVIDENCE CONCERNING JP.

APSELL 69 PRL 23 884 + (BRANDEIS, MARYLAND, SYRACUSE, TUFTS)
 SUPERSEDED BY BMST 70.

SCHMIDT 73 PURDUE CONF. 363 SCHMIDT (BRANDEIS)
 BRIEFEL 75 PRD 12 1859 +GOUREVITCH, KIRSCH+ (BRAN+UMD+SYR+TUFT)

$\Xi(1940)$

52 XI*1/2(1940, JP=) I=1/2

WE LIST UNDER XI(1940) EVERYTHING REPORTED IN THE MASS RANGE 1875-2300 MEV.

52 XI*1/2(1940) MASS (MEV)

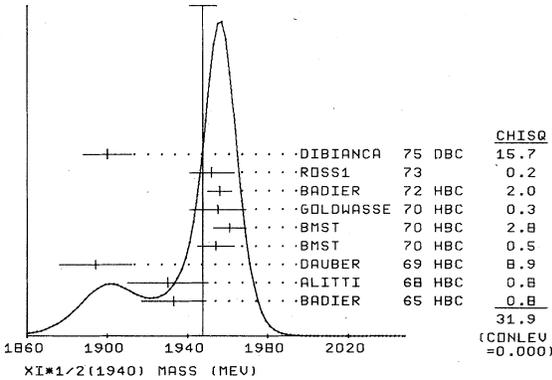
M	35 1933.0	16.0	BADIER	65 HBC	0 XI- PI*	
M	27 1930.0	20.0	ALITTI	68 HBC	0 XI- PI*	11/68
M	66 1894.0	18.0	DAUBER	69 HBC	- XI PI	11/68
M	110 1954.0	9.0	BMST	70 HBC	0 XI-PI+ (2.9 K-P)	7/70
M	40 1961.0	8.0	BMST	70 HBC	XI(1530) PI	7/70
M	21 1955.0	14.0	GOLDWASSE	70 HBC	- XI PI	10/70
M	29 1956.0	6.0	BADIER	72 HBC	XI PI, XI2PI, K Y	10/71
M	25 1952.0	11.0	ROSSI	73	(XI PI)-	2/74
M	1900.0	12.0	DIBIANCA	75 DBC	XI PI	1/76*
M					AVERAGE MEANINGLESS (SCALE FACTOR = 2.0)	
M					(SEE IDEOGRAM BELOW)	

Baryons

$\Xi(1940)$, $\Xi(2030)$, $\Xi(2250)$

Data Card Listings

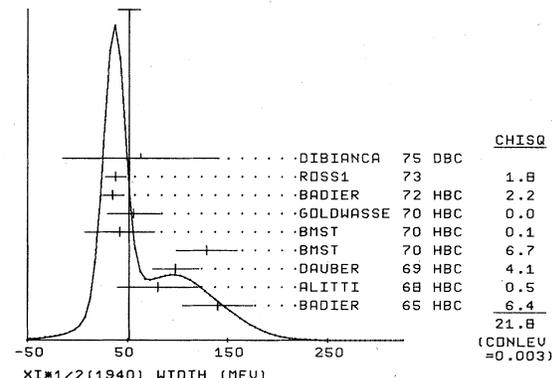
For notation, see key at front of Listings.



52 $\Xi^*1/2(1940)$ WIDTH (MEV)

W	Mass (MeV)	Width (MeV)	Experiment	Notes	CHISQ
W	35	140.0	BADIER	65 HBC	0 XI- PI+
W	27	80.0	ALITTI	68 HBC	0 XI- PI+
W	66	98.0	DAUBER	69 HBC	- XI PI
W	110	129.0	BMST	70 HBC	0 XI-PI+ (2.9 K-P)
W	40	42.0	BMST	70 HBC	XI(1530) PI
W	21	56.0	GOLDWASSE	70 HBC	- XI PI
W	29	35.0	BADIER	72 HBC	XI PI, XI2PI, K Y
W	38	10.0	ROSSI	75	(XI PI)-
W	63	78.0	DIBIANCA	75 DBC	XI PI

AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)
(SEE IDEOGRAM BELOW)



52 $\Xi^*1/2(1940)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^*1/2(1940)$ INTO XI PI	1321+ 139
P2	$\Xi^*1/2(1940)$ INTO XI*(1530) PI	153+ 135
P3	$\Xi^*1/2(1940)$ INTO XI PI PI (EXCLUDING P2)	1321+ 139+ 139
P4	$\Xi^*1/2(1940)$ INTO XI0 PI-	1314+ 139
P5	$\Xi^*1/2(1940)$ INTO XI- P10	1321+ 134

52 $\Xi^*1/2(1940)$ BRANCHING RATIOS

THE $\Xi(1940)$ IS SEEN MAINLY IN XI PI AND SOME IN XI(1530) PI. IT HAS BEEN LOOKED FOR IN OTHER CHANNELS BUT NOT SEEN.

R	Decay Mode	Branching Ratio	Experiment	Notes	CHISQ
R1	$\Xi^*1/2(1940)$ INTO (XI PI)/(XI*(1530) PI)	(P1)/(P2)			
R1		2.8	0.7	0.6 APSELL 70 HBC	0
R2	$\Xi^*1/2(1940)$ INTO (XI PI PI)/(XI*(1530) PI)	(P3)/(P2)			
R2		0.0	0.3	APSELL 70 HBC	0
R3	$\Xi^*1/2(1940)$ INTO (XI0 PI-)/(XI- P10)	(P4)/(P5)			
R3		25	2.6	0.6	1.6 ROSSI 73 (XI PI)-
R3	1 THIS BR IS 2.0(0.5) FOR AN I=1/2(I=3/2) XI*(1940).				2/74

REFERENCES FOR $\Xi^*1/2(1940)$

BADIER	45 PL 16 171	+DEMULIN, GELDBERG, + (EPOL, SACLAY, AMST) I
ALITTI	68 PRL 21 1119	+FLAMINIO, METZGER, RADOJICIC, + (BNL, SYRACUSE) I
DAUBER	69 PR 179 1262	+BERGE, HUBBARD, MERRILL, MULLER (LRL) I
APSELL	70 PRL 24 777	+ (BRANDEIS, MARYLAND, SYRACUSE, TUFTS) I
BMST	70 DUKE 317	BRANDEIS+MARYLAND+SYRACUSE+TUFTS COLLABOR.
GOLDWASSE	70 PR 1D 1960	E L GOLDWASSER, P F SCHULTZ (ILLINOIS)
BADIER	72 NP 837, 429	+BARRELET, CHARLTON, VIDEAU (EPOL)
ROSSI	73 PURDUE CONF. 345	RCSS, LLOYD, RADOJICIC (OXFORD)
DIBIANCA	75 NP 898 137	DIBIANCA, ENDORFR (CARN)

PAPERS NOT REFERRED TO IN DATA CARDS

APSELL 69 PRL 23 884 + (BRANDEIS, MARYLAND, SYRACUSE, TUFTS)
 SUPERSEDED BY BMST 70.
 SCHMIDT 73 PURDUE CONF. 363 + SCHMIDT (BRANDEIS)
 BRIEFEL 75 PRD 12 1859 +GOUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)

$\Xi(2030)$

68 $\Xi^*1/2(2030)$, JP=) I=1/2
 THE STATE AT 2129 MEV REPORTED BY BLOKZIJL 75 IS REALLY INTERMEDIATE BETWEEN XI*(2130) AND XI*(2250). WE LIST IT HERE TEMPORARILY.

68 $\Xi^*1/2(2030)$ MASS (MEV)

M	Mass (MeV)	Width (MeV)	Experiment	Notes	CHISQ
M	42	2030.0	10.0	ALITTI 69 HBC	K-P 3.9-5 BEV/C
M	40	2058.0	17.0	BARTSCH 69 HBC	K-P 10GEV/C
M	15	2019.0	7.0	ROSSI 73	SIGMA KBAR
M	25	2129.0	5.0	BLOKZIJL 75 HBC	K- P AT 4.2 GEV
M	2044.0	8.0	DIBIANCA 75 DBC	XI 2PI, XI* PI	1/76*

AVERAGE MEANINGLESS (SCALE FACTOR = 7.4)

68 $\Xi^*1/2(2030)$ WIDTH (MEV)

W	Mass (MeV)	Width (MeV)	Experiment	Notes	CHISQ
W	45.0	40.0	20.0	ALITTI 69 HBC	-
W	57.0	30.0		BARTSCH 69 HBC	-
W	15	35.0	17.0	ROSSI 73	SIGMA KBAR
W	25	21.0	8.0	BLOKZIJL 75 HBC	K- P AT 4.2 GEV
W	60.0	24.0		DIBIANCA 75 DBC	XI 2PI, XI* PI

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

68 $\Xi^*1/2(2030)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^*1/2(2030)$ INTO XI PI	1321+ 139
P2	$\Xi^*1/2(2030)$ INTO LAMBDA KBAR	1115+ 497
P3	$\Xi^*1/2(2030)$ INTO SIGMA KBAR	1197+ 497
P4	$\Xi^*1/2(2030)$ INTO XI*(1530) PI	153+ 139
P5	$\Xi^*1/2(2030)$ INTO LAMBDA (OR SIGMA) KBAR PI	1115+ 497+ 139

68 $\Xi^*1/2(2030)$ BRANCHING RATIOS

R	Decay Mode	Branching Ratio	Experiment	Notes	CHISQ
R1	$\Xi^*1/2(2030)$ INTO (XI PI)/(MODES P1 TO P4)	(P1)/(P1+P2+P3+P4)			
R1		0.30	OR LESS	ALITTI 69 HBC	- 1 STD DEV LIMIT
R2	$\Xi^*1/2(2030)$ INTO (LAM KBAR)/(MODES P1 TO P4)	(P2)/(P1+P2+P3+P4)			
R2		0.25	0.15	ALITTI 69 HBC	-
R3	$\Xi^*1/2(2030)$ INTO (SIG KBAR)/(MODES P1 TO P4)	(P3)/(P1+P2+P3+P4)			
R3		0.75	0.20	ALITTI 69 HBC	-
R4	$\Xi^*1/2(2030)$ INTO (XI* PI)/(MODES P1 THRU P4)	(P4)/(P1+P2+P3+P4)			
R4		0.15	OR LESS	ALITTI 69 HBC	- 1 STD DEV LIMIT
R5	$\Xi^*1/2(2030)$ INTO LAMBDA (OR SIGMA) KBAR PI (P5)	(P5)			
R5		SEEN	BARTSCH 69 HBC		9/69

REFERENCES FOR $\Xi^*1/2(2030)$

ALITTI	69 PRL 22 79	+BARNES, FLAMINIO, METZGER, + (BNL, SYRACUSE) I
BARTSCH	69 PL 288 439	+ (AACHEN, BERLIN, CERN, LOIC, VIENNA)
ROSSI	73 PURDUE CONF. 345	RCSS, LLOYD, RADOJICIC (OXFORD)
BLOKZIJL	75 ANL-HEP-CP-75-58	+DEGROOT, HOOGLAND+ (AMST+CERN+NIJH+OXF)
BRIEFEL	75 PALERMO CONF.	+DEGROOT, HOOGLAND+ (AMST+CERN+NIJH+OXF)
DIBIANCA	75 NP 898 137	DIBIANCA, ENDORFR (CARN)

$\Xi(2250)$

22 $\Xi^*1/2(2250)$, JP=)
 THE EVIDENCE FOR THIS STATE IS WEAK. BARTSCH 69 SEE A BUMP OF NOT MUCH STATISTICAL SIGNIFICANCE IN LAMBDA-KBAR-PI, SIGMA-KBAR-PI, AND XI-PI-PI MASS SPECTRA. GOLDWASSE 70 SEE A NARROWER BUMP IN XI-PI-PI AT A HIGHER MASS. PERHAPS THEY ARE THE SAME STATE, PERHAPS THEY ARE NOT, BUT SEE ALSO MORRIS 75.

22 $\Xi^*1/2(2250)$ MASS (MEV)

M	Mass (MeV)	Width (MeV)	Experiment	Notes	CHISQ
M	35	2244.0	52.0	BARTSCH 69 HBC	K-P 10 GEV/C
M	18	2295.0	15.0	GOLDWASSE 70 HBC	K-P 5.5 GEV/C

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

22 $\Xi^*1/2(2250)$ WIDTH (MEV)

W	Mass (MeV)	Width (MeV)	Experiment	Notes	CHISQ
W	130.0	80.0		BARTSCH 69 HBC	-
W	LESS THAN	30.0		GOLDWASSE 70 HBC	-

K-P 5.5 GEV/C

22 $\Xi^*1/2(2250)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^*1/2(2250)$ INTO XI PI PI	1321+ 139+ 139
P2	$\Xi^*1/2(2250)$ INTO LAMBDA KBAR PI	1115+ 497+ 139
P3	$\Xi^*1/2(2250)$ INTO SIGMA KBAR PI	1197+ 497+ 139

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Xi(2250)$, $\Xi(2500)$, Ω^-

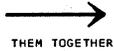
REFERENCES FOR $\Xi(2500)$

BARTSCH 69 PL 288 439 * (AACHEN, BERLIN, CERN, LOIC, VIENNA)
 GOLDWASS 70 PR 10 1960 E L GOLDWASSER, P F SCHULTZ (ILLINOIS)

PAPERS NOT REFERRED TO IN DATA CARDS.

MORRIS 75 ANL-HEP-CP-75-58 MORRIS, CH, PARKER, SMITH, WHITMORE (MSU)

$\Xi(2500)$



99 $\Xi(2500)$, JP=) I=1/2

IT IS QUITE POSSIBLE THAT THE REASON THE EXPERIMENTS DISAGREE ABOUT THE MASS AND WIDTH IS THAT THEY ARE SEEING DIFFERENT Ξ 'S. FOR NOW, HOWEVER, WE GROUP THEM TOGETHER.

99 $\Xi(2500)$ MASS (MEV)

M	30	2430.0	20.0	ALITTI	69 HBC	-	K-P 4.6-5 GEV/C	9/69
M	45	2500.0	10.0	BARTSCH	69 HBC	-0	K-P 10 GEV/C	9/69
M		2392.	27.	DIBIANCA	75 DBC		XI 2PI	1/76*
M	AVERAGE MEANINGLESS (SCALE FACTOR = 3.2)							

99 $\Xi(2500)$ WIDTH (MEV)

W	150.0	60.0	40.0	ALITTI	69 HBC	-		9/69
W	59.0	27.0		BARTSCH	69 HBC	-0		9/69
W	75.	69.		DIBIANCA	75 DBC		XI 2PI	1/76*
W	AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)							

99 $\Xi(2500)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\Xi(2500)$ INTO XI PI	1321+ 139
P2	$\Xi(2500)$ INTO LAMBDA KBAR	1115+ 497
P3	$\Xi(2500)$ INTO SIGMA KBAR	1197+ 457
P4	$\Xi(2500)$ INTO $\Xi(1530)$ PI	1533+ 139
P5	$\Xi(2500)$ INTO LAMBDA (OR SIGMA) KBAR PI	1115+ 497+ 139
P6	$\Xi(2500)$ INTO XI PI PI	1321+ 139+ 139

99 $\Xi(2500)$ BRANCHING RATIOS

R1	$\Xi(2500)$ INTO (XI PI)/(MODES P1 THRU P4)	(P1)/(P1+P2+P3+P4)	
R1	(0.5) OR LESS	ALITTI 69 HBC	1 STD DEV LIMIT 9/69
R2	$\Xi(2500)$ INTO (LAM KBAR)/(MODES P1 THRU P4)	(P2)/(P1+P2+P3+P4)	
R2	0.5 0.2	ALITTI 69 HBC	- 9/69
R3	$\Xi(2500)$ INTO (SIG KBAR)/(MODES P1 THRU P4)	(P3)/(P1+P2+P3+P4)	
R3	0.5 0.2	ALITTI 69 HBC	- 9/69
R4	$\Xi(2500)$ INTO (XI* PI)/(MODES P1 THRU P4)	(P4)/(P1+P2+P3+P4)	
R4	(0.2) OR LESS	ALITTI 69 HBC	1 STD DEV LIMIT 9/69
R5	$\Xi(2500)$ INTO (LAMBDA (OR SIGMA) KBAR PI)/TOTAL	(P5)	
R5	SEEN	BARTSCH 69 HBC	-0 9/69
R6	$\Xi(2500)$ INTO (XI PI PI)/TOTAL	(P6)	
R6	SEEN	BARTSCH 69 HBC	-0 9/69

REFERENCES FOR $\Xi(2500)$

ALITTI 69 PRL 22 79 +BARNES, FLAMINIO, METZGER, + (BNL, SYRACUSE) I
 BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LOIC, VIENNA)
 DIBIANCA 75 NP 898 137 DIBIANCA, ENDORFER (CERN)

S=-3 I=0 HYPERON STATE (Ω)

Ω^-

24 OMEGA-(1675, JP=3/2+) I=0

SEE STABLE PARTICLE DATA CARD LISTINGS

Appendix I

TEST OF $\Delta I=1/2$ RULE FOR K DECAYS

The quantities of interest for making tests of theoretical predictions regarding the $\Delta I=1/2$ rule for K decay are usually partial decay rates for single channels or special sums of channels. It is not possible to compute the errors on sums, differences, and ratios of partial decay rates from the information given in the Table of Stable Particles because of the presence of off-diagonal terms in the error matrix. For this reason we give some of these quantities in Table I. Throughout this Appendix, italics are used to indicate that a quantity has changed by more than one (old) standard deviation since our previous edition, and S gives the scale factor included in the quoted error because of inconsistencies in the data (see footnote at end of Stable Particle Table for definition of S).

Table I. (000) and (+-0) refer to the sign of the pions into which the K_L decays.

$\Gamma_{K_{L3}^+} = \Gamma_{K_{e3}^+} + \Gamma_{K_{\mu 3}^+}$	$= (6.483 \pm 0.090) 10^6 \text{ sec}^{-1}$
$\Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}$	$= 0.663 \pm 0.018 \quad S=1.7^*$
$\Gamma_{K_{\mu 3}^+} / \Gamma_{K_{\tau}^+}$	$= 3.227 \pm 0.083$
$\Gamma_{K_{L3}^0} = \Gamma_{K_{e3}^0} + \Gamma_{K_{\mu 3}^0}$	$= (12.75 \pm 0.15) 10^6 \text{ sec}^{-1} \quad S=1.1^*$
$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0}$	$= 0.696 \pm 0.017$
$\Gamma_{K^0(000)} / \Gamma_{K^0(+ - 0)}$	$= 1.747 \pm 0.070 \quad S=1.1^*$

1. Leptonic decay rates

The $\Gamma_{K_{L3}}$ rates are useful in testing the leptonic $\Delta I=1/2$ rule in the way suggested by Trilling.¹ The predictions are

$$\Gamma_{K_{L3}^0} / 2\Gamma_{K_{L3}^+} = 1.012, \text{ a phase-space factor,}^2$$

and

$$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0} = \Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}$$

From Table 1,

$$\Gamma_{K_{L3}^0} / 2\Gamma_{K_{L3}^+} = 0.983 \pm 0.018$$

and

$$\frac{\Gamma_{K_{\mu 3}^0}}{\Gamma_{K_{e3}^0}} \left[\frac{\Gamma_{K_{\mu 3}^+}}{\Gamma_{K_{e3}^+}} \right]^{-1} = 1.050 \pm 0.038 .$$

These results seem to show a less than 2σ disagreement with the predictions, but the errors should be regarded with caution in view of the internal disagreements in the data. (Note the ideograms in the Data Listings for the charged K meson.)

2. Three-pion decays

We follow here the tests done by Mast et al.,³ based on the general analysis of K decays suggested by Zemach.⁴ Both decay rates (Γ) and slopes (g , the energy dependence of the Dalitz plot distributions) are used. The $\Delta I=1/2$ rule predicts that the following test quantities are all equal to zero:

$$\text{Test 1} = \frac{2}{3} \frac{\Gamma_{K^0(000)}}{\phi_1} \left[\frac{\Gamma_{K^0(+ - 0)}}{\phi_2} \right]^{-1} - 1 ,$$

$$\text{Test 2} = \frac{1}{4} \frac{\Gamma_{K_{\tau}^+}}{\phi_3} \left[\frac{\Gamma_{K_{\tau}^+}}{\phi_4} \right]^{-1} - 1 ,$$

$$\text{Test 3} = \frac{1}{2} \frac{\Gamma_{K_{\tau}^+}}{\phi_3} \left[\frac{\Gamma_{K^0(+ - 0)}}{\phi_2} \right]^{-1} - 1 ,$$

$$\text{Test 4} = \frac{1}{2} g_{K_{\tau}^+} + g_{K_{\tau}^+} ,$$

$$\text{Test 5} = g_{K^0(+ - 0)} + g_{K_{\tau}^+} - \frac{1}{2} g_{K_{\tau}^+} .$$

The ϕ_i are phase-space factors which have been calculated as described in Mast et al.³ by use of a relativistic formulation and the masses and slopes from this edition. The factors labeled UDP are the relative areas of the Dalitz plots, assuming a uniform distribution. The NU DP include the observed slopes (see below). The CNU DP have been calculated by including the final-state Coulomb interaction.

The values are:

	Method		
	UDP	NU DP	CNU DP
$\phi_1(000) =$	1.489	1.489	1.444
$\phi_2(+ - 0) =$	1.221	1.300	1.284
$\phi_3(++ -) =$	1.000	1.000	1.000
$\phi_4(+00) =$	1.247	1.180	1.144

For convenience, we repeat the slope parameters tabulated in the Stable Particle Table. They are as follows:

$$g_{K_{\tau}^+} = -0.214 \pm 0.005 \quad S=1.7^*$$

$$g_{K_{\tau}^-} = -0.214 \pm 0.007 \quad S=2.7^*$$

$$\bar{g}_{K_{\tau}^+} = -0.214 \pm 0.004$$

$$g_{K_{\tau}^0} = 0.550 \pm 0.020 \quad S=1.6^*$$

$$g_{K^0(+ - 0)} = 0.646 \pm 0.014 \quad S=2.5^*$$

A difference in the τ^+ and τ^- slopes would be an indication of CP violation in this decay. Since no difference is observed at this time, we average the two and use this value in Test 4 and Test 5.

We use the CNUDP factors and the rates and slopes reported in this edition to compute the five test quantities which the $\Delta I=1/2$ rule predicts to be zero. The results are:

$$\begin{aligned}\text{Test 1} &= 0.036 \pm 0.042 \\ \text{Test 2} &= -0.077 \pm 0.024 \\ \text{Test 3} &= 0.227 \pm 0.021 \\ \text{Test 4} &= 0.061 \pm 0.011 \\ \text{Test 5} &= 0.157 \pm 0.018\end{aligned}$$

The three-pion final state can be in isospin states $I = 1, 2, 3$. Tests 1 and 2 test the existence of isospin $I = 3$ in the final state. Since the rate tests (Tests 1, 2, and 3) could differ from zero by as much as 0.1 owing to the mass differences and the occurrence of big slopes⁵, no evidence for $I=3$ is found. Test 4 is related to the $I=2$ amplitude in the final state and indicates the presence of $I=2$. Tests 3 and 5 give information on the $\Delta I=3/2$ part of the $I=1$ amplitude relative to the $\Delta I=1/2$ part. Both tests indicate the presence of $\Delta I=3/2$.

References

1. G. Trilling, K-Meson Decays, UCRL-16473, (updated from Argonne Conference Proceedings, 1965, p. 115).
2. N. Brene (CERN), private communication. In our Jan. 1968 edition we had erroneously used 1.04.
3. T. S. Mast, L. K. Gershwin, M. Alston-Garnjost, R. O. Bangerter, A. Barbaro-Galtieri, J. J. Murray, F. T. Solmitz, and R. D. Tripp, Phys. Rev. **183**, 1200 (1969).
4. C. Zemach, Phys. Rev. **133**, B1201 (1964).
5. C. Bouchiat and M. Veltman, Topical Conference on Weak Interactions, CERN 69-7 (1969), p. 225.

Appendix II

A. SU(3) CLASSIFICATION OF BARYON RESONANCES

It is established that a symmetry higher than SU(3) is necessary to classify the known baryon resonances. However, many higher-symmetry schemes have been proposed, and even for SU(6) various versions exist (for a review see Dalitz¹). Since it is not clear which one of these schemes best fits the data, we do not review them here, but we report once again fits of baryon states into SU(3) multiplets.

For the reader's convenience, we collect here the relevant formulae.

Exact SU(3) symmetry predicts that all the members of a multiplet should have the same mass and the same couplings for decays into other multiplets. It has been found, however, that the members of the octet of stable baryons lie within 20% of their mean mass; therefore a symmetry-breaking interaction has been introduced by Gell-Mann and Okubo independently.² In addition, for the isospin-0 vector mesons (ω and ϕ), an additional symmetry-breaking interaction has been introduced by Sakurai³ to take care of octet-singlet mixing. The relevant formulae for masses and decay rates are given below.

Mass Formulae

Broken SU(3) gives:

$$\text{Decuplet} \quad \Delta - \Sigma = \Sigma - \Xi^* = \Xi^* - \Omega \quad \text{GMO} \quad (1)$$

$$\text{Octet} \quad 2(N + \Xi) = 3\Lambda + \Sigma \quad \text{GMO} \quad (2)$$

$$\text{Octet-Singlet} \left\{ \begin{array}{l} \sin^2 \theta = \frac{\Lambda - M_B}{\Lambda - \Lambda'} \\ M_B = \frac{2(N + \Xi) - \Sigma}{3} \end{array} \right. \quad \begin{array}{l} \text{Mixing} \\ \text{angle}^{\text{18}} \end{array} \quad (3)$$

$$\text{mixing} \left\{ \begin{array}{l} \sin^2 \theta = \frac{\Lambda - M_B}{\Lambda - \Lambda'} \\ M_B = \frac{2(N + \Xi) - \Sigma}{3} \end{array} \right. \quad \text{GMO} \quad (4)$$

Here GMO stands for the Gell-Mann-Okubo formula; the particle symbol indicates its mass. The formulae would be the same if squared masses were used. For the nonet case, Λ is the "mostly-octet" particle, Λ' is the "mostly-singlet" particle.

Decay Rates

In terms of a relativistically invariant matrix element T , the decay rate for two-body decay of a resonance of mass M_R is

$$\Gamma \propto \frac{|T|^2 R_2}{M_R}, \quad (5)$$

where $R_2 = k/M_R$ is the two-body phase space factor. Since the numerator is an invariant, and since Γ must transform as $1/E$, we introduce the denominator M_R .⁴

For meson decays (see below) the rates are calculated according to Eq. (5); for baryon resonance decays into $1/2^+$ baryons and 0^- mesons, one next takes into account the fact that spin sums in $|T|^2$ introduce another factor M_R , cancelling the $1/M_R$. We are then left with

$$\Gamma = \frac{|T|^2 k}{M_R} M_N, \text{ for baryons} \quad (5')$$

$$= \frac{|T|^2 k}{M_R^2} M_N^2, \text{ for mesons.} \quad (5'')$$

The powers of the nucleon mass M_N or M_N^2 have been introduced so that we can treat $|T|$ as dimensionless.

$|T|^2$ contains centrifugal barrier factors, which we call B_l . We then have

$$\left. \begin{array}{l} \text{Decuplet} \\ \text{Singlet} \end{array} \right\} \Gamma = (c_g)^2 B_l(k) \frac{M_N}{M_R} k \quad (6)$$

$$\text{Octet} \quad \Gamma = (c_D g_D + c_F g_F)^2 B_l(k) \frac{M_N}{M_R} k \quad (7)$$

$$\left. \begin{array}{l} \text{Octet-Singlet} \\ \text{mixing} \end{array} \right\} \begin{array}{l} \Lambda = G_8 \cos \theta + G_1 \sin \theta \\ \Lambda' = -G_8 \sin \theta + G_1 \cos \theta \end{array} \quad (8)$$

$$\text{with} \quad \begin{array}{l} G_8 = c_D g_D + c_F g_F \\ G_1 = c_1 g_1. \end{array} \quad (9)$$

Here c_i are the SU(3) coefficients with the sign convention adopted in this article [see note in the Table of SU(3) Isoscalar Factors and Fig. 2 in the text]. M_N is the nucleon mass, M_R is the resonance mass for which Γ is calculated, k is the center-of-mass momentum for the channel being considered, and g_i are the relevant couplings. For the case of singlet-octet mixing, formula (8) has to be used in conjunction with (6) and (7). G_8 and G_1 represent the couplings for the multiplet, and Λ and Λ' represent the couplings for the physical states.

The relation between g_D , g_F , and the parameter α is

$$\alpha = \left[1 + \frac{\sqrt{5}}{3} \frac{g_F}{g_D} \right]^{-1}. \quad (10)$$

Exact SU(3) predicts that the couplings g_i for all the members of a multiplet are the same; however, since the symmetry is broken for the masses, it is probably broken for the widths. In the case of the $3/2^+$ decuplet, for broken SU(3) a sum rule has been derived by Becchi⁵ and by Gupta⁶ independently. It relates the g_i for the members of the decuplet by the relation

$$2(\Delta + \Xi) = 3\Sigma^*(\Lambda\pi) + \Sigma^*(\Sigma\pi), \quad (11)$$

where $\Sigma^*(\Lambda\pi)$ is the coupling for the $\Sigma(1385) \rightarrow \Lambda\pi$ decay and $\Sigma^*(\Sigma\pi)$ is the coupling for the decay $\Sigma(1385) \rightarrow \Sigma\pi$.

As mentioned in the text (Sec. IV B) the determination of the relative signs of resonant amplitudes can be useful in making an SU(3) assignment of resonances. In fact the resonant amplitude $T \propto \sqrt{\kappa_e \kappa_i} \propto G_e G_i$, where the subscript e refers to the elastic channel and the G_e , G_i are the couplings of Eqs. (6) through (9). Assuming that all g_i are positive, the sign of the G_i are dependent upon the sign of the Clebsh-Gordon coefficients c_i . Once a sign convention is adopted (we use the Levi-Setti⁷ convention, see Fig. 2 in the text) and the signs for a Σ state ($I=1$) and a Λ state ($I=0$) of known SU(3) assignment have been chosen for reference, the signs of all the other amplitudes can be useful in determining multiplet assignments. For exact SU(3) all the decays of members of a decuplet have the same sign. For octets the relative sign depends upon the value of g_D/g_F and the mixing angle, as seen from Eqs.(7) through (9).

Fits to the Data

Fits of baryon decay rates within SU(3) can be found in, among others, papers by Tripp,^{8,9} Levi-Setti,⁷ Samios,¹⁰ and Plane.¹¹ The most recent fits were made by Barbaro-Galtieri¹² and Samios.¹³

In fitting the data a choice for B_l has to be made. Plane¹¹ tried two forms for B_l :

(a) The form $B_l = (kr)^{2l} D_l(kr)$, r being the radius of interaction and D_l the polynomials in kr given by Blatt and Weisskopf.¹⁴ Usually r is taken to be 1 fermi.⁸

(b) The form $B_l = k^{2l}$.

However, for final results form (b) was chosen. A discussion of the differences among these two forms has been given by Barbaro-Galtieri.¹⁵ It turns out that not only the values of the couplings, g_i , depend upon the form used for B_f , but also the value obtained for the mixing angle. For the $3/2^-$ singlet, $\Lambda(1520)$, and the isospin-0 member of the octet, $\Lambda(1690)$, the mixing angles obtained in the two cases are

$$\theta_a = (-16.1^{+1.4}_{-1.3})^\circ, \theta_b = (-27.5^{+3.6}_{-3.4})^\circ,$$

in disagreement by a few standard deviations. It turns out that if a radius of interaction of $r = 0.15$ fermi is used for form (a), the two values of θ agree. This value of r does not fit resonance shapes when used in the Breit-Wigner resonant form.

Samios¹³ used form (b) for B_f .

Table I is a summary of the fits made by us (update of Barbaro-Galtieri¹²) using the barrier factor form (a) and exact SU(3). A few comments follow.

1/2⁻ Nonet (Baryon-Eta Resonances)

For this nonet Eq. (7) was multiplied by the factor

$$\left[\frac{M_R - M_B}{\bar{M}_R - \bar{M}_B} \right]^2,$$

where M_B is the decay baryon and $\bar{M}_R - \bar{M}_B = 564$ MeV is the difference of the mean $1/2^-$ and $1/2^+$ baryon octet masses. This kinematic factor comes from PCAC

arguments (i.e., the assumption that the axial vector current remains an octet in the presence of symmetry breaking) and it was advocated by Graham.¹⁶ For the $1/2^-$ nonet it was used in this form first by Gell-Mann.¹⁷

3/2⁺ Decuplet

The agreement among the coupling constants obtained for the four rates in this decuplet is very bad. The fit made using form (b) for B_f has $\chi^2=50$ for 3 degrees of freedom; the one made with form (a) for B_f has $\chi^2/DF=24/3$. The broken SU(3) relation (11), however, is very well satisfied.

B. SU(3) CLASSIFICATION OF MESON RESONANCES

All of the discussion above applies, except that for bosons the GMO formula is usually applied to the square of the masses, as opposed to the first power for fermions. Thus for example, Eq. (2) becomes

$$4\hat{K} = 3\hat{\eta} + \hat{\pi}. \tag{2'}$$

The symbol \hat{K} was introduced by Glashow and Socolow¹⁸ for the square of the K mass, etc.

Because of the difference between Eqs. (5') and (5''), there is also an extra factor of (M_N/M_R) in Eqs. (6) and (7). The three established nonets ($0^-, 1^-, 2^+$) and their mixing angles are listed at the bottom of the Meson Table.

Table I. SU(3) baryon multiplets with two or more known members. Values of θ and α [defined by Eqs. (8) and (10)] are the result of fits made to all the measured two-body decay rates of each multiplet.

J^P	Octet members ^a				Singlet	$\theta(\text{deg})^b$	α
$1/2^-$	N(1535)	$\Lambda(1670)$	$\Sigma(1750)$	$[\Xi(1825)]$	$\Lambda(1405)$	3 ± 5	$0.93 \pm .11$
$3/2^-$	N(1520)	$\Lambda(1690)$	$\Sigma(1670)$	$[\Xi(1815)]$	$\Lambda(1520)$	-23 ± 4	$0.31 \pm .05$
$5/2^-$	N(1670)	$\Lambda(1830)$	$\Sigma(1765)$				$1.17 \pm .04$
$5/2^+$	N(1688)	$\Lambda(1815)$	$\Sigma(1915)$				$0.65 \pm .03$
Decuplet members ^c					g_{10}		
$3/2^+$	$\Delta(1232)$	$\Sigma(1385)$	$\Xi(1530)$	Ω^-	$1.78-2.29$	$\chi^2/DF=50/3$	
$7/2^+$	$\Delta(1950)$	$\Sigma(2030)$					

^aMasses in parentheses are the nominal masses used in the Baryon Table. The Ξ members have masses as calculated by using formulae (1) and (2) with the mixing angle θ derived from the decay widths.

^bSee text for a discussion of the $3/2^-$ mixing angle.

^cCoupling constants from Ref. 12.

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Appendix III

TEST OF $\Delta I=1/2$ RULE FOR HYPERON DECAYS

O. E. Overseth
University of Michigan

1. *Nonleptonic decay Amplitudes*

In this edition we again use the new convention for the amplitudes A and B adopted in 1973. Some theorists have suggested that dimensionless amplitudes are more useful to them than the ones appearing in the literature. Berge¹ used a convention with A and B in units of $\text{sec}^{-1/2}$. Samios² used a convention which gave A and B in units of $(\text{MeV}\cdot\text{sec})^{-1/2}$. Following is the convention suggested by Jackson³, which gives dimensionless A and B.

The effective Lagrangian density for nonleptonic hyperon decays ($B_1 \rightarrow B_2 + \pi$) can be written

$$L_{\text{eff}} = G\mu_c^2 [\bar{\psi}_2(A+B\gamma_5)\psi_1]\phi_\pi,$$

where $G=10^{-5}m_p^{-2}$ is a coupling constant characteristic of first-order weak decays, μ_c is the charged pion mass, and A and B are dimensionless complex numbers giving the relative amplitudes of the parity-violating and parity-conserving decays, respectively. The matrix γ_5 is to be taken in the Pauli form, $\gamma_5 = \begin{pmatrix} 0 & -I \\ -I & 0 \end{pmatrix}$. The invariant amplitude for the decay is

$$M = G\mu_c^2 [\bar{u}(p)(A+B\gamma_5)u(P)],$$

where P is the 4-momentum of the decaying hyperon of mass M, and p is the 4-momentum of the baryon decay product of mass m. With the normalization convention, $\bar{u}_i u_i = 2m_i$, the Pauli form of the matrix element in the rest frame of the decaying hyperon is

$$M = G\mu_c^2 (\chi_2 | \sqrt{2M(E+m)}A + \sqrt{2M(E-m)}B\vec{\sigma}\cdot\hat{q} | \chi_1),$$

where E is the total energy of the final baryon and \hat{q} is a unit vector in the direction of motion of the final baryon. Comparison with Sec. VI D of the text shows that the amplitudes s and p defined there are proportional to A and B:

$$\frac{p}{s} = \left(\frac{E-m}{E+m} \right)^{1/2} \frac{B}{A} = \left[\frac{(M-m)^2 - \mu^2}{(M+m)^2 - \mu^2} \right]^{1/2} \frac{B}{A}.$$

Here μ is the mass of the pion entering the decay. The parameters α , β , and γ can therefore be expressed in terms of A and B, rather than s and p, if desired.

The decay rate for $B_1 \rightarrow B_2 + \pi$ is

$$\Gamma = \frac{G^2 \mu_c^4}{8\pi q} \left\{ \left[\frac{(M+m)^2 - \mu^2}{M^2} \right] |A|^2 + \left[\frac{(M-m)^2 - \mu^2}{M^2} \right] |B|^2 \right\},$$

where q is the c.m. momentum of the decay products. For reference, the dimensionless constant in this expression has the value $(G^2 \mu_c^4 / 8\pi) = 1.9488 \times 10^{-15}$.

Table I summarizes the amplitudes A and B for the nonleptonic decays of the Λ , Σ , and Ξ hyperons. These amplitudes have been calculated by using the experimental data for mean lives, branching ratios, and the decay asymmetry α given in the Stable Particle Table of this Review. Time-reversal invariance is assumed and final-state interactions are neglected, so A and B are taken to be relatively real. The subscript on the hyperon refers to the sign of the decaying pion. The statistical correlation coefficient

$$C_{AB} = \frac{\langle \Delta A \Delta B \rangle}{\sqrt{\langle \Delta A^2 \rangle \langle \Delta B^2 \rangle}}$$

is also given. The absolute signs of A and B have been assigned, using the following convention. Taking $A(\Lambda^0)$ as positive, the other S -wave decay amplitudes are chosen to give an approximate fit to the triangular relationships

$$\sqrt{2}A(\Sigma_0^+) + A(\Sigma_+^+) = A(\Sigma^-) \text{ and } \sqrt{3}A(\Sigma_0^+) + A(\Lambda^0) = 2A(\Xi^-).$$

The signs of the B amplitudes relative to those of the corresponding A amplitudes are determined by the sign of the appropriate α decay parameter.

Table I

$M \rightarrow m + \mu$	A	B	C_{AB}
$\Lambda^0 \rightarrow p + \pi^-$	1.48 ± 0.01	10.17 ± 0.24	-0.272
$\Lambda_0^0 \rightarrow n + \pi^0$	-1.08 ± 0.02	-7.28 ± 0.59	-0.747
$\Sigma_+^+ \rightarrow n + \pi^+$	0.06 ± 0.02	19.06 ± 0.16	0.003
$\Sigma_0^+ \rightarrow p + \pi^0$	1.48 ± 0.05	-12.04 ± 0.59	0.918
$\Sigma^- \rightarrow n + \pi^-$	1.93 ± 0.01	-0.65 ± 0.08	-0.024
$\Xi_0^0 \rightarrow \Lambda + \pi^0$	1.53 ± 0.03	-5.90 ± 1.11	0.347
$\Xi^- \rightarrow \Lambda + \pi^-$	2.04 ± 0.02	-6.71 ± 0.38	0.198

2. Tests of the $\Delta I=1/2$ Rule

(a) Λ Decay

For Λ decay the $\Delta I=1/2$ rule predicts that $\Gamma_0/\Gamma_- = 0.50$ and $\alpha_0 = \alpha_-$. In order to determine the magnitude of possible $\Delta I=3/2$ amplitudes present we write the linear expressions⁴ for the $\Delta I=3/2$ A- and B-wave amplitudes in terms of $\Delta\alpha$, where $\Delta\alpha$ is the measured value of α_0/α_- minus the predicted value, and in terms of $\Delta\Gamma$ similarly defined. Evaluating these we find

$$\Delta\alpha = -1.53 (A_3/A_1) + 1.60 (B_3/B_1),$$

$$\Delta\Gamma = 1.83 (A_3/A_1) + 0.26 (B_3/B_1).$$

Here the $\Delta I=3/2$ amplitudes are expressed relative to the $\Delta I=1/2$ amplitudes. The numerical values of the coefficients depend on the ratio B/A . The uncertainties in the coefficients are small compared to the uncertainties in $\Delta\alpha$ and $\Delta\Gamma$. Final-state πN interactions have been included in these relations but have a very small effect. From the Stable Particle Table,

$$\Delta\alpha = 0.006 \pm 0.066, \quad \Delta\Gamma = 0.058 \pm 0.012,$$

and hence

$$(A_3/A_1) = 0.027 \pm 0.008$$

and

$$(B_3/B_1) = 0.030 \pm 0.037.$$

The possible 3% $\Delta I=3/2$ A-wave amplitude is due to the disagreement of decay rates with prediction. At this level the results are sensitive to electromagnetic corrections. However, in Λ decay the phase space correction and the other radiative corrections appear to be about equal in magnitude and have opposite signs,^{5,6} and hence cancel each other in the correction to the decay rates.

(b) Ξ Decay

The analysis for Ξ decay is very similar to that for Λ decay. If the $\Delta I=1/2$ rule is valid, $\Gamma_0(\Xi^0)/\Gamma_-(\Xi^-) = 0.50$ and $\alpha_0 = \alpha_-$. For this case the expressions linear in $\Delta I=3/2$ A- and B-wave amplitudes are⁴

$$\Delta\alpha = 1.38 (A_3/A_1) - 1.38 (B_3/B_1),$$

$$\Delta\Gamma = -1.44 (A_3/A_1) - 0.06 (B_3/B_1).$$

From the Stable Particle Table,

$$\Delta\alpha = 0.12 \pm 0.21, \quad \Delta\Gamma = 0.058 \pm 0.024,$$

and we find

$$(A_3/A_1) = -0.035 \pm 0.017$$

and

$$(B_3/B_1) = -0.13 \pm 0.15.$$

(c) Σ Decay

The traditional test of the $\Delta I=1/2$ rule in Σ decay is that the amplitudes satisfy the relationship

$$\sqrt{2} \Sigma_0^+ + \Sigma_+^+ - \Sigma^- = 0.$$

Graphically this is equivalent to closing the Σ triangle when the amplitudes are plotted on A, B axes. Including $\Delta I \geq 3/2$ amplitudes in Σ decay analysis, the " Σ triangle" relationship becomes

$$\sqrt{2} A_0 + A_+ - A_- = -3\sqrt{2/5} A_3 + \frac{2}{\sqrt{15}} A_5,$$

where A_3 and A_5 are $\Delta I=3/2$ and $\Delta I=5/2$ amplitudes, respectively. There is a similar equation for the B amplitudes. From Table I,

$$\sqrt{2} A_0 + A_+ - A_- = 0.22 \pm 0.09$$

and

$$\sqrt{2}B_0 + B_+ - B_- = 2.7 \pm 1.1 .$$

If we neglect the $\Delta I=5/2$ amplitudes and assume all amplitudes to be real we can solve for possible $\Delta I=3/2$ amplitudes. The result is

$$\frac{A_3}{A_-} = -0.060 \pm 0.026$$

and

$$\frac{B_3}{B_+} = -0.074 \pm 0.030 .$$

Thus for hyperon decay, present experimental data limit $\Delta I=3/2$ amplitudes to less than about 5%.

3. The Lee-Sugawara Relation

From Table I the Lee-Sugawara relation,^{7,8} $\sqrt{3}\Sigma_0^+ + \Lambda^0 - 2\Sigma^- = 0$, is satisfied to -0.05 ± 0.12 for the A amplitudes, and to 2.7 ± 2.0 for the B amplitudes.

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$$\frac{\Gamma_0}{\Gamma_-} \approx \frac{1}{2} \left\{ 1 + 3\sqrt{2} \times \left[\frac{S_{11}S_{33}\cos(\delta_1 - \delta_3) + P_{11}P_{33}\cos(\delta_{11} - \delta_{31})}{S_{11}^2 + P_{11}^2} \right] \right\} .$$

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Appendix IV

GROWTH OF INFORMATION

From time to time we have presented figures demonstrating the amount of experimental work which has gone into spectroscopy, and the amount of new information available as a result. The 1976 versions of these figures are shown as Figs. 1, 2, and 3.

Figure 1 is a simple count of the number of meson resonances listed in the Tables, categorized as those "understood" -- i.e., all quantum numbers are believed known -- and those simply "listed". For the 1976 edition, there is an increase in both of these categories because of the discovery of the J/ψ and related particles.

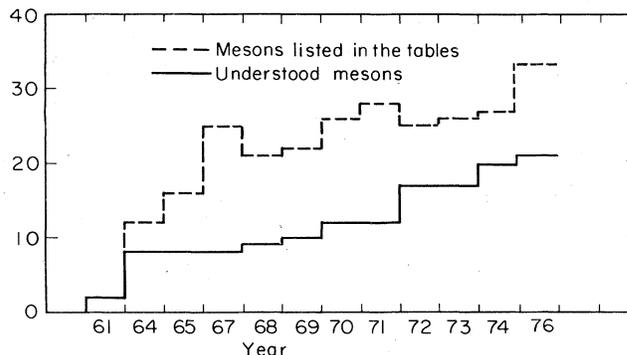


Fig. 1. Number of meson resonances listed in the Tables (dashed line) and those for which all quantum numbers are known (solid line), as a function of year of publication of the Review of Particle Properties. Note abscissa omits years in which no publication occurred.

In Figure 2 we present similar information for the baryon resonances, but concentrating here on the "growth of understanding". That is, the number of known baryons (we include for this figure only those with known J^P) has grown only very slowly with time (dashed line); the real progress has been in the measurement of the properties of those baryons. Therefore we show as the solid curve a count of the number of baryonic properties -- mass, width, and branching ratios. Most of these results are from partial-wave analyses.

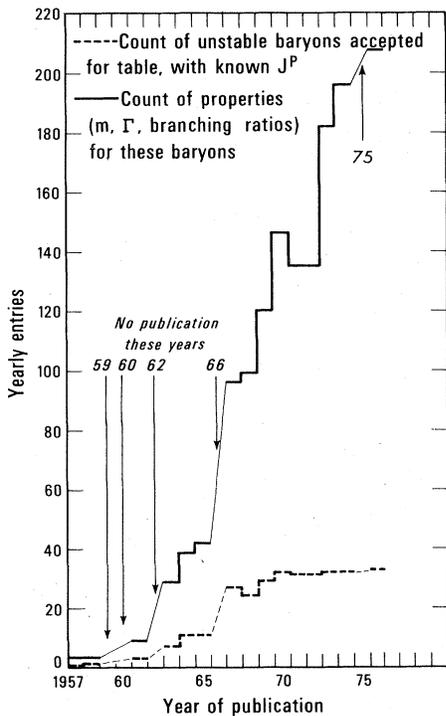


Fig. 2. Total amount of information (mass + width + branching ratios) on baryon resonances listed in the Tables, restricted to those with well-established J^P (solid line). Dashed line shows numbers of such resonances listed. Abscissa shows year of publication of Review of Particle Properties.

Finally, in Figure 3 we show a count of the number of new results put in the Listings each year, shown according to the type of detector that was used in obtaining the result, and according to the type of particle. N^* and Z^* particles have been omitted from this figure, because for these we use mainly the results of partial-wave analyses, rather than the primary data.

The fall in productivity in recent years doubtless reflects the tight budget situation and the declining emphasis on spectroscopy. Intense activity on the J/ψ and related particles is evident in the curve for meson resonances.

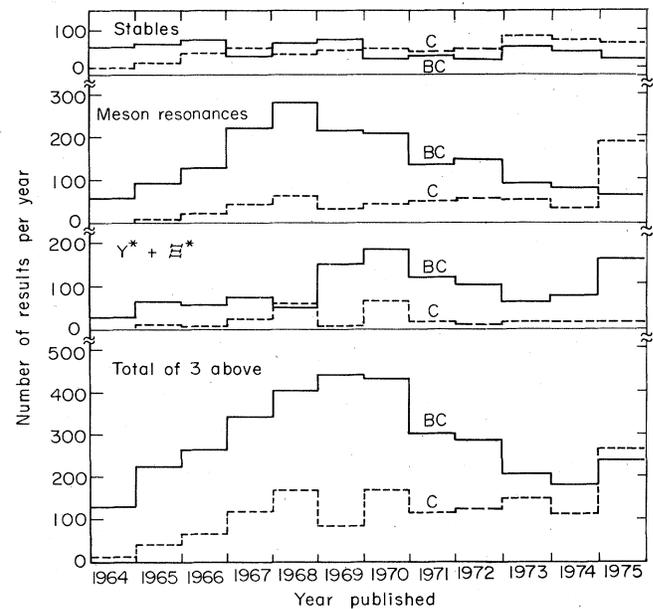


Fig. 3. The rate of production of data on particle properties, as a function of year of publication of the original result, based on data cards added to the Listings. Solid line shows bubble chamber results, dashed line counter or other electronic systems results.

