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Effect of a Form Factor on dE/dx from Close Collisions J. D. Jackson

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Point 1:

Close collisions and distant collisions contribute approximately equally to dE/dx. A form factor is possibly relevant only for the close collision half. We consider these close collisions in the "free electron" approximation.

Point 2:

We parameterize the pion-electron interaction with a form factor showing a ρ -meson propagator:

$$F(Q^2) = \frac{m_{\rho}^2}{m_{\rho}^2 + Q^2}$$
, where $m_{\rho} = 0.77$ GeV

Calculation:

$$s = m_{\pi}^2 + m_e^2 + 2m_e E_{\pi} \text{ (lab)}$$
$$Q_{\max}^2 = s - 2 \left(m_{\pi}^2 + m_e^2\right) + \frac{\left(m_{\pi}^2 - m_e^2\right)^2}{s}$$
$$T_{\max} = Q_{\max}^2 / 2m_e \quad (\text{maximum kinetic energy transfer})$$

The form factor multiplies the cross section given by Rossi [1], p. 16, Eq. 2.3.6, which in the notation we are using then reads

$$\frac{d\sigma}{dT} = \frac{A}{T^2} \left(1-\beta^2 \frac{T}{T_{\rm max}}\right) \left|F\left(2m_eT\right)\right|^2 \label{eq:dstar}$$

Upon multiplication by T and integration over T, we obtain the equivalent of Rossi's Eq. 2.5.4:

$$\frac{dE}{dx}(\text{close}) = A \int_{T_{\min}}^{T_{\max}} \frac{dT}{T} \left(1 - \beta^2 \frac{T}{T_{\max}} \right) \left| F\left(2m_e T\right) \right|^2 \equiv A \left(\mathcal{I}_1 - \mathcal{I}_2 \right)$$

where T_{\min} is the dividing energy between close and distant collisions (called η by Rossi) Following Rossi, I put $T_{\min} = 10^5 \text{ eV} = 10^{-4} \text{ GeV}$. Now

$$\mathcal{I}_1 = \int_{T_{\min}}^{T_{\max}} \frac{dT}{T} \ \frac{m_{\rho}^4}{\left(m_{\rho}^2 + 2m_e T\right)^2}$$

My CRC Tables yield

$$\mathcal{I}_{1} = \ln\left(\frac{T_{\max}}{T_{\min}}\right) - \ln\left(\frac{m_{\rho}^{2} + 2m_{e}T_{\max}}{m_{\rho}^{2} + 2m_{e}T_{\min}}\right) + \frac{m_{\rho}^{2}}{m_{\rho}^{2} + 2m_{e}T_{\max}} - \frac{m_{\rho}^{2}}{m_{\rho}^{2} + 2m_{e}T_{\min}} .$$

For $m_{\rho}^2 \to \infty$, $\mathcal{I}_1 \to \ln\left(\frac{T_{\max}}{T_{\min}}\right)$ [the usual result].

Similarly,

$$\begin{split} \mathcal{I}_{2} &= \frac{\beta^{2} m_{\rho}^{4}}{T_{\max}} \int_{T_{\min}}^{T_{\max}} \frac{dT}{\left(m_{\rho}^{2} + 2m_{e}T\right)^{2}} \\ &= \frac{\beta^{2}}{T_{\max}} \frac{m_{\rho}^{4}}{2m_{e}} \left[\frac{1}{m_{\rho}^{2} + 2m_{e}T_{\min}} - \frac{1}{m_{\rho}^{2} + 2m_{e}T_{\max}} \right] \\ &= \beta^{2} \left(\frac{m_{\rho}^{2}}{2m_{e}T_{\max}} \right) \left[\frac{1}{\left(1 + 2m_{e}T_{\min}/m_{\rho}^{2}\right)} - \frac{1}{\left(1 + 2m_{e}T_{\max}/m_{\rho}^{2}\right)} \right] \end{split}$$

In the limit $m_{\rho}^2 \to \infty$, $\mathcal{I}_2 \to \beta^2 (1 - T_{\min}/T_{\max})$,

$$\beta^2 = 1 - 1/\gamma^2 = 1 - m_\pi^2/E_\pi^2$$

Results:

I wrote a short program on my Macintosh (using THINK Pascal) to compute the various quantities:

$$Q_{\max}^2$$
, T_{\max} ,
 $\ln (T_{\max}/T_{\min}) - \beta^2 (1 - T_{\min}/T_{\max})$ = usual close term,
and $R = rac{(\mathcal{I}_1 - \mathcal{I}_2)}{[\text{usual close term}]}$

Everything is in GeV or GeV² or dimensionless. Inputs: $m_{\rho} = 0.77$, $m_e = 5.11 \times 10^{-4}$, $T_{\rm min}$ cutoff = 10⁻⁴/ GeV. Attached is a page of output,* showing the effect for pions. The quantity called "dE/dx ratio" is R, the ratio of dE/dx(close) with the form factor to dE/dx(close) with no form factor.

We see that for pions of 850 GeV there is a 6% reduction with the form factor, corresponding to roughly 3% reduction in dE/dx. The reduction is even smaller for protons.

Addendum:

It is useful to summarize the result as the usual close-collision term minus a correction. Using the approximations $2m_e T_{\min}/m_{\rho}^2 \ll 1$ and $T_{\min}/T_{\max} \ll 1$, and defining x = $2m_e T_{\rm max}/m_{\rho}^2$, I find

$$\frac{dE}{dx}(\text{close}) = A \left[\ln \left(T_{\text{max}}/T_{\text{min}} \right) - \beta^2 - f(x) \right]$$

where $f(x) = \ln (1+x) + \frac{1}{\gamma^2} \frac{x}{1+x}$

Reference:

1. B. Rossi, *High Energy Particles*, Prentice-Hall, Inc., Englewood Cliffs, NJ, 1952.

^{*} Not available in postscript for web listing. Request copy by mail from the PDG, LBNL.